

Hot matter at RHIC: hot physics at Kent, STAR, PHENIX

Given at the Physics Department, Kent State University,
Ohio.

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April 10, 2007

1 Outline

What:

- Motivation
- Data and future directions
- Conclusions

1 Outline

What:

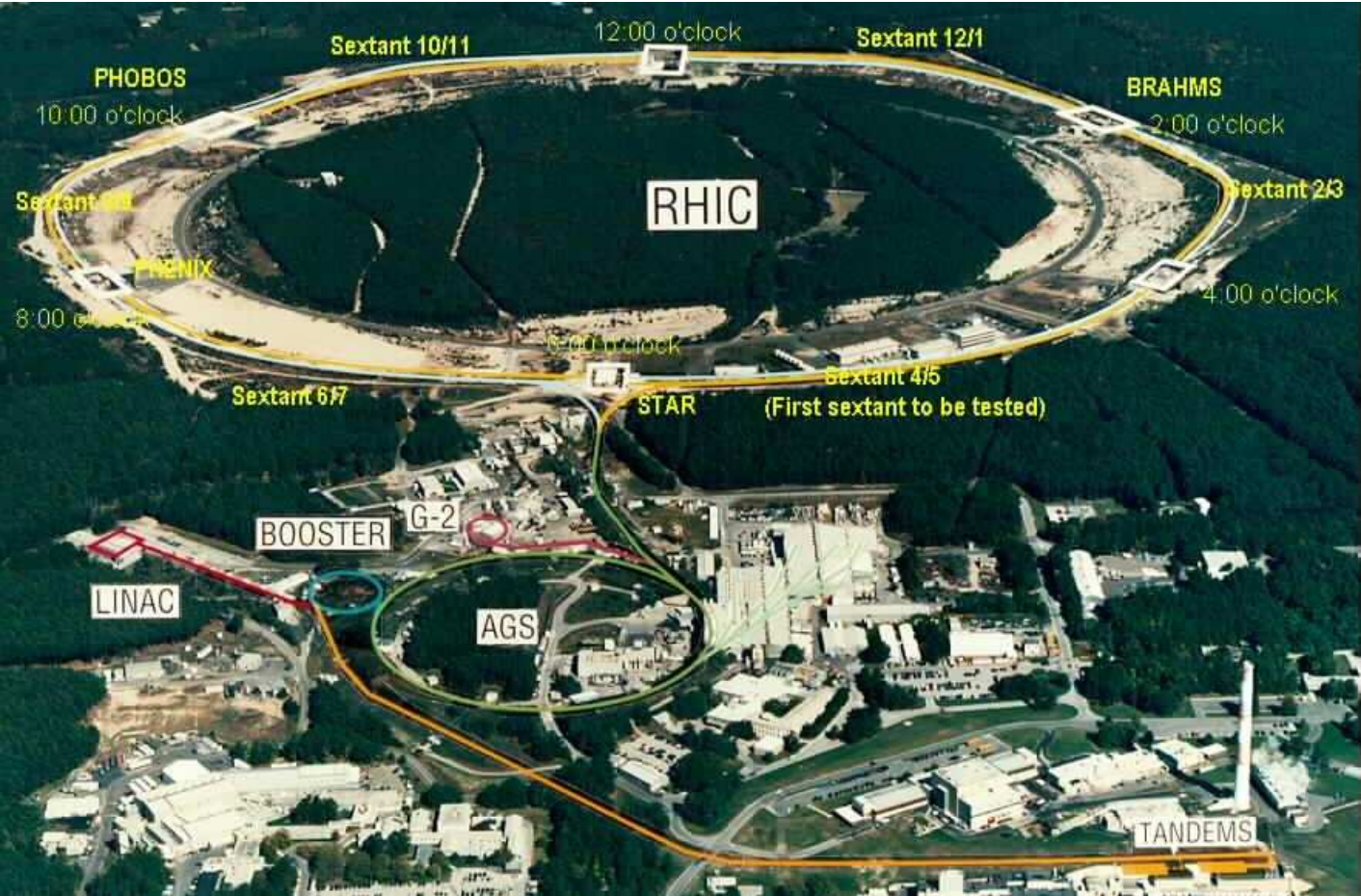
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How:

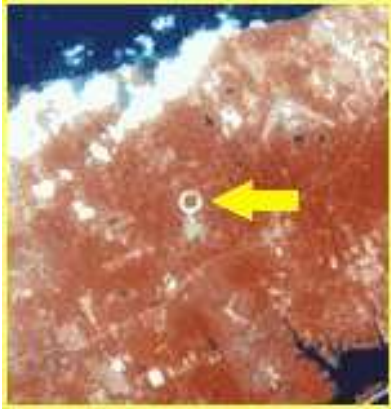
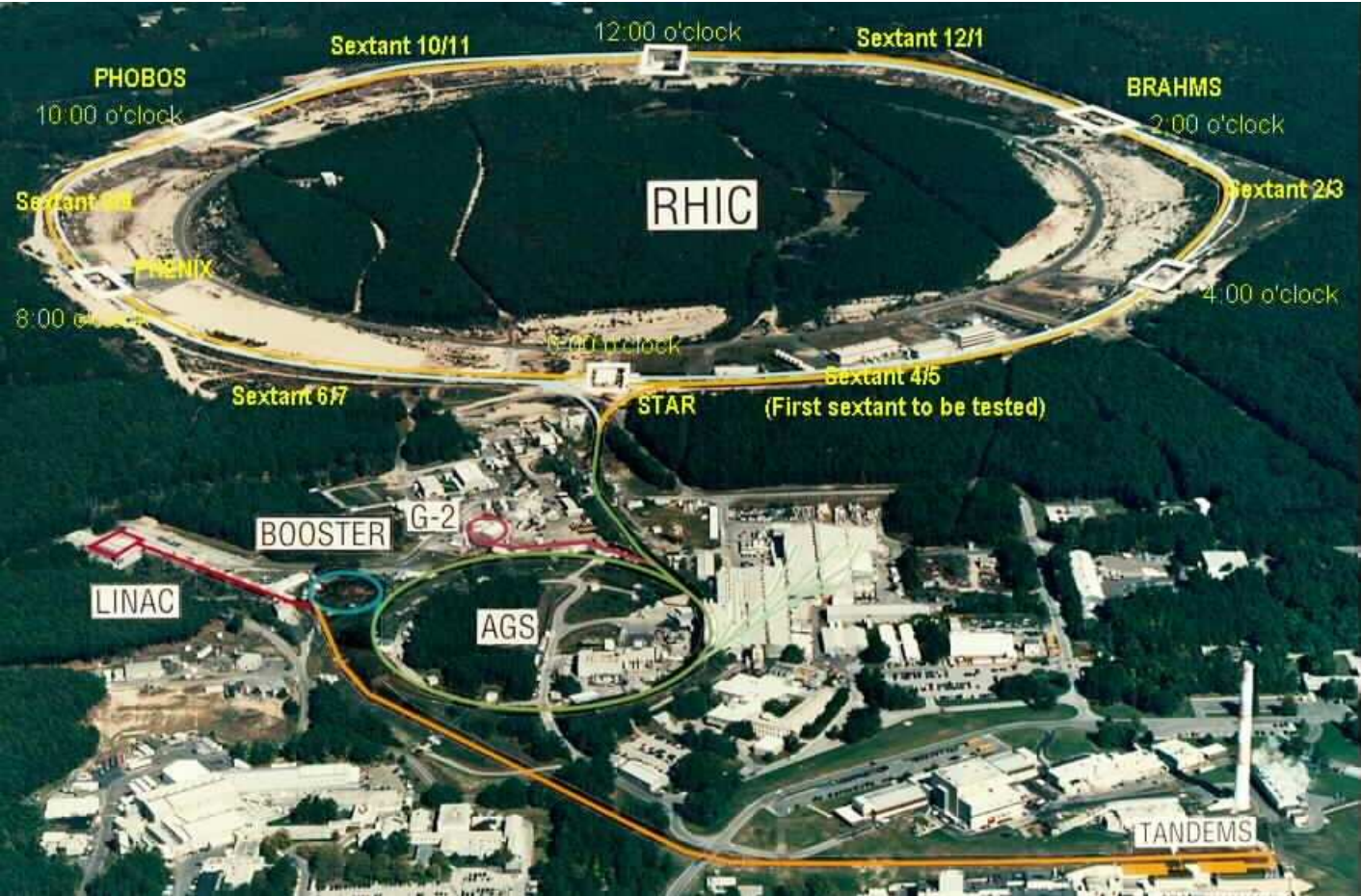
- Focus on recent developments (2005 →)
- Place them in historical and conceptual context

2 The RHIC

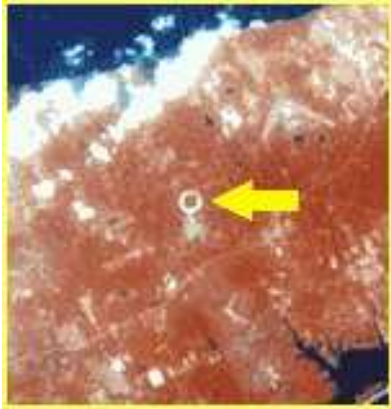
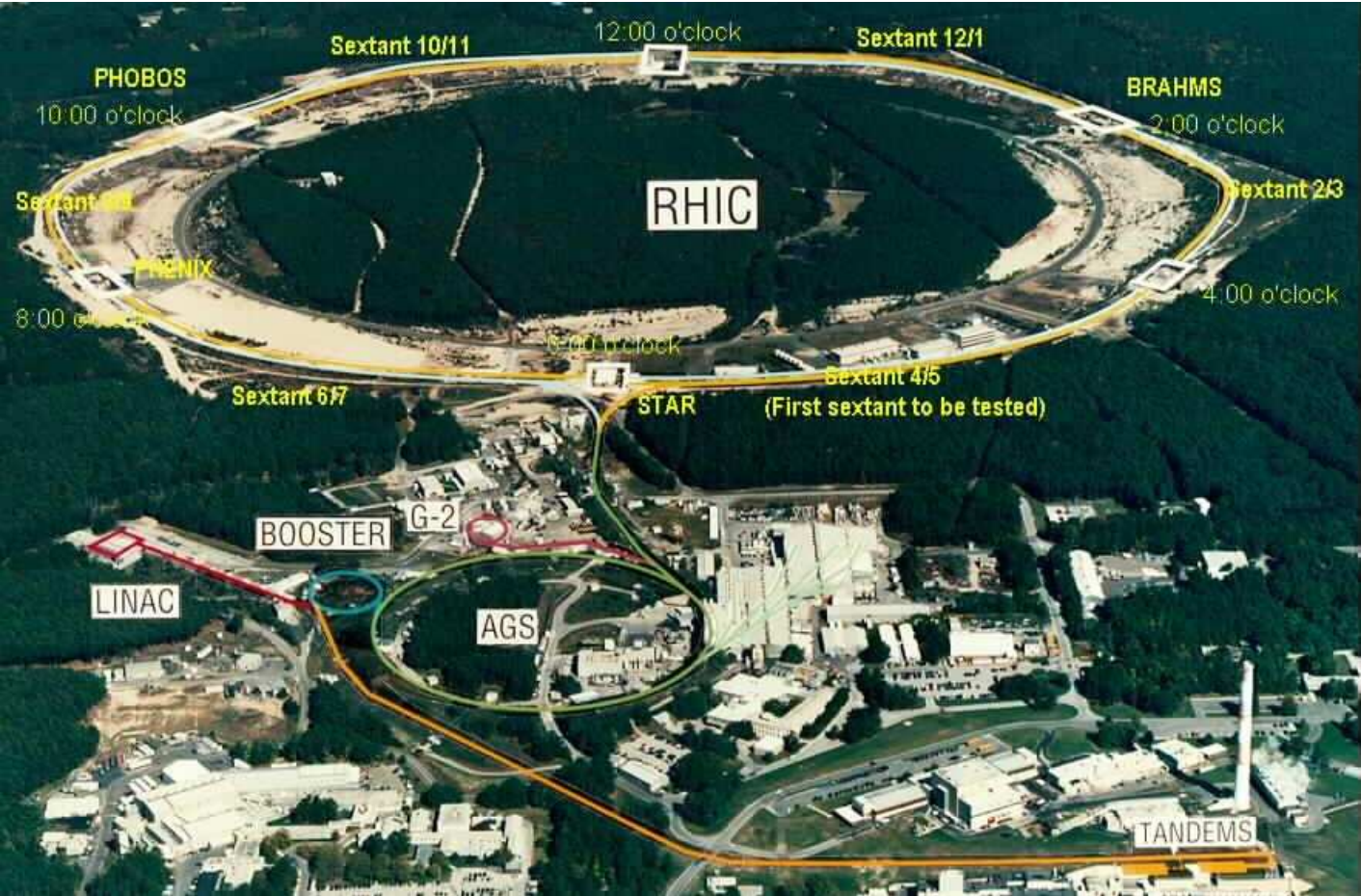
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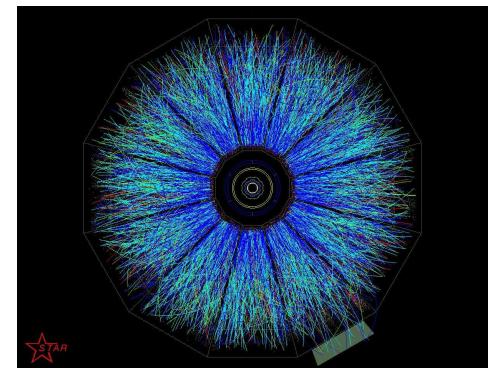
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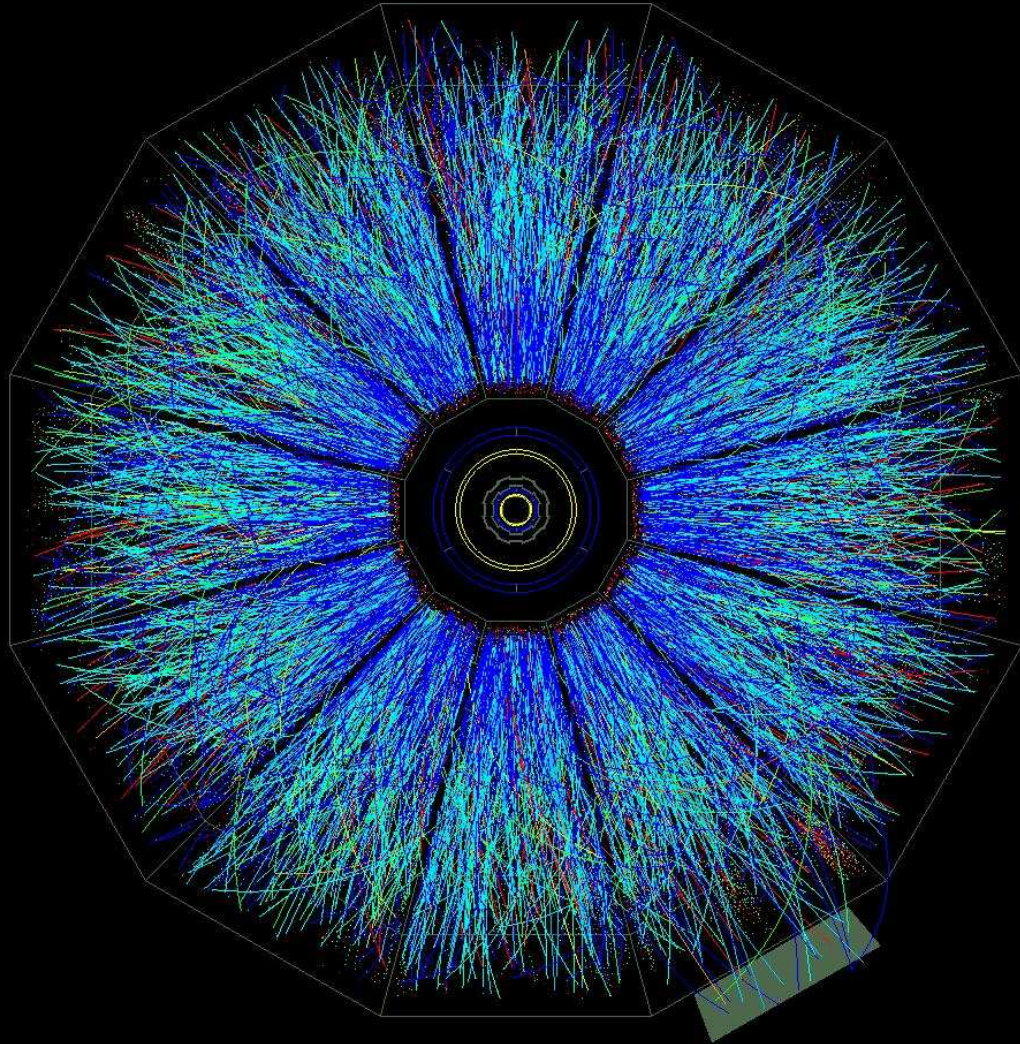
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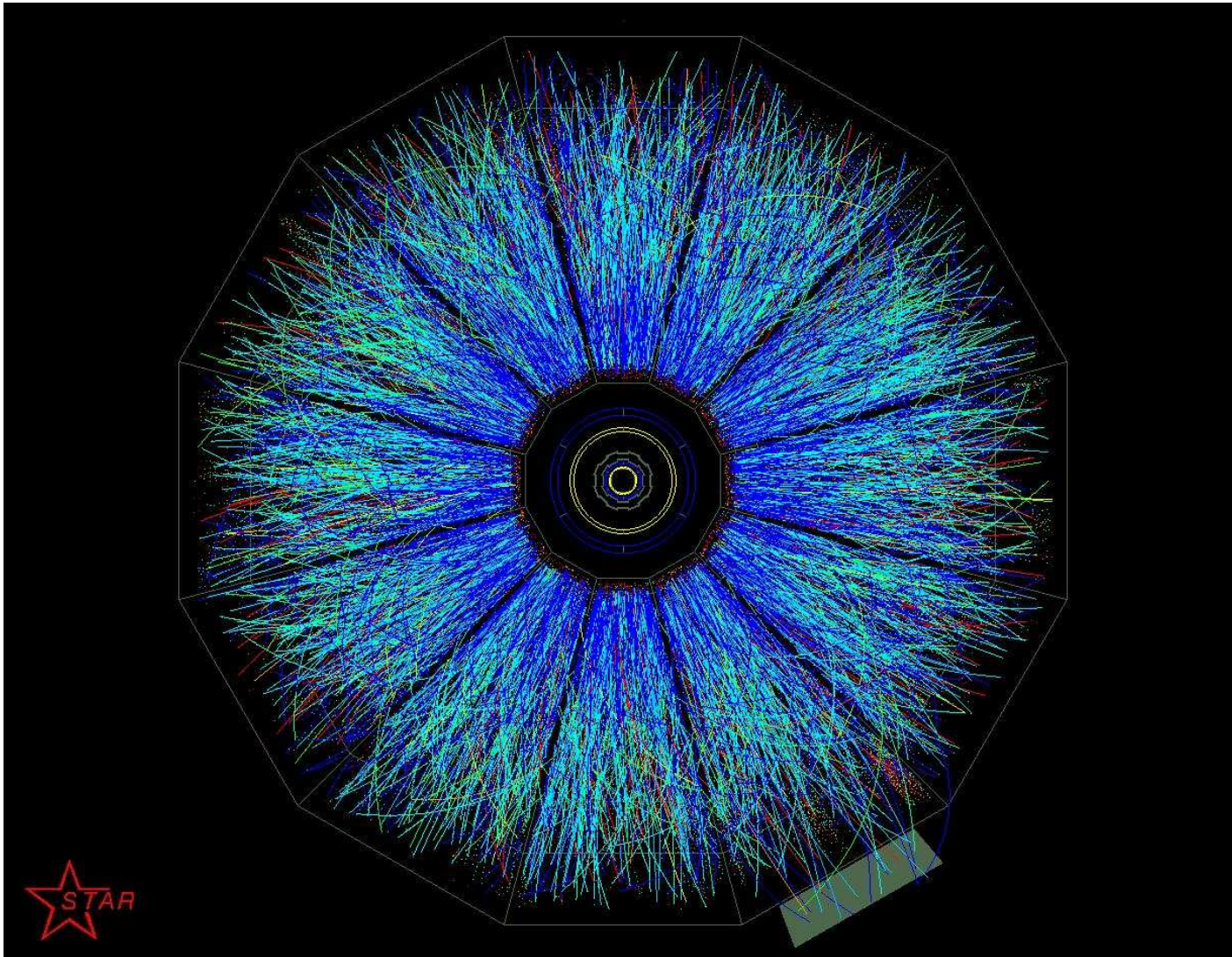
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3 Au+Au collisions at RHIC...

