

Early stage of RHIC collisions and its equilibration: a story told by correlations

BNL Nuclear Physics Seminar

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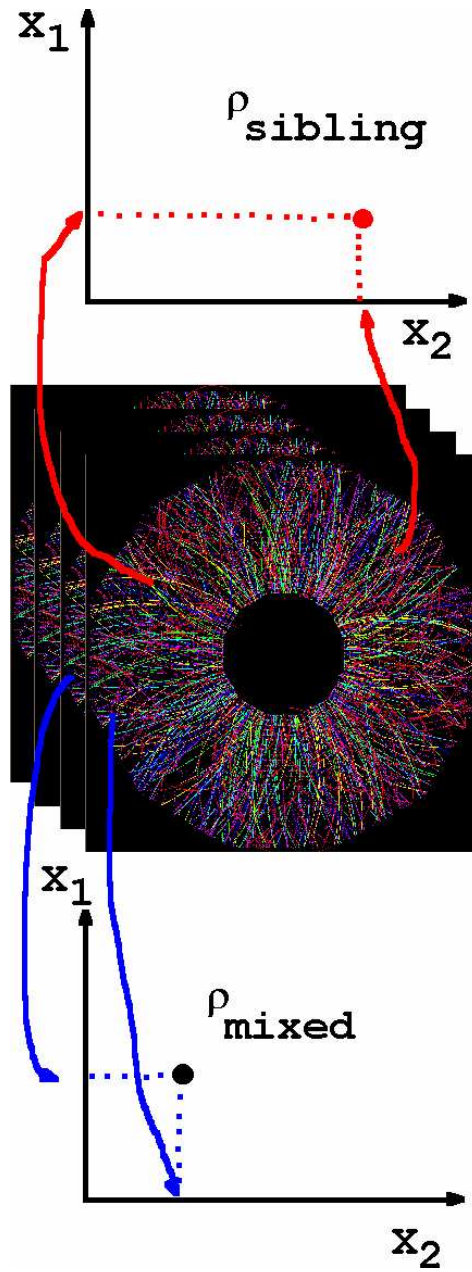
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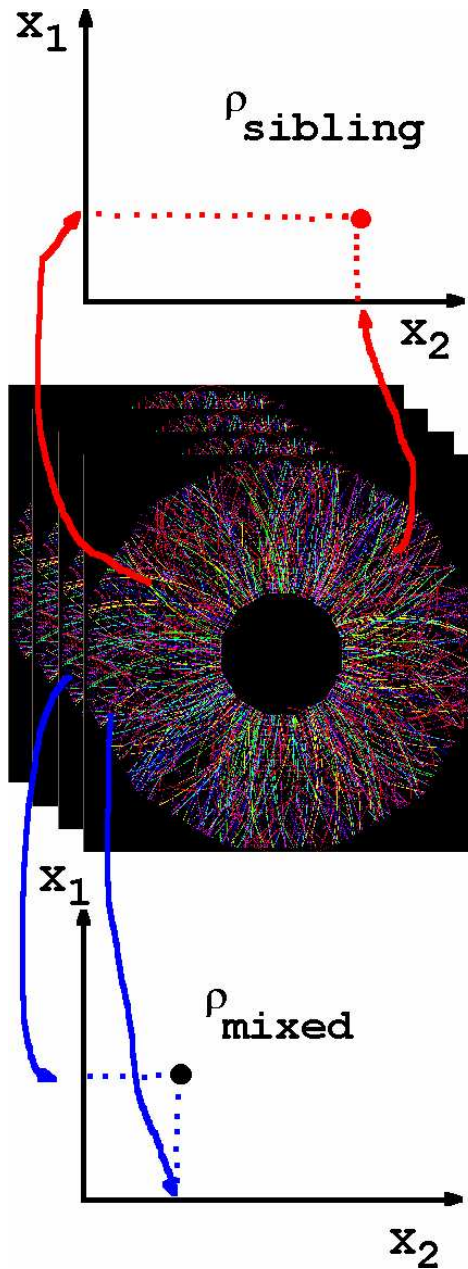
- Conclusions

2 Autocorrelation

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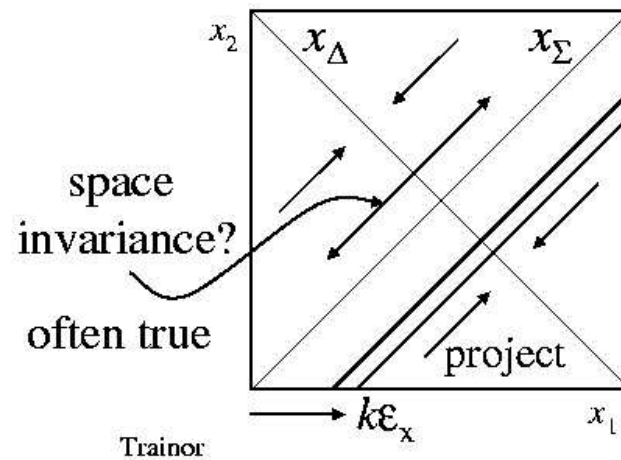


$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix} \rightarrow \begin{pmatrix} x_{\Sigma} \equiv x_1 + x_2 \\ x_{\Delta} \equiv x_1 - x_2 \end{pmatrix},$$

always a lossless transformation of data.

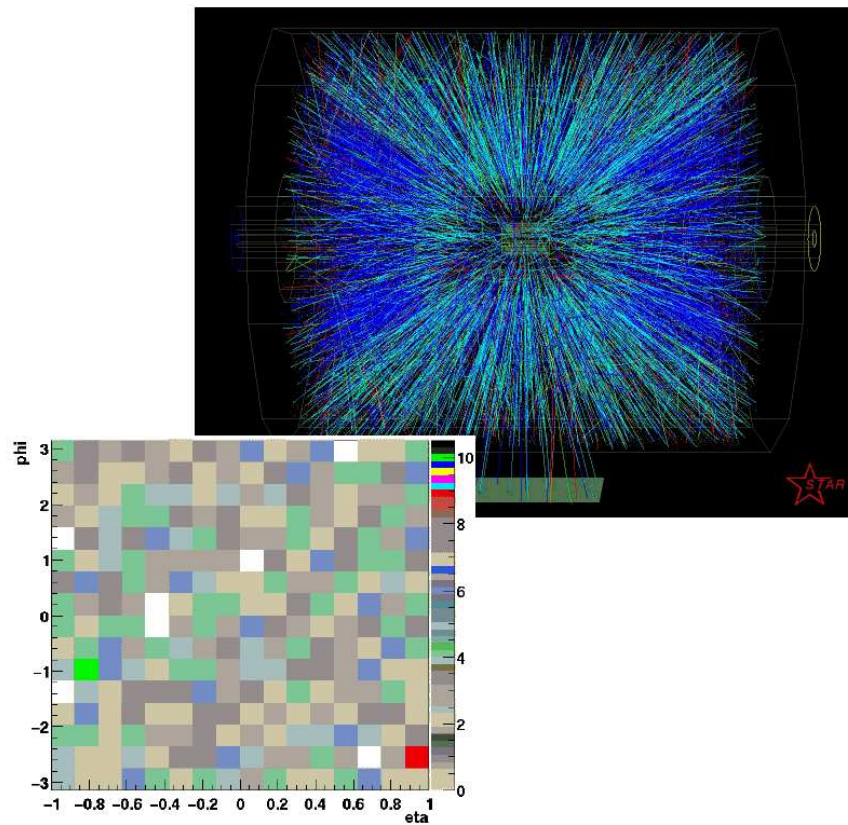
Autocorrelation A is a projection of a two-point distribution onto difference variable(s) x_{Δ} , lossless for x_{Σ} -invariant (homogenous, stationary) problems.

$$\Delta R(x_1, x_2) = \frac{\rho_{\text{sibling}}(x_1, x_2)}{\rho_{\text{mixed}}(x_1, x_2)} - 1$$



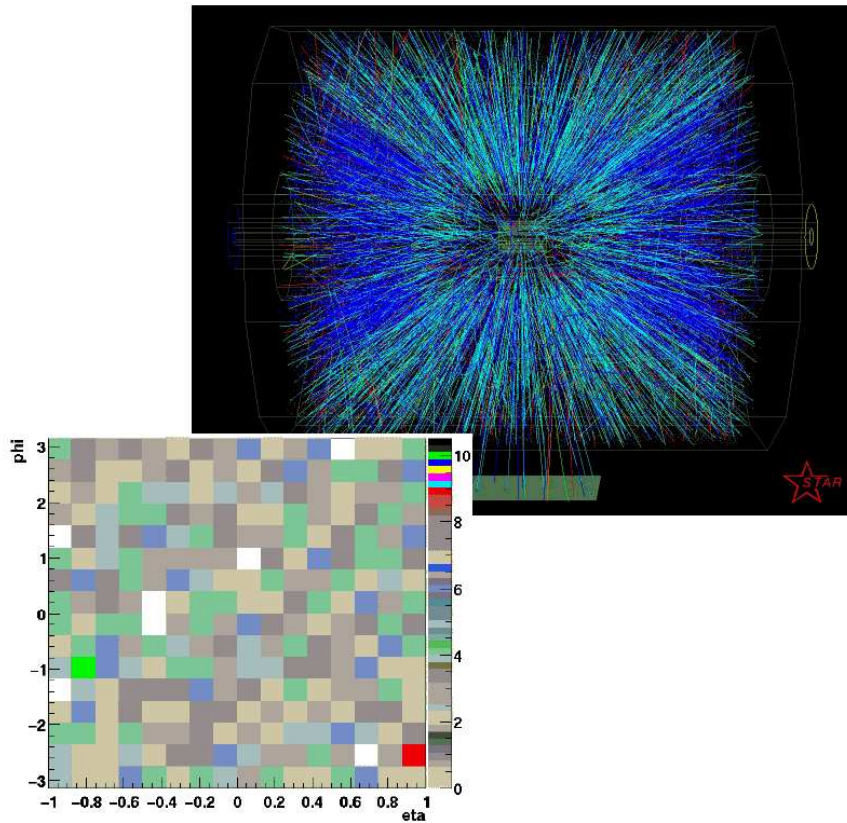
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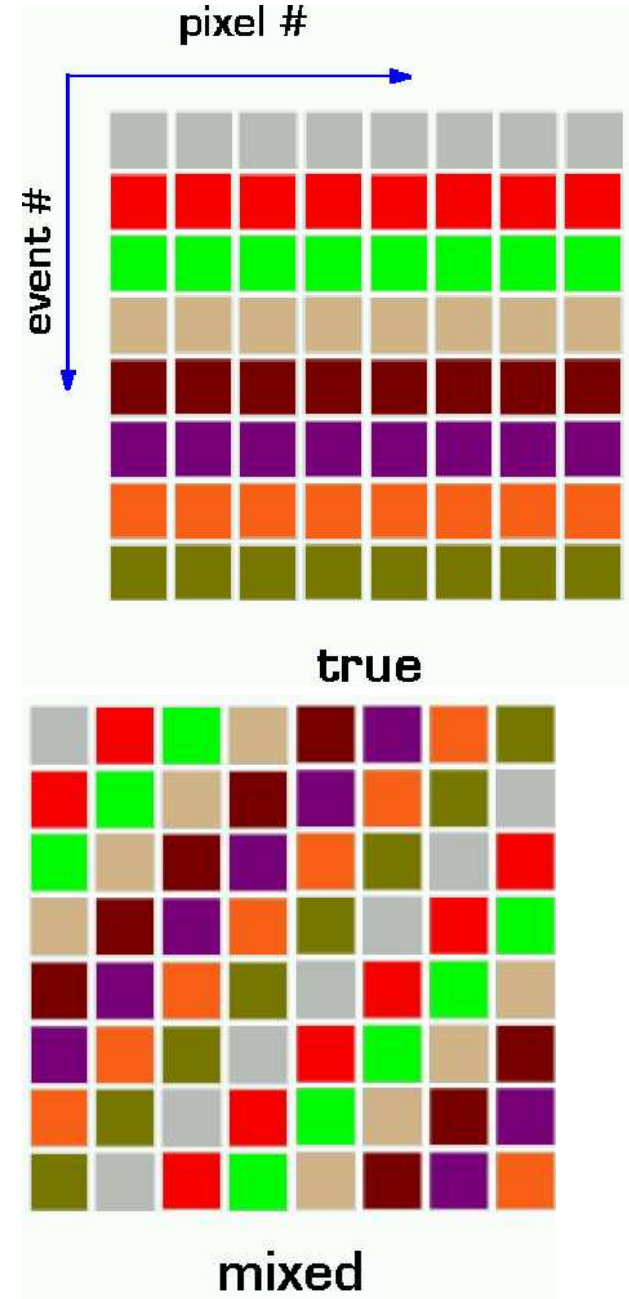


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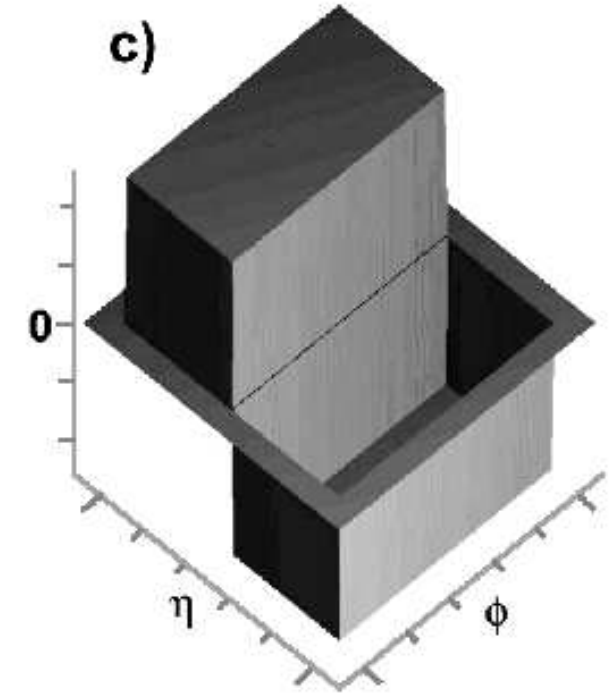
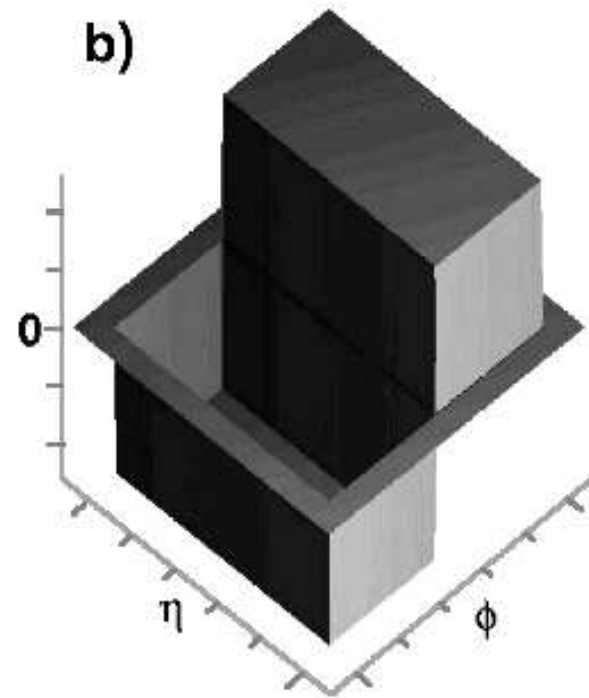
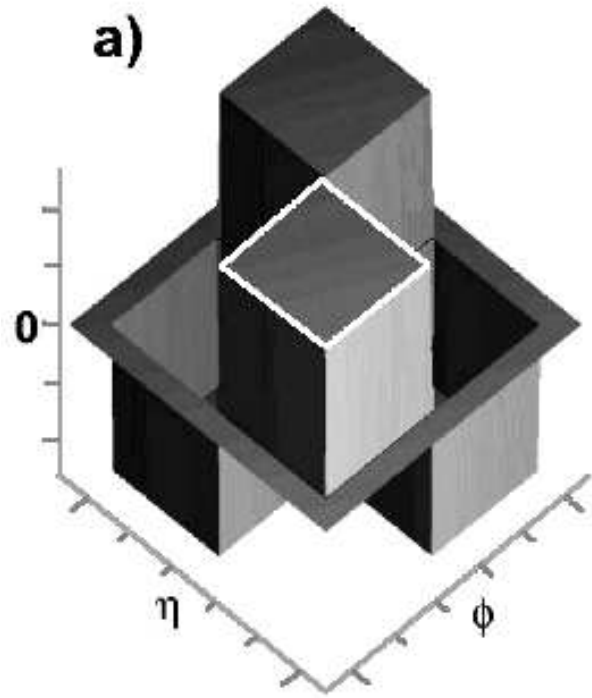