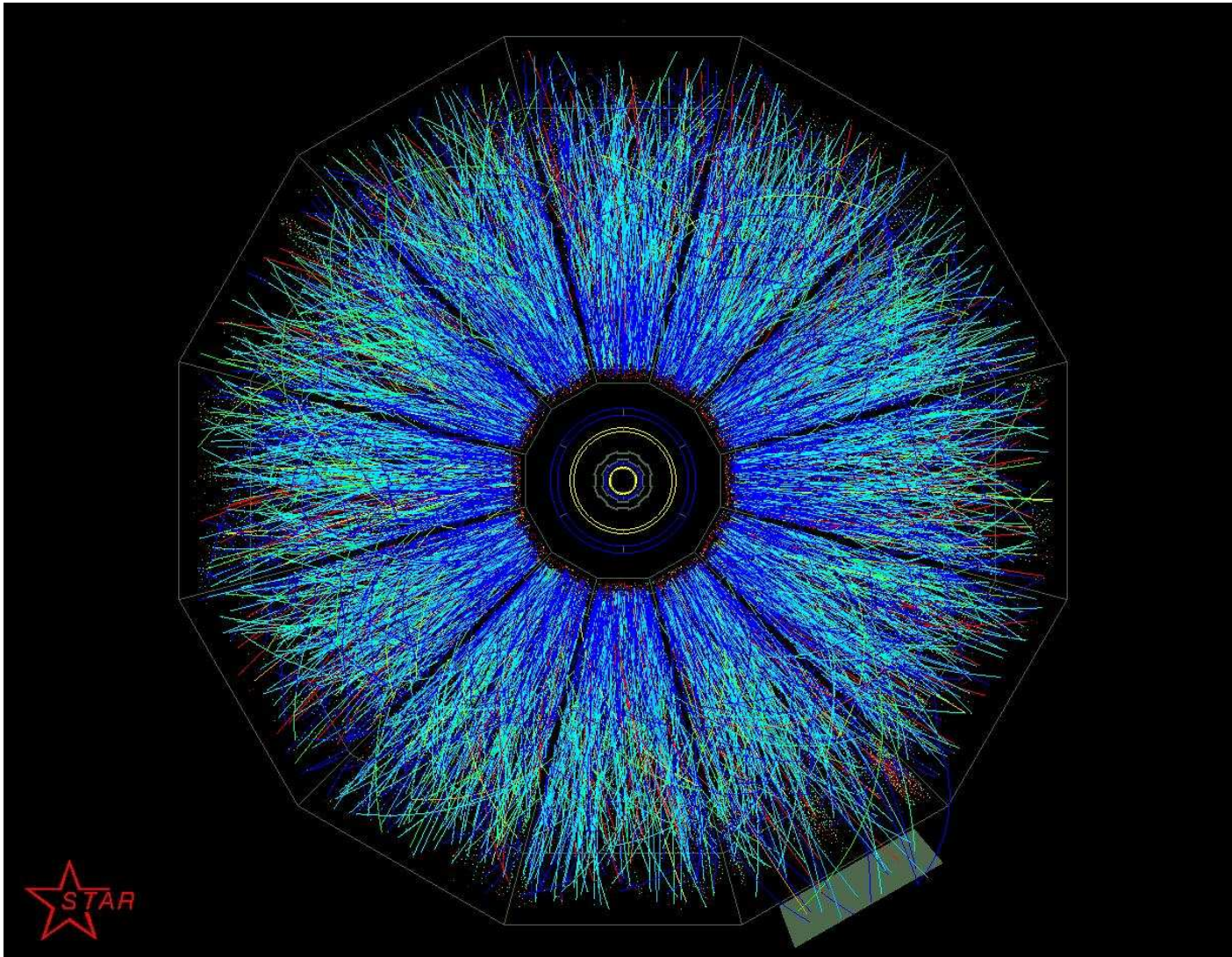


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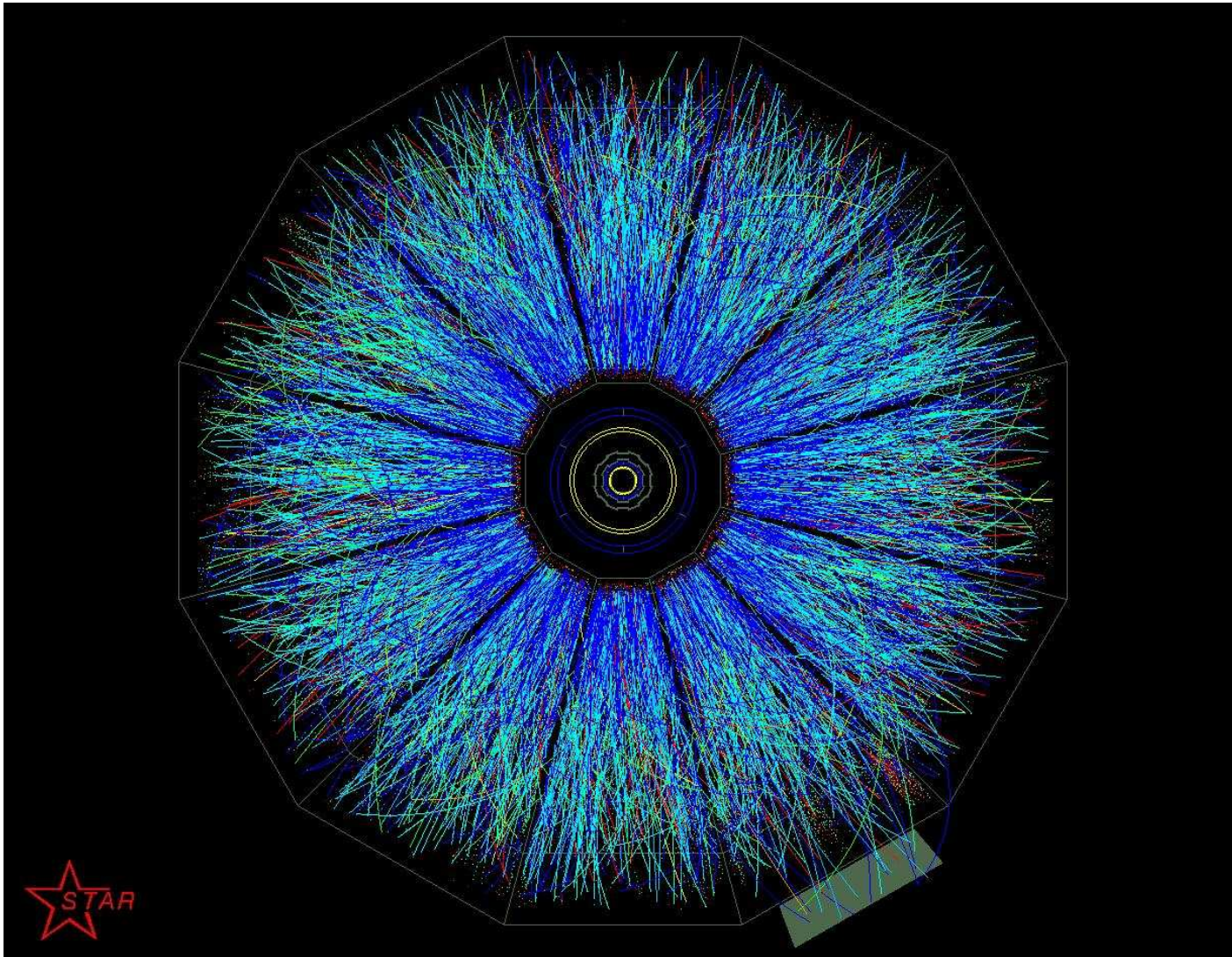
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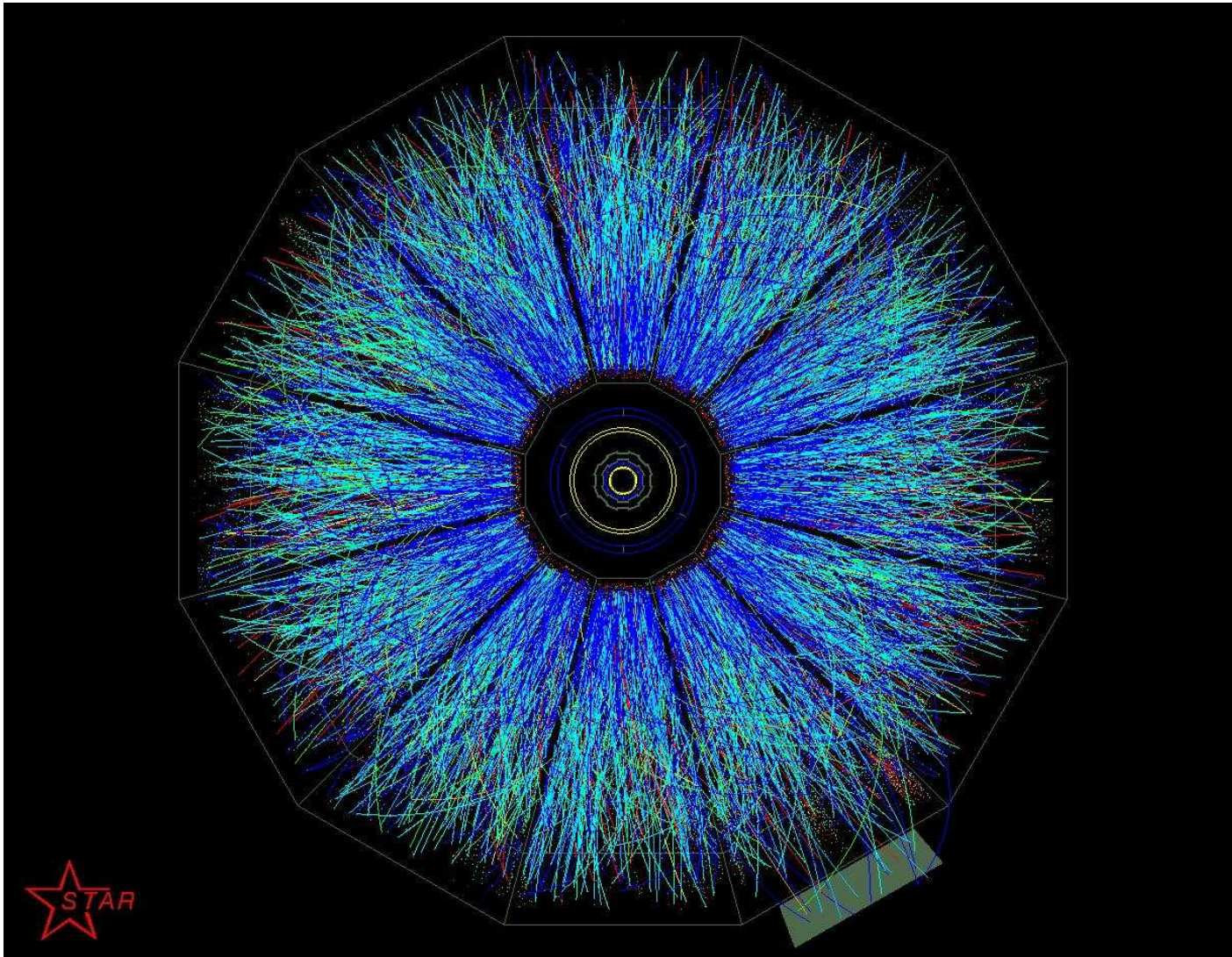
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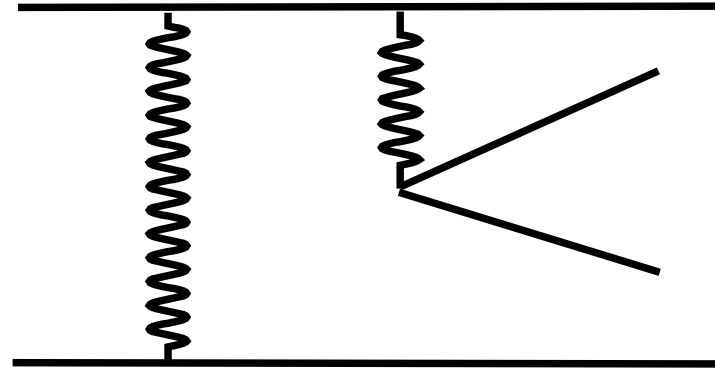
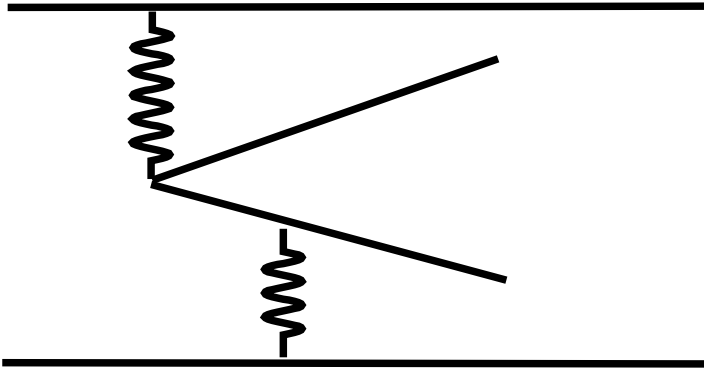
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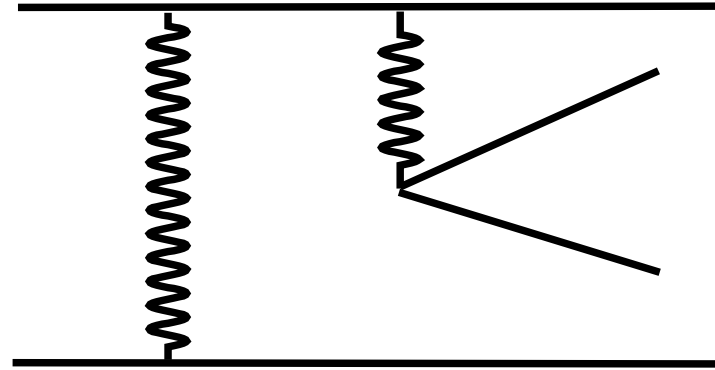
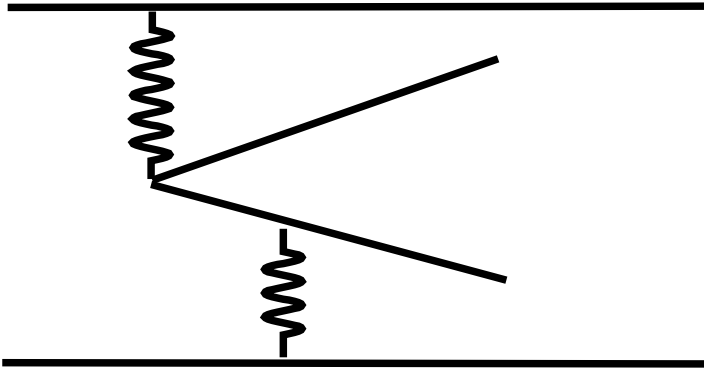
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4 A perfect theory approach may work in QED/electro-weak realm...

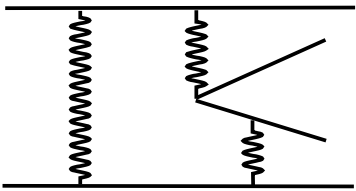


Expansion in powers of $\alpha \approx 1/137$, higher orders matter less...

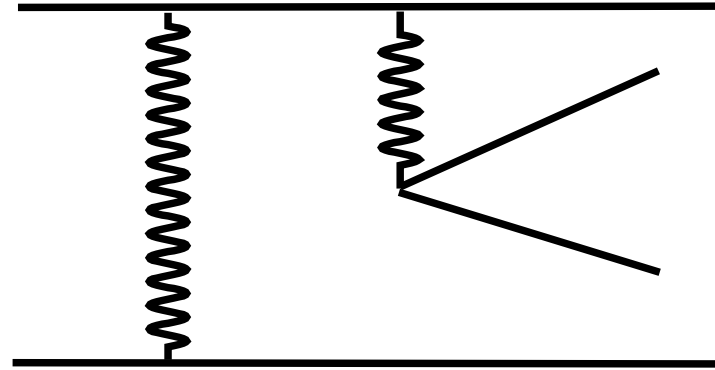
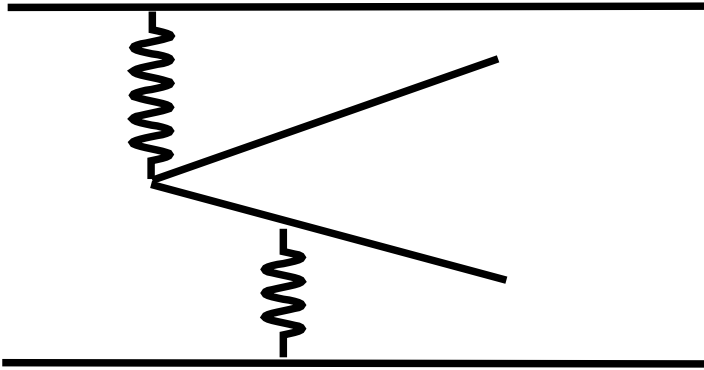
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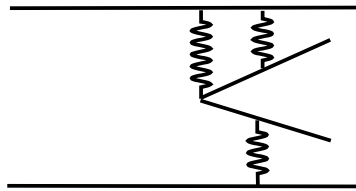
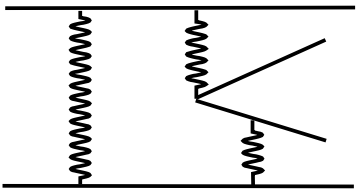
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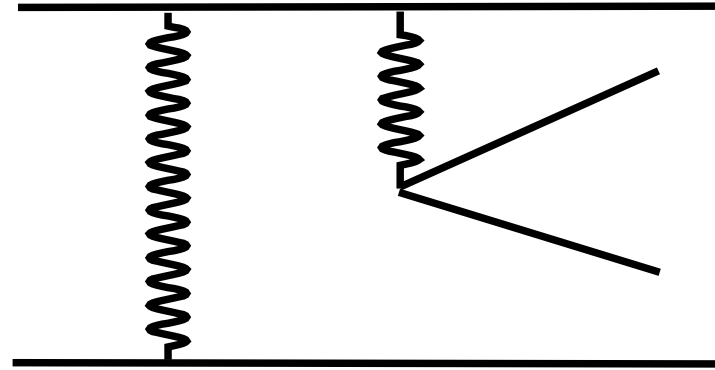
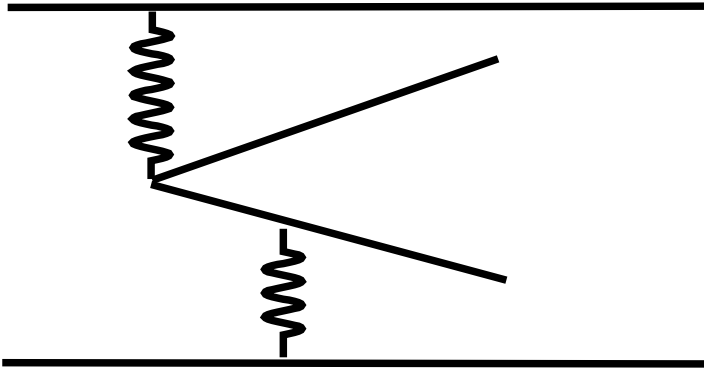
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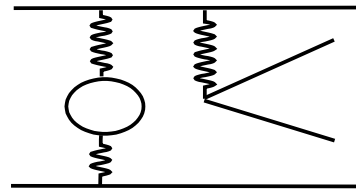
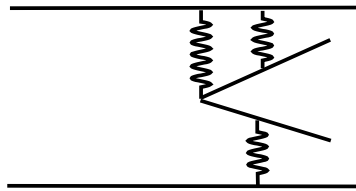
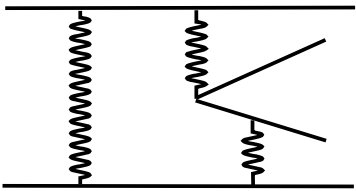
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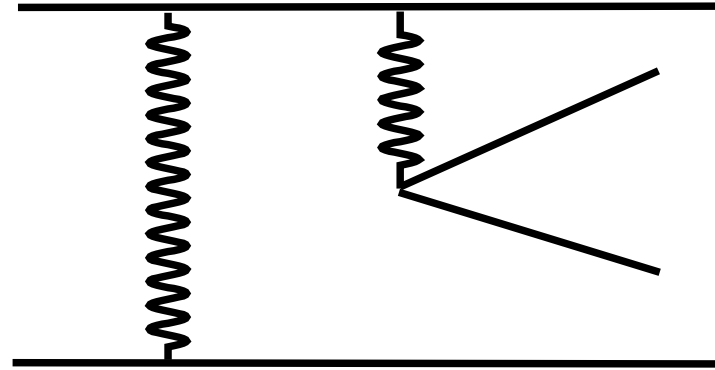
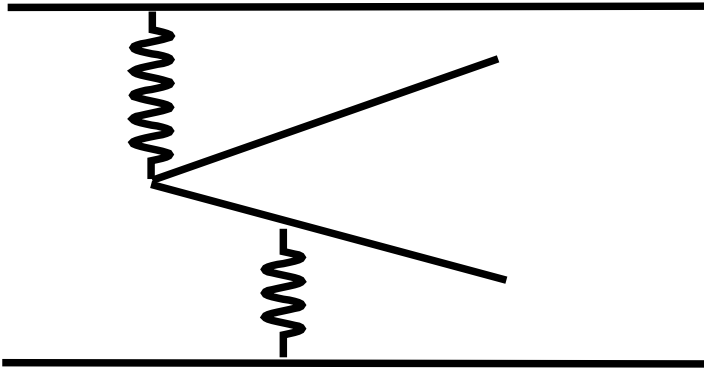
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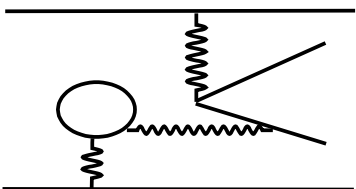
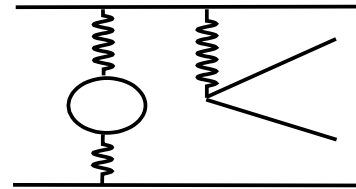
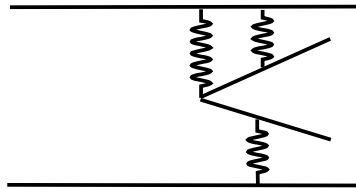
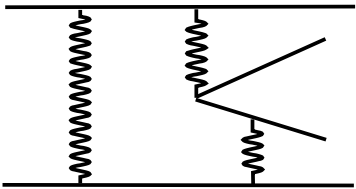
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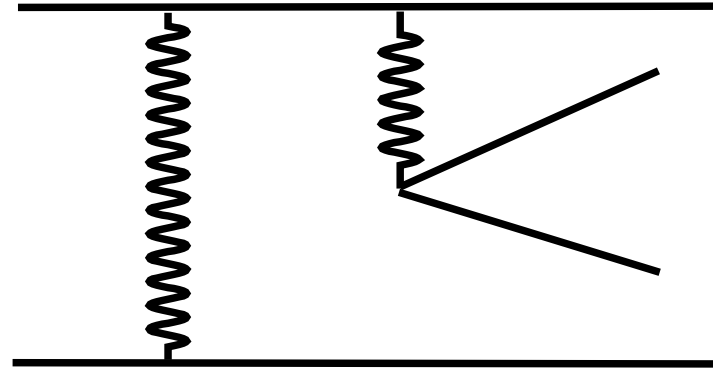
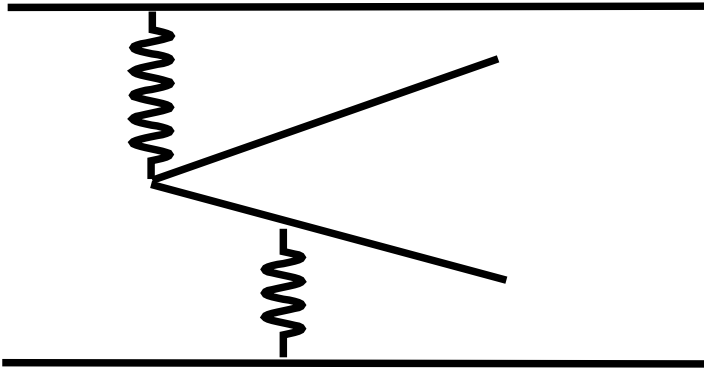
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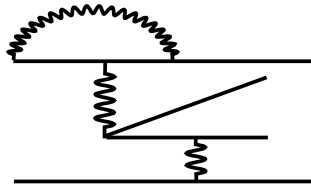
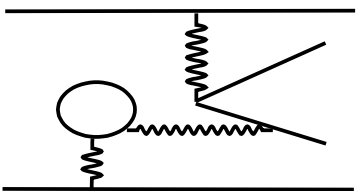
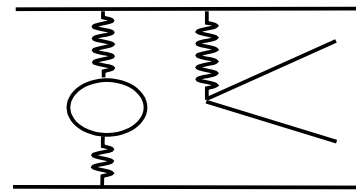
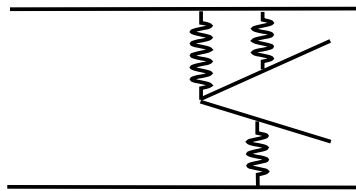
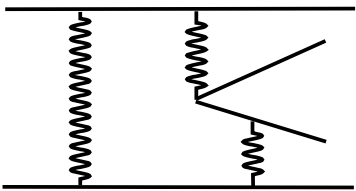
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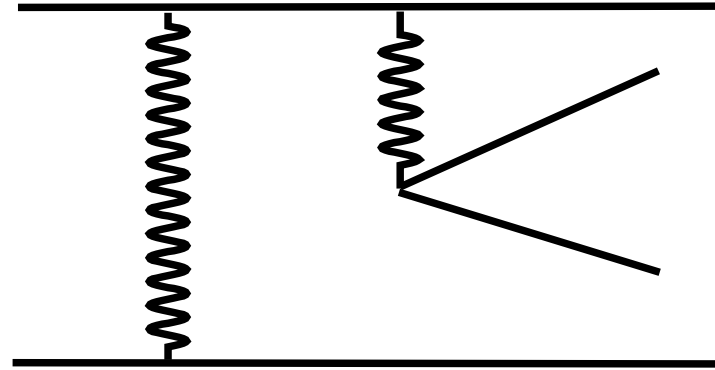
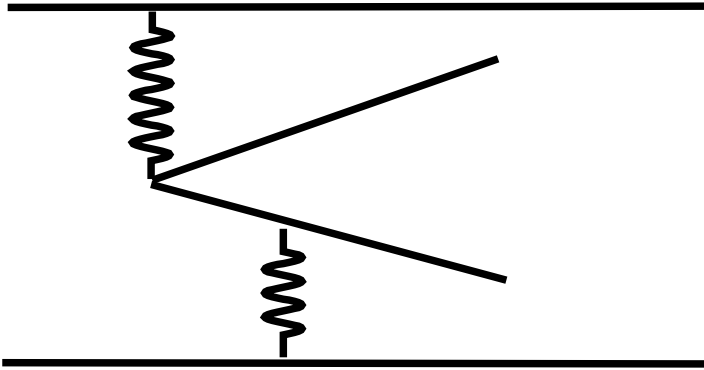
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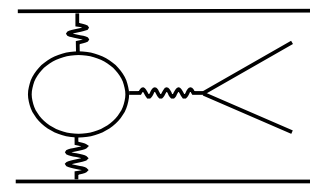
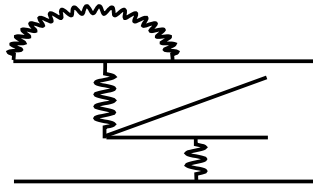
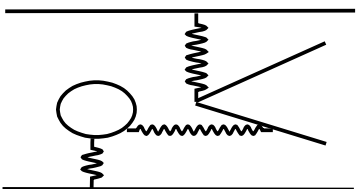
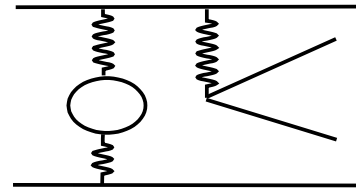
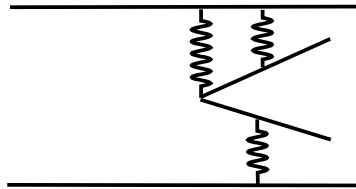
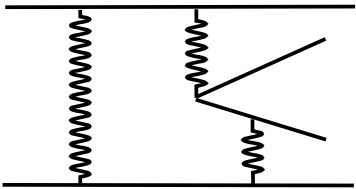
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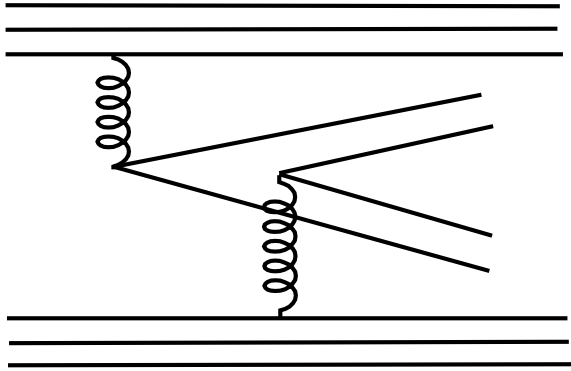


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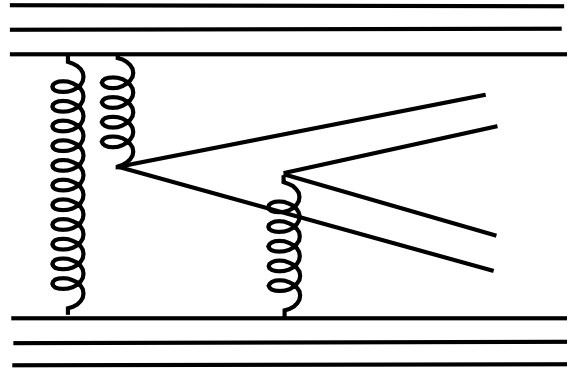
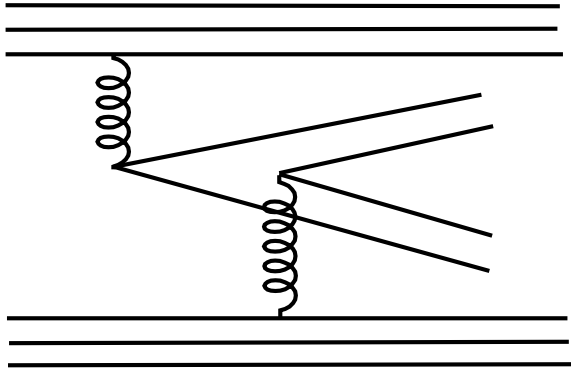


5 Coming to QCD...

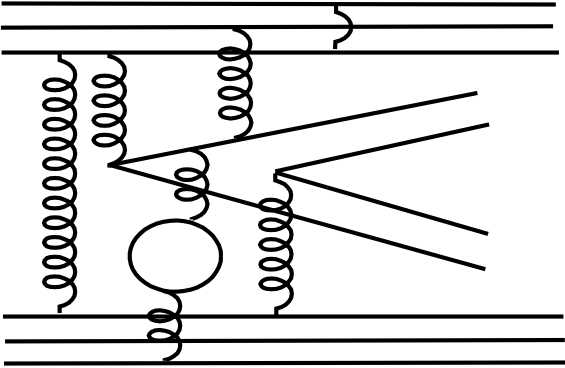
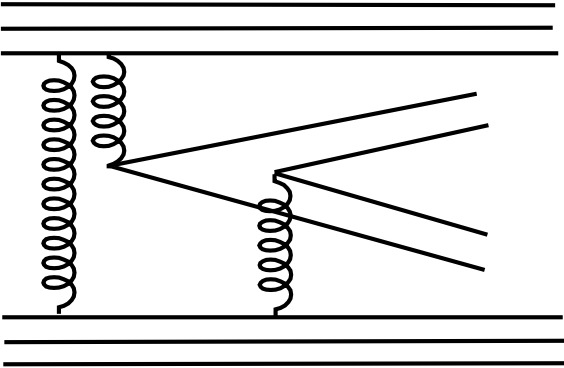
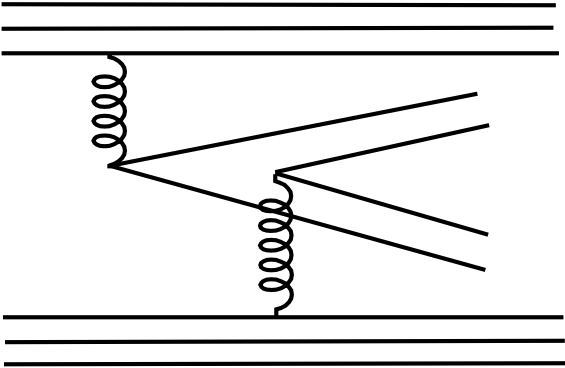
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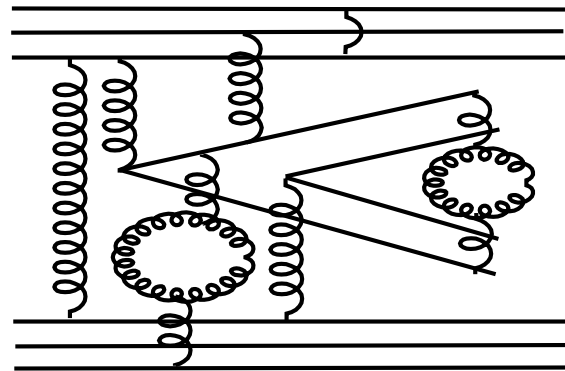
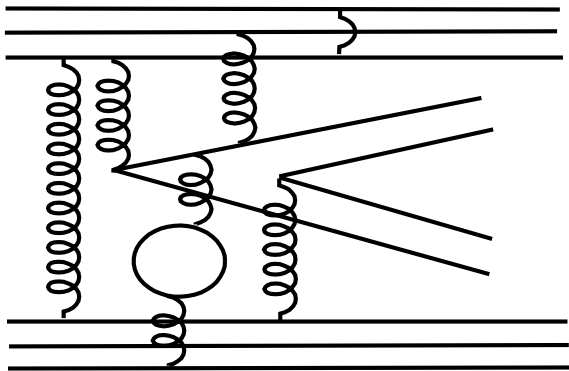
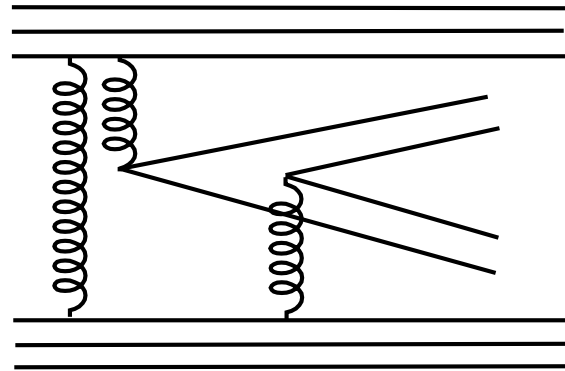
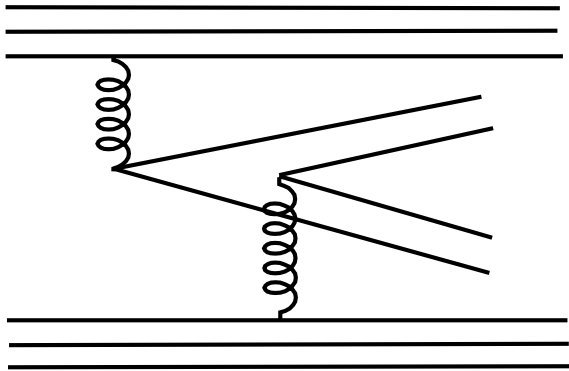
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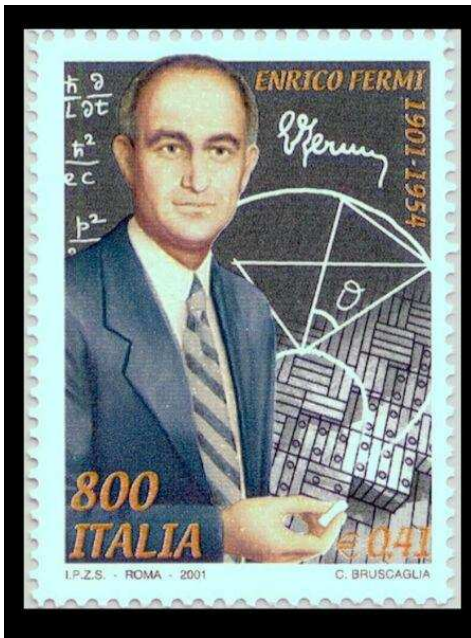


Hard to get matrix elements:

- with $\alpha_S \sim 1$, comparable contributions in all orders, series converge slowly (if at all)

6 E.Fermi – an extreme view: forget about matrix elements!

E.Fermi, "High Energy Nuclear Events", Progr. Theor. Phys. 5, No.4, 1950
(Yukawa theory, no QCD!)

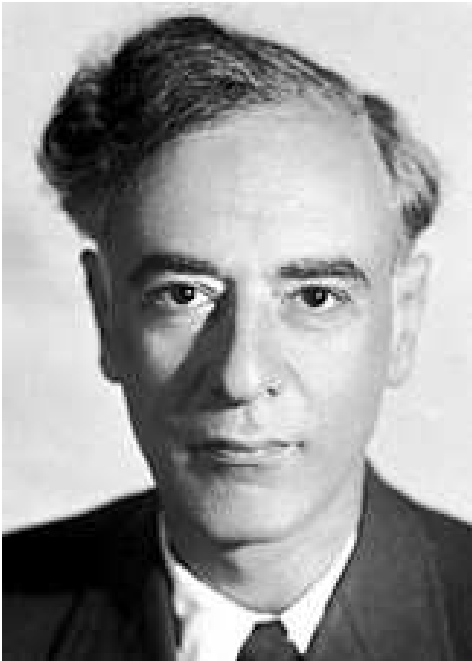


"When two nucleons collide with very great energy in their center of mass system this energy will be suddenly released in a small volume surrounding the two nucleons. $\langle \dots \rangle$ Since the interactions of the pion field are strong we may expect that rapidly this energy will be distributed among the various degrees of freedom $\langle \dots \rangle$ according to **statistical laws**. $\langle \dots \rangle$ It is realized that this description of the phenomenon is probably **as extreme**, although in the opposite direction, as is the **perturbation theory approach**."

7 I.Ya.Pomeranchuk (1951), L.D.Landau (1953) – forget about "individual" particles!

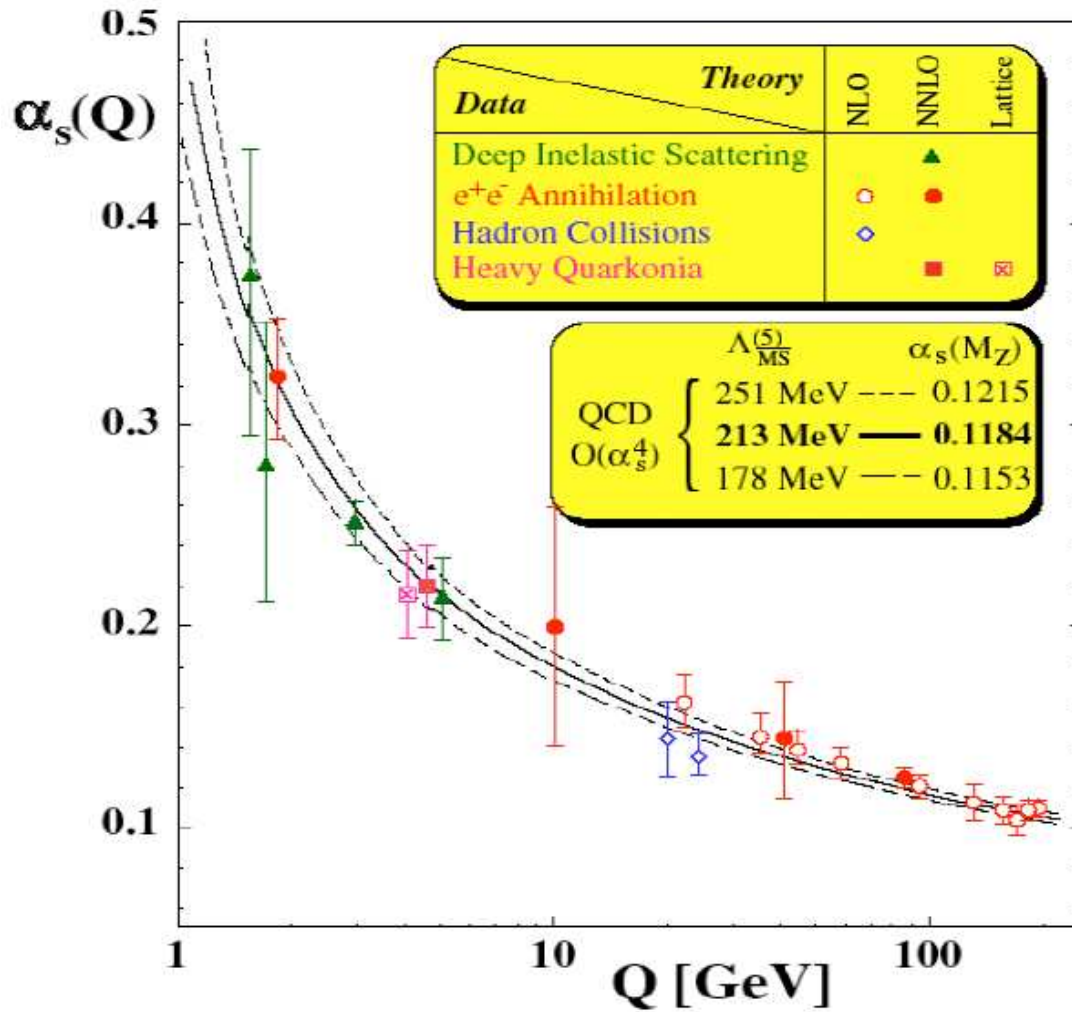


Time evolution of non-viscous hydro with freeze-out.



- "Hot and dense" phase – no "particles", mean free path $\lambda \ll L \Rightarrow$ relativistic hydrodynamics of an ideal (non-viscous and non-heat-conducting) liquid is applicable.
- Free separation at temperature $T \sim m_\pi$ and $\lambda \sim L$, particles reappear.

8 QCD running coupling: infrared slavery and asymptotic freedom

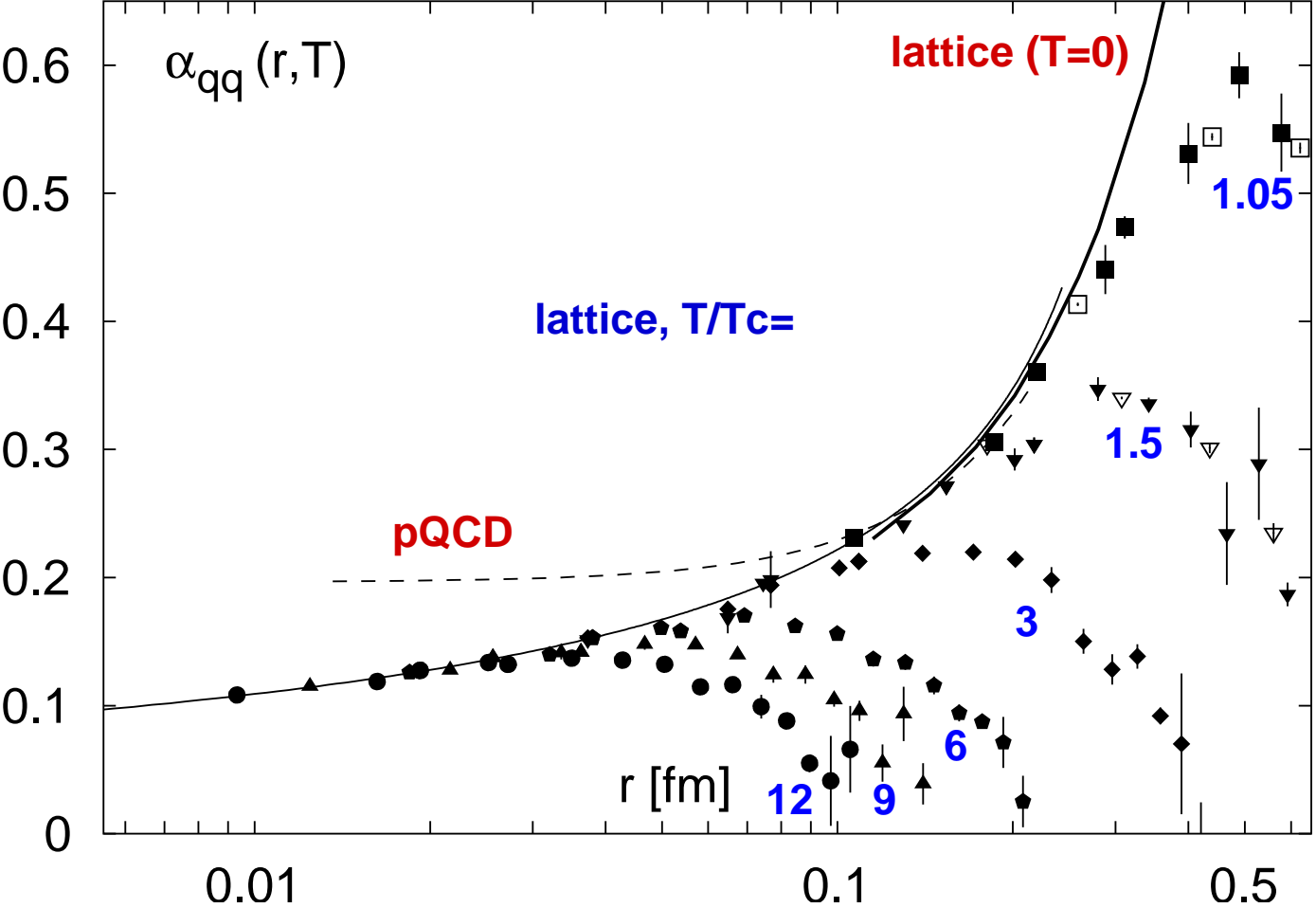


Gross, Politzer, Wilczek

Remember $\Delta p_x \Delta x \approx \hbar$! Asymptotic freedom as seen in particle physics experiments (F.Wilczek's Nobel Lecture 2004)

9 Deconfinement

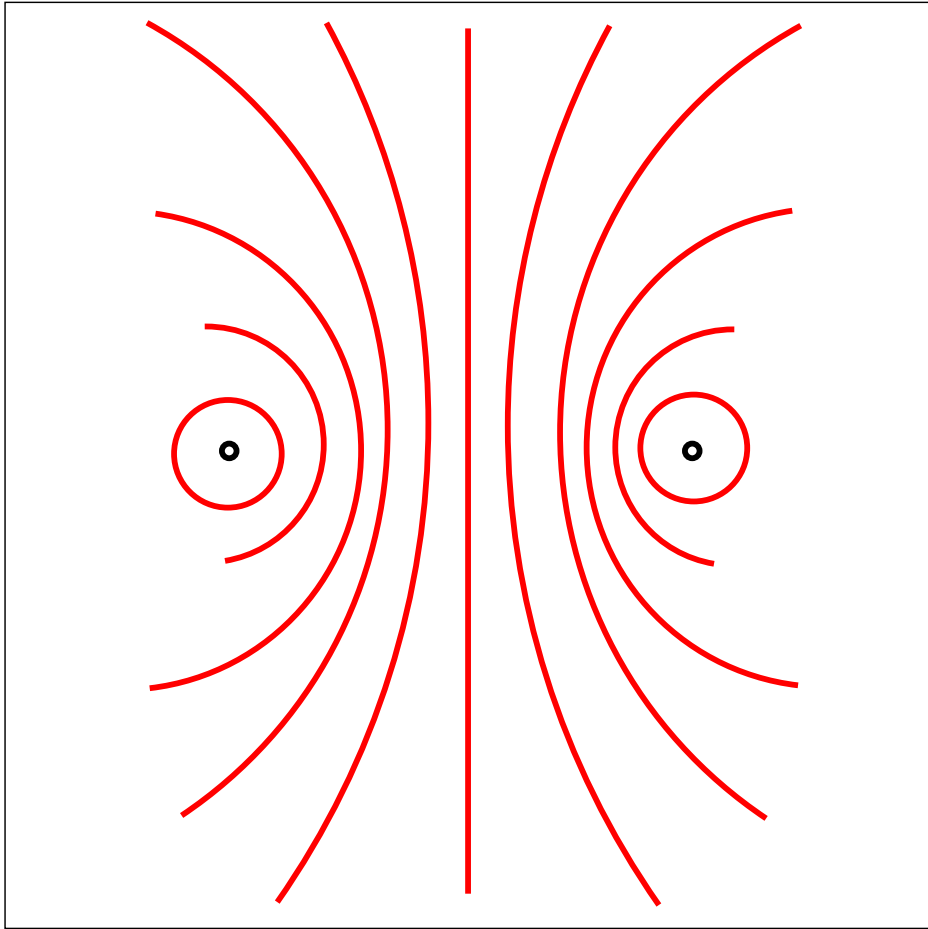
Kaczmarek, Karsch, Zantow, Petreczky PRD70(074505) 2004



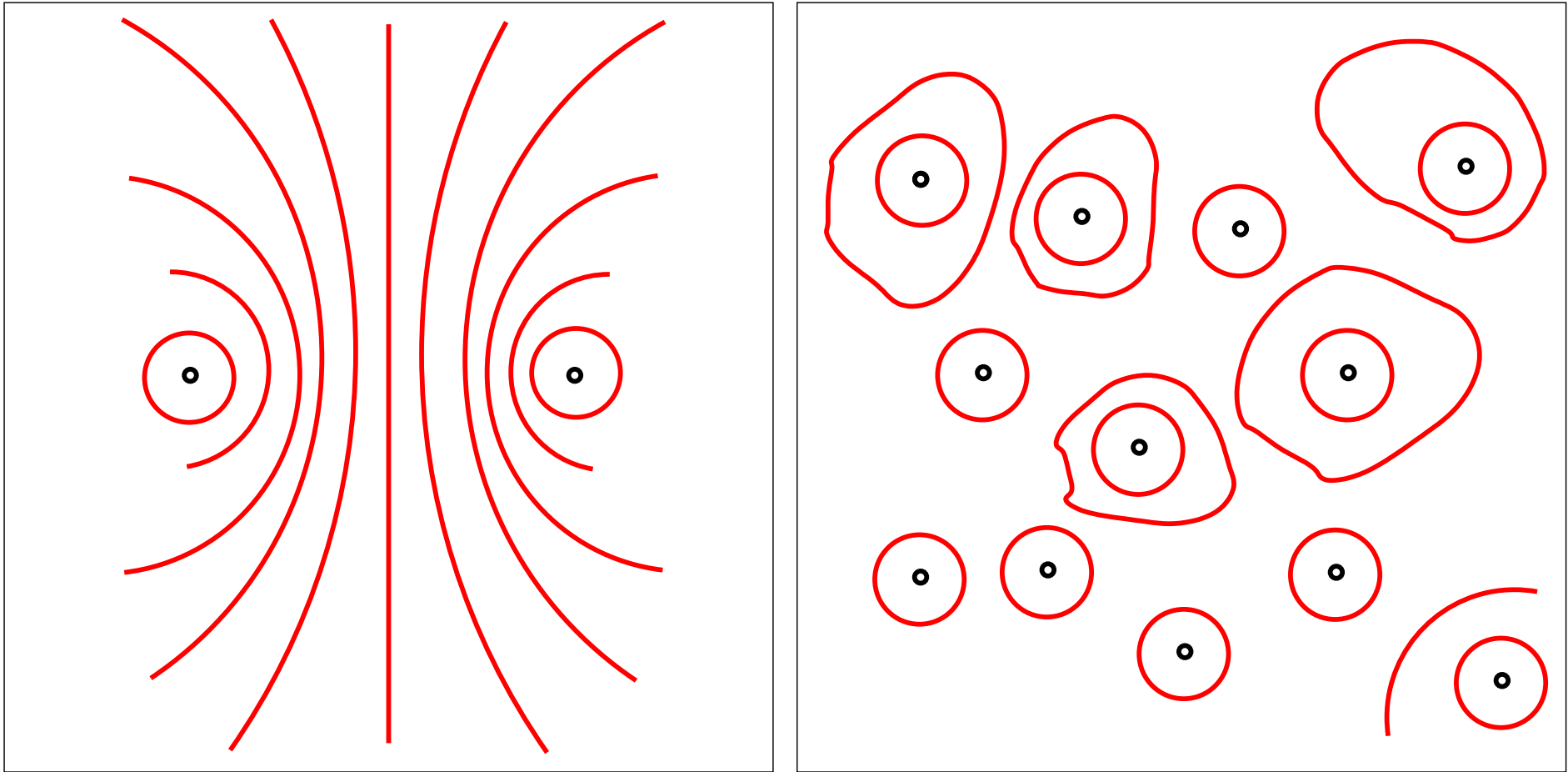
QCD running coupling at $T > T_c$, showing screening of strong force, albeit at relatively large distances.

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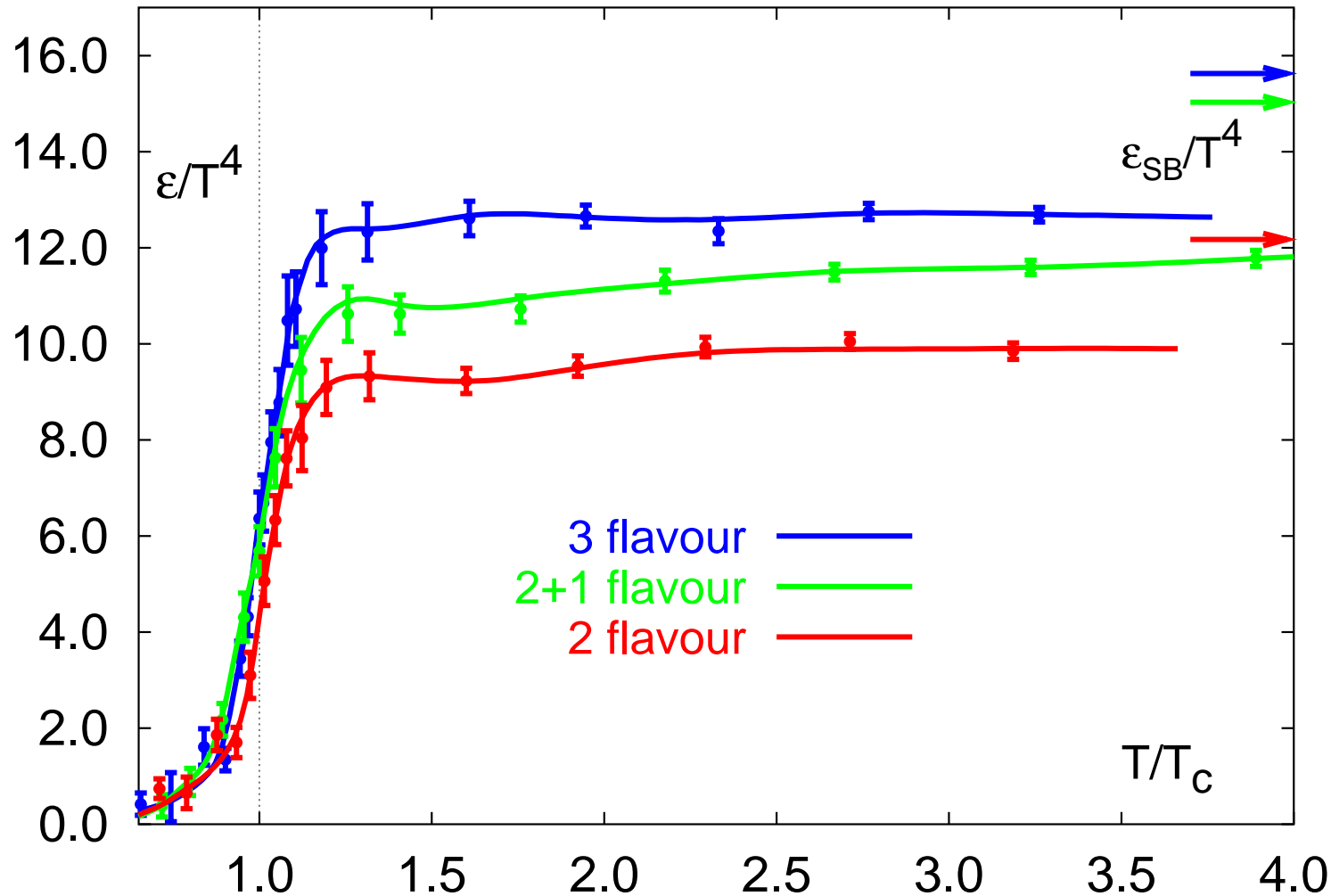


10 A more intuitive picture...



Screening leads to deconfinement at high density or temperature. Analogous to Debye screening in ordinary plasma, there is r_D .

11 Phase transition \Rightarrow increase in the number of degrees of freedom, EOS change



F. Karsch Lect. Notes Phys. 583 (2002) 209 Pressure becomes excessive after phase transition back to ordinary matter \Rightarrow fireball may "explode".

12 Approaches in theory

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- "Freedom": test the nature of the medium by falsifying perturbative predictions. High p_t , "jet tomography", photons, leptons. Work with a subset of specific particles or even expect QGP itself to be "asymptotically free" \Rightarrow perturbative.