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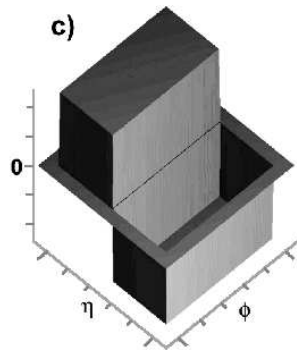
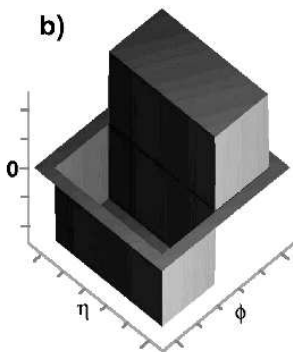
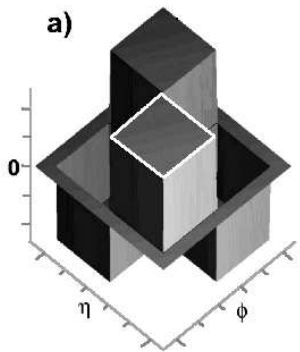
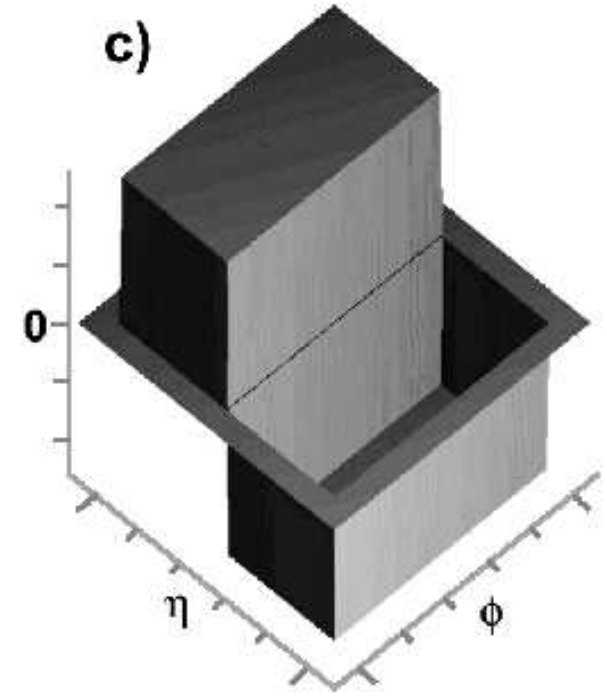
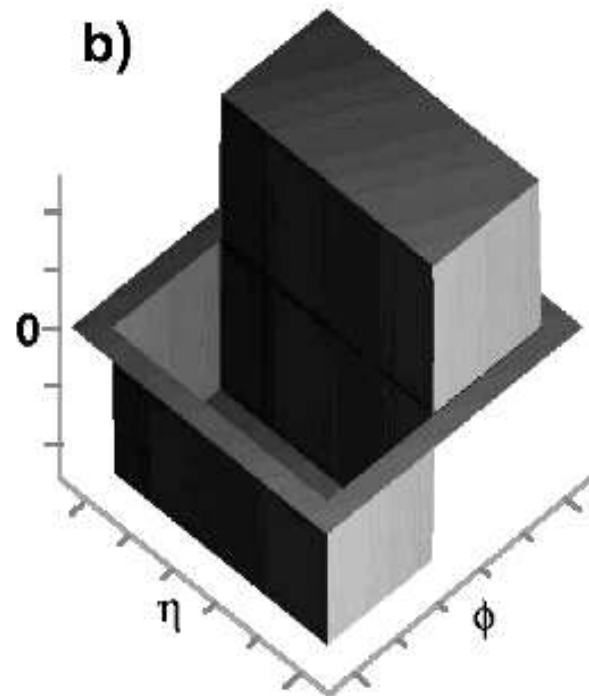
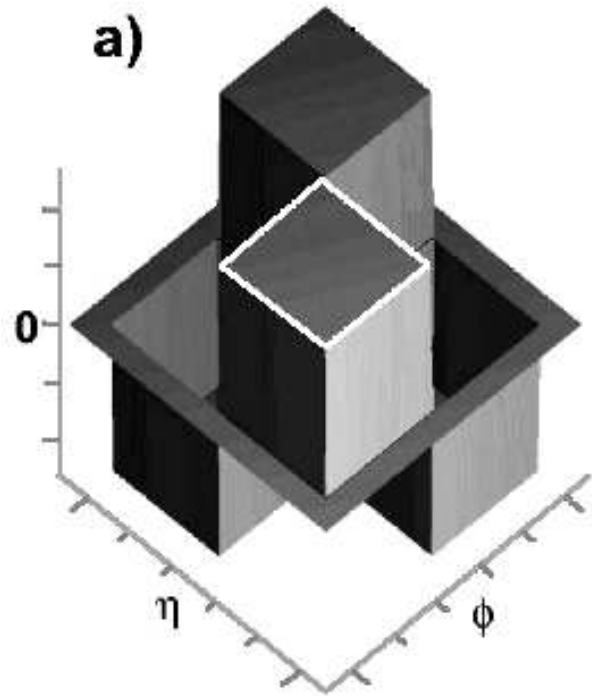