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- **Lattice QCD** is not considered a "paradigm", it's heavy artillery, ultima ratio regum

13 Coming next...

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- Perturbative diagnostics
 - high p_t spectra (hadrons, γ)
 - charm

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- Experimental strategies
- Perturbative diagnostics
 - high p_t spectra (hadrons, γ)
 - charm
- Quasi-classical diagnostics
 - flow (hydro)
 - mini-jets in the medium (dissipation ?!)
 - hadro-chemistry

14 Strategy dilemmas in experiment

identify many particle IDs including leptons and gamma

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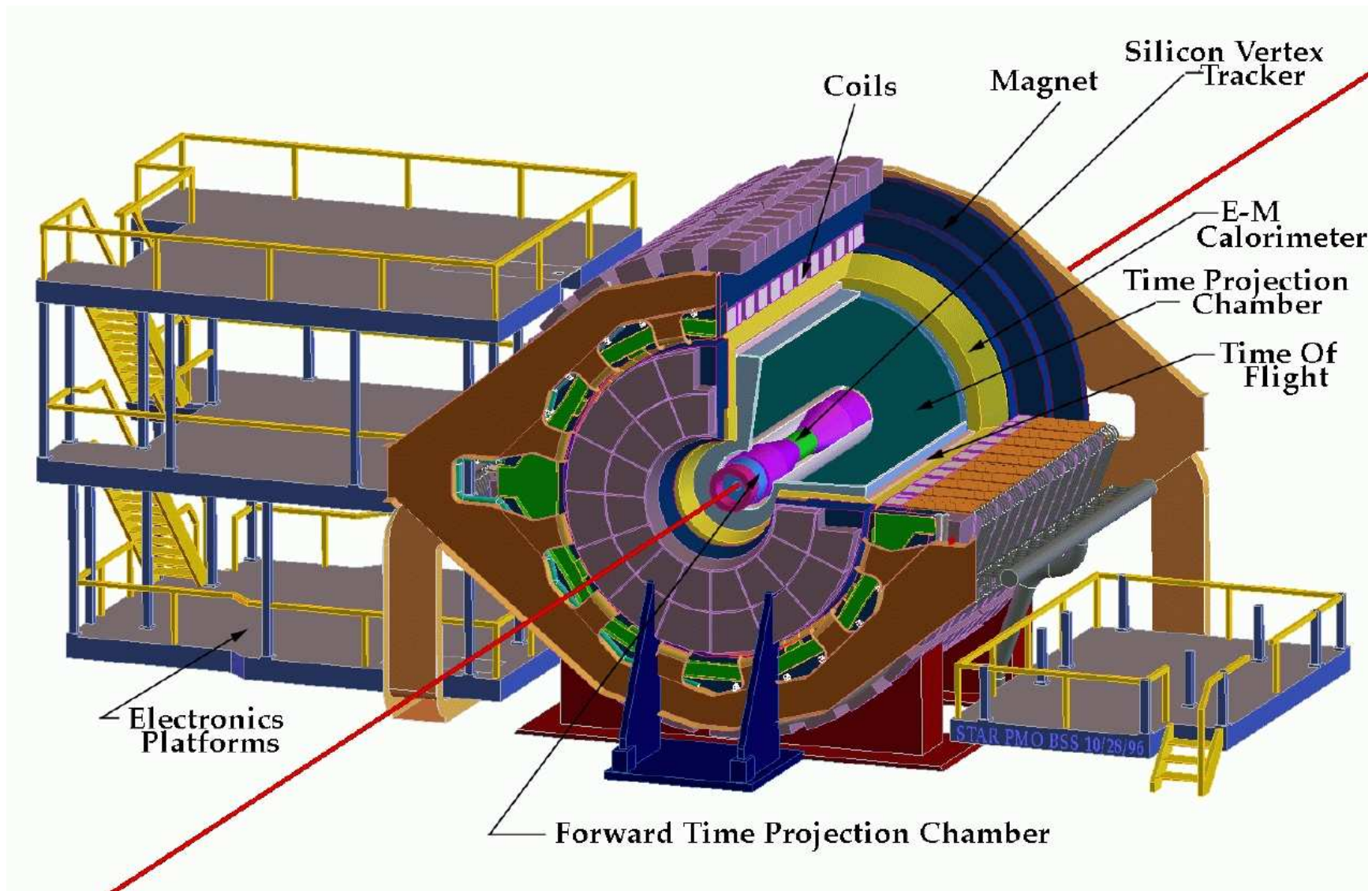
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15 STAR — subsystems



recent Kent contributions: EEMC, ZDC SMD, computing infrastructure

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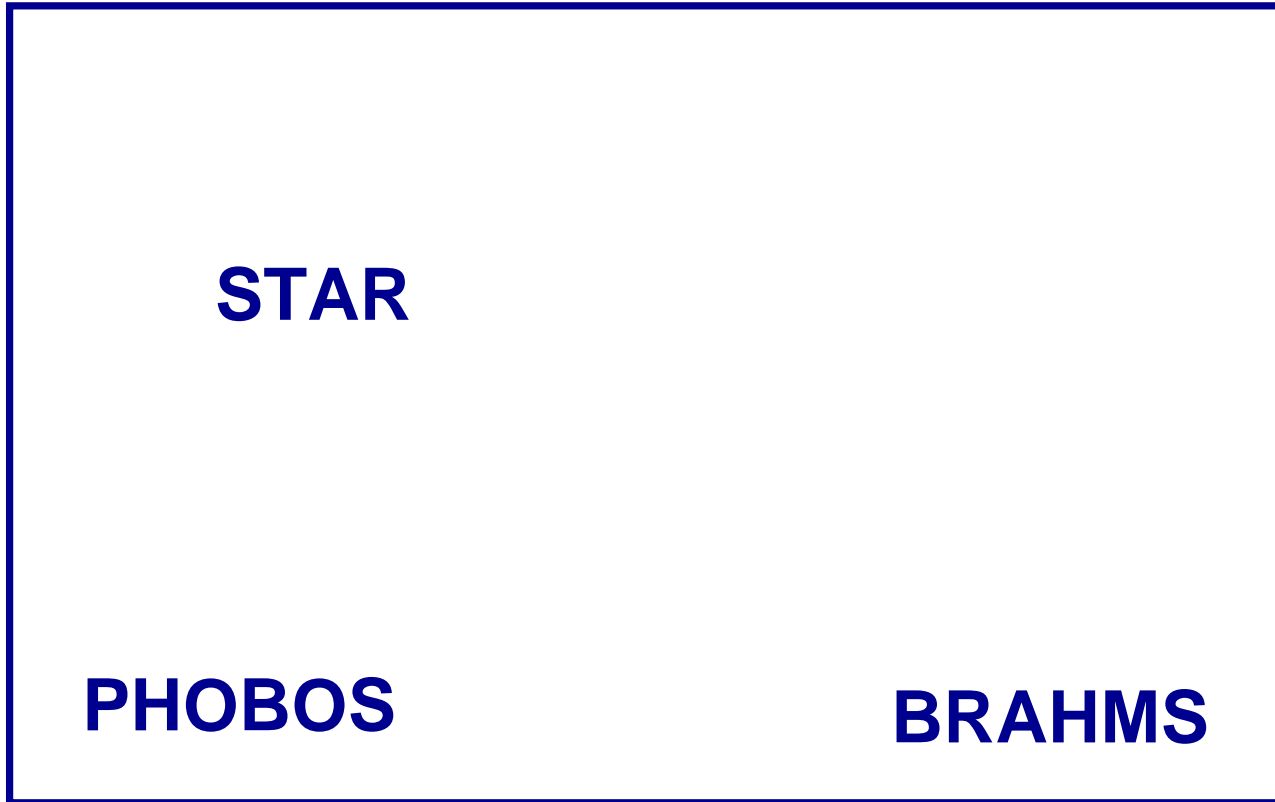
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TPC is the key; relatively infrequent large events

17 STAR — growth

Concept for an Experiment on Particle and Jet Production at Midrapidity

J.W. Harris,⁷ M. Bloomer,⁷ P. Brady,¹ J. Carroll,² S.I. Chase,⁷ W. Christie,⁷
J. Cramer,¹² E. Friedlander,⁷ D. Greiner,⁷ C. Gruhn,⁷ M. Gyulassy,⁷ T. Hallman,⁴
E. Hjort,¹⁰ G. Igo,² P. Jacobs,⁷ K. Kadija,¹³ D. Keane,⁵ L. Madansky,⁴ C. Naudet,⁷
D. Nygren,⁷ G. Odyniec,⁷ D. Olson,⁷ G. Paic,¹³ A. Poskanzer,⁷ G. Rai,⁷ H.G. Ritter,⁷
R. Scharenberg,¹⁰ L.S. Schroeder,⁷ P. Seidl,⁷ P. Seyboth,⁸ D. Shy,⁷ R. Stock,³
T.J. M. Symons,⁷ L. Teitelbaum,⁷ M.L. Tincknell,⁹ H. van Hecke,⁶ X.N. Wang,⁷
R. Welsh,⁴ W. Wenzel,⁷ H. Wieman,⁷ and K.L. Wolf¹¹

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 - 5 Kent State University, Kent Ohio 44242, U.S.A.
 - 6 Los Alamos National Laboratory, Los Alamos, New Mexico, U.S.A.
 - 7 Lawrence Berkeley Laboratory, Berkeley, California 94720, U.S.A.
 - 8 Max Planck Institute for Physics, Munich, West Germany
 - 9 Oak Ridge National Laboratory, Oak Ridge, Tennessee, U.S.A.
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LBL-29488

Concept for an Experiment on Particle and Jet Production at Midrapidity

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Scaling Properties of Hyperon Production in Au+Au Collisions at $\sqrt{s_{NN}} = 200$ GeV

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Diwan,¹⁶ S.M. Dogra,¹⁶ W.J. Dong,¹⁶ X. Dong,¹⁶ J.E. Draper,¹⁶ F. Du,¹⁶ V.B. Duin,¹⁶ J.C. Dunlop,¹⁶ M.R. Datta Munshi,¹⁶ Y. Ekaev,¹⁶ W.R. Edwards,¹⁶ L.G. Eklund,¹⁶ V. Emelianov,¹⁶ J. Engelage,¹⁶ G. Eppley,¹⁶ B. Erman,¹⁶ M. Estabrook,¹⁶ P. Fackin,¹⁶ R. Fatemi,¹⁶ J. Fedorin,¹⁶ K. Filimonov,¹⁶ P. Filin,¹⁶ E. Finch,¹⁶ V. Finn,¹⁶ Y. Fisyak,¹⁶ J. Fu,¹⁶ C.A. Gagliardi,¹⁶ L. Gallard,¹⁶ J. Gao,¹⁶ M.S. Gauger,¹⁶ Y. Gaurichvane,¹⁶ P. Ghosh,¹⁶ J.E. Gomez,¹⁶ Y.G. Gong,¹⁶ R. Goswami,¹⁶ D. Greiner,¹⁶ S.M. Grunin,¹⁶ K.S.P. Guzman,¹⁶ Y. Guo,¹⁶ N. Gupta,¹⁶ T.D. Guizoneau,¹⁶ B. Haag,¹⁶ T.J. Hallman,⁴ A. Hamad,¹⁶ F.W. Harin,¹⁶ W. He,¹⁶ M. Heine,¹⁶ T.W. Henry,¹⁶ S. Hogganwood,¹⁶ R. Hoggarty,¹⁶ A. Hirsch,¹⁶ E. Hjort,¹⁰ G.W. Hoffmann,¹⁶ M.J. Horne,¹⁶ H.Z. Huang,¹⁶ S.L. Huang,¹⁶ E.W. Hughes,¹⁶ T.J. Humanic,¹⁶ G. Igo,² P. Jacak,¹⁶ W.W. Jacobs,⁷ P. Jale,¹⁶ H. Jiang,¹⁶ P.G. Jones,¹⁶ E.G. Judd,¹⁶ S. Kabanov,¹⁶ K. Kang,¹⁶ J. Kapitan,¹⁶ M. Kapcia,¹⁶ D. Keane,⁵ A. Kechin,¹⁶ V.V. Khlopov,¹⁶ R.J. Kim,¹⁶ J. Kiryluk,¹⁶ A. Kisel,¹⁶ E.M. Kislov,¹⁶ S.R. Klein,¹⁶ D.D. Koethe,¹⁶ T. Kollegger,¹⁶ M. Kopylov,¹⁶ L. Kouchard,¹⁶ V. Kouchal,¹⁶ K.L. Kovalek,¹⁶ M. Krauss,¹⁶ P. Kravtsov,¹⁶ V.I. Kravtsov,¹⁶ K. Krueger,¹⁶ C. Kuhn,¹⁶ A.I. Kulkov,¹⁶ A. Kumar,¹⁶ A.A. Kurepin,¹⁶ M.A.C. Lacerda,¹⁶ J.M. Lando,¹⁶ S. Lange,¹⁶ S. LePoutre,¹⁶ F. Lee,¹⁶ J. Leiser,¹⁶ S.A. Leisher,¹⁶ B. Letichev,¹⁶ C.H. Lee,¹⁶ S. Lehar,¹⁶ M.J. LeVine,¹⁶ C. Li,¹⁶ Q. Li,¹⁶ Y. Li,¹⁶ G. Lin,¹⁶ S.J. Lindenbaum,¹⁶ S.A. Liu,¹⁶ F. Liu,¹⁶ H. Liu,¹⁶ L. Liu,¹⁶ Z. Liu,¹⁶ T. Ljubicic,¹⁶ W. J. Llope,¹⁶ R. Long,¹⁶ R.S. Longacre,¹⁶ M. Lopez-Noreña,¹⁶ W.A. Love,¹⁶ Y. Lu,¹⁶ T. Ludwik,¹⁶ D. Lynn,¹⁶ G.L. Ma,¹⁶ J.G. Ma,¹⁶ Y.G. Ma,¹⁶ D. Magagnoli,¹⁶ D.P. Malabarba,¹⁶ R. Majka,¹⁶ L.K. Mangotra,¹⁶ R. Marwede,¹⁶ S. Margutti,¹⁶ C. Markert,¹⁶ B. Marot,¹⁶ H.S. Matis,¹⁶ Y.A. Matushkov,¹⁶ C.J. McDiarmid,¹⁶ T.K. Mehta,¹⁶ N.A. McIntosh,¹⁶ A. Meshkin,¹⁶ M.L. Miller,¹⁶ N.G. Mishev,¹⁶ S. Mischarewski,¹⁶ C. Misner,¹⁶ A. Mischke,¹⁶ D.K. Mishra,¹⁶ J. Mitchell,¹⁶ M. Miñano,¹⁶ L. Miskewicz,¹⁶ C.P. Moore,¹⁶ D.A. Moore,¹⁶ M.G. Muchnik,¹⁶ B.K. Nandi,¹⁶ C. Nattar,¹⁶ T.K. Nayak,¹⁶ I.M. Nigmatov,¹⁶ N. Nigmatov,¹⁶ V.A. Nisimov,¹⁶ L.V. Nizich,¹⁶ S.B. Nurushev,¹⁶ G. Odyniec,⁷ A. Ogawa,¹⁶ V. Okunov,¹⁶ M. Odrzyasik,¹⁶ D. Olson,⁷ M. Pacha,¹⁶ S.K. Pal,¹⁶ Y. Panferov,¹⁶ S.Y. Panikell,¹⁶ A.I. Pankratov,¹⁶ T. Panzini,¹⁶ V. Pereslavskiy,¹⁶ V. Pervez,¹⁶ S. Piatek,¹⁶ J. Pochay,¹⁶ S.C. Poon,¹⁶ R. Picha,¹⁶ M. Planinic,¹⁶ J. Piatek,¹⁶ N. Pijak,¹⁶ N. Pivko,¹⁶ J. Pister,¹⁶ A.M. Piskunov,¹⁶ M. Piskunov,¹⁶ E. Piskunov,¹⁶ B.V.K.S. Prasad,¹⁶ D. Prindle,¹⁶ C. Prunus,¹⁶ J. Putschke,¹⁶ G. Raban,¹⁶ B. Radicevic,¹⁶ R.E. Raabe,¹⁶ S.Y. Raabe,¹⁶ S.J. Rattner,¹⁶ D. Rautava,¹⁶ F. Reiter,¹⁶ A. B. Ridge,¹⁶ H.G. Ritter,⁷ J.B. Roberts,¹⁶ O.V. Ropchuk,¹⁶ J.L. Romero,¹⁶ A. Rose,¹⁶ C. Roy,¹⁶ L. Ruan,¹⁶ M.J. Ruscak,¹⁶ R. Sahoo,¹⁶ I. Sakrejda,¹⁶ S. Saito,¹⁶ J. Sandhu,¹⁶ M.J. Sarantsev,¹⁶ P.S. Sathya,¹⁶ J. Schambach,¹⁶ R.P. Scharenberg,¹⁰ N. Schmitz,¹⁶ K. Schmitz,¹⁶ J. Schuch,¹⁶ I. Selyuzhko,¹⁶ P. Seyboth,⁸ A. Shabanov,¹⁶ F. Shalunov,¹⁶ M. Shao,¹⁶ M. Sharma,¹⁶ W.Q. Shen,¹⁶ S.S. Shennikoff,¹⁶ E. Shkhrin,¹⁶ F. Simon,¹⁶ R. Simons,¹⁶ S. Simons,¹⁶ G. Sissakian,¹⁶ G. Sioli,¹⁶ P. Sitar,¹⁶ J. Sittler,¹⁶ J. Sjostrand,¹⁶ H.M. Spinka,¹⁶ B. Srivastava,¹⁶ A. Stadnik,¹⁶ T.D.S. Stasidis,¹⁶ B. Stock,¹⁶ A. Stolpovsky,¹⁶ M. Strikhanov,¹⁶ R. Strongfellow,¹⁶ A.P. Svanlov,¹⁶ E. Suvorov,¹⁶ M. Sundeen,¹⁶ Z. Sun,¹⁶ B. Surrow,¹⁶ M. Swasey,¹⁶ T.J. Symons,⁷ A. Szabo,¹⁶ Z. Tang,¹⁶ A. T. Tikhonchuk,¹⁶ A.H. Tang,¹⁶ T. Tarnovsky,¹⁶ D. Thies,¹⁶ J.H. Thomas,¹⁶ A.R. Timmins,¹⁶ S. Timonen,¹⁶ M. Tokarev,¹⁶ T.A. Trainor,¹⁶ B. Trzeciak,¹⁶ R.E. Tribble,¹⁶ C.D. Tsai,¹⁶ J. Urey,¹⁶ J. Ullrich,¹⁶ D.G. Underwood,¹⁶ G. Van Buren,¹⁶ N. van der Kolk,¹⁶ M. van Leeuwen,¹⁶ A.M. VandeMolen,¹⁶ B. Varet,¹⁶ C.M. VandenWyngaert,¹⁶ A.N. VandenVliet,¹⁶ B. Vorer,¹⁶ S.E. Vugot,¹⁶ V.P. Vysotskiy,¹⁶ S. Vukobratovic,¹⁶ S.A. Voloshin,¹⁶ W.T. Wong,¹⁶ F. Wang,¹⁶ C. Wang,¹⁶ J.S. Wang,¹⁶ X. Wang,¹⁶ J.W. Wang,¹⁶ J.C. Wang,¹⁶ G.D. Westfall,¹⁶ A. Weisberg,¹⁶ C. White,¹⁶ H. Wieman,⁷ S.W. Wissiak,¹⁶ R. Witt,¹⁶ J. Wood,¹⁶ C. Wu,¹⁶ X. Xu,¹⁶ Q.H. Xu,¹⁶ Z. Xu,¹⁶ P. Vepko,¹⁶ L.K. You,¹⁶ V.I. Yurievich,¹⁶ W. Zhan,¹⁶ H. Zhang,¹⁶ W.M. Zhang,¹⁶ Y. Zhang,¹⁶

Z.P. Zhang,²⁹ Y. Zhan,³⁰ C. Zhang,³⁰ R. Zankovskiy,¹⁶ Y. Zankovskiy,¹⁶ A.N. Zaslavski,¹⁶ and J.X. Zhu,³⁰ (STAR Collaboration)

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- ⁴California Institute of Technology, Pasadena, California 91125
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- ⁶University of California, Davis, California 95616
- ⁷University of California, Los Angeles, California 90024
- ⁸Cornell University, Ithaca, New York 14853
- ⁹Cornell University, Ithaca, New York 14853
- ¹⁰Laboratory for High Energy Physics, Tsinghua University, Beijing, China
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- ¹³Lawrence Berkeley Laboratory, Berkeley, California 94720
- ¹⁴Lawrence Livermore National Laboratory, Livermore, California 94550
- ¹⁵Lawrence Berkeley Laboratory, Berkeley, California 94720
- ¹⁶Lawrence Berkeley Laboratory, Berkeley, California 94720
- ¹⁷Lawrence Berkeley Laboratory, Berkeley, California 94720
- ¹⁸Lawrence Berkeley Laboratory, Berkeley, California 94720
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- ²⁸Lawrence Berkeley Laboratory, Berkeley, California 94720
- ²⁹Lawrence Berkeley Laboratory, Berkeley, California 94720
- ³⁰Lawrence Berkeley Laboratory, Berkeley, California 94720

We present the scaling properties of Λ , Σ , Ξ , and Ξ' meson anti-particles produced at mid-rapidity in Au+Au collisions at RHIC at $\sqrt{s_{NN}} = 200$ GeV. The yield of anti-strange baryons per participant nucleon increases from peripheral to central collision centrality more rapidly than expected to correspond to an increasing strange quark density of matter produced. The value of the strange phase quark occupancy factor, \bar{K} , obtained from a statistical model fit, approaches unity for the most central collisions. We also show that the nuclear modification factors, R_{pA} , of Λ and Ξ are consistent with each other and with that of protons in the same centrality range. This $R_{pA} < 1$ behavior is consistent with a scenario of hadron formation from coexisting quark degrees of freedom through quark recombination or coalescence.

18 STAR — computing challenges

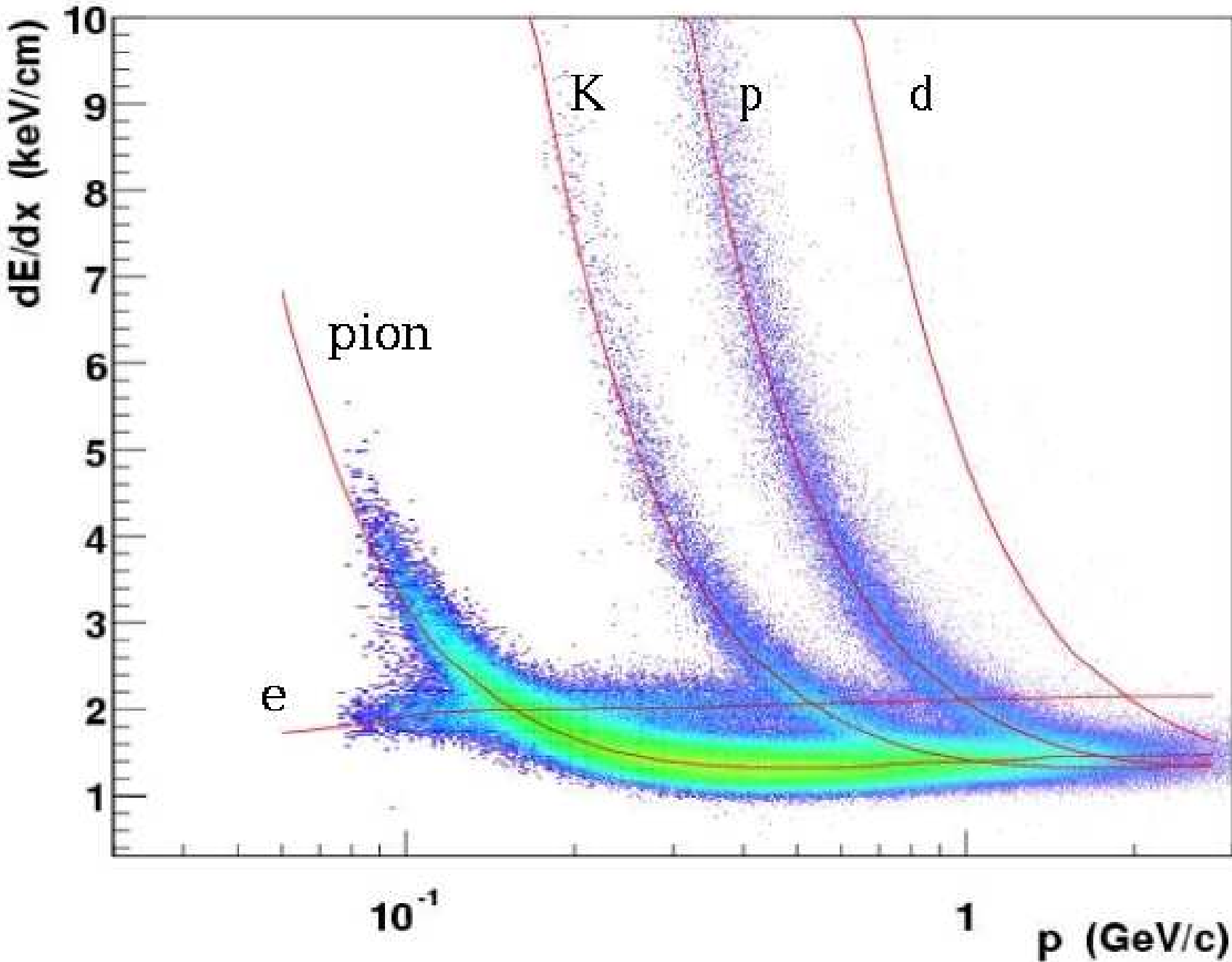
- high volume of data → 10^9 events next year
- distributed resources and users

Solution: distributed computing (grid).

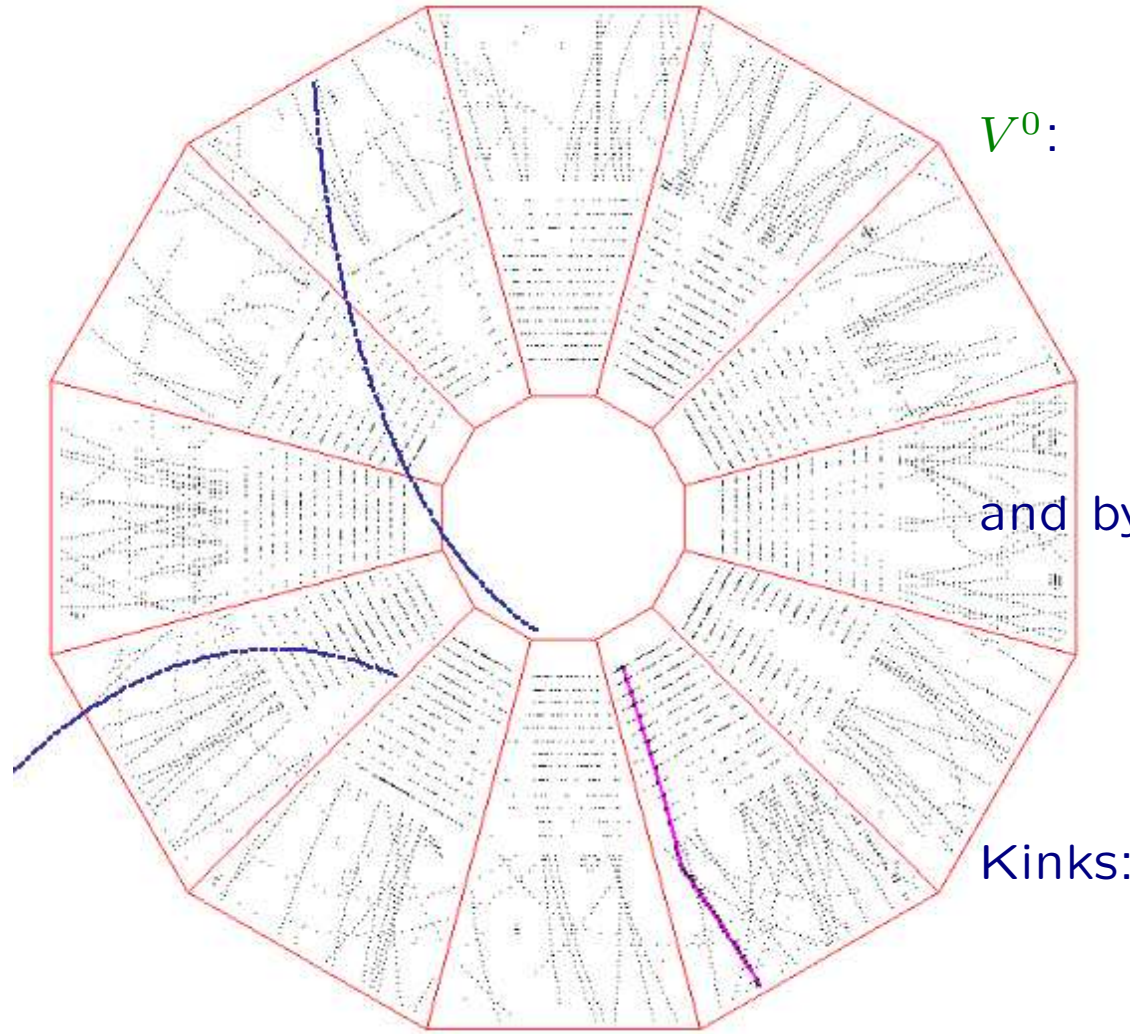
Projects with Kent contribution:

- grid collector (event catalog) – Wei-Ming Zhang
- database API, load-balancer – Mikhail Kopytine