4 Discrete Wavelet Transform (DWT)

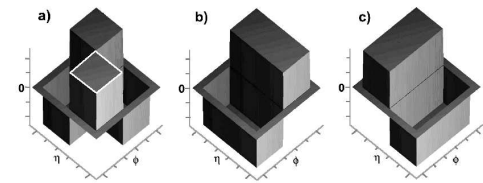
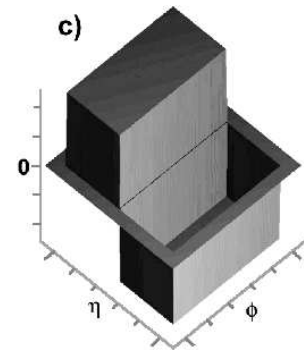
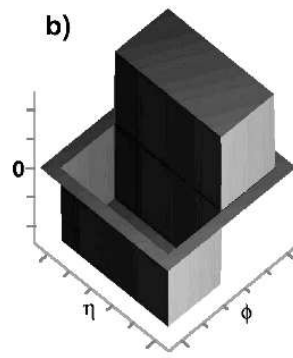
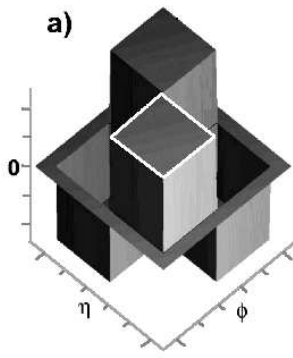
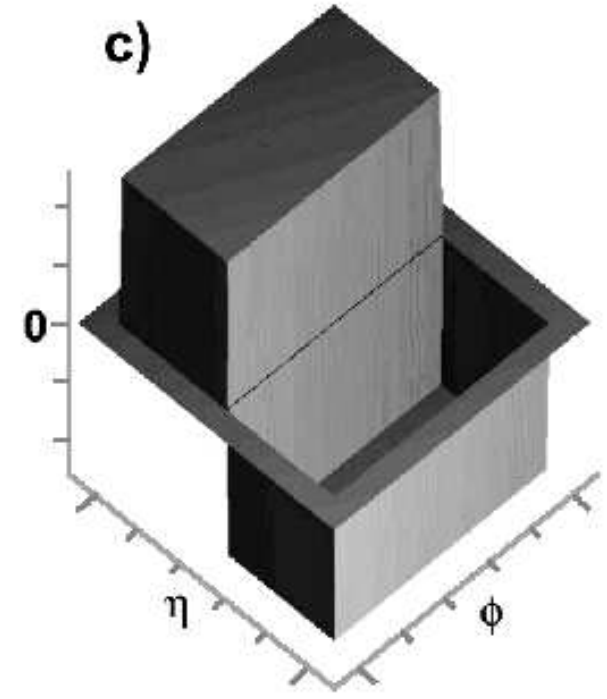
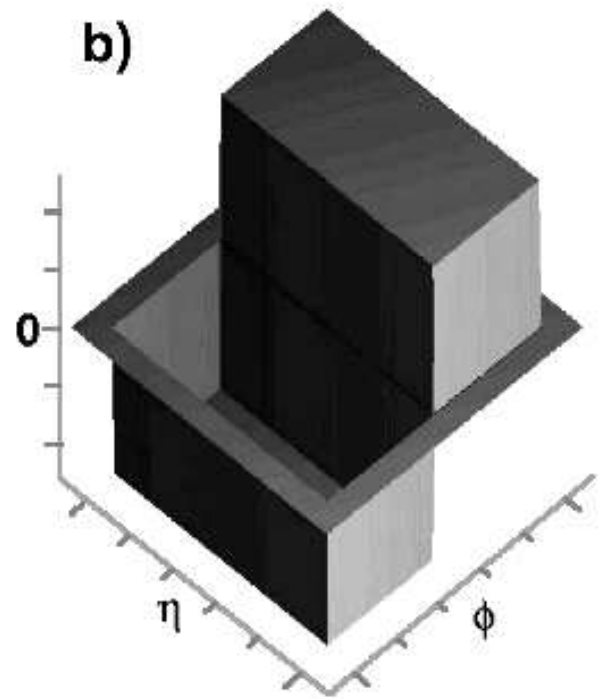
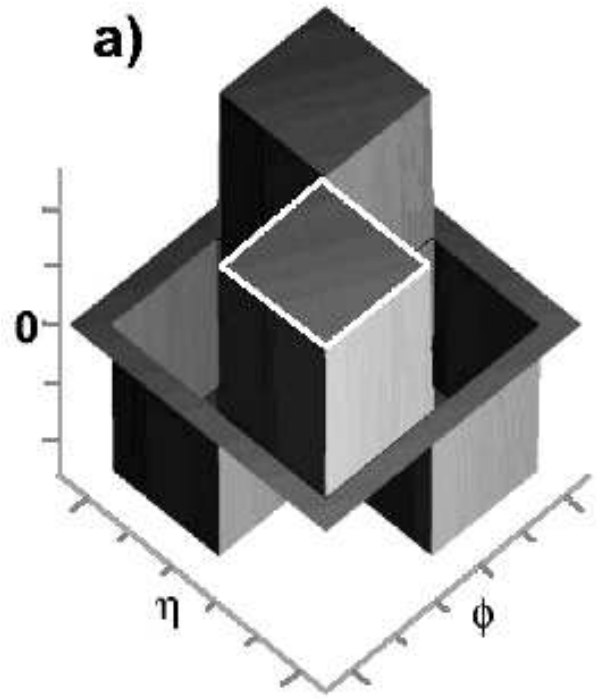
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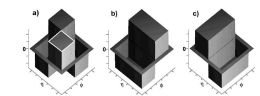
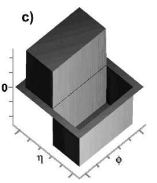
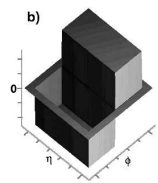
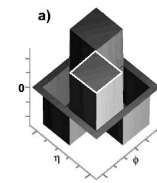
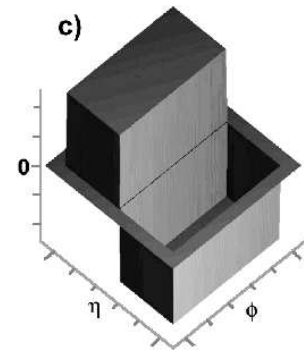
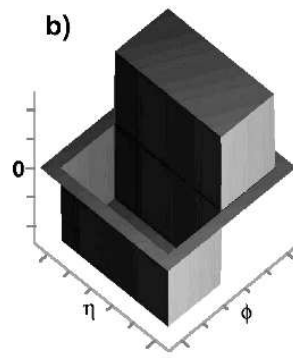
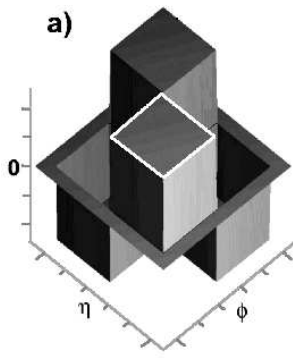
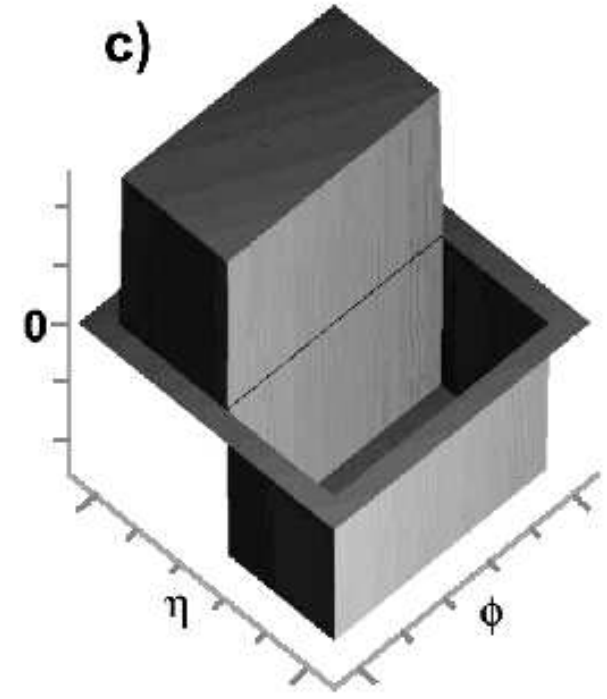
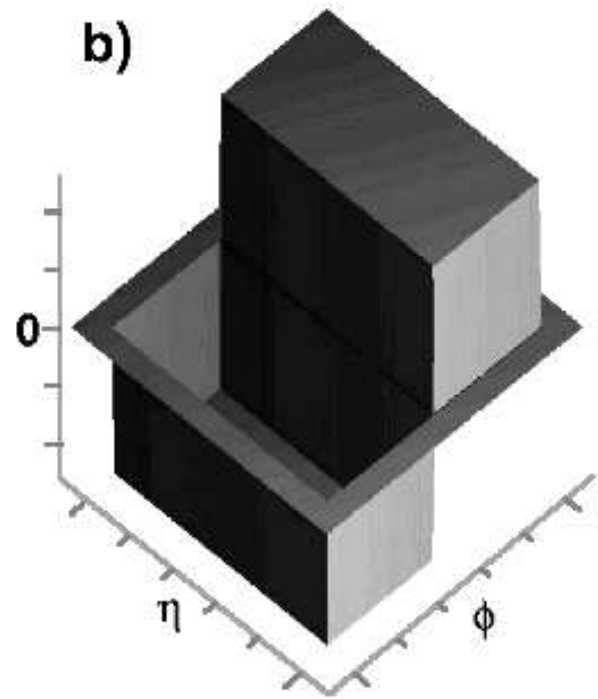
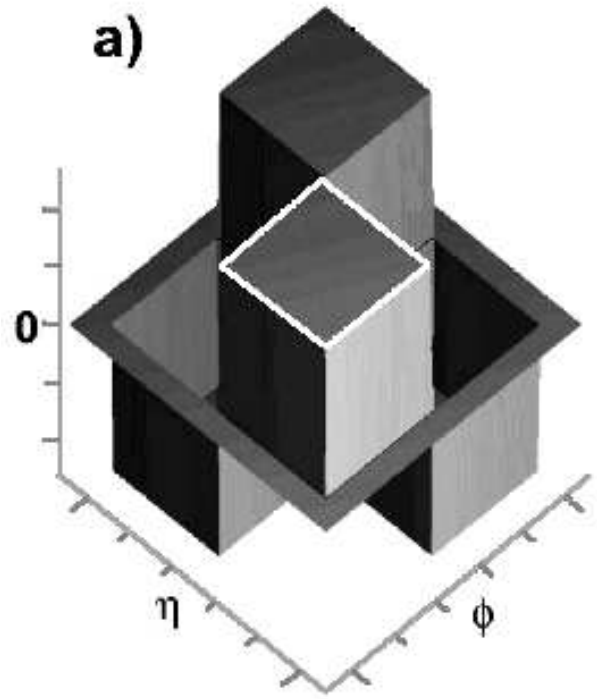
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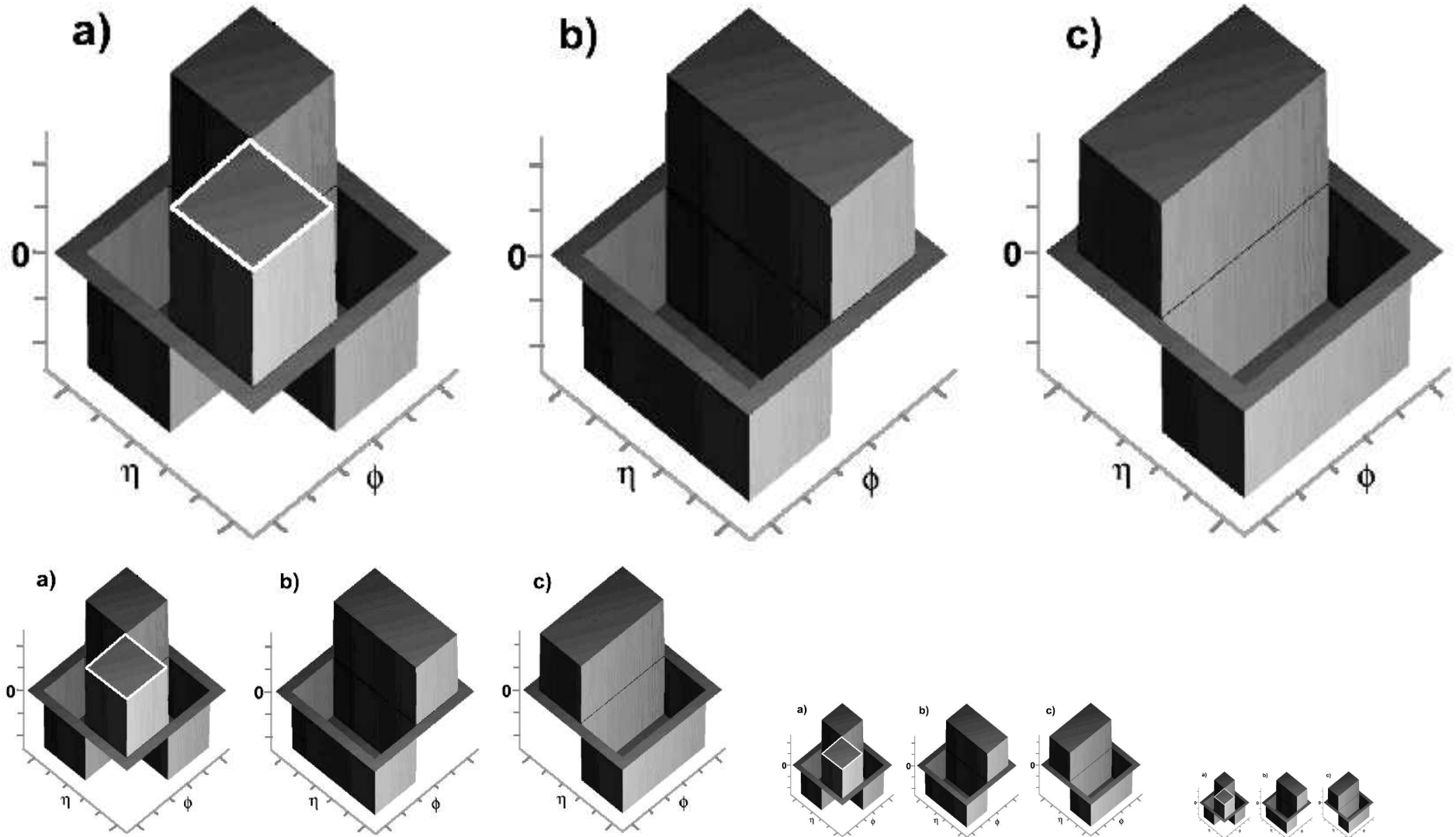
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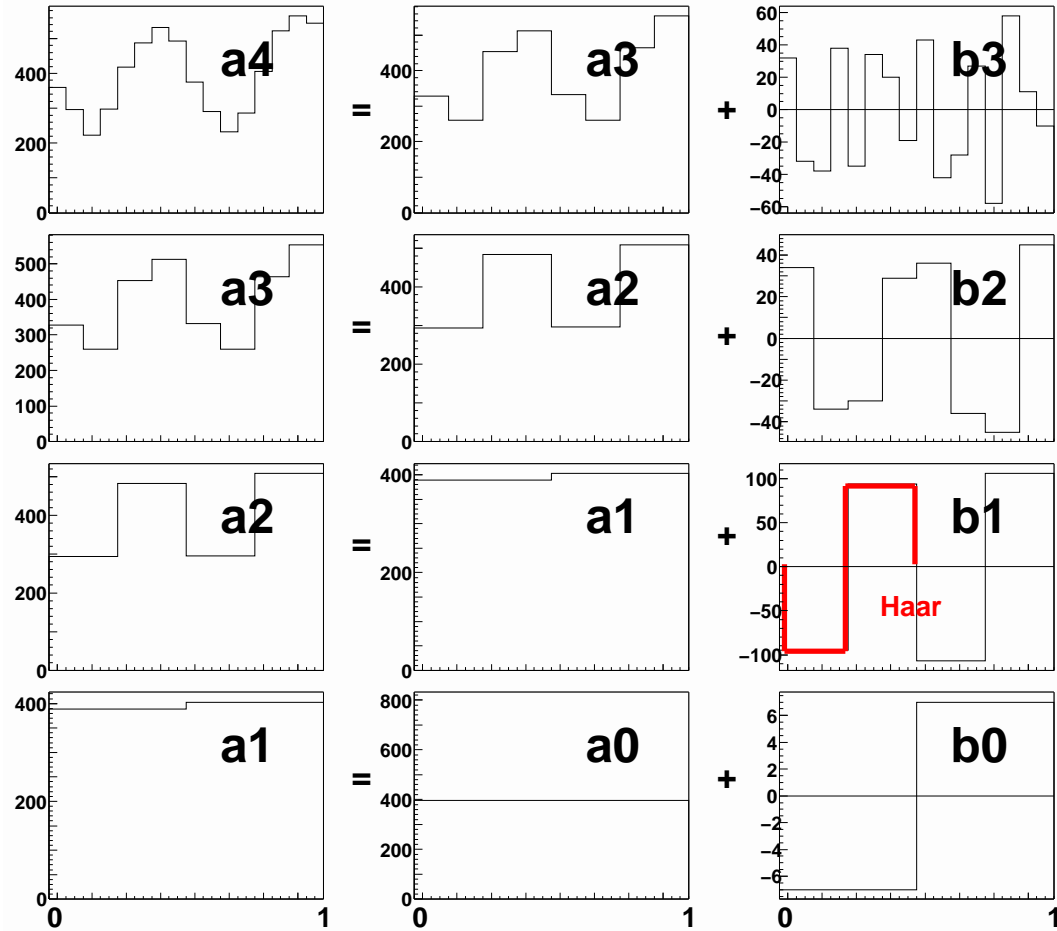
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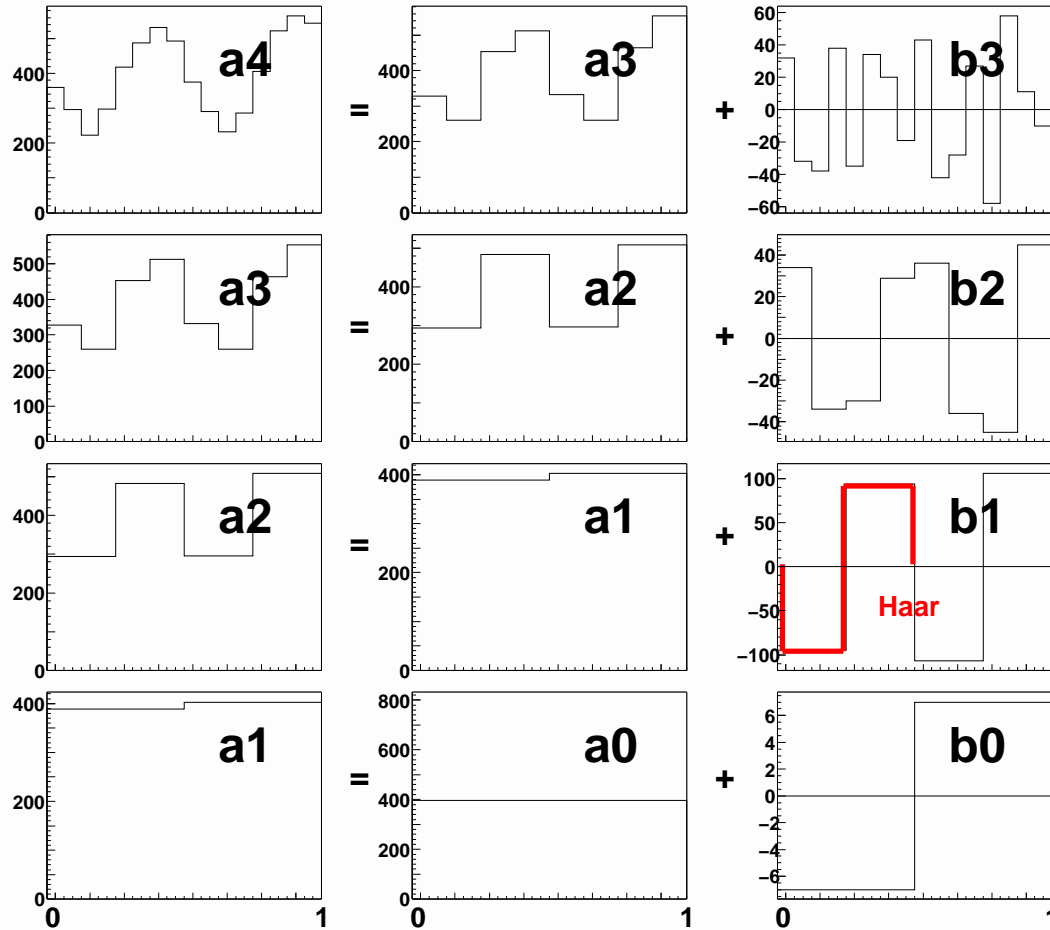
$F_{m,i,j}^\lambda(\phi, \eta)$ —Haar wavelet **orthonormal basis** in (ϕ, η) : scale fineness (m), directional modes of sensitivity (λ), track density $\rho(\eta, \phi, p_T)$, locations in 2D (i, j) . **DWT** is an expansion in this basis.

5 A flow-inspired example

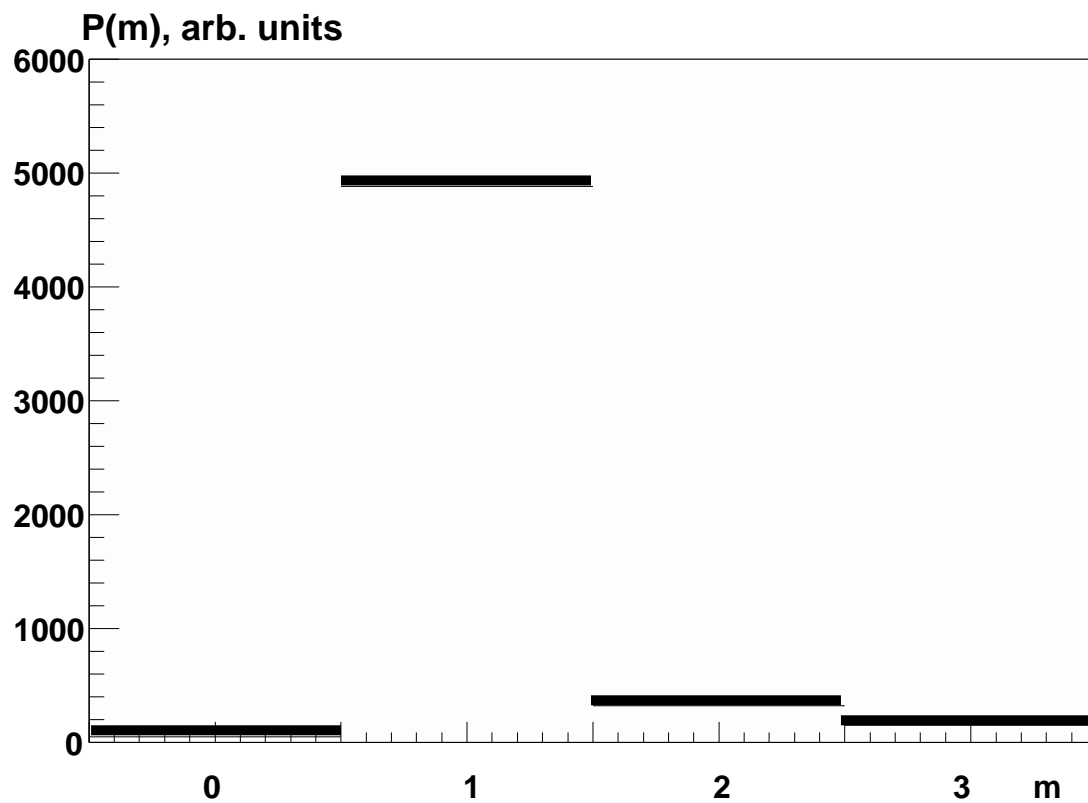
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Elliptic flow-inspired example:
 x axis – an angle in “natural units” ($2\pi = 1$), y axis – multiplicity. The multiresolution theorem: $a_4 = a_0 + b_0 + b_1 + b_2 + b_3$, can have better fineness.



Power spectrum of that flow event as a function of “fineness” m . The dominant contribution is $m = 1$ (the “ v_2 ” harmonic, **b1**). Statistical fluctuations also contribute.

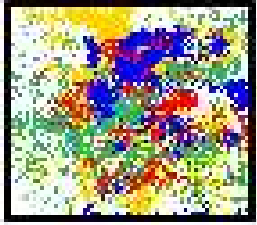
$$P(m) = 2^{-m} \sum_i \langle \rho, F_{m,i} \rangle^2.$$

Computational complexity $O(N)$!

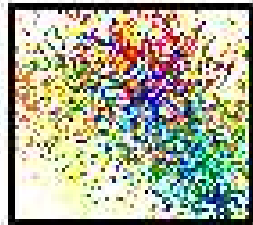
6 DWT Power Spectra and the Hilbert Space

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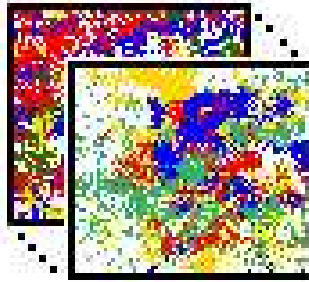
real



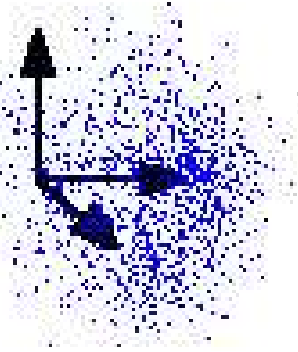
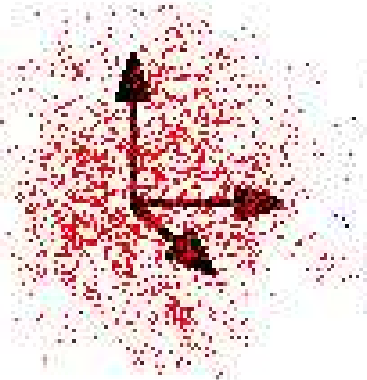
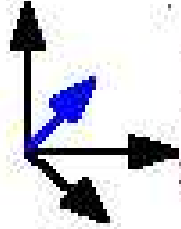
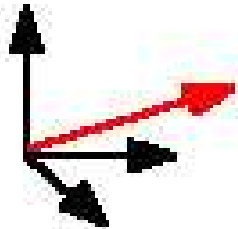
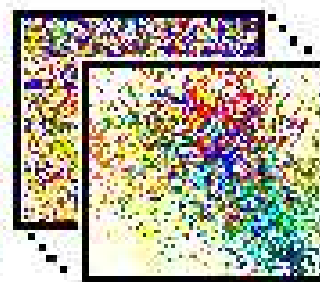
mixed



many real

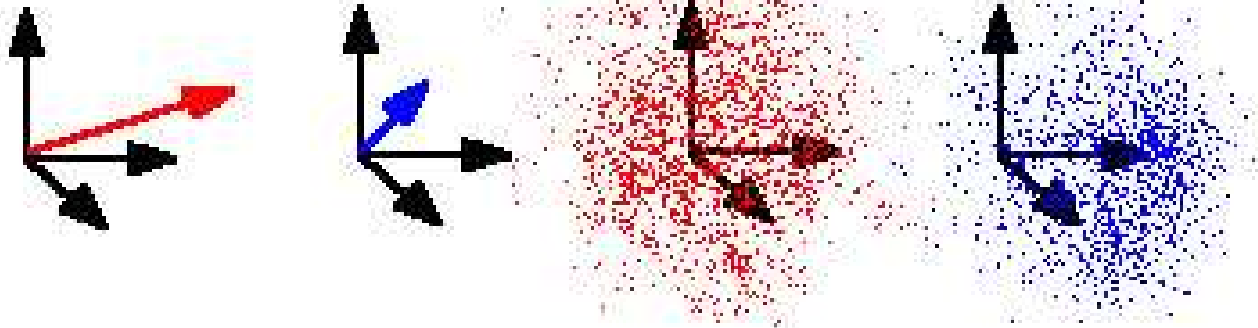
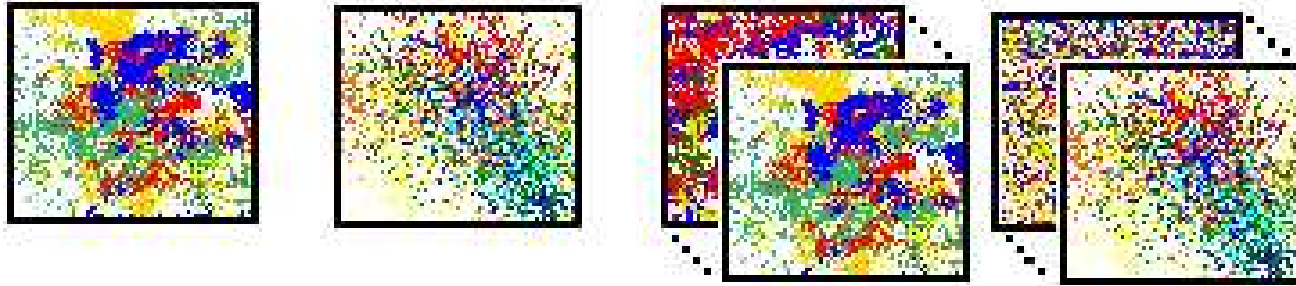


many mixed



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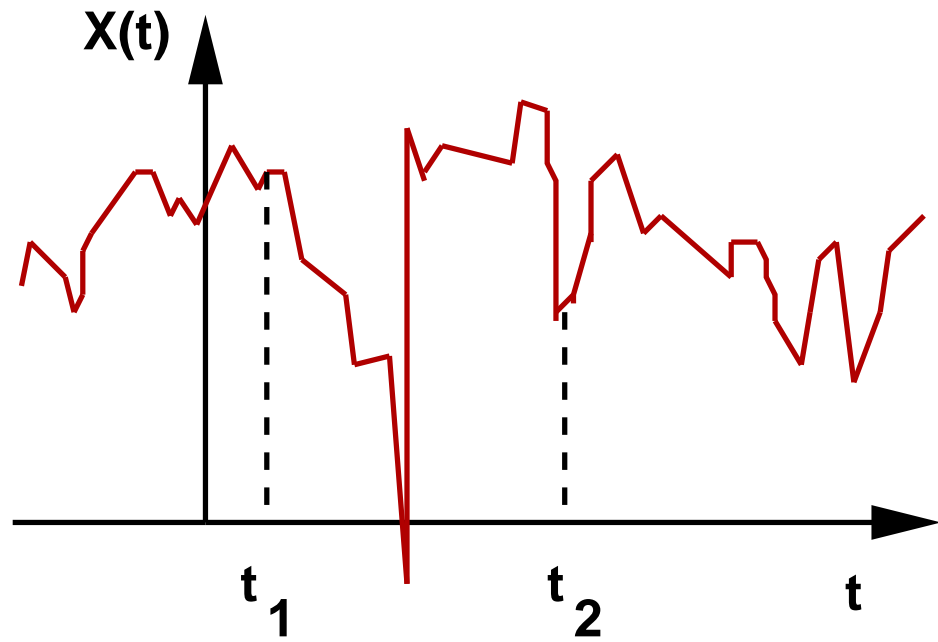
real mixed many real many mixed



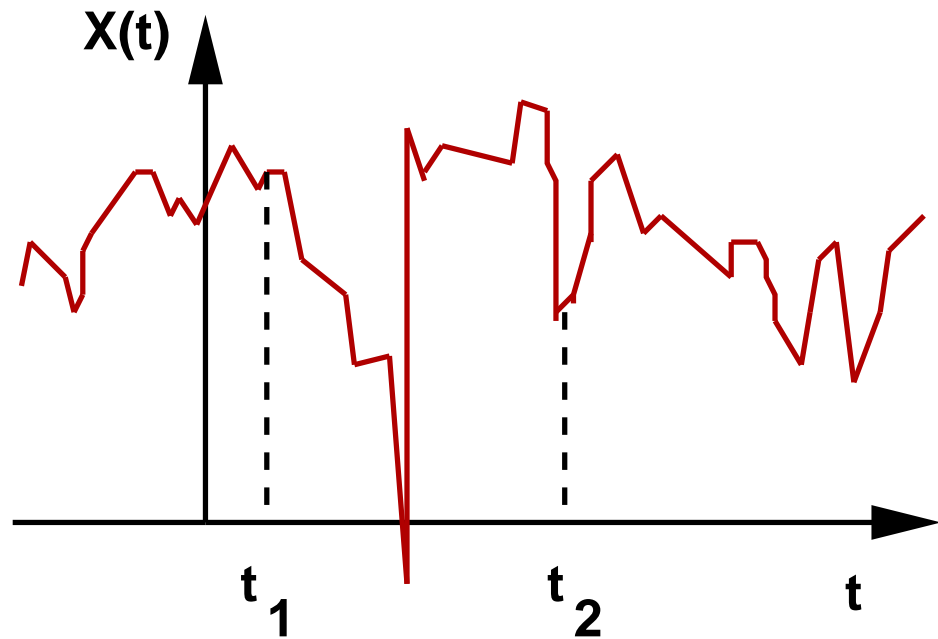
Power of local fluctuations, mode λ : $P^\lambda(m) = 2^{-2m} \sum_{i,j} \langle \rho, F_{m,i,j}^\lambda \rangle^2 \propto \text{norm}^2$
in the DWT subspace. “Dynamic texture” $P_{dyn}^\lambda(m) \equiv P_{true}^\lambda(m) - P_{mix}^\lambda(m)$.

7 DWT Power Spectra and Correlations

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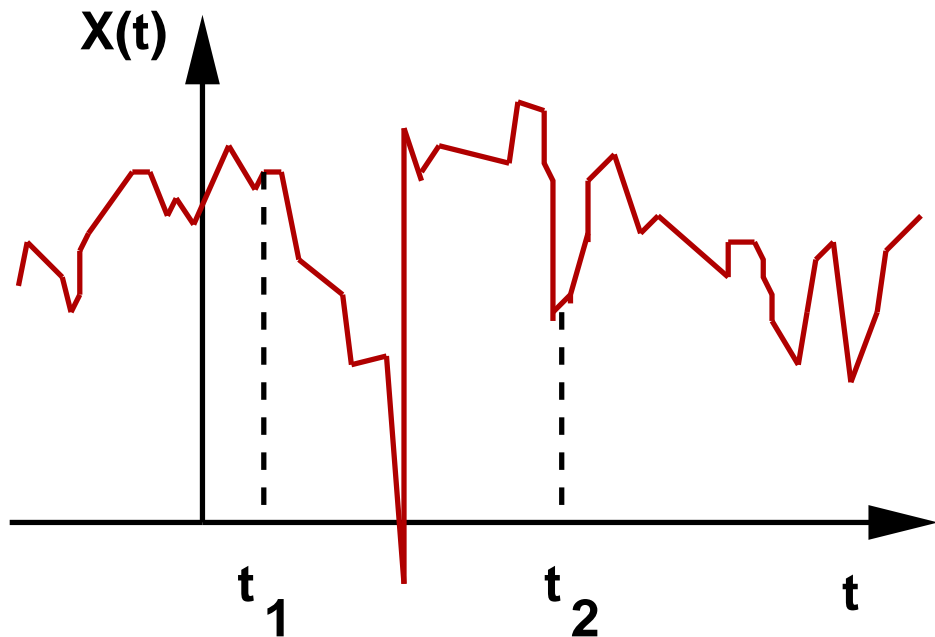
Autocorrelation

$$A(\tau) = \int_{-\infty}^{\infty} X(t)X(t + \tau) dt$$

(1)

where $\tau = t_2 - t_1$, and $X(t)$ is
“homogeneous random field”.

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Wiener-Khinchin theorem relates $A(\tau)$ with the local fluctuation power spectrum $P(\omega)$ via Fourier transform \mathcal{F} .

$$\mathcal{F}_{\omega \rightarrow \tau}(P(\omega)) = A(\tau) \quad (2)$$

$$P(\omega) = \mathcal{F}_{\tau \rightarrow \omega}(A(\tau)) \quad (3)$$

$$O(N) \rightarrow O(N^2) \quad (4)$$

Prefer $O(N)$ for initial data processing for CPU reasons.

In the **Haar discrete wavelet** basis, the integral equation 3 looks different:

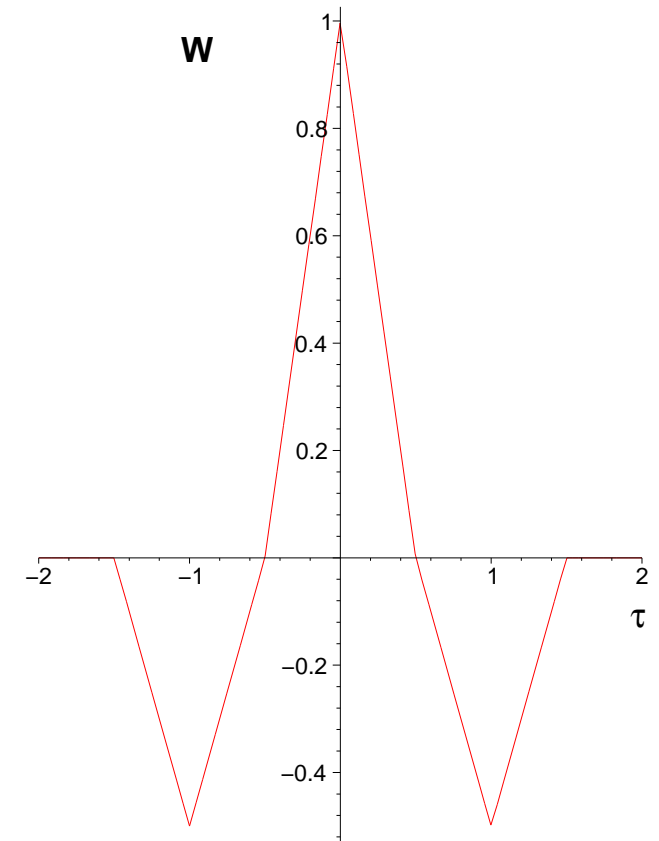
$$P(m) = \overline{\int_{-\infty}^{\infty} X(\tau/2)X(-\tau/2)W(\tau, m) d\tau},$$

where W is the weight function for the Haar wavelet. $P(m)$ reflects differential structure on scale m . See example:

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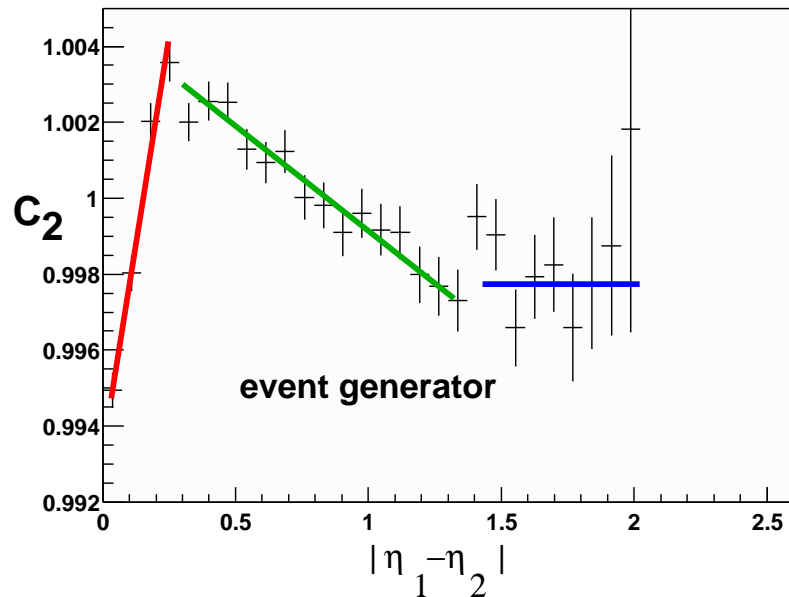
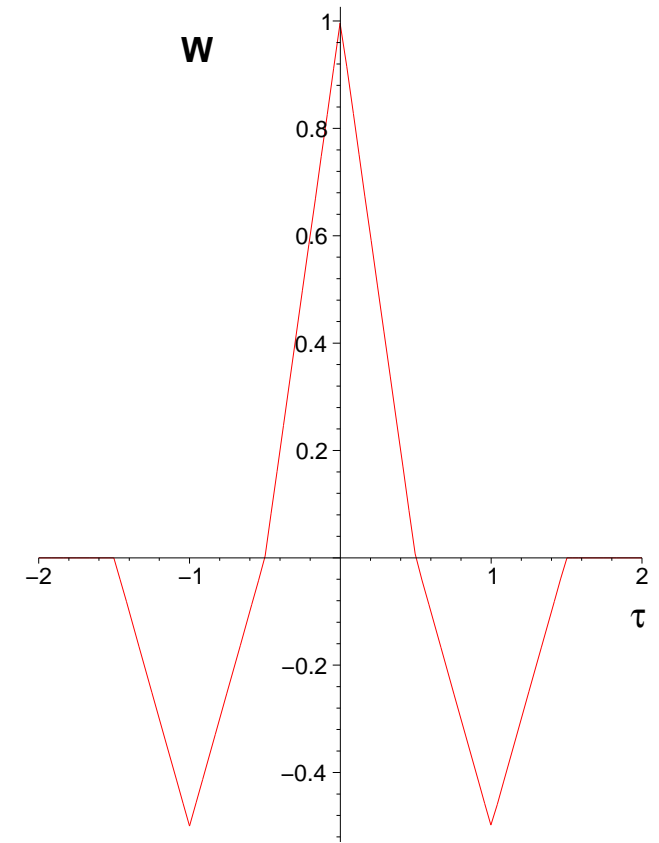
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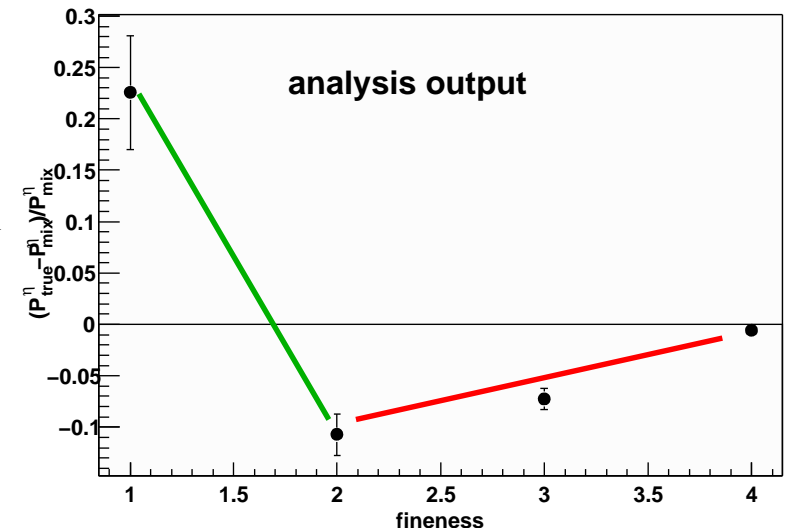
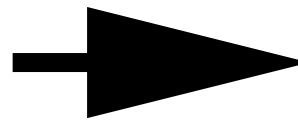
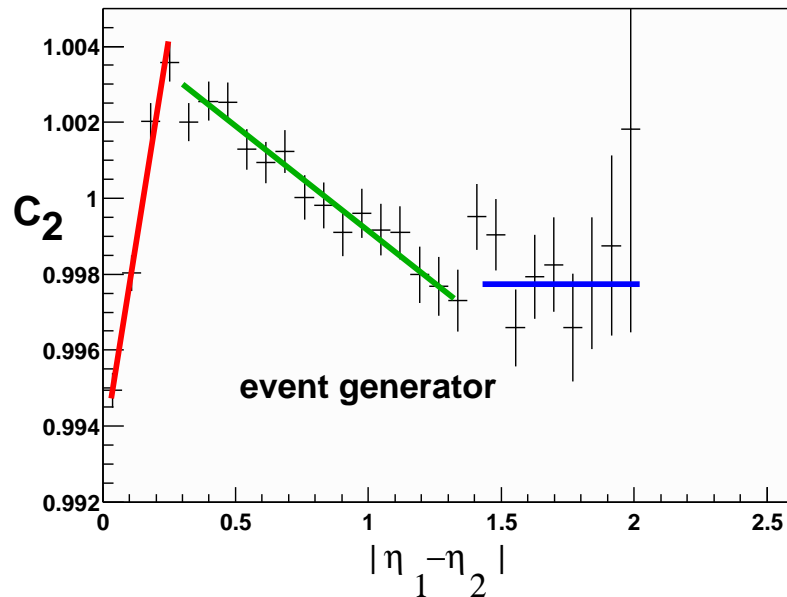
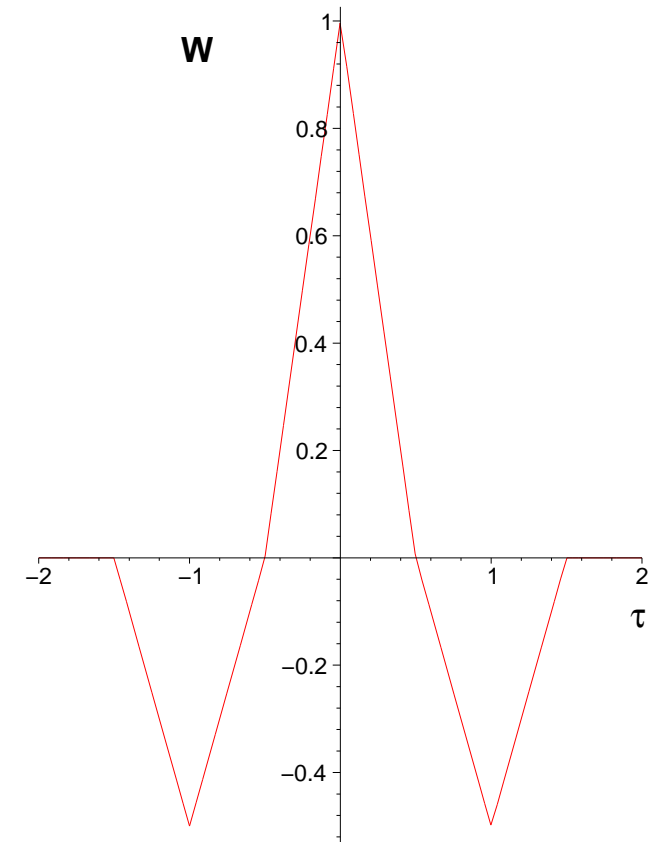
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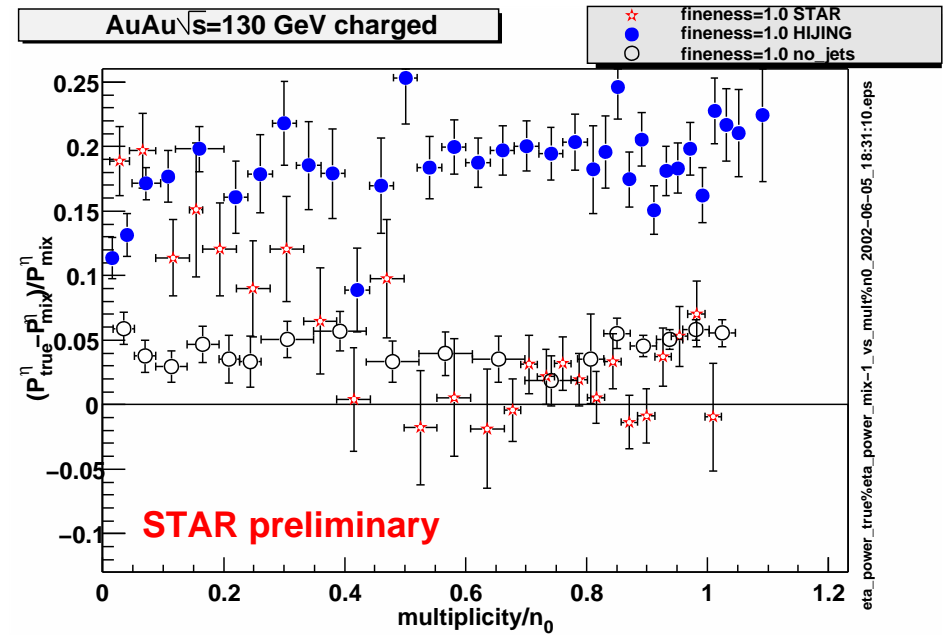
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8 Centrality dependence

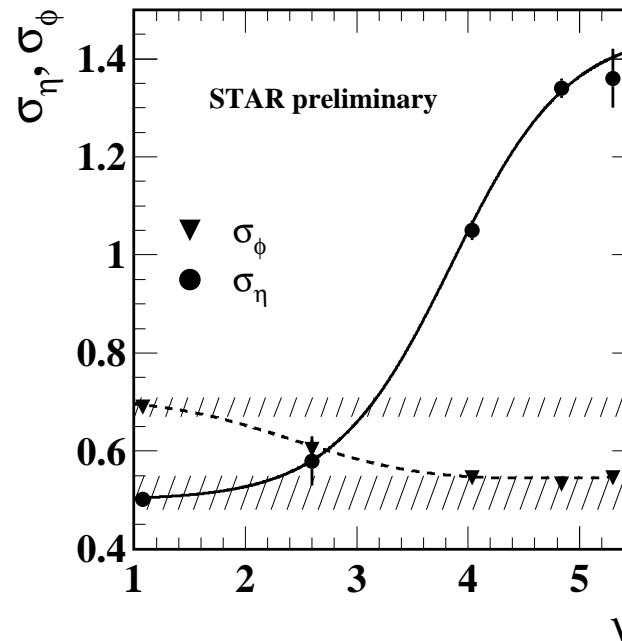
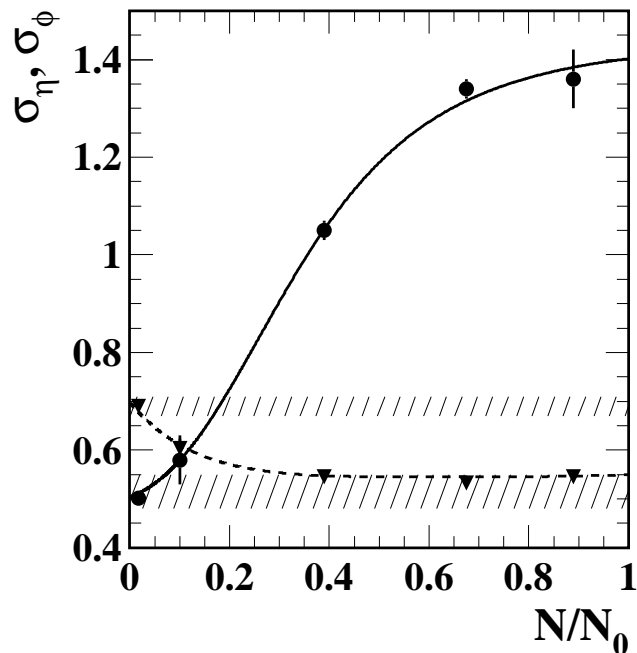
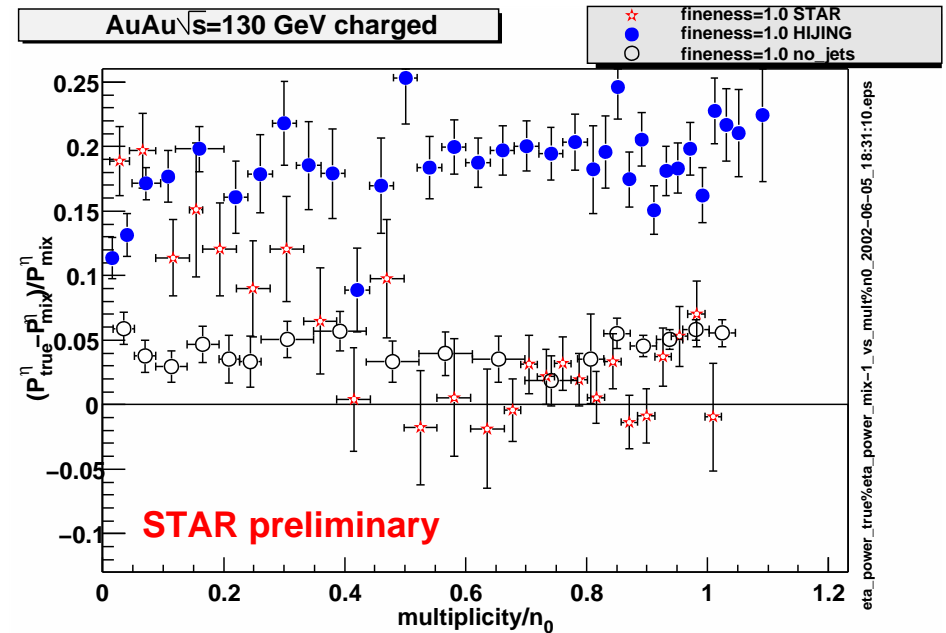
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Istron Bay, Greece, June 2002,
nucl-ex/0211015. STAR DWT
analysis of charge-independent
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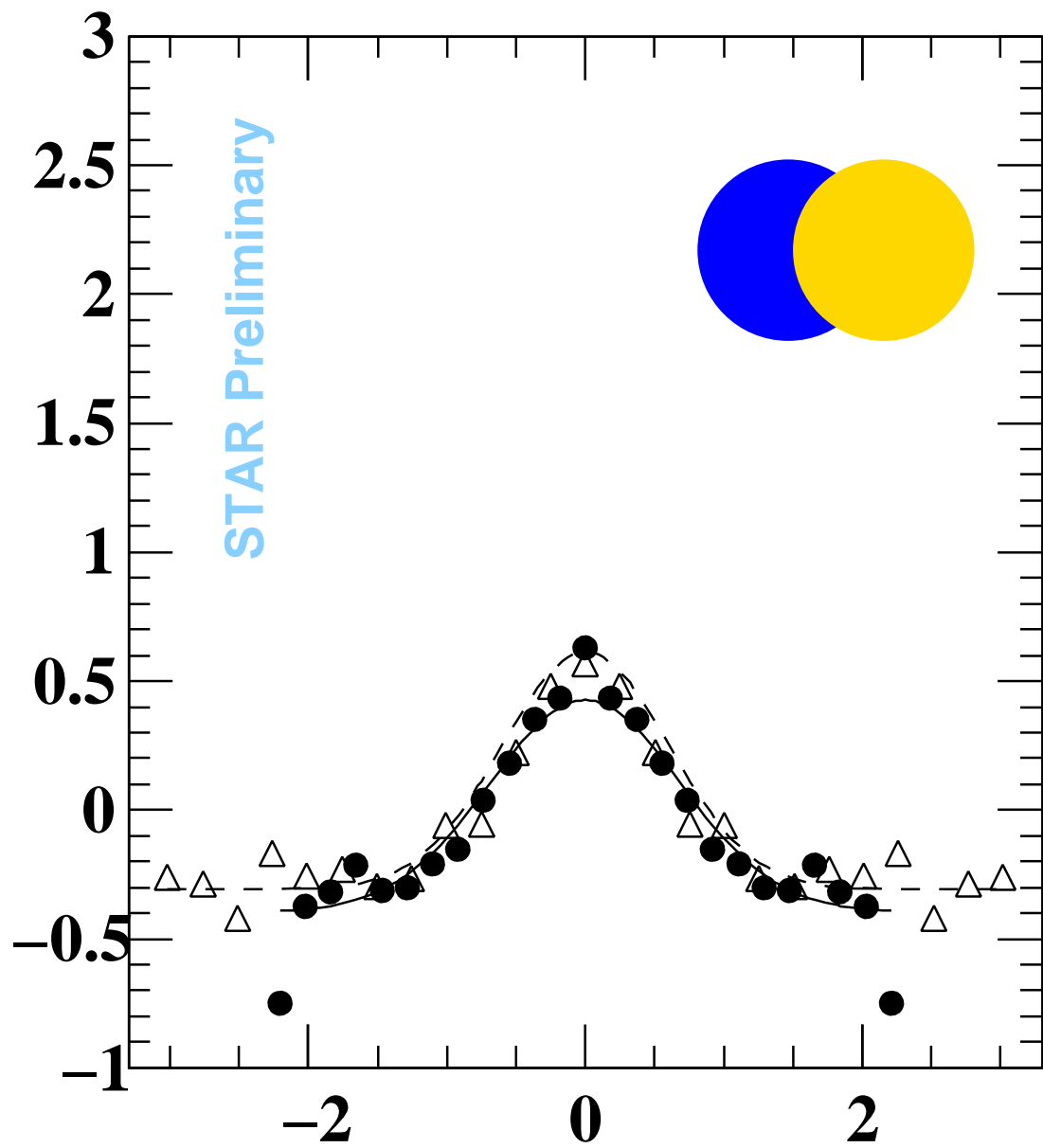