19 STAR — Particle identification — by dE/dx



20 STAR — Particle identification — by topology



$$K^0 \rightarrow \pi^+ \pi^- \quad (1)$$

$$\Lambda \rightarrow p \pi^- \quad (2)$$

$$\bar{\Lambda} \rightarrow \bar{p} \pi^+ \quad (3)$$

$$\Sigma^- \rightarrow \Lambda \pi^- \quad (4)$$

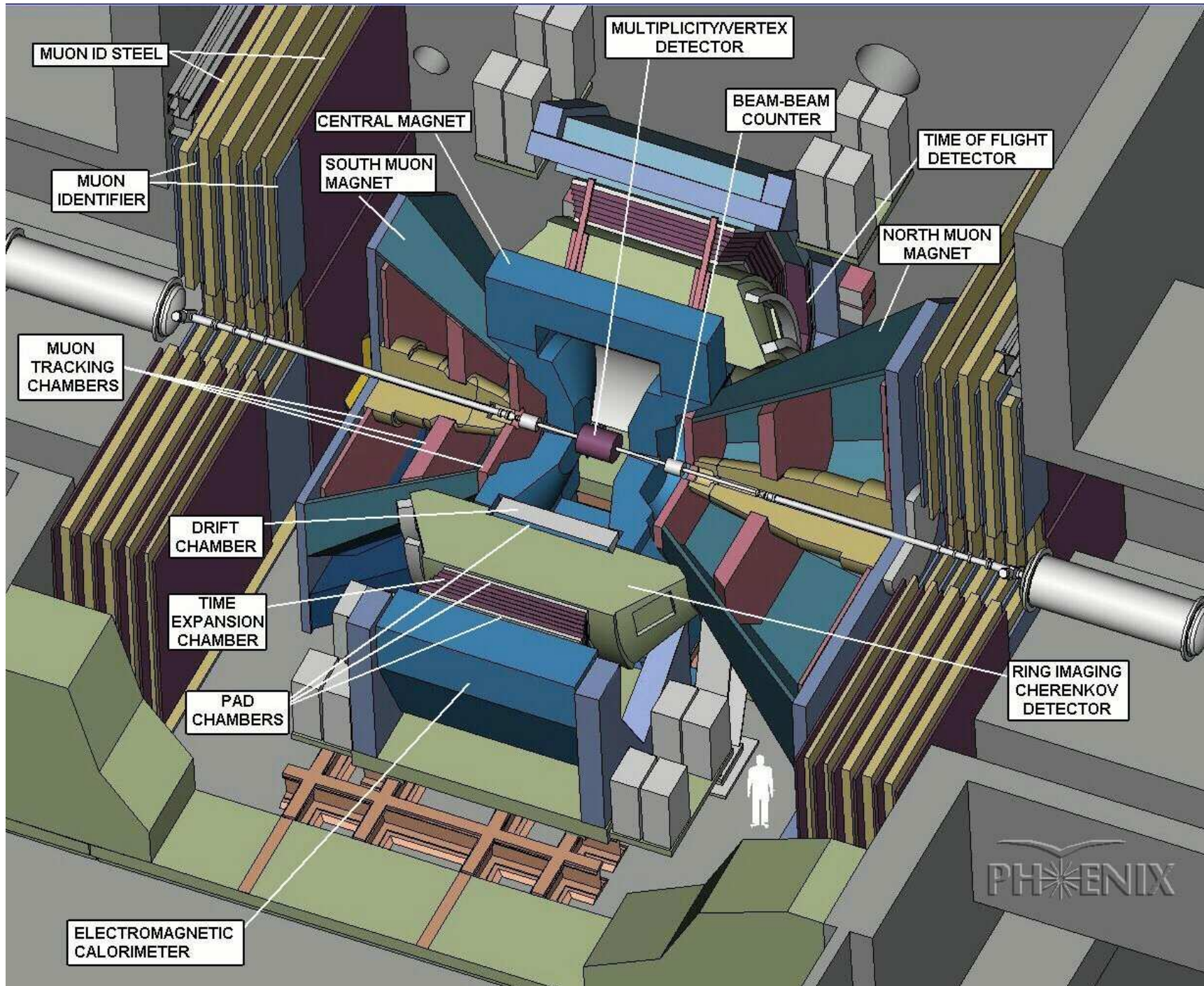
$$\Omega^- \rightarrow \Lambda K^- \quad (5)$$

$$K^\pm \rightarrow \mu^\pm \nu \quad (6)$$

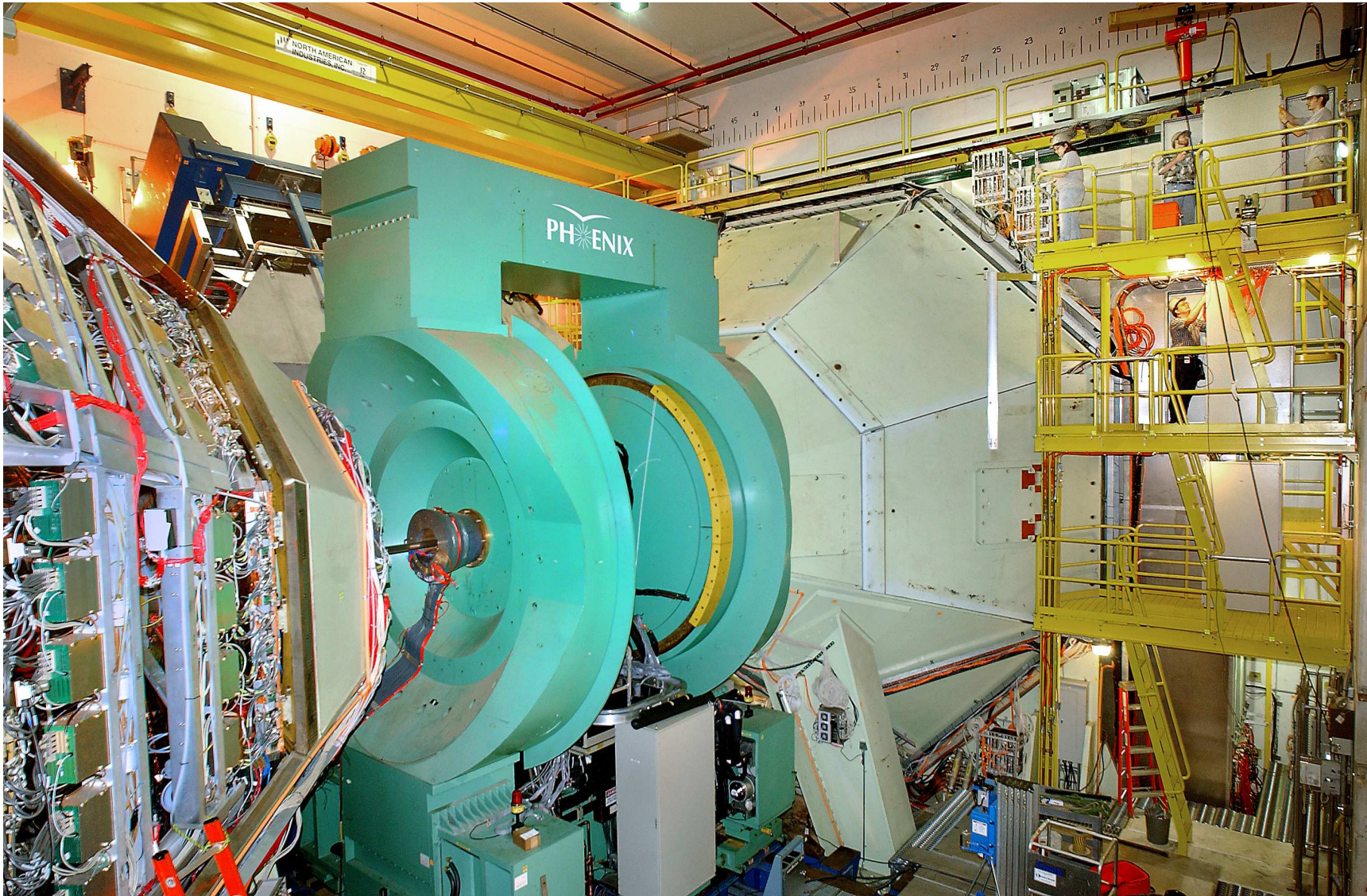
$$K^\pm \rightarrow \pi^\pm \pi^0 \quad (7)$$

About 10% of a central event.

21 PHENIX — subsystems



22 PHENIX — photograph



Accepts $\theta = (90 \pm 20)^\circ$, highly inhomogeneous \vec{B} -field up to 0.8T.

23 PHENIX — strategy

23 PHENIX — strategy

identify many particle IDs including leptons and gamma

want
large
acceptance
to study
collective
effects
and
correlations

STAR

can
sacrifice
acceptance
to improve
quality of
particle ID
and
tracking

PHOBOS

BRAHMS

physics is in hadrons, rare = irrelevant

23 PHENIX — strategy

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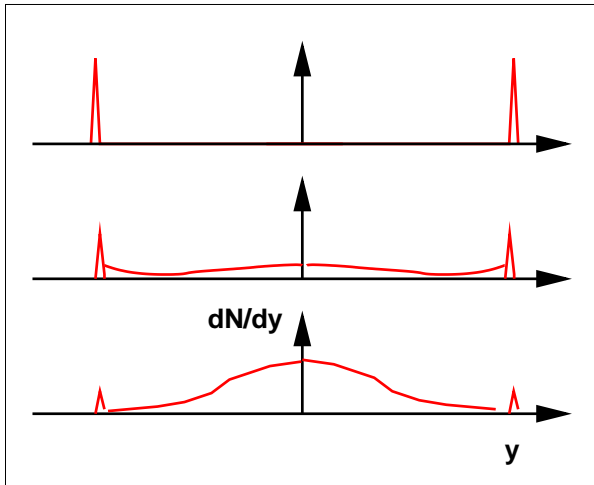
can
sacrifice
acceptance
to improve
quality of
particle ID
and
tracking

physics is in hadrons, rare = irrelevant

Technologically heterogeneous, high rate, limited acceptance, + leptons and γ

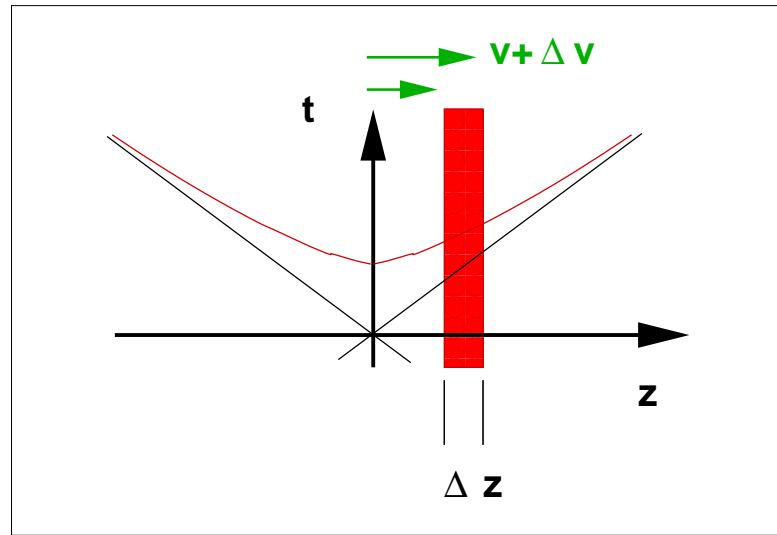
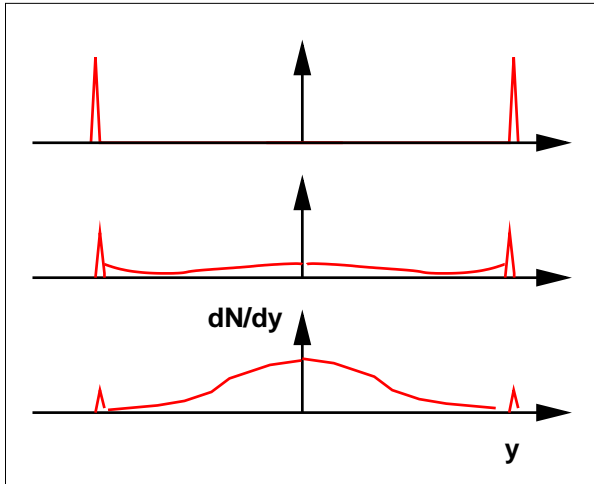
24 Stopping — Longitudinal expansion — Energy density

24 Stopping — Longitudinal expansion — Energy density



$$y = \frac{1}{2} \ln \left(\frac{1+v}{1-v} \right) = v + \mathcal{O}(v^3) \quad (8)$$

24 Stopping — Longitudinal expansion — Energy density

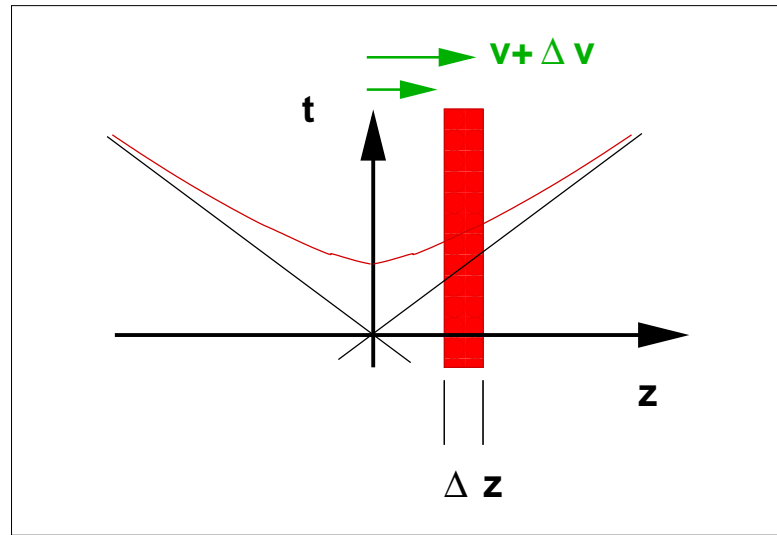
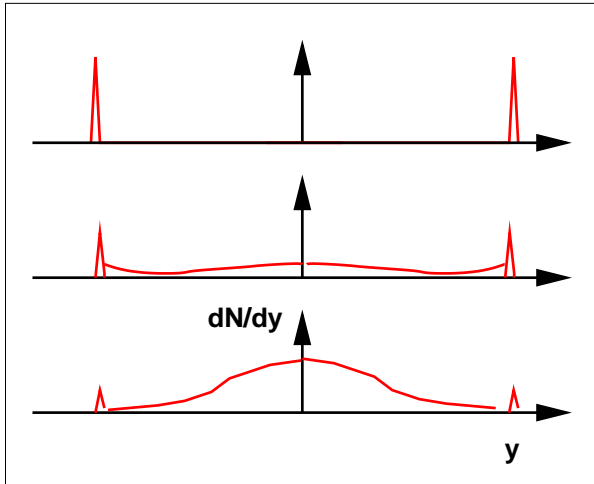


$$y = \frac{1}{2} \ln \left(\frac{1+v}{1-v} \right) = v + \mathcal{O}(v^3) \quad (8)$$

Bjorken:

$$\Delta y \approx \Delta v = \frac{\Delta z}{t} \quad (9)$$

24 Stopping — Longitudinal expansion — Energy density



$$y = \frac{1}{2} \ln \left(\frac{1+v}{1-v} \right) = v + \mathcal{O}(v^3) \quad (8)$$

Bjorken:

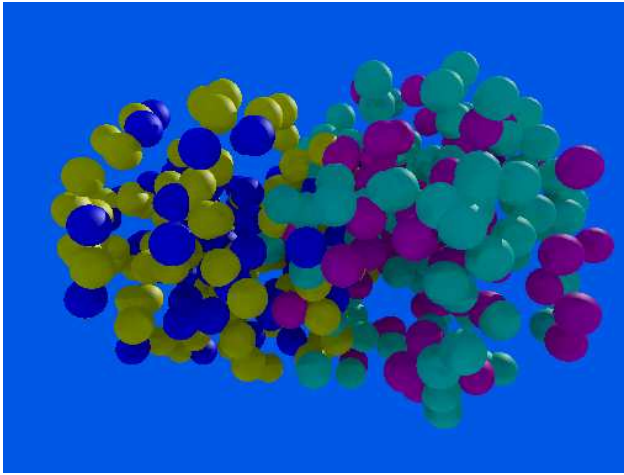
$$\Delta y \approx \Delta v = \frac{\Delta z}{t} \quad (9)$$

$$E = N \frac{d\langle E \rangle}{dy} \Delta y = N \frac{d\langle E \rangle}{dy} \frac{\Delta z}{t} \quad (10)$$

$$\epsilon(t = t_{\text{form}}) = \frac{E}{S \Delta z} = \frac{N}{S t_{\text{form}}} \frac{d\langle E \rangle}{dy} \approx \frac{dN}{dy} \frac{\langle m_t \rangle}{S t_{\text{form}}} \quad (11)$$

25 Energy density

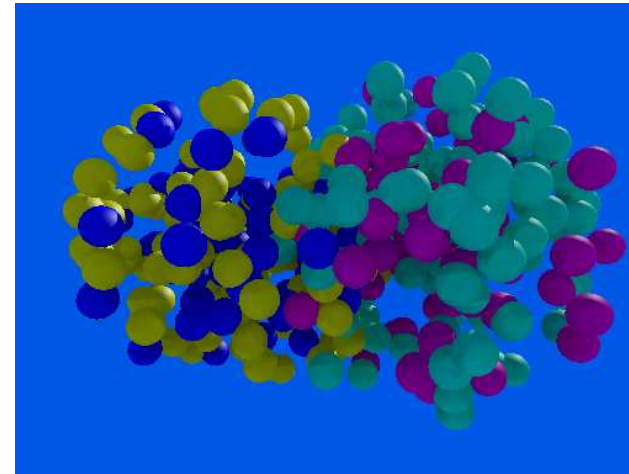
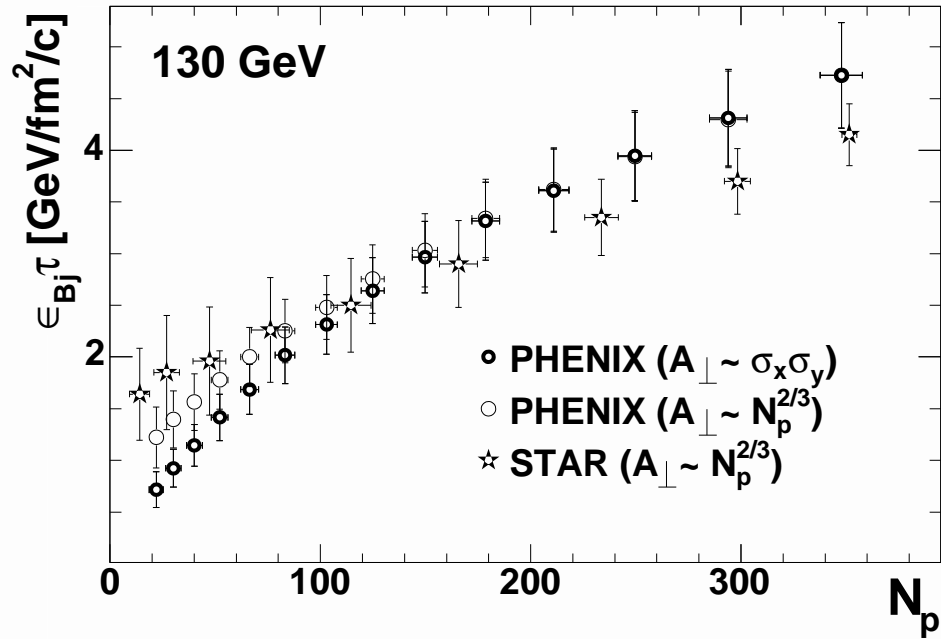
25 Energy density



N_p is number of participants. Energy density $> 1\text{GeV}/\text{fm}^3$ is believed to be adequate for the phase transition.

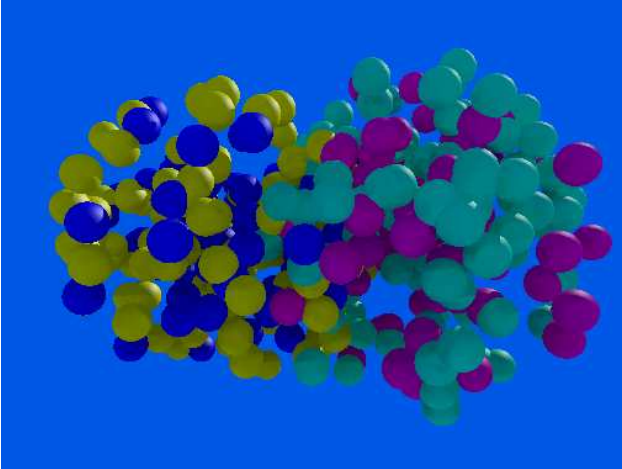
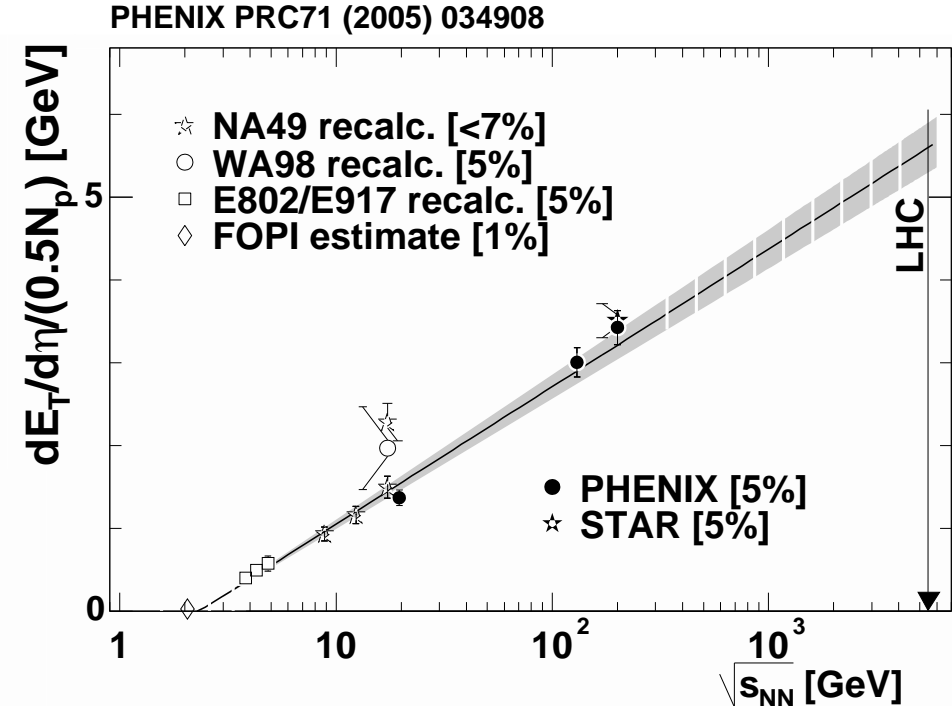
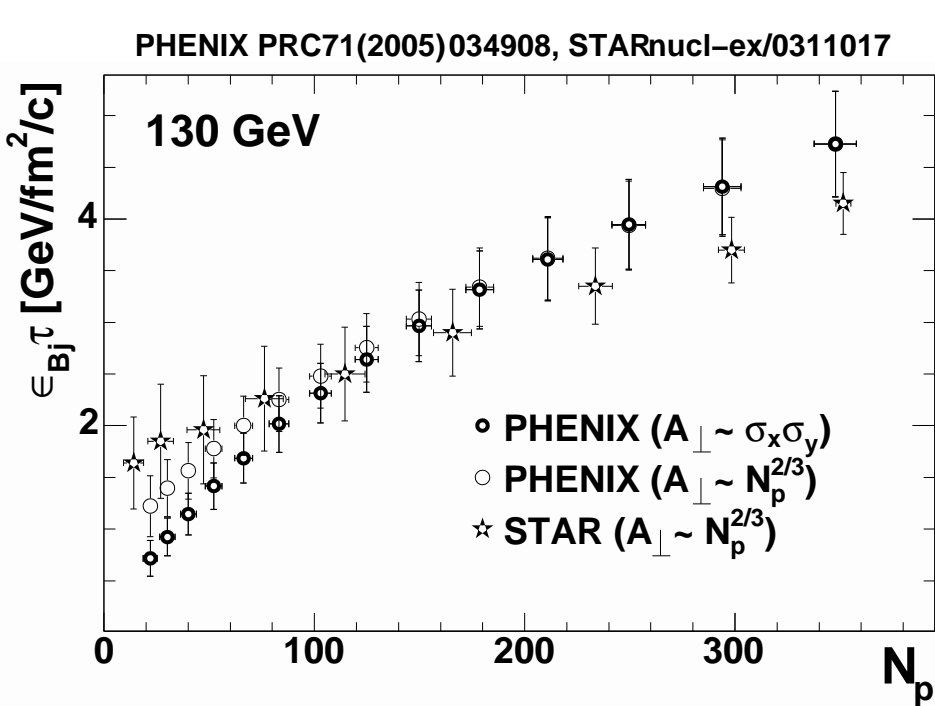
25 Energy density

PHENIX PRC71(2005)034908, STARnucl-ex/0311017



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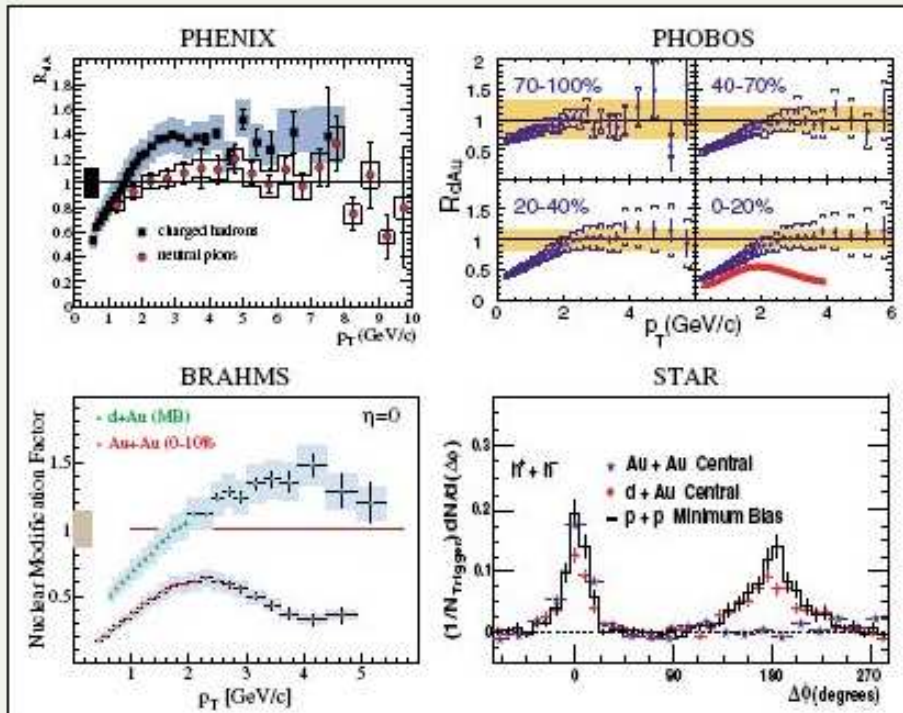
26 High p_t suppression — hadrons

PHYSICAL REVIEW LETTERS

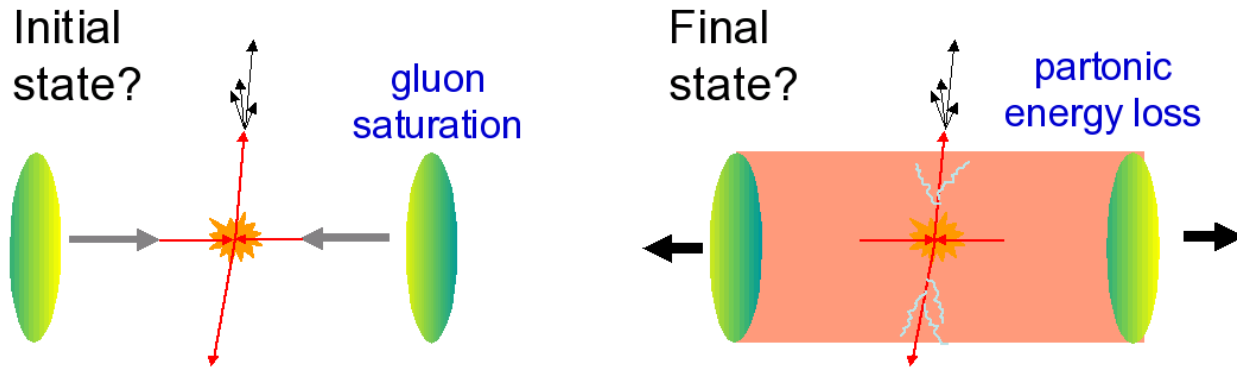
Articles published week ending

15 AUGUST 2003

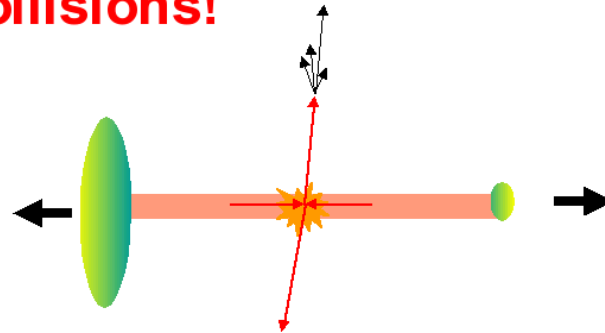
Volume 91, Number 7



27 High p_t suppression — alternatives



How to discriminate? Turn off final state \Rightarrow
d+Au collisions!