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Istron Bay, Greece, June 2002, nucl-ex/0211015. STAR DWT analysis of charge-independent correlations, AuAu, $\sqrt{s_{NN}} = 130$ GeV



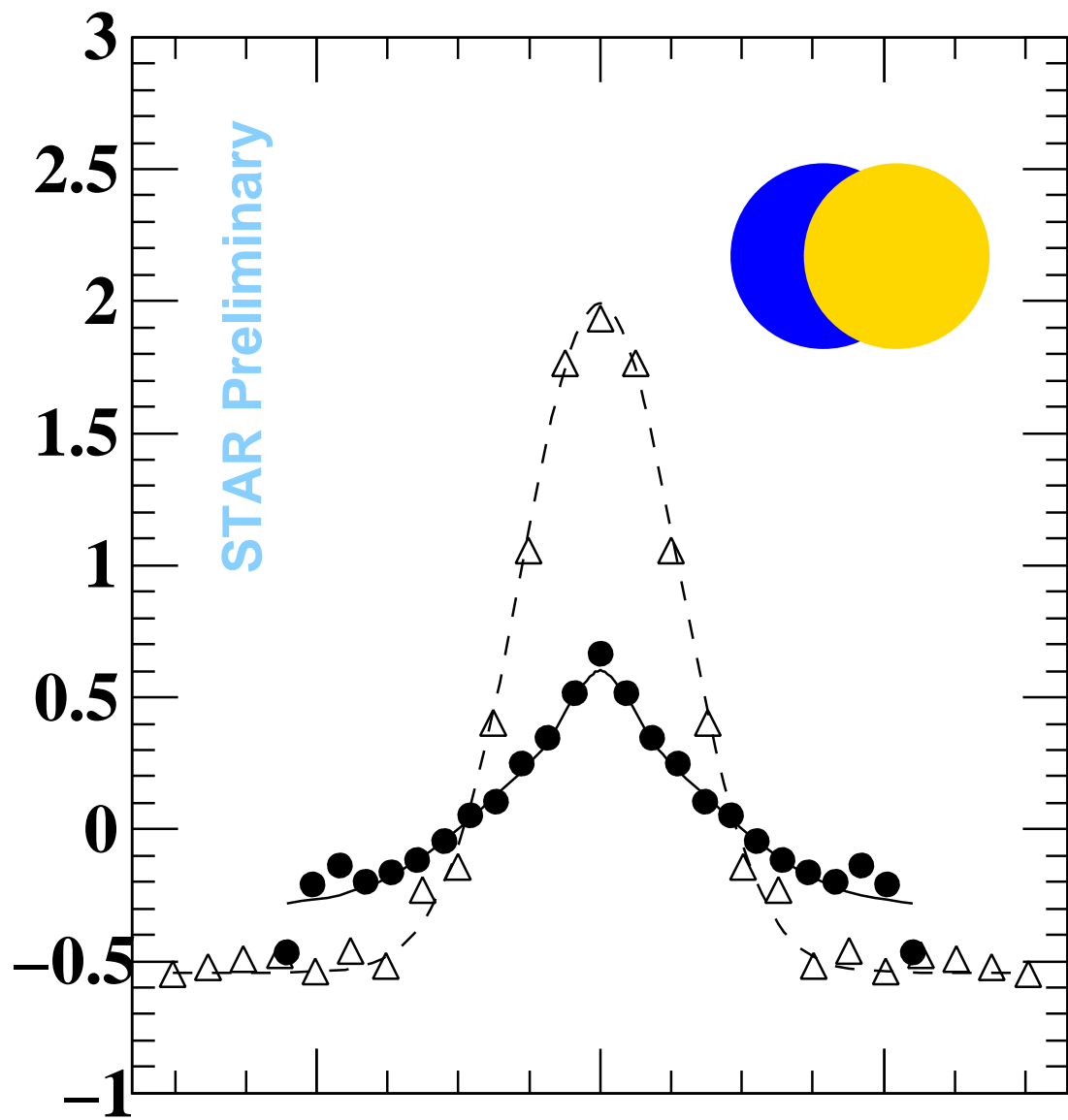
[STAR two-particle correlation analysis, to be submitted to PRL]
The step-like character of the centrality dependence is elucidated using $\nu = (N_{part}/2)^{1/3}$



(d)

\mathbf{x}_Δ

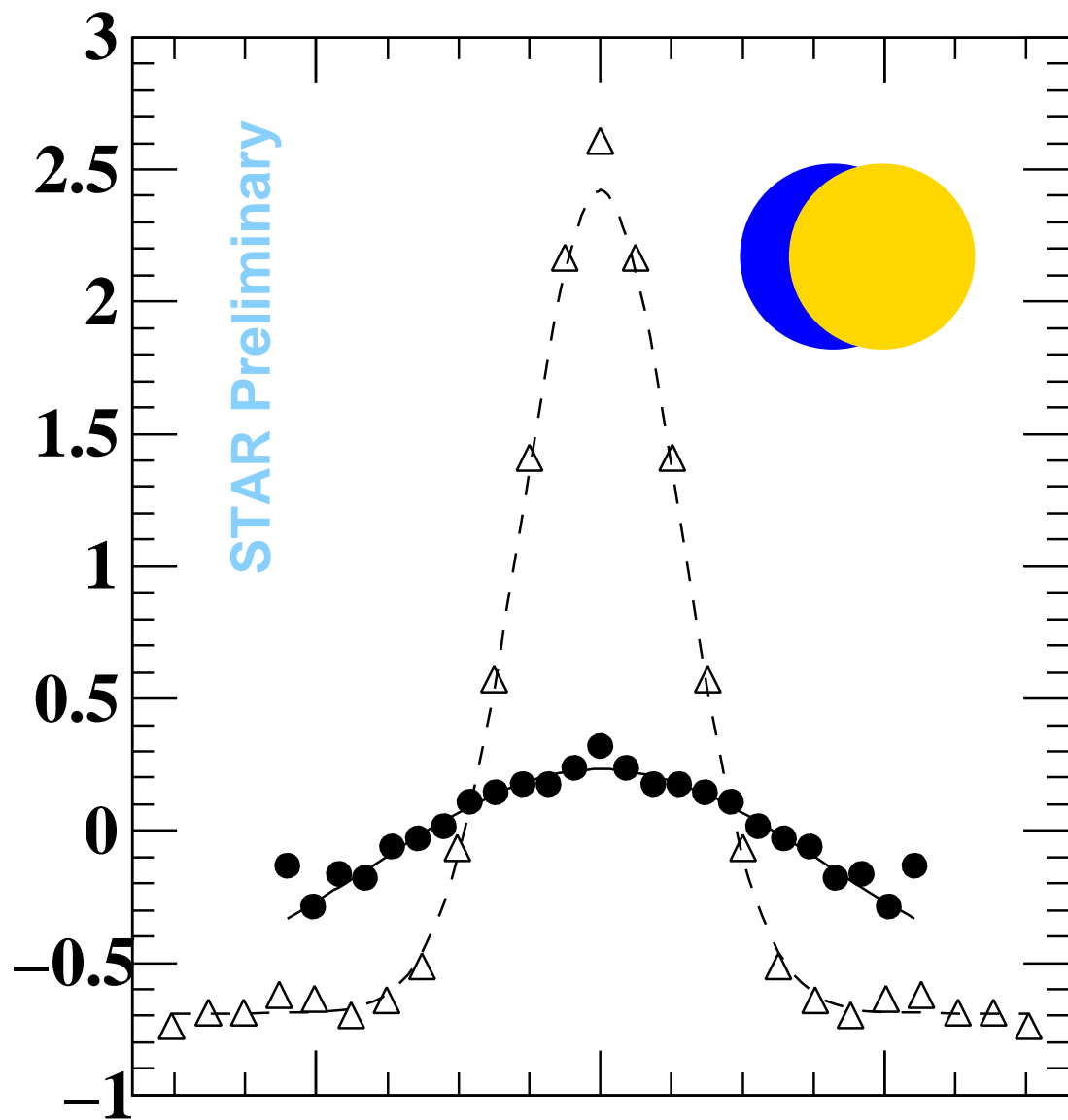
Per-particle-normalized charge-independent autocorrelations $N\Delta R$ in difference variables η_Δ (solid dots) and ϕ_Δ (open triangles).



(c)

x_{Δ}

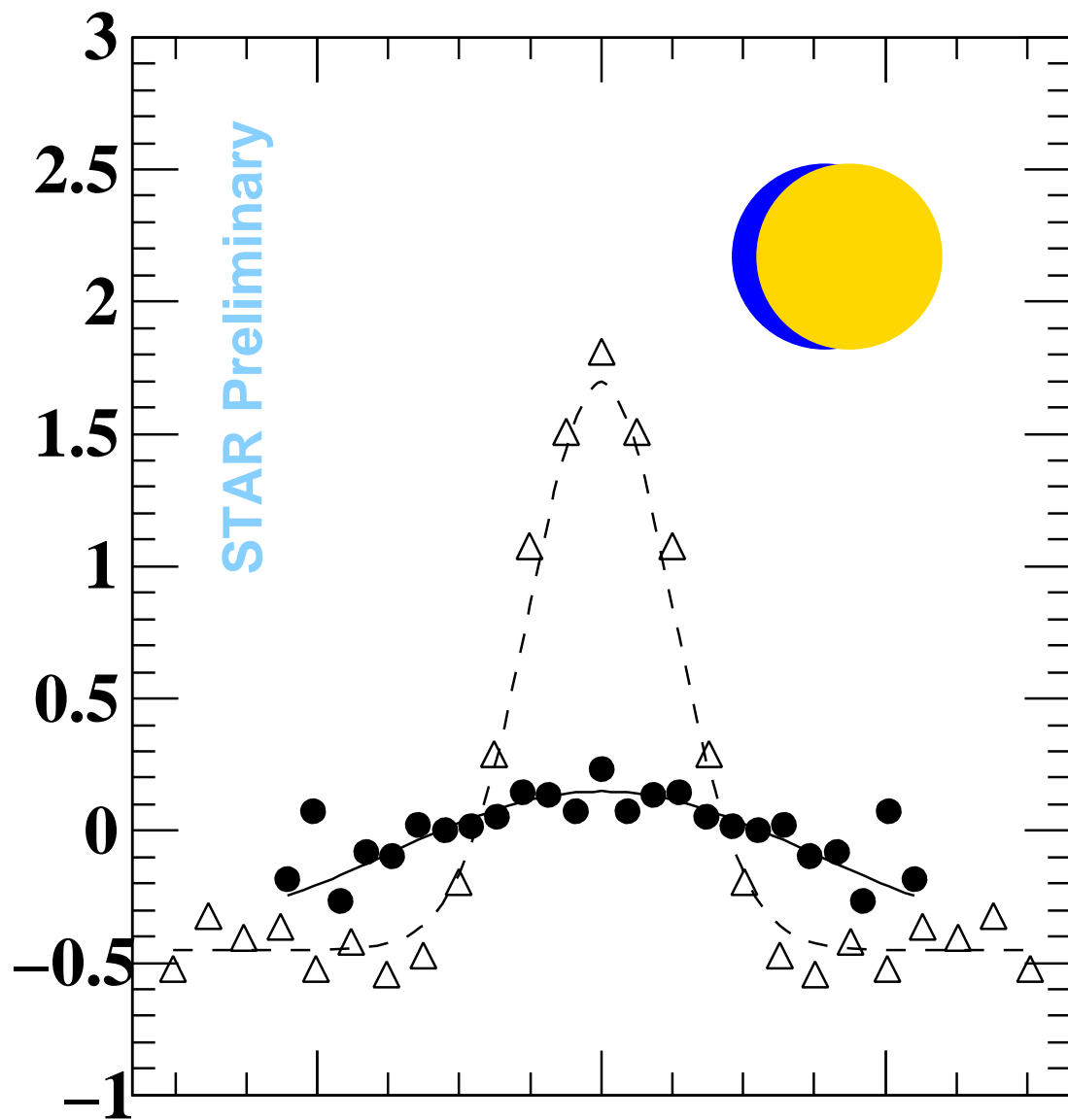
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(b)

x_{Δ}

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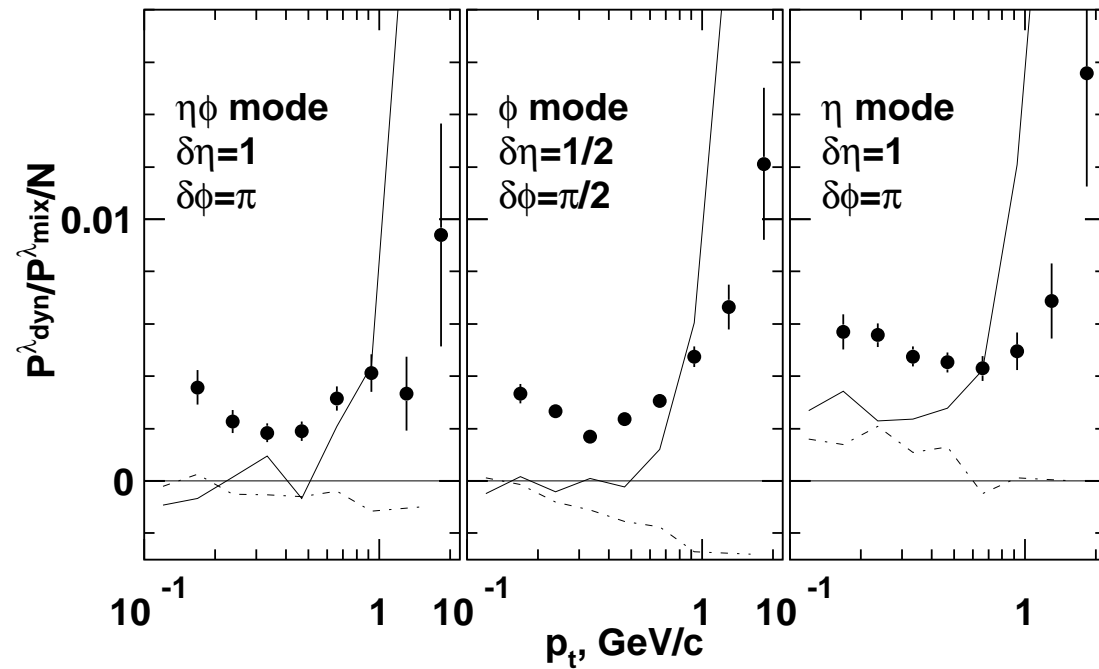