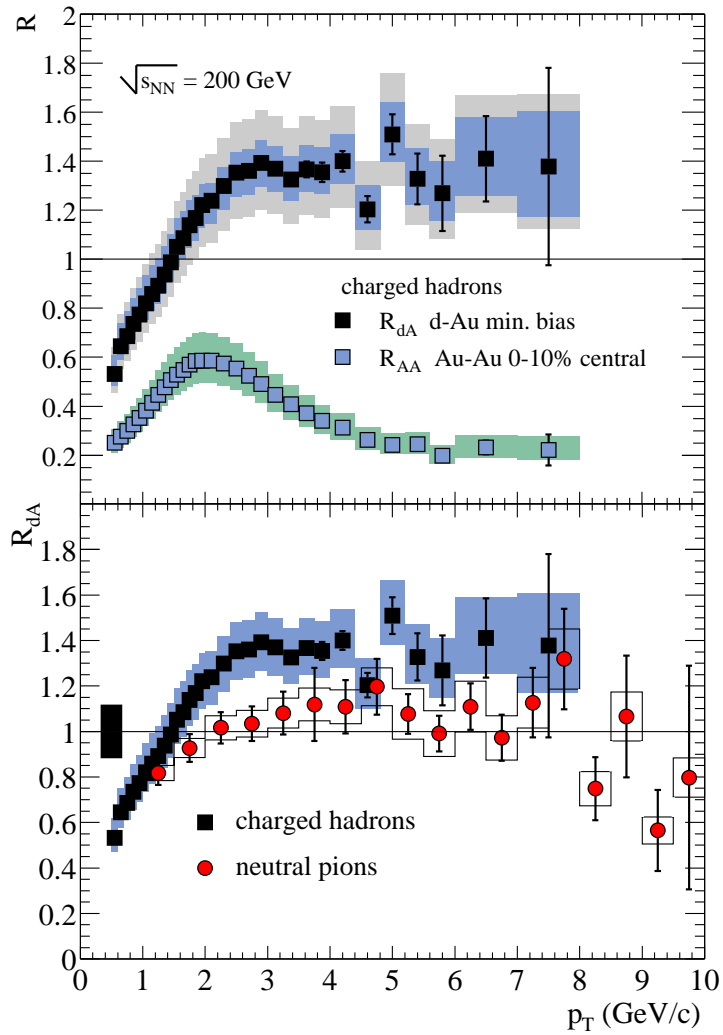
Carl Gagliardi – d+A at RHIC

7

$$R_{AA}(p_t) = \frac{\frac{1}{n_{\text{evt}}} \frac{d^2 N_{AA}}{dp_t dy}}{\frac{\langle N_{\text{coll}} \rangle}{\sigma_{pp}^{\text{inel}}} \frac{d^2 \sigma_{pp}}{dp_t dy}} \quad (12)$$

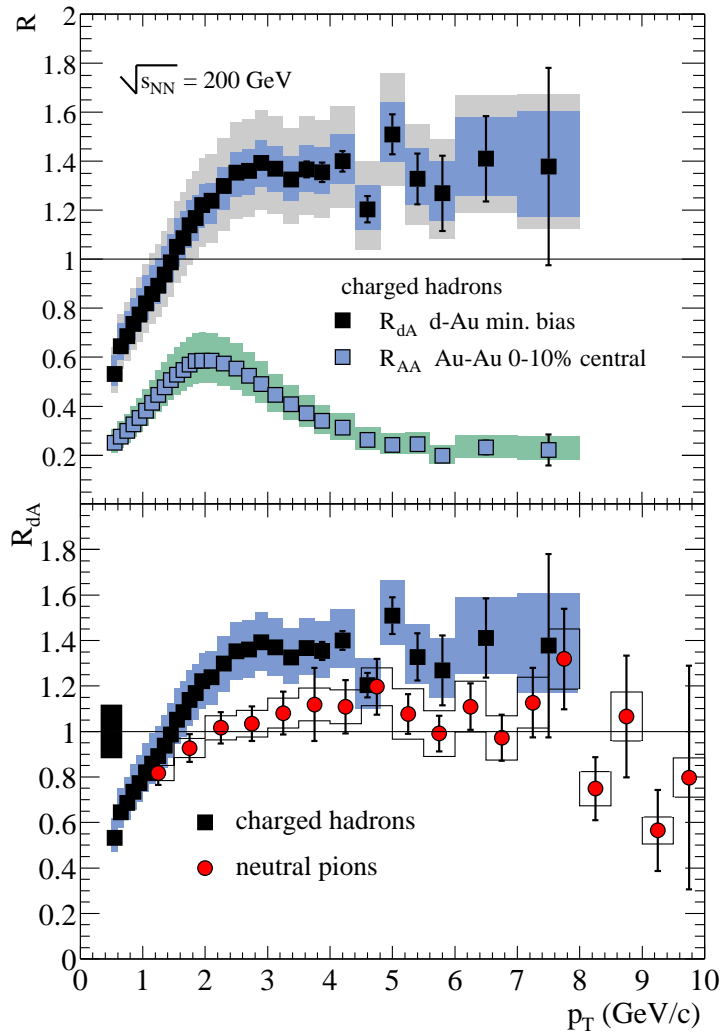
28 High p_t suppression — hadrons — STAR and PHENIX

28 High p_t suppression — hadrons — STAR and PHENIX

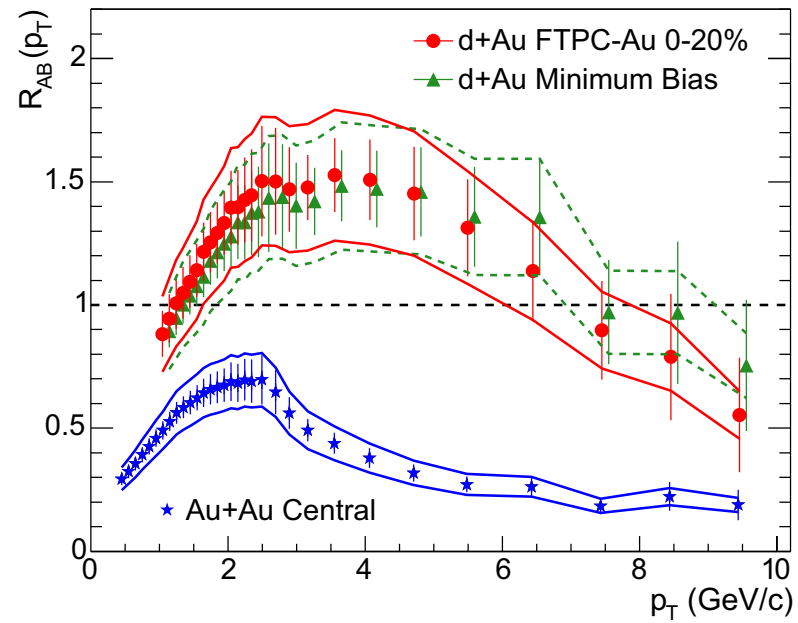


PHENIX

28 High p_t suppression — hadrons — STAR and PHENIX

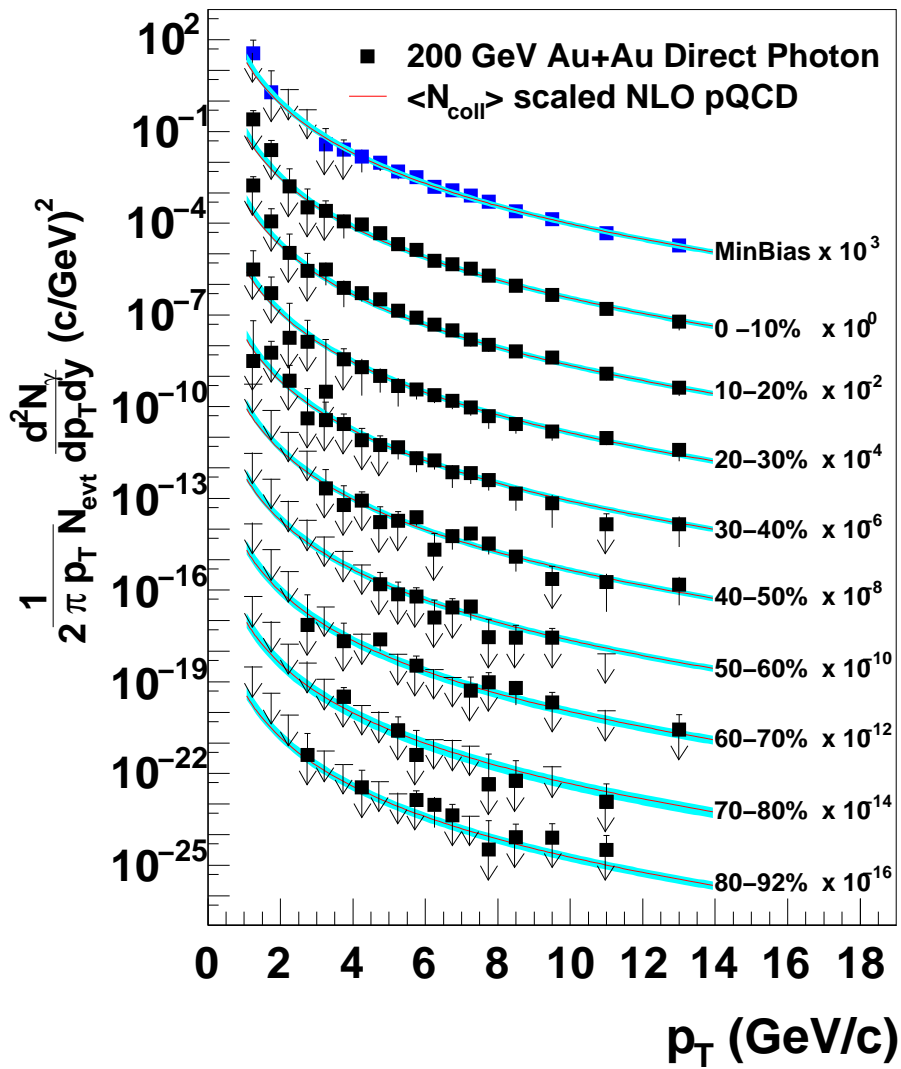


PHENIX

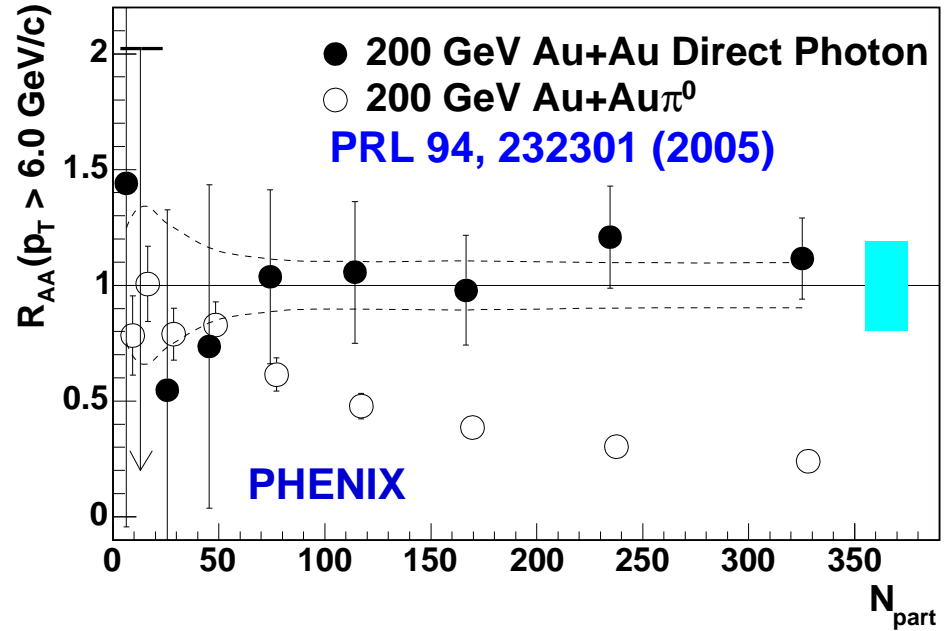


STAR

29 High p_t suppression — photons — spectra



PHENIX PRL 94, 232301 (2005)

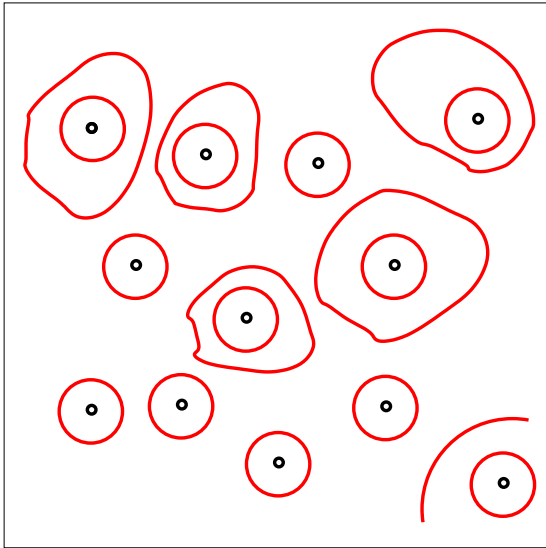
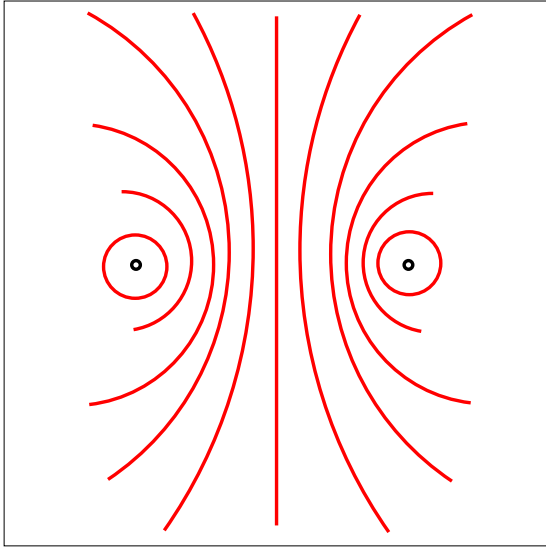


Photons are produced in participant NN collisions with no initial state modification. The high p_t suppression in hadronic sector is not an initial state effect.

Binary-collision scaled pQCD gets it right

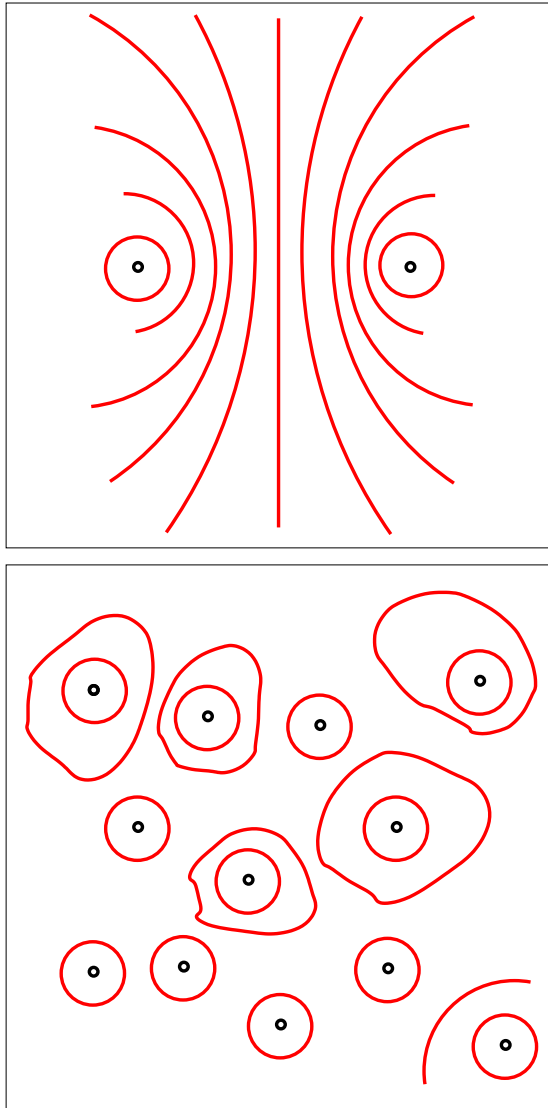
30 Charm

30 Charm

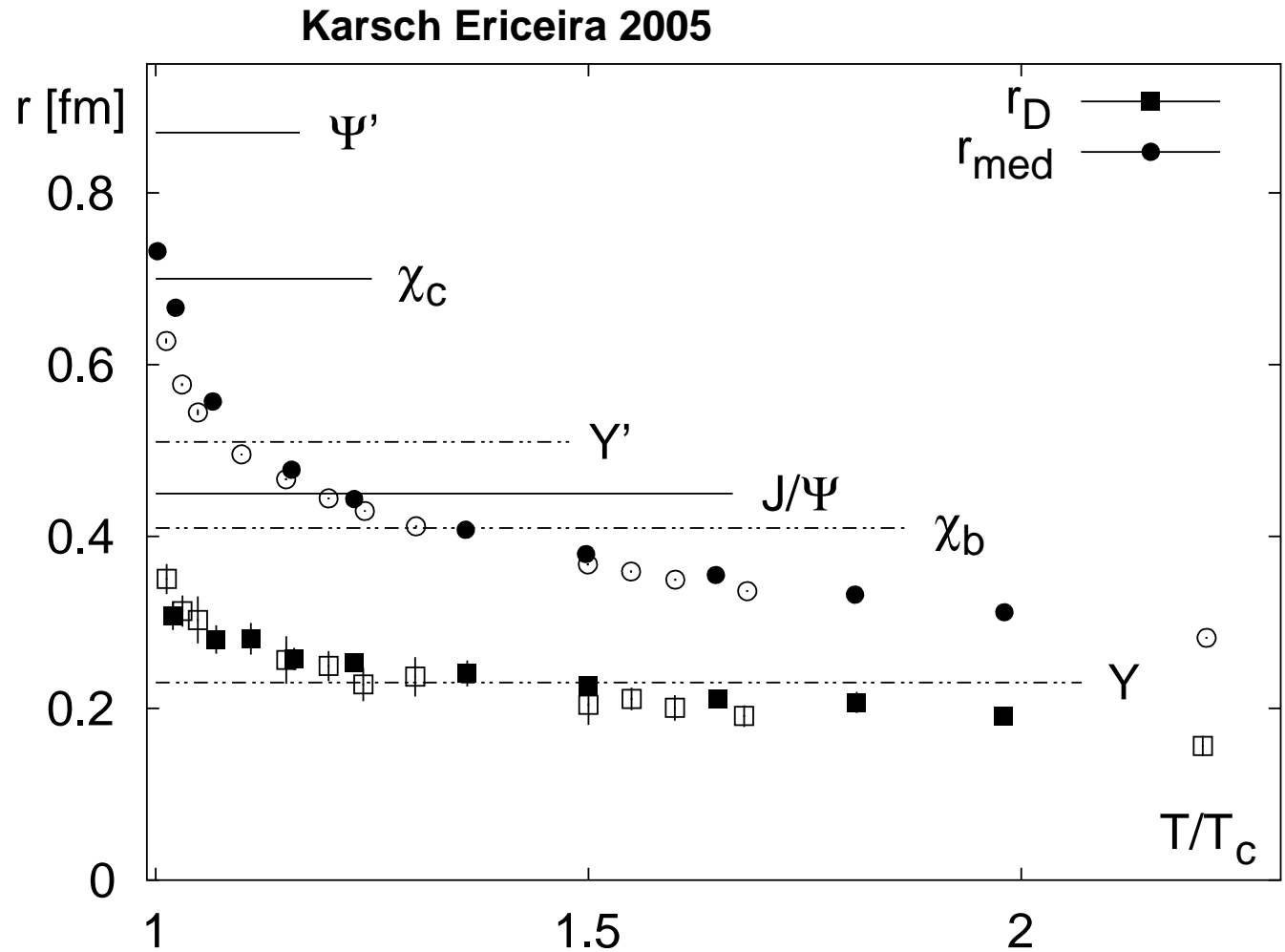


Matsui, Satz
mechanism of J/ψ
suppression. PLB
178:416, 1986.

30 Charm



Matsui, Satz
mechanism of J/ψ
suppression. PLB
178:416, 1986.



Open/closed SU(3)/SU(2); for $r > r_{med}$ the $q\bar{q}$ force is strongly modified by the colored medium; r_D is the Debye screening radius. Horizontal lines are $\sqrt{r_{q\bar{q}}^2}$ for charmonium states.

31 PHENIX — sources of electrons

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photonic (calibrated out using γ converter of known thickness):

$$\pi^0 \rightarrow \gamma e^+ e^- \rightarrow \gamma - \text{conversion} \quad (13)$$

$$\eta, \eta', \rho, \omega, \phi \rightarrow \pi^0 \rightarrow \gamma e^+ e^- \rightarrow \gamma - \text{conversion} \quad (14)$$

31 PHENIX — sources of electrons

photonic (calibrated out using γ converter of known thickness):

$$\pi^0 \rightarrow \gamma e^+ e^- \rightarrow \gamma - \text{conversion} \quad (13)$$

$$\eta, \eta', \rho, \omega, \phi \rightarrow \pi^0 \rightarrow \gamma e^+ e^- \rightarrow \gamma - \text{conversion} \quad (14)$$

non-photonic (wanted):

$$D^+ \rightarrow e^+ X \quad (15)$$

$$D^0 \rightarrow e^+ X \quad (16)$$

$$B \rightarrow \quad (17)$$

non-photonic (unwanted):

$$K \rightarrow \pi e \nu \quad (18)$$

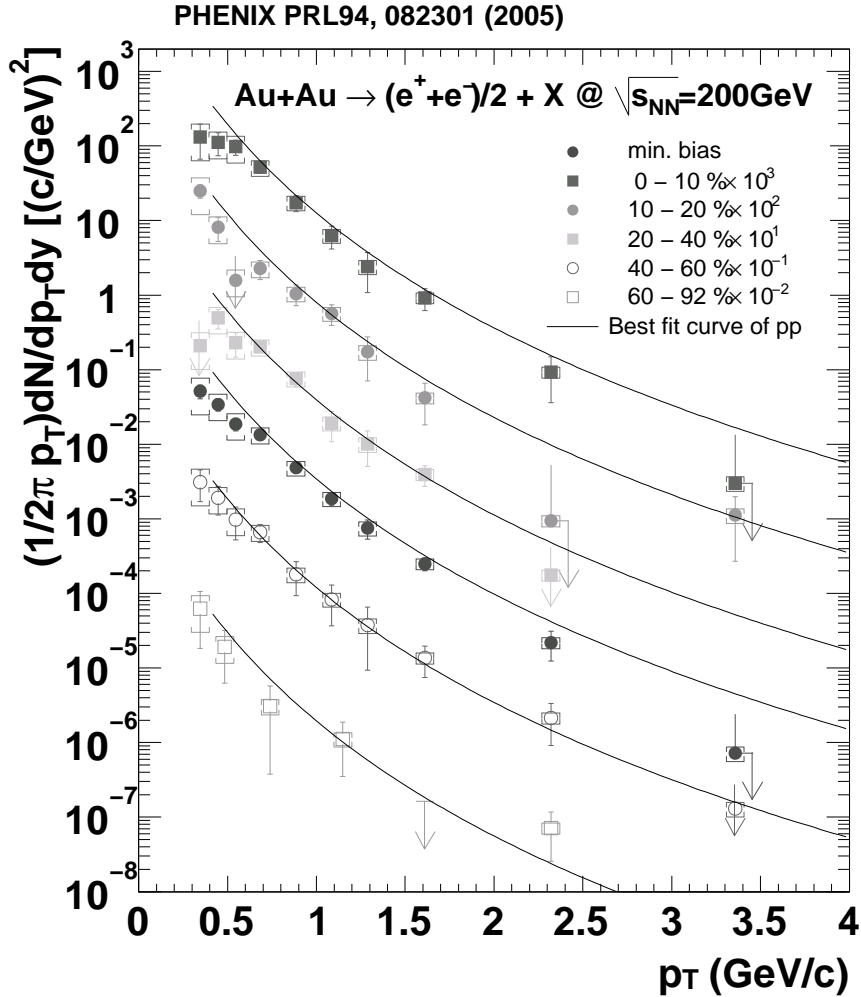
(GEANT)

$$\rho, \omega, \phi \rightarrow e^+ e^- \quad (19)$$

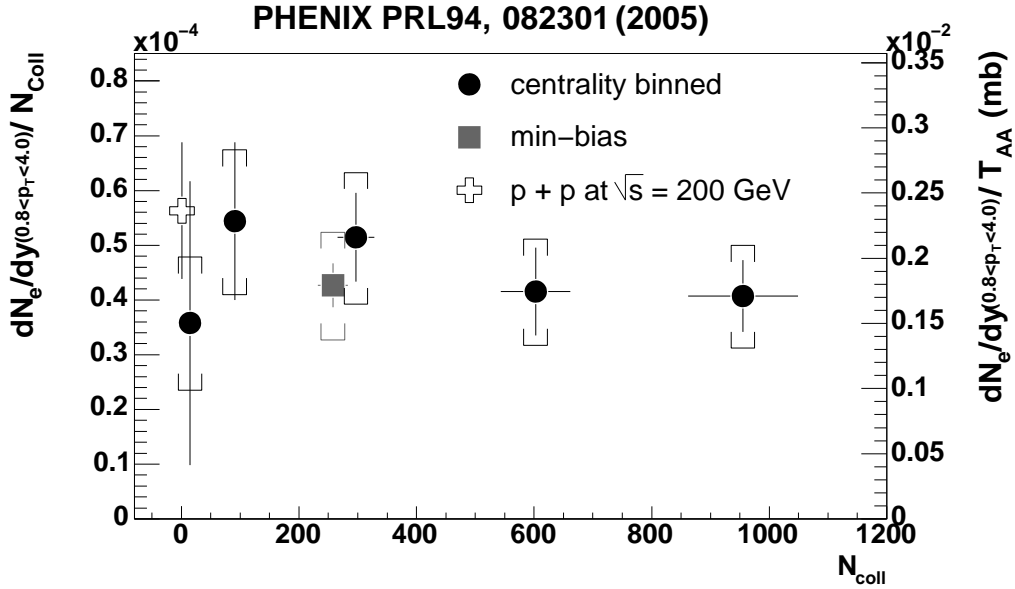
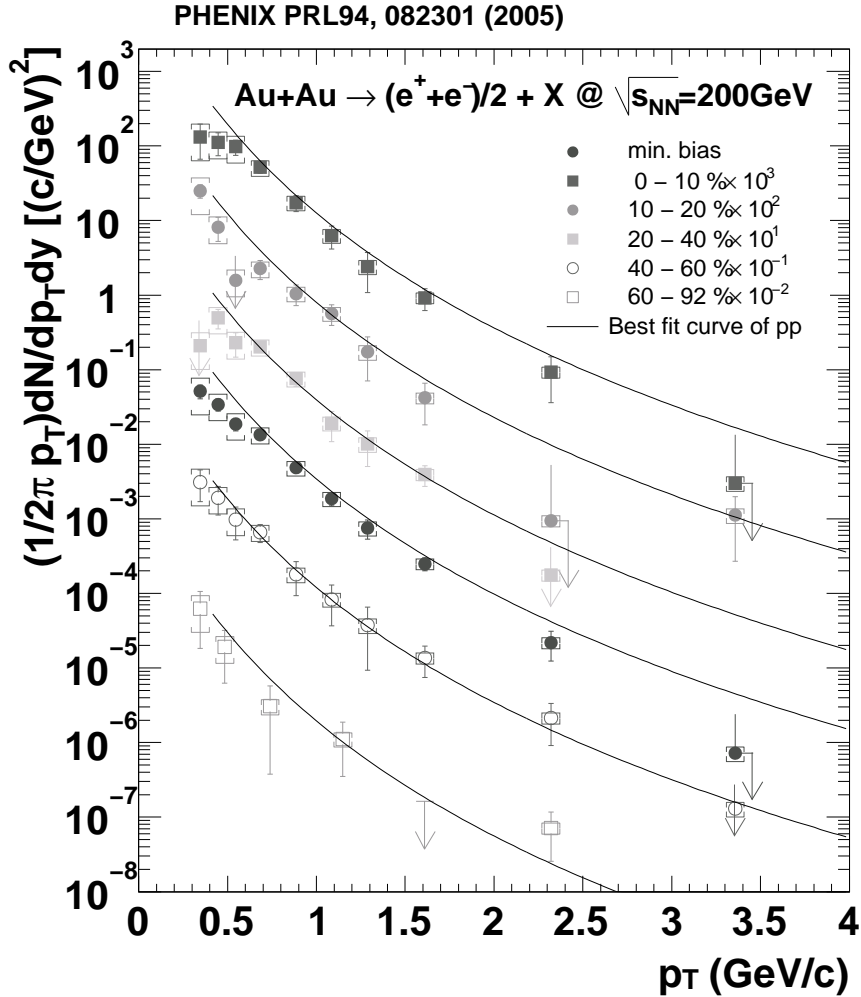
m_t scaling of measured π^0 spectra