(a)

\mathbf{x}_{Δ}

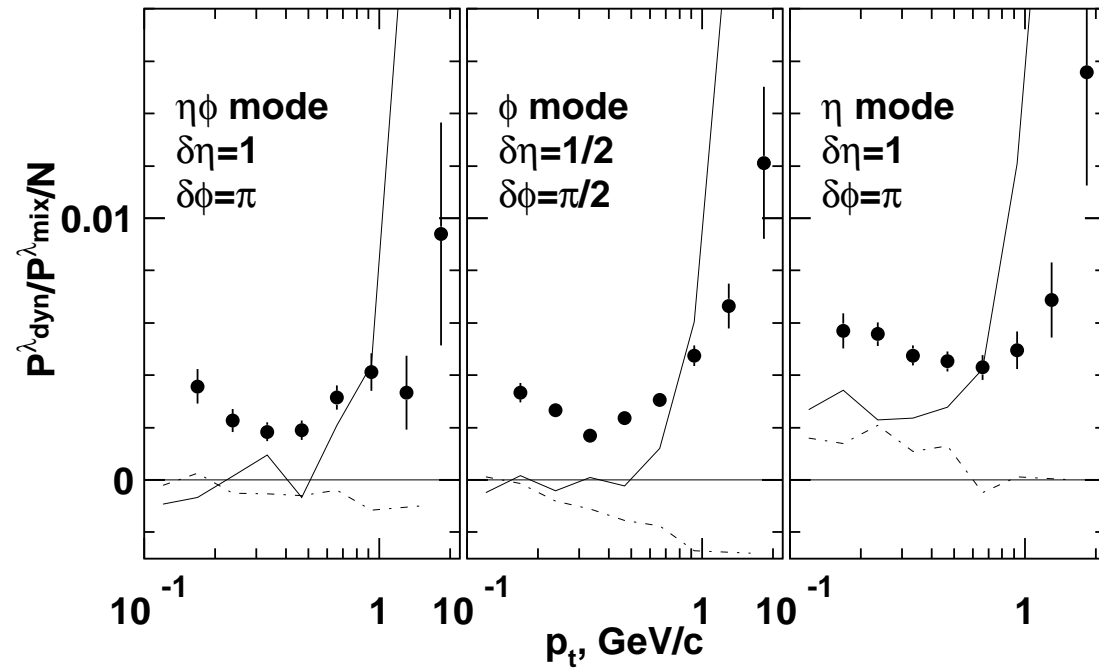
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9 p_t dependence: peripheral events, $\sqrt{s_{NN}} = 200$ GeV



Peripheral events (60-84%): normalized dynamic texture for fineness scales $m = 0, 1, 0$ from left to right panels, respectively, as a function of p_t . ● – STAR data; solid line – standard HIJING, dash-dotted line – HIJING without jets.

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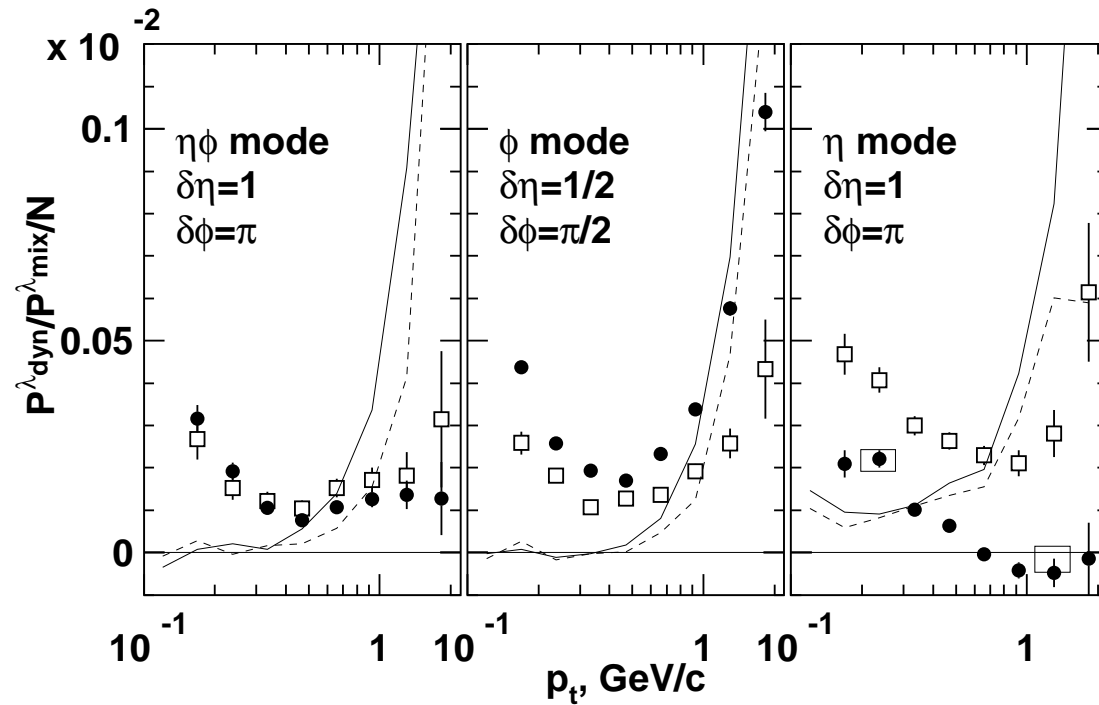


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Qualitative trends in peripheral data are as expected. What signal to expect in the central data, if correlation does not change ?

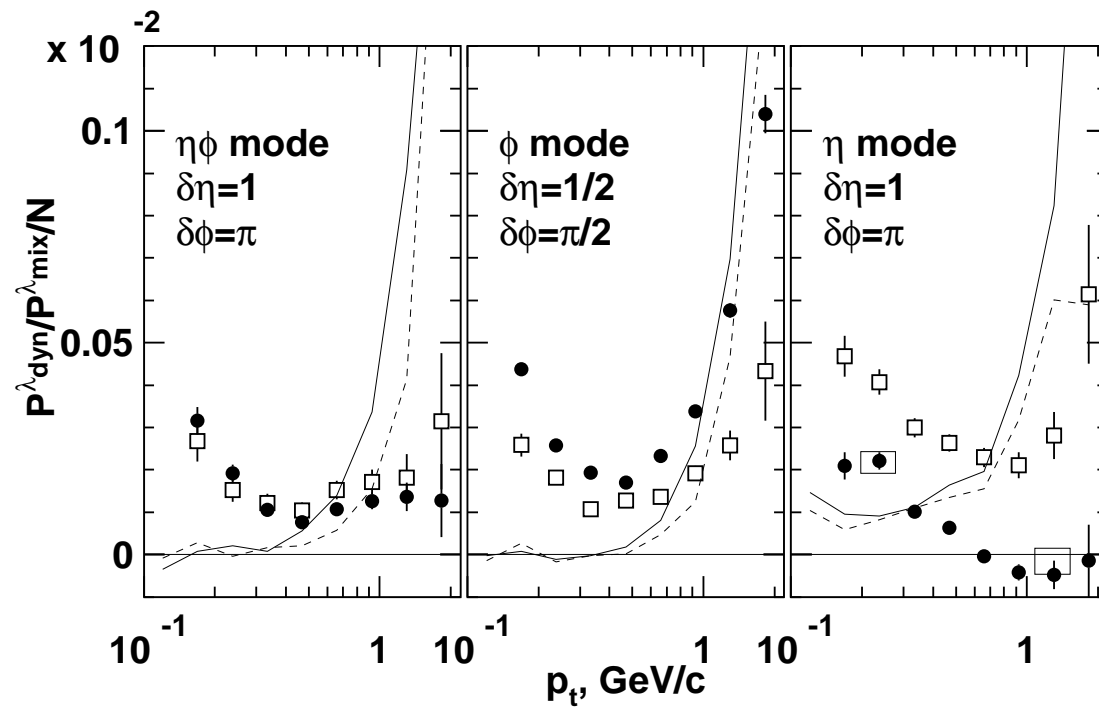
$$\left(\frac{P_{true}}{P_{mix}} - 1 \right) \frac{1}{N} \Big|_{central} = \left(\frac{P_{true}}{P_{mix}} - 1 \right) \Big|_{peripheral} \frac{1}{N_{central}} \quad (5)$$

10 p_t dependence: central events, $\sqrt{s_{NN}} = 200$ GeV



Central (top 4%) events:
 normalized dynamic texture
 for fineness scales $m = 0, 1, 0$
 from left to right panels,
 respectively, as a function of
 p_t . ● – STAR data; solid line –
 standard HIJING; dashed line –
 HIJING with jet quenching;
 boxes – peripheral STAR data
 just shown, renormalized as
 just described.

10 p_t dependence: central events, $\sqrt{s_{NN}} = 200$ GeV



Central (top 4%) events: normalized dynamic texture for fineness scales $m = 0, 1, 0$ from left to right panels, respectively, as a function of p_t . ● – STAR data; solid line – standard HIJING; dashed line – HIJING with jet quenching; boxes – peripheral STAR data just shown, renormalized as just described.

We are observing a modification of the minijet structure predominantly in the longitudinal, η direction. Longitudinal expansion of the hot and dense medium formed early in the collision makes this direction special and is likely to be part of the modification mechanism.

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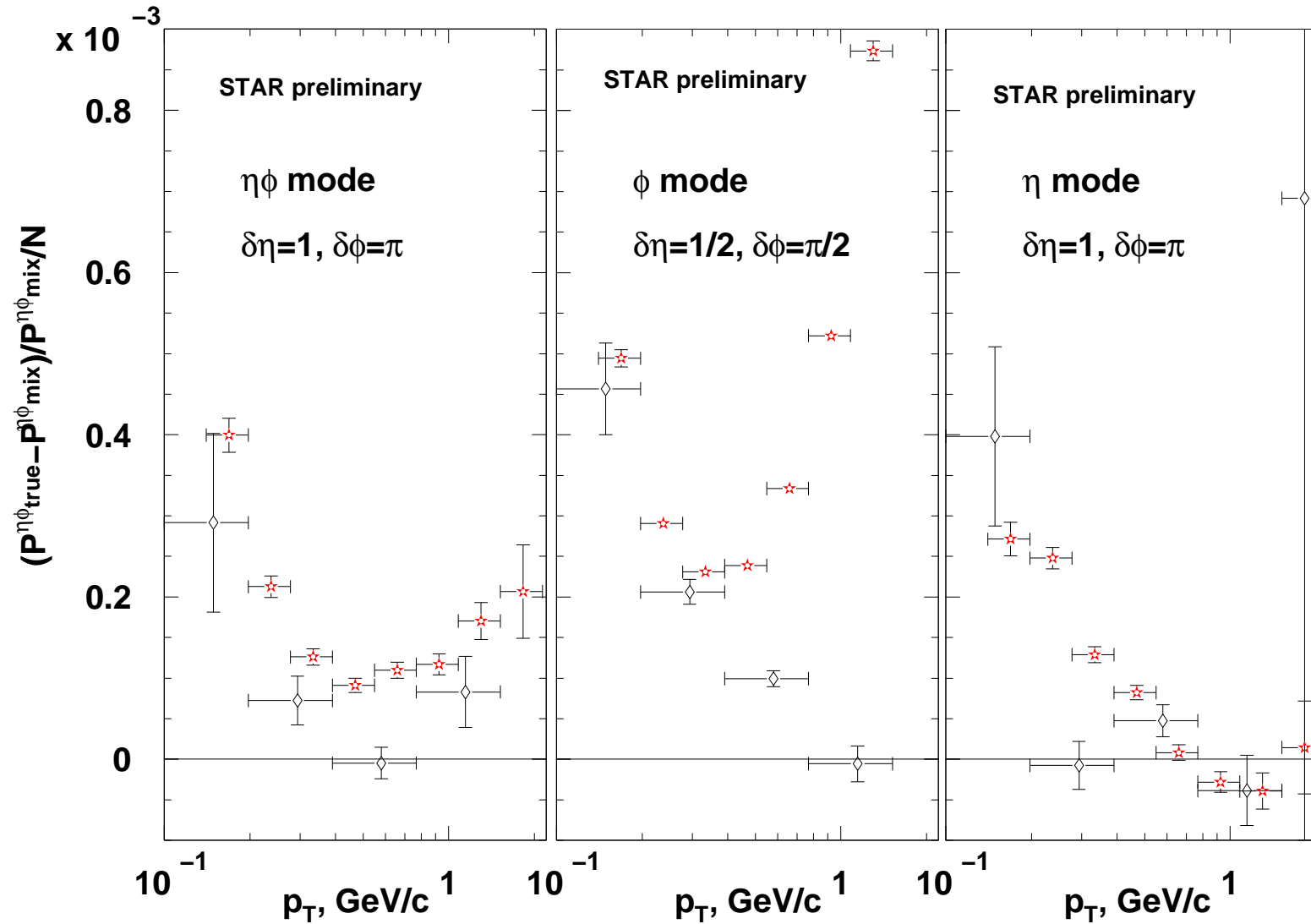
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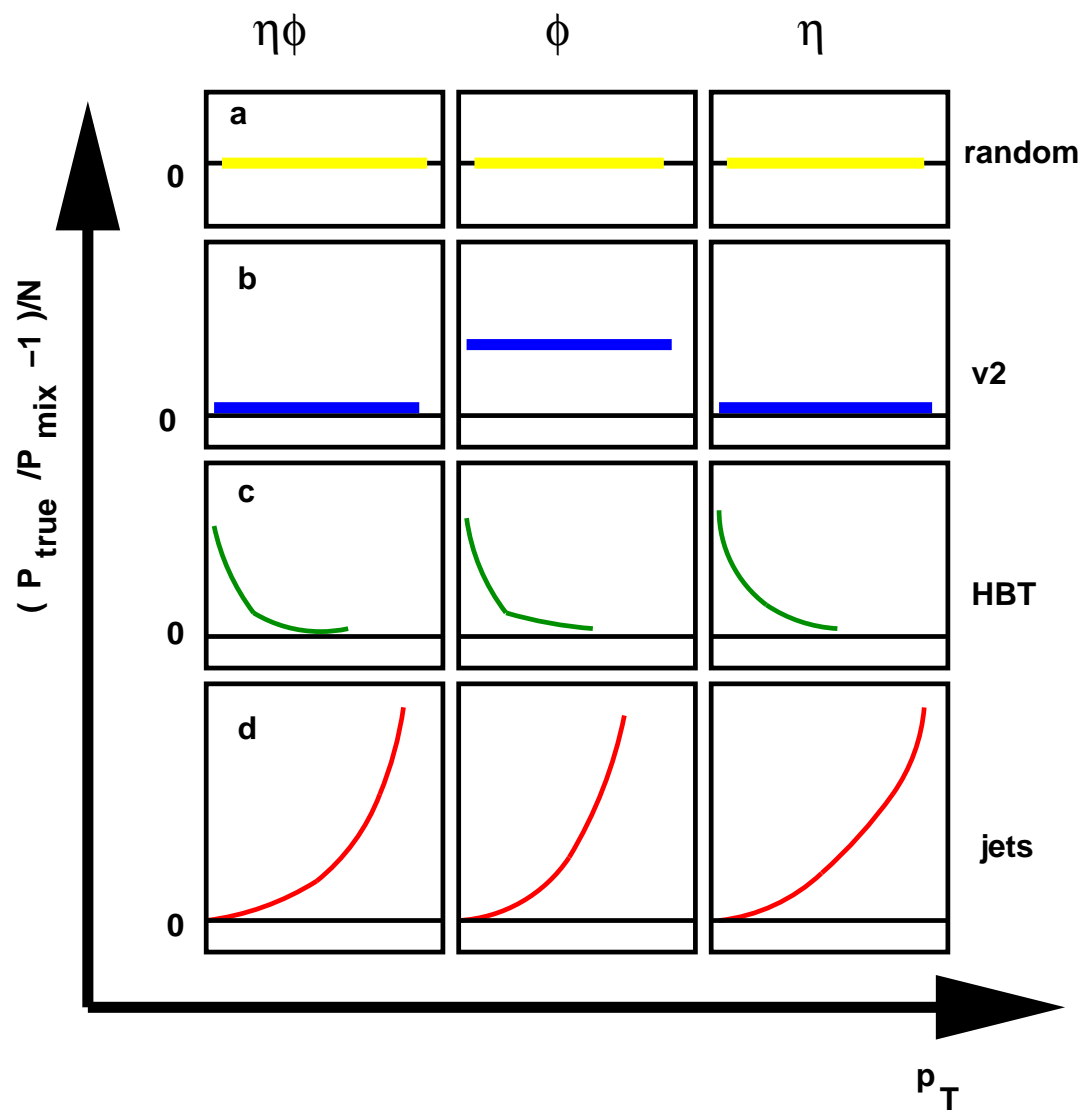
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$$P(N|1, 2, \dots, N - 1) \approx P(N)\prod_{i=1}^{N-1} C(i, N)$$

12 Understanding low p_t



An event generator tuned to reproduce like- and unlike-sign correlations in Q_{inv} , reproduces the low p_t trends in the data. HBT, Coulomb and string fragmentation physics contribute at low p_t .