32 PHENIX — non-photonic single electrons

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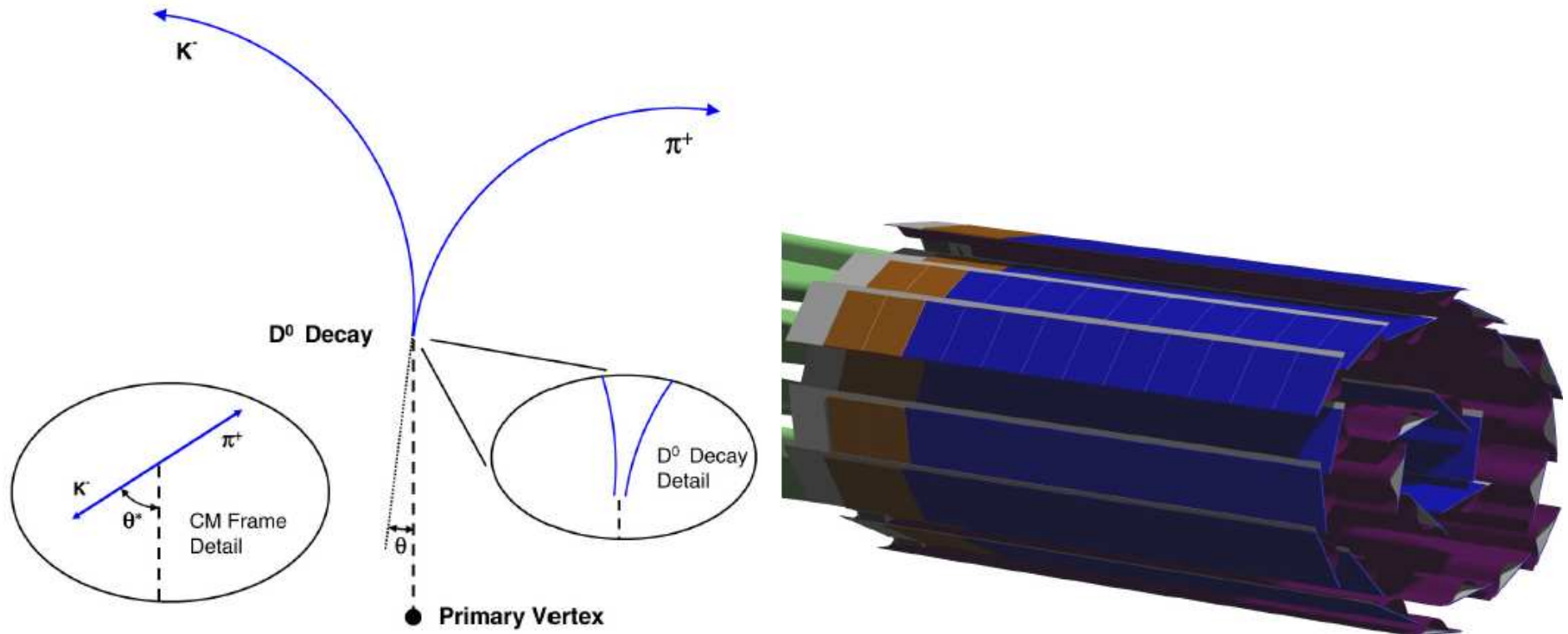
32 PHENIX — non-photonic single electrons



Believed to come from open heavy flavor decays, non-photonic electrons do not show high p_t suppression: the spectra scale with binary collisions.

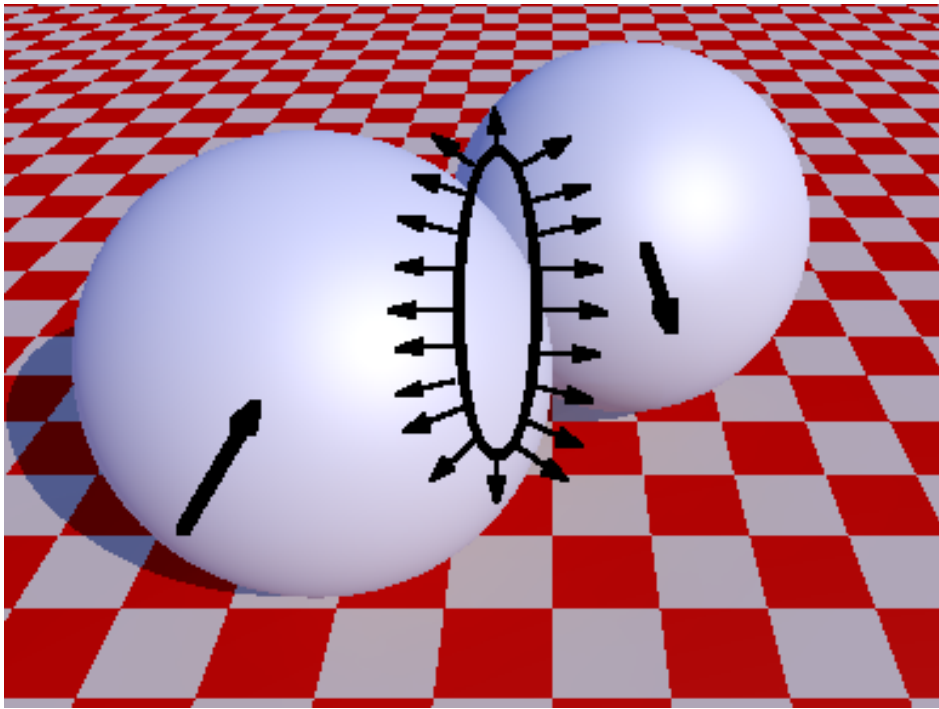
33 Future of charm: STAR Heavy Flavor Tracker

- charm production for J/Ψ signature
- clarify charm R_{AA} story
- elliptic flow of D mesons – is flow partonic ? is c part of it ?



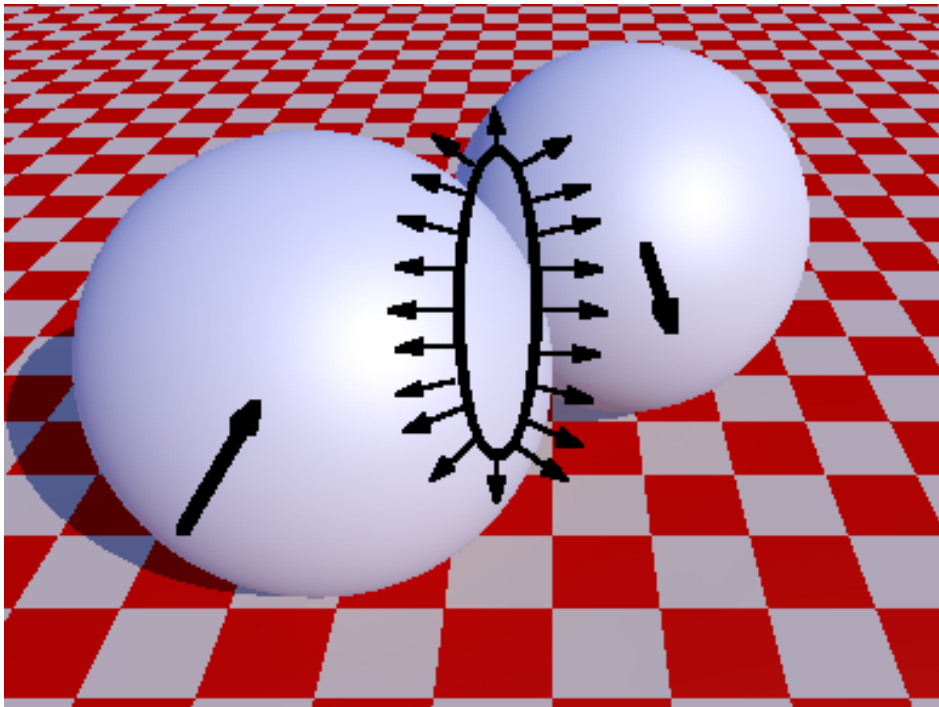
$c\tau \propto 100 - 300 \mu m$. Reduce combinatorics by vertex-finding. $10 \mu m$ resolution.

34 Azimuthal asymmetry — flow



(x, y) anisotropy \rightarrow rescattering \rightarrow
 (p_x, p_y) anisotropy

34 Azimuthal asymmetry — flow



(x, y) anisotropy \rightarrow rescattering \rightarrow
 (p_x, p_y) anisotropy

$$E \frac{d^3 N}{d^3 p} = \frac{1}{2\pi} \frac{d^2 N}{p_t dp_t dy} \left\{ 1 + \sum_{m=1}^{\infty} 2v_m \cos[m(\phi - \Psi_r)] \right\} \quad (20)$$

- flow starts early – perhaps before hydro is applicable (stopping stage)
- testifies to equilibration
- sensitive to pressure and density gradients
- flow is a multiparticle effect; there is “non-flow”

35 Directed flow — Importance of impact vector

35 Directed flow — Importance of impact vector

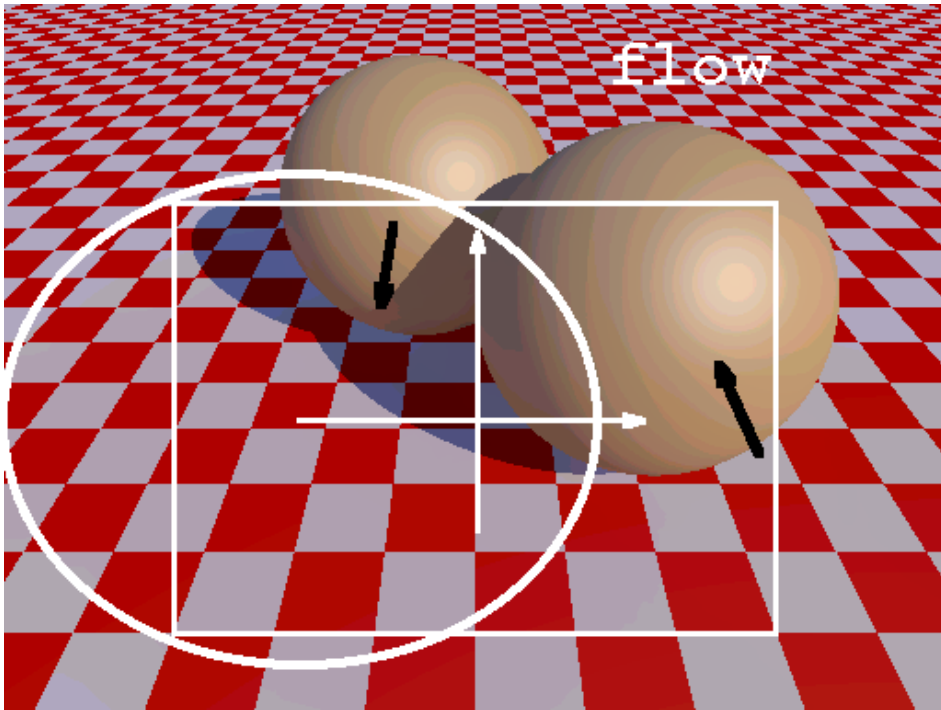
$$E \frac{d^3 N}{d^3 p} = \frac{1}{2\pi} \frac{d^2 N}{p_t dp_t dy} \left\{ 1 + \sum_{m=1}^{\infty} 2v_m \cos[m(\phi - \Psi_r)] \right\} \quad (21)$$

for v_1 , need to know $0 \leq \Psi_r < 2\pi$:

35 Directed flow — Importance of impact vector

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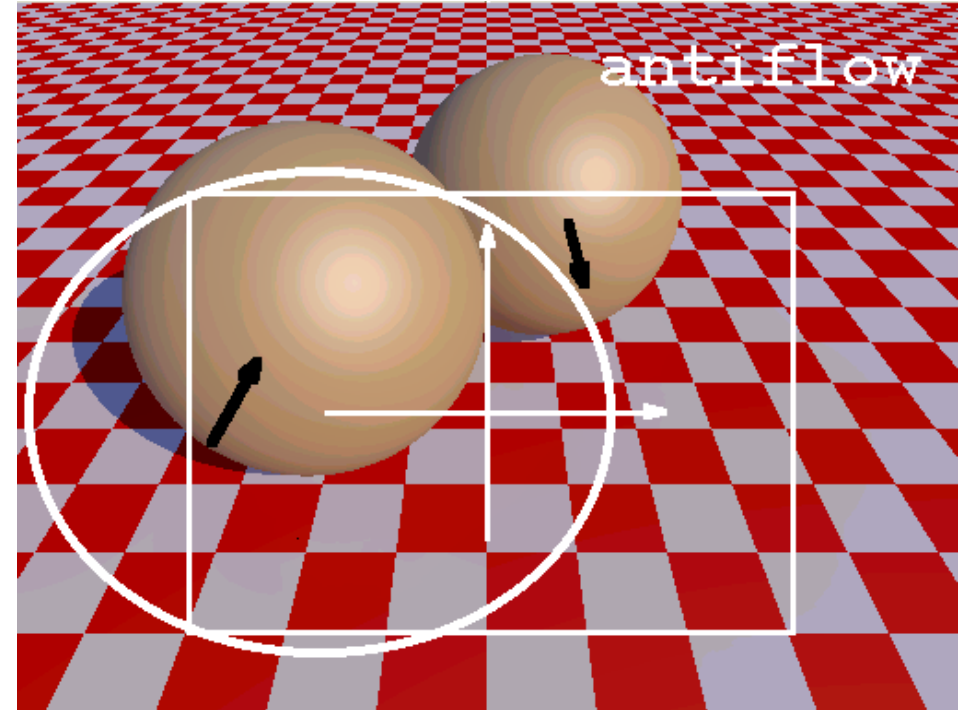
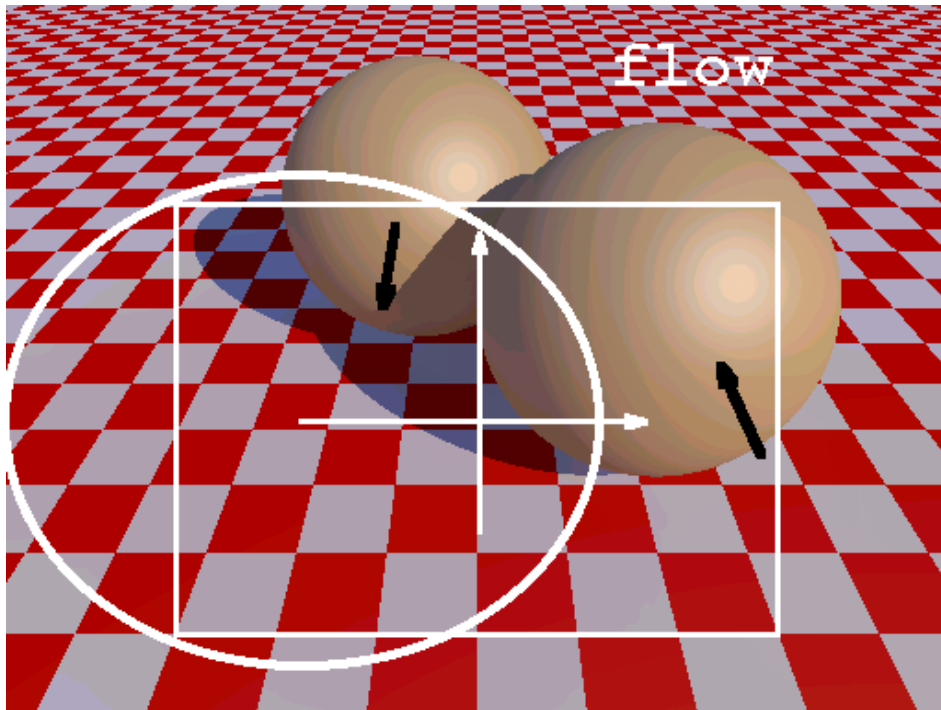
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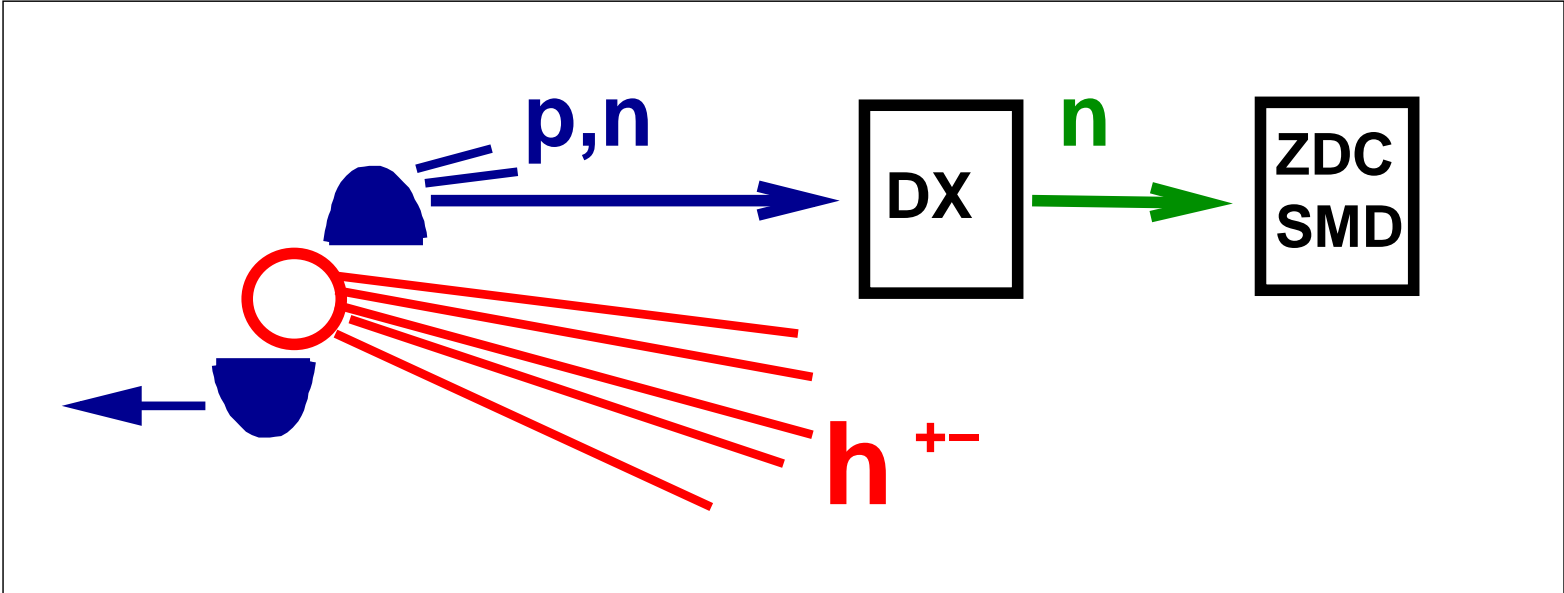
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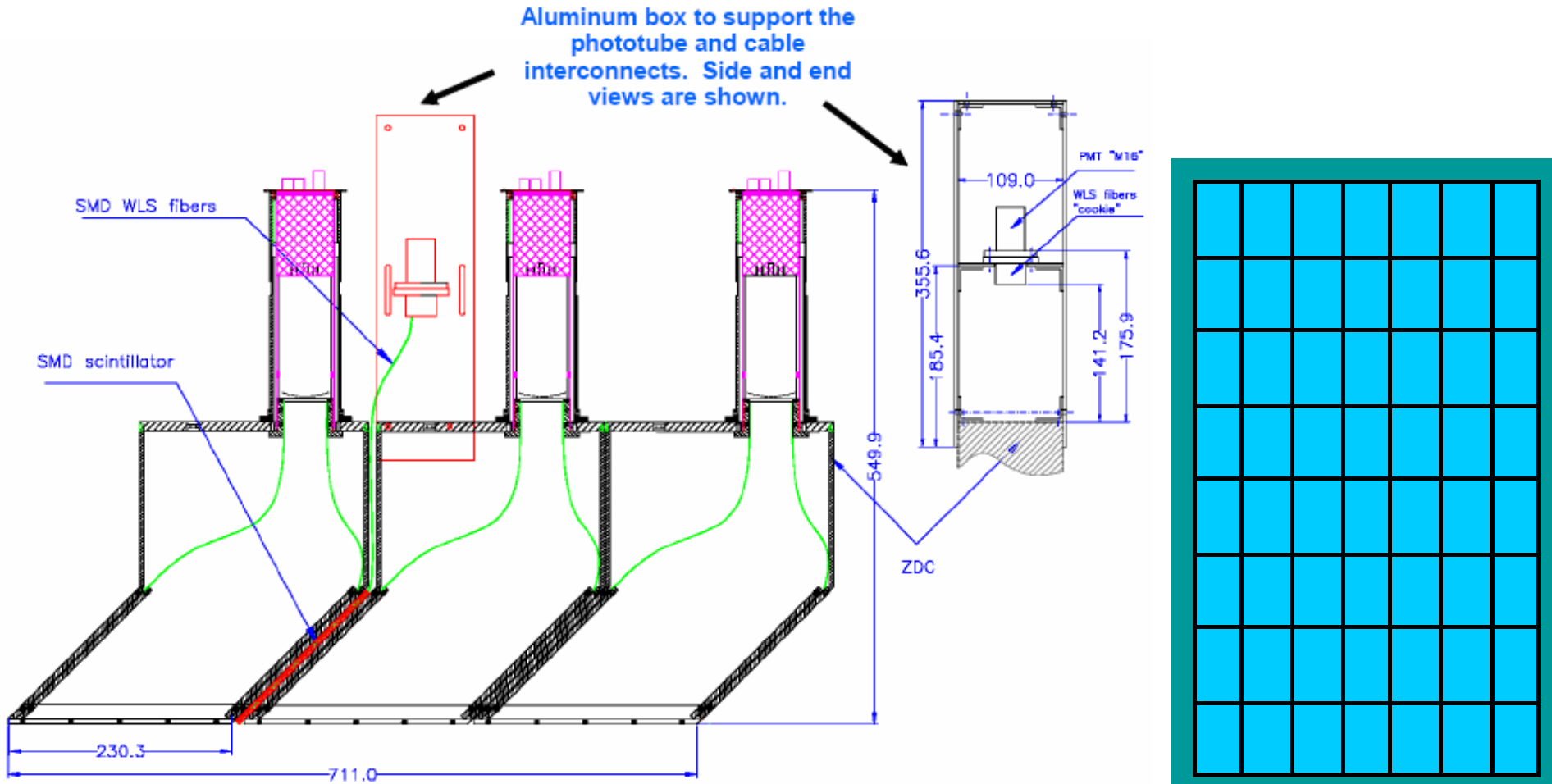
need to be able to distinguish the two !