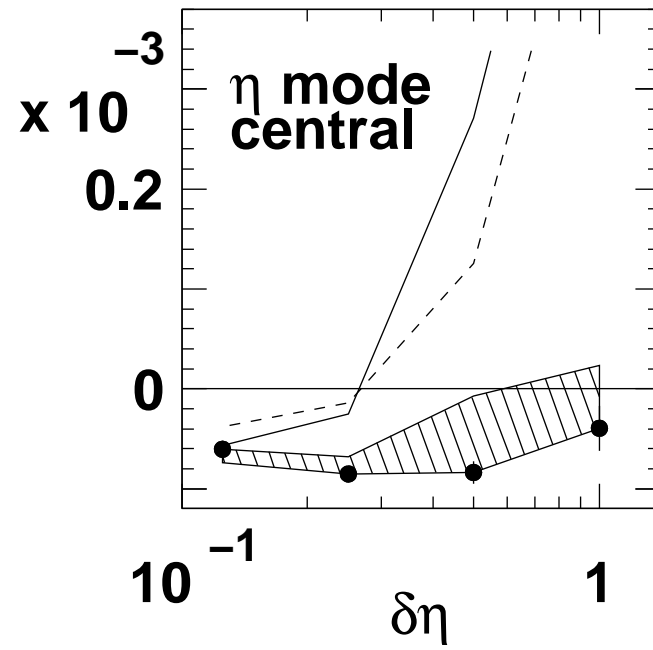
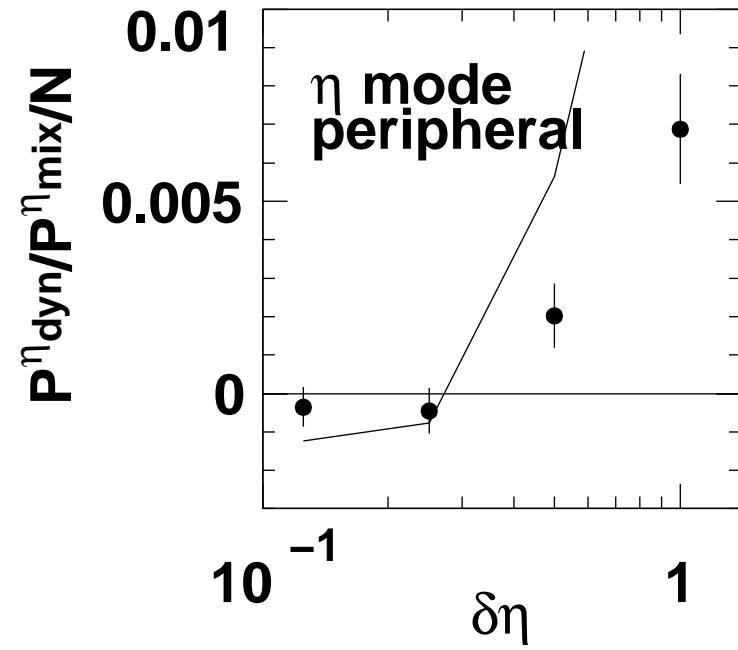


Dynamic texture response in various idealized situations (showing only one scale):

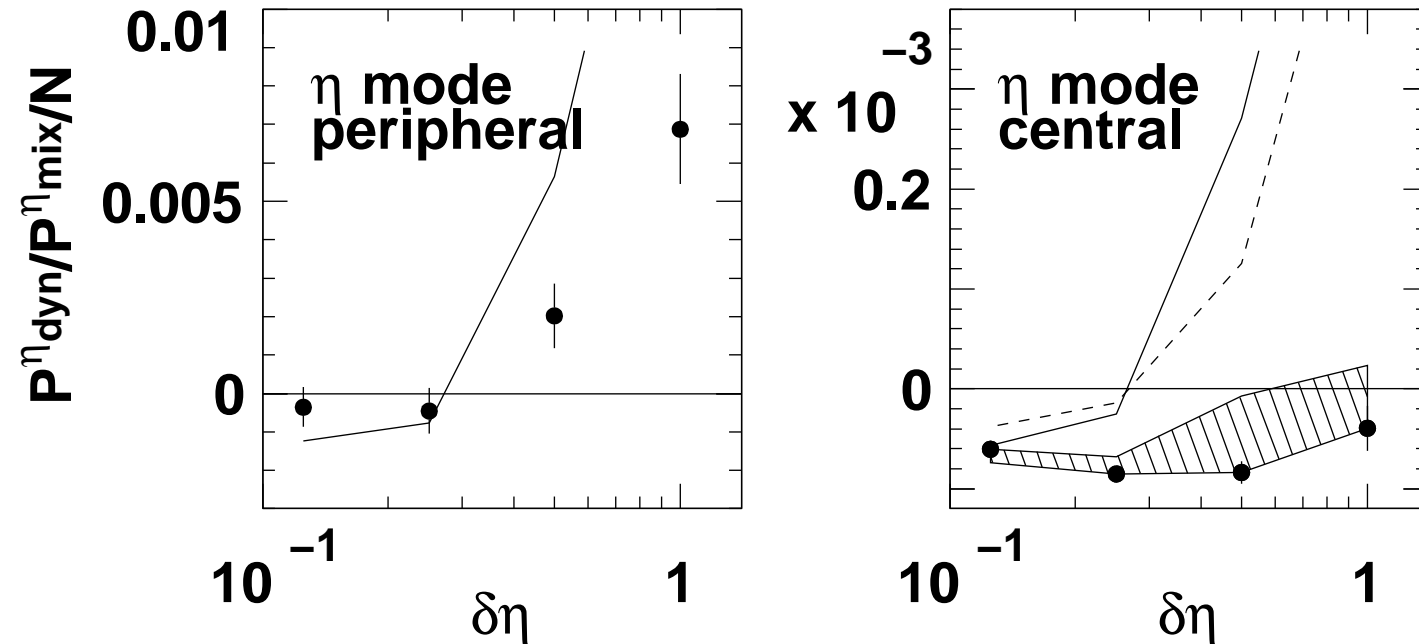
- (a) events of random (uncorrelated) particles
- (b) p_t -independent elliptic flow
- (c) Correlations at low Q_{inv} (Bose-Einstein correlations and Coulomb effect)
- (d) HIJING jets

13 Scale dependence

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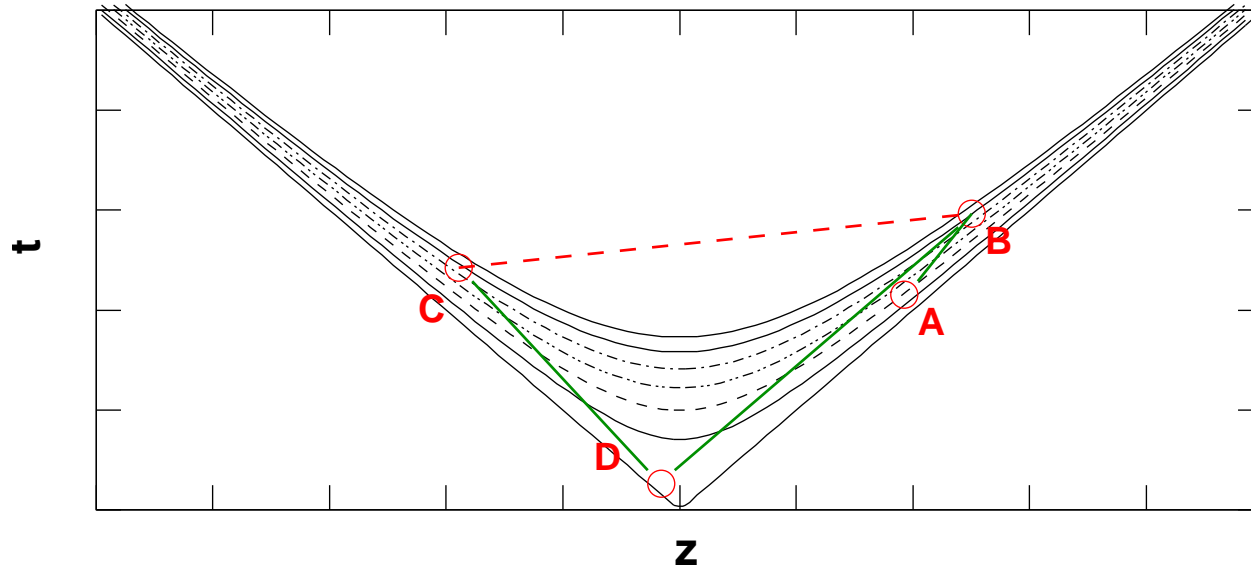
13 Scale dependence



Scale dependence of the dynamic texture measure in peripheral and central events for $1.1 < p_t < 1.5$ GeV/ c . ● – STAR data; solid line – standard HIJING; dashed line – HIJING with jet quenching. An estimate of systematic error, mainly due to track merging, is shown as a hatched area.

14 Rapidity scale and collision history

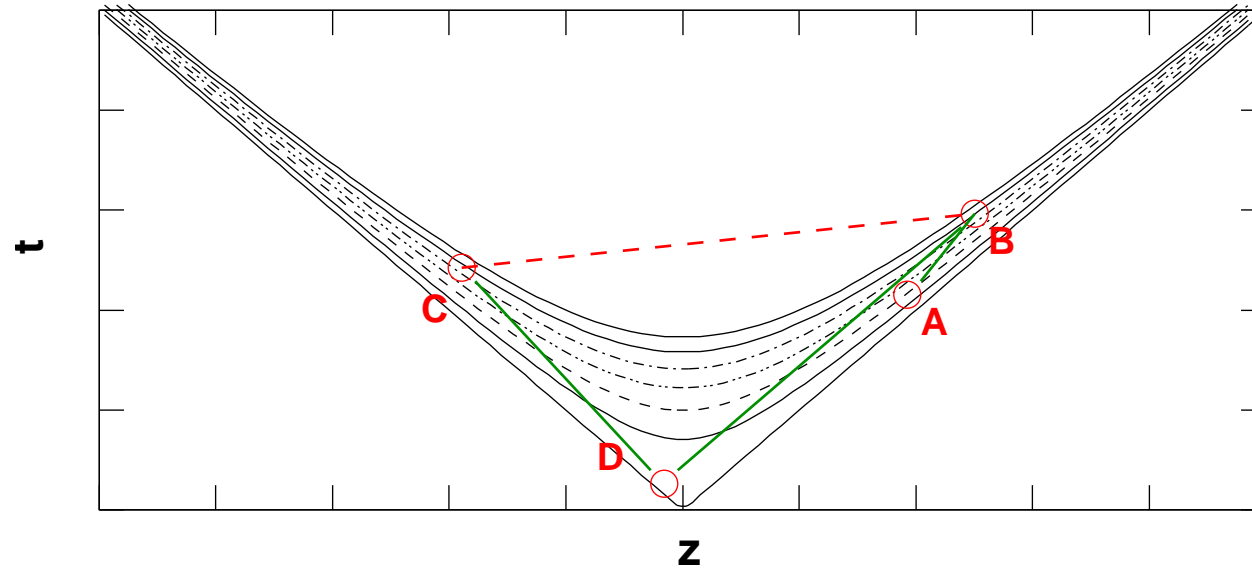
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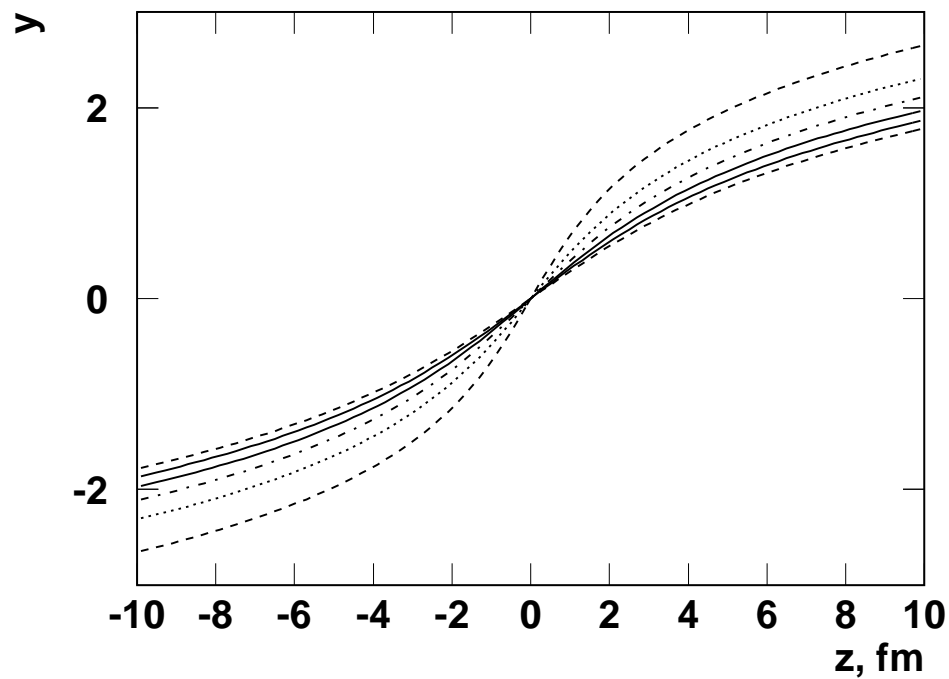
Bjorken expansion:

$t = (z^2 + \tau^2)^{1/2}$, CB is space-like; DC, DA, DB, AB contacts possible.

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$y = \frac{1}{2} \ln \frac{t+z}{t-z}$ Rapidity and causality:
 large $\delta y \iff$ large δz . Large δy correlations reflect early state, otherwise acausal.

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- More info: nucl-ex/0407001