36 Directed flow — STAR ZDC SMD

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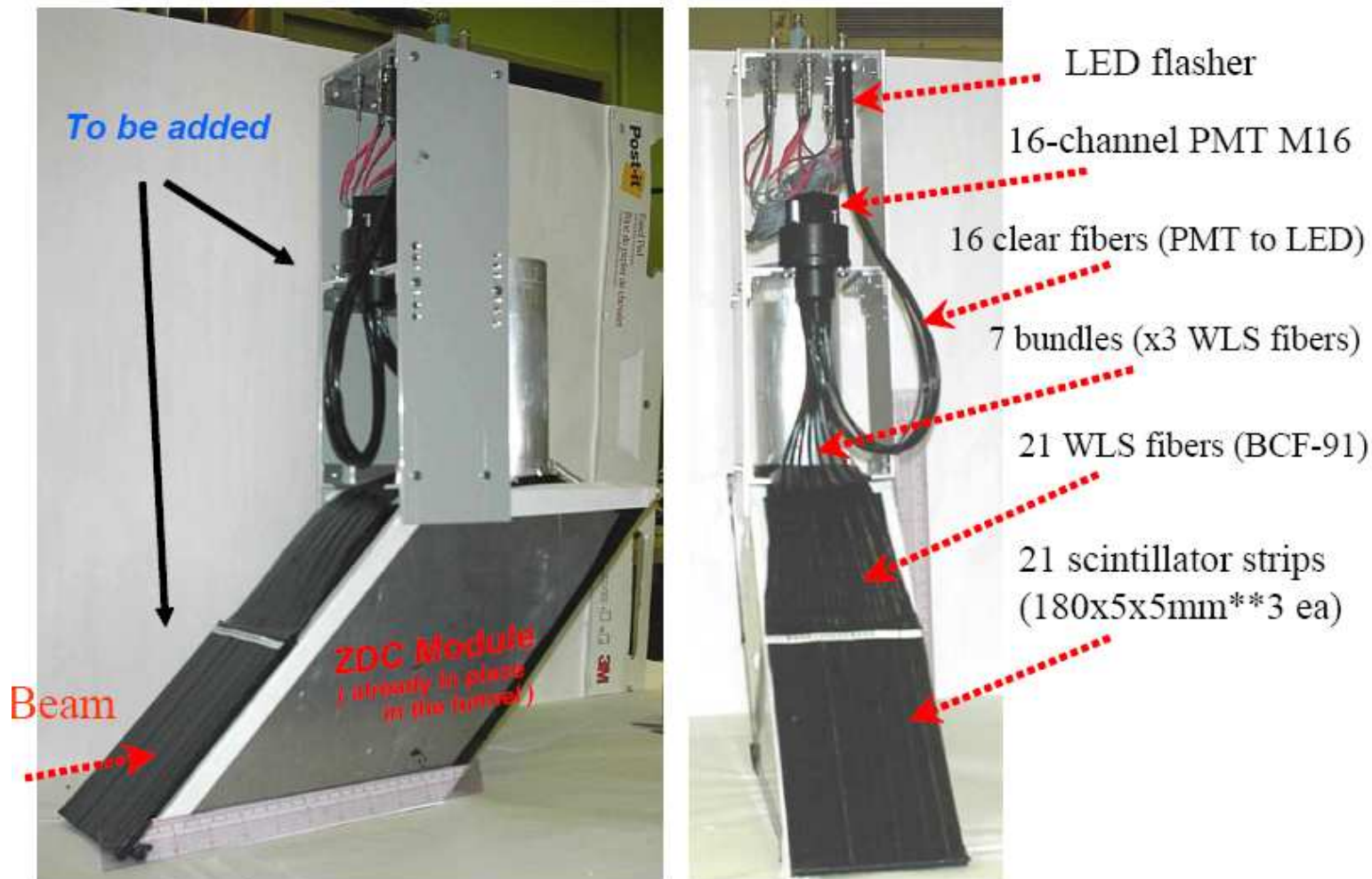


36 Directed flow — STAR ZDC SMD



7 vertical and 8 horizontal SMD slats at 1/3 of ZDC depth measure transverse asymmetry of spectator neutrons

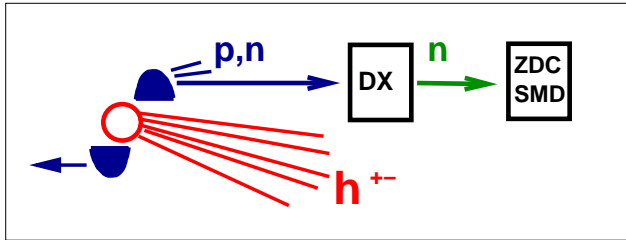
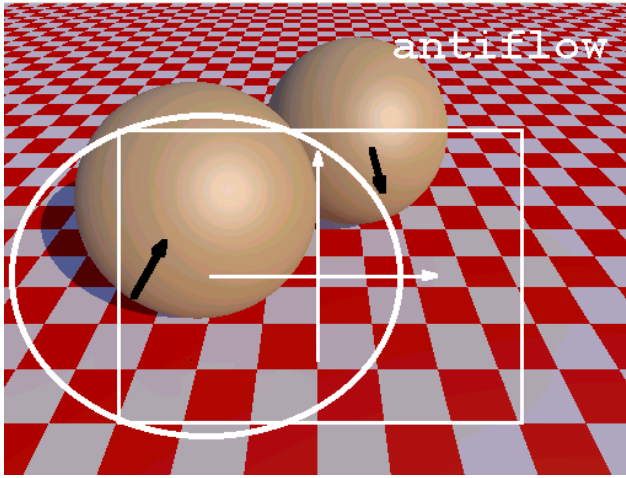
36 Directed flow — STAR ZDC SMD



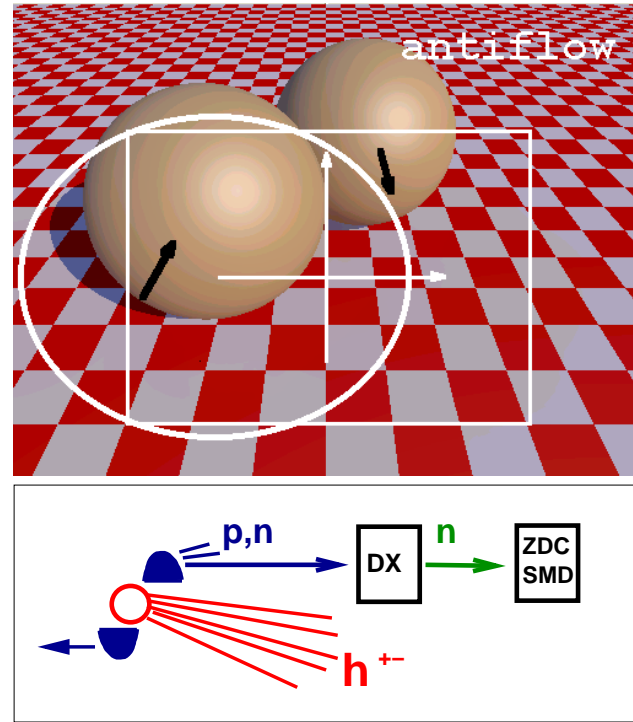
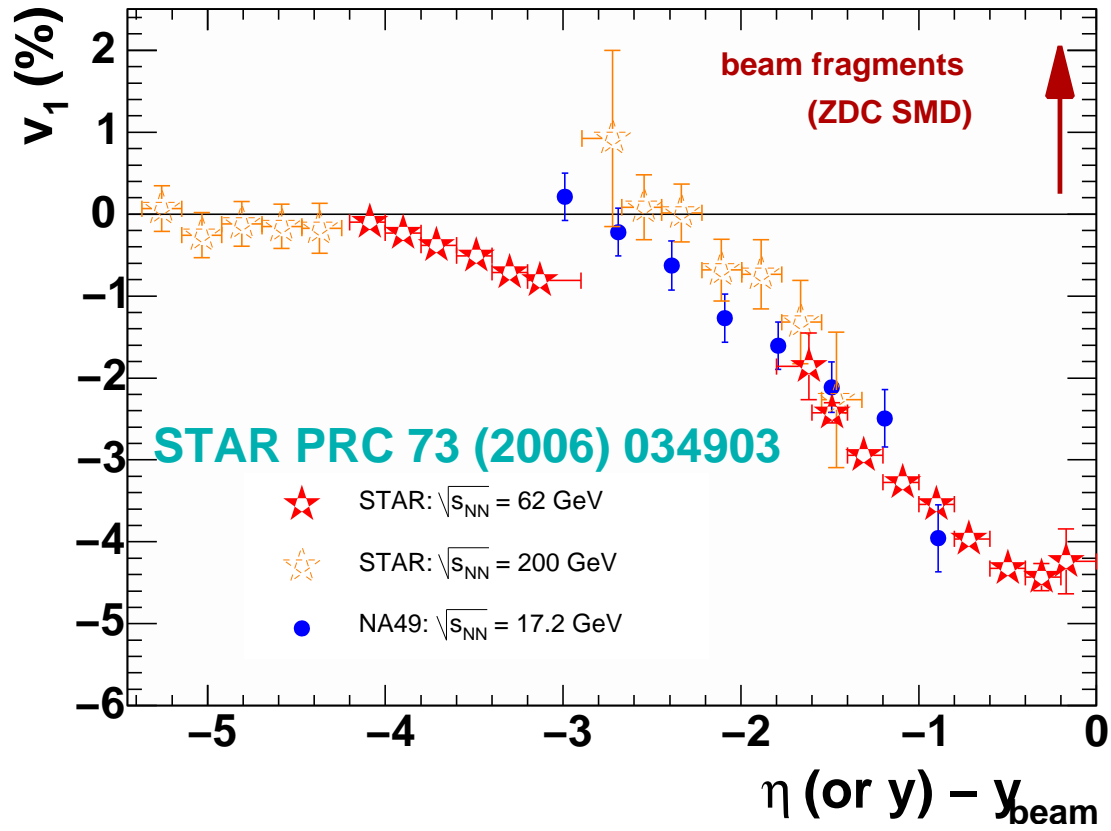
- effort spearheaded by Kent people
- besides flow, contributes to strangelet search, UPC and spin physics

37 Directed flow

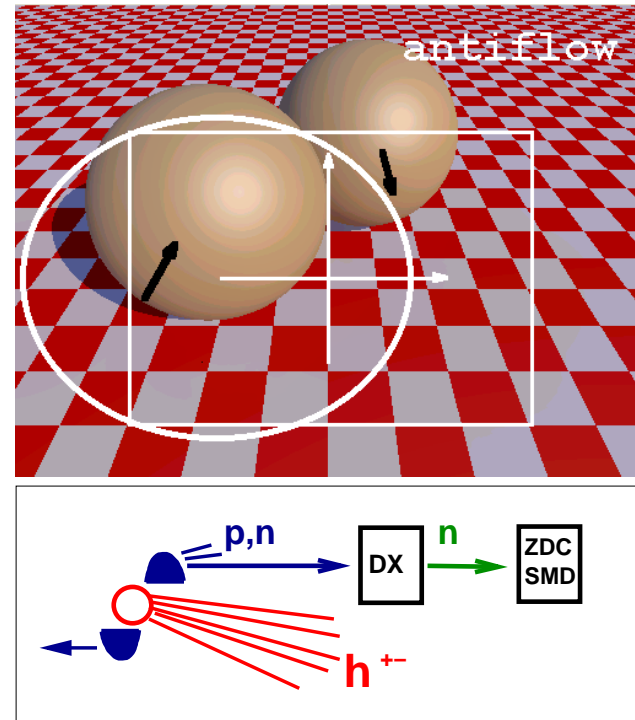
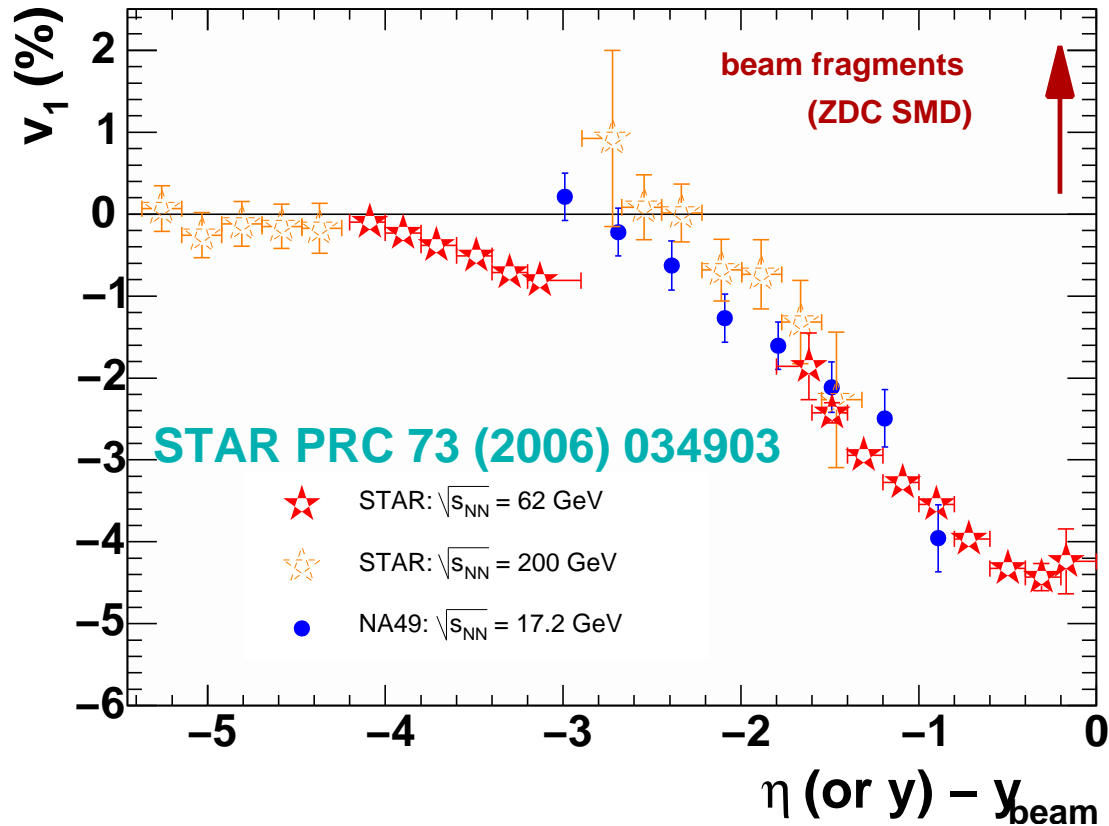
37 Directed flow



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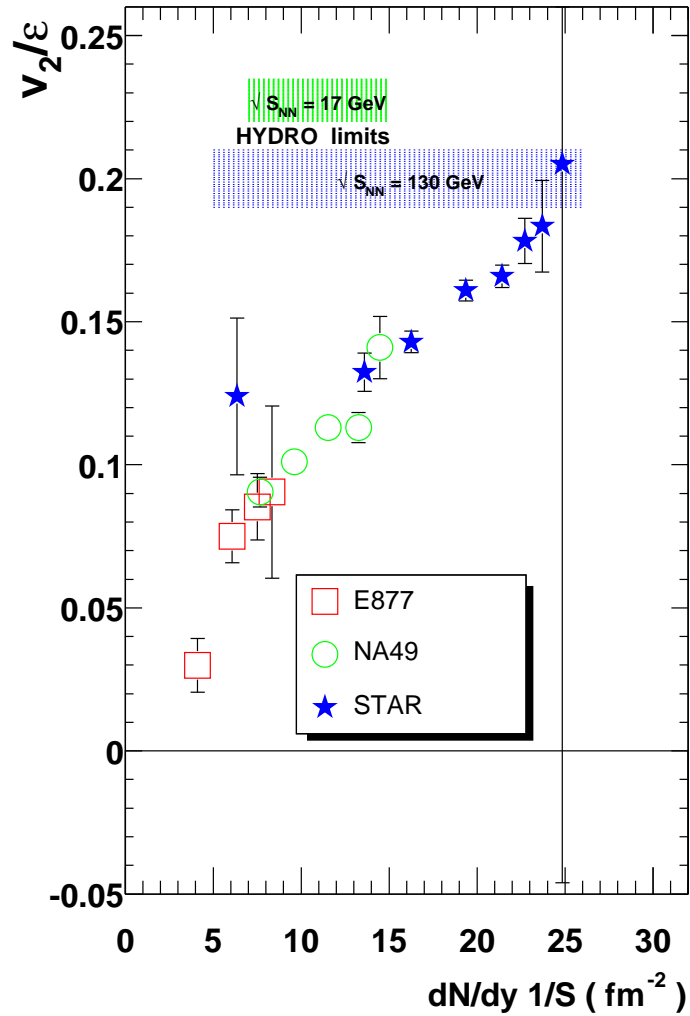


Charged-particles v_1 from 3-particle cumulants in the projectile frame.

- monotonic around midrapidity
- Supports limiting fragmentation
- Antiflow !

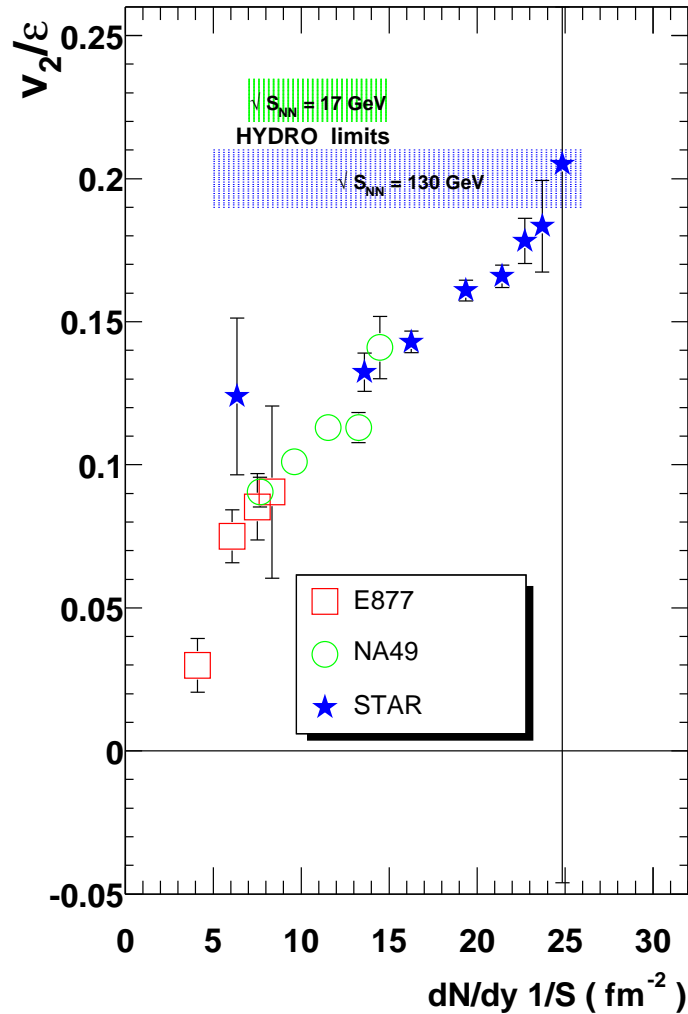
38 Elliptic flow and hydro fluidity

38 Elliptic flow and hydro fluidity



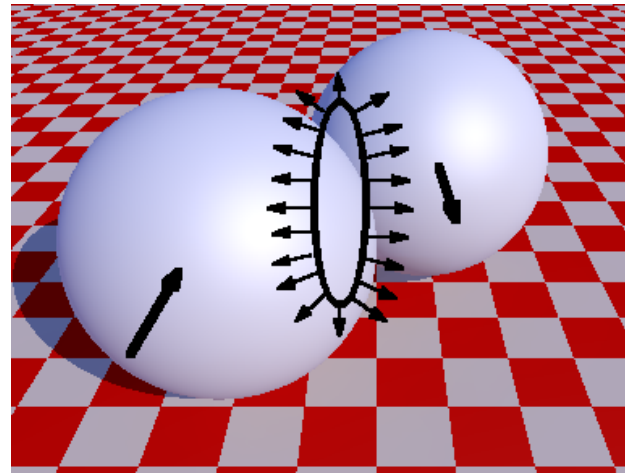
STAR PRC66 (2002) 034904
(cumulant v_2). Hydro limits by
Kolb, Sollfrank, Heinz, PRC62
(2000) 054909.

38 Elliptic flow and hydro fluidity



- v_2 is a response to excentricity

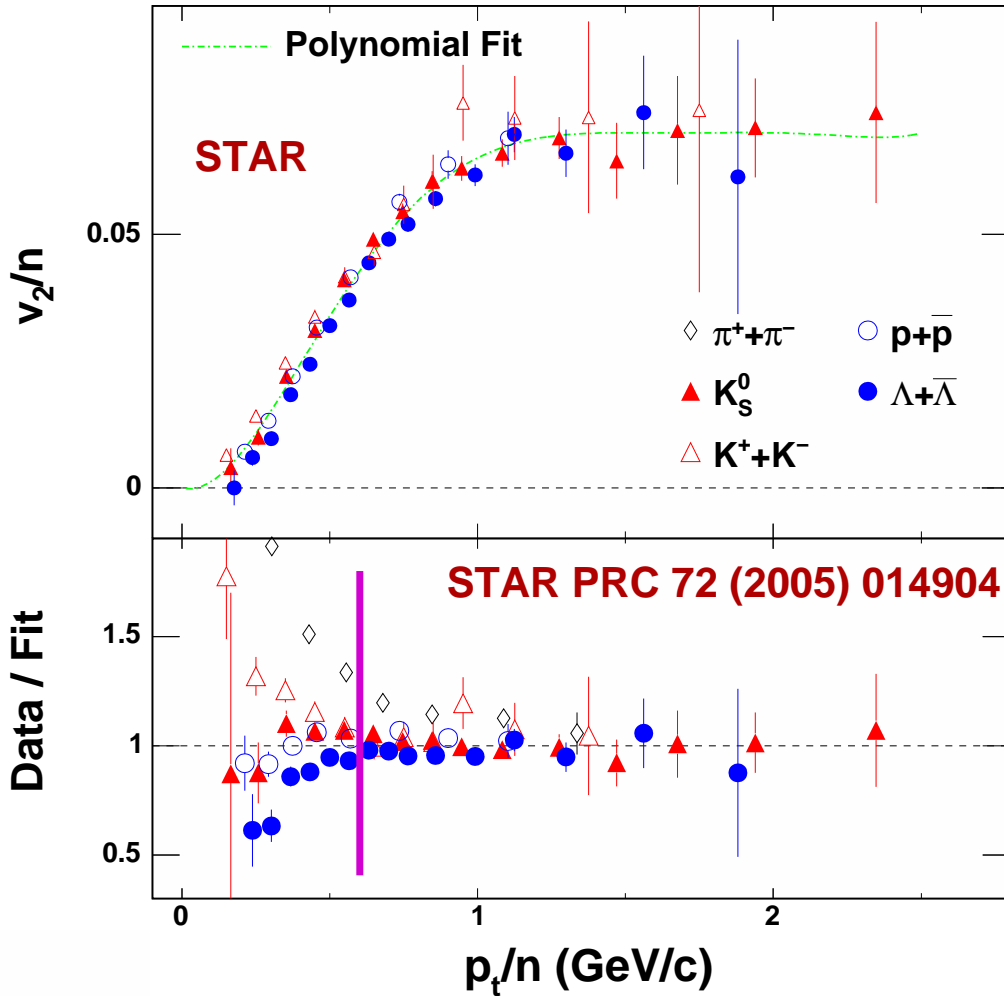
$$\epsilon = \frac{(y^2 - x^2)}{(y^2 + x^2)}$$



- low viscosity \iff high cross-sections ! "sQGP".

STAR PRC66 (2002) 034904
 (cumulant v_2). Hydro limits by
 Kolb, Sollfrank, Heinz, PRC62
 (2000) 054909.

39 Elliptic flow and quark coalescence



$$\frac{dN}{d\phi} \propto 1 + 2v_2 \cos(2\phi) \quad (22)$$

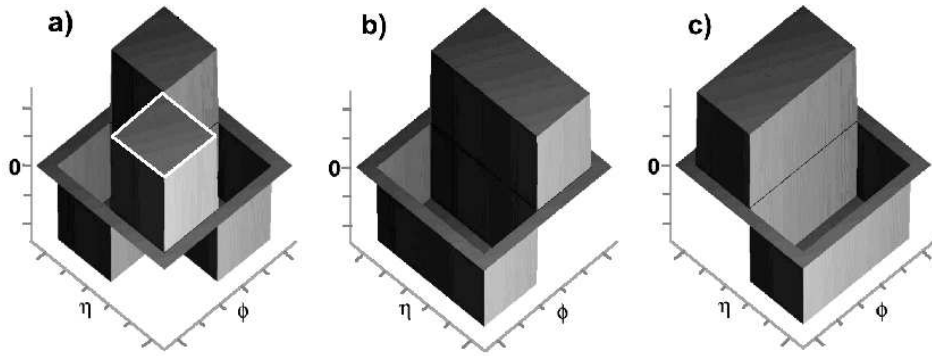
$$\frac{dN_{\text{clscnc},n}}{d\phi}(p_t) \propto \left(\frac{dN(\frac{p_t}{n})}{d\phi} \right)^n \quad (23)$$

$$(1 + 2v_2 \cos(2\phi))^n = 1 + 2v_2 n \cos(2\phi) + \mathcal{O}(v_2^2) \quad (24)$$

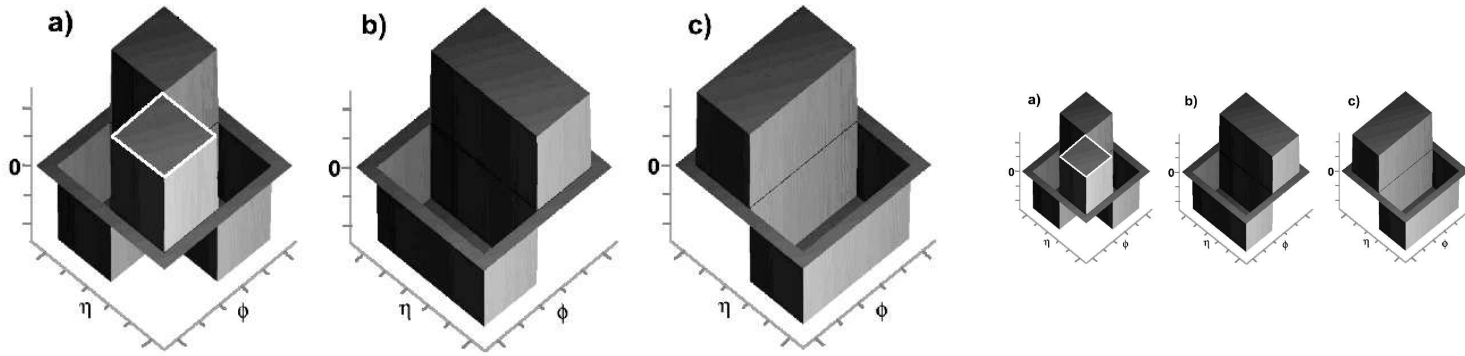
STAR AuAu 200 GeV minbias; n is number of constituent quarks. Expect universality if quark coalescence dominates hadronization after the universal flow sets in. Valid at $p_t/n > 0.6$ GeV/c for $K_S^0, K^\pm, p, \bar{p}, \Lambda, \bar{\Lambda}$.

40 Local hadron density fluctuations and Discrete Wavelet Transform (DWT)

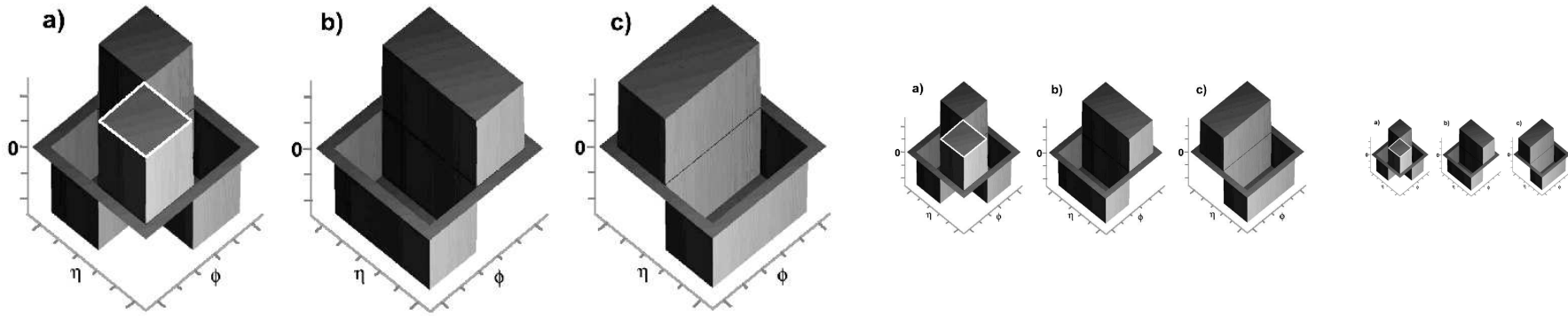
40 Local hadron density fluctuations and Discrete Wavelet Transform (DWT)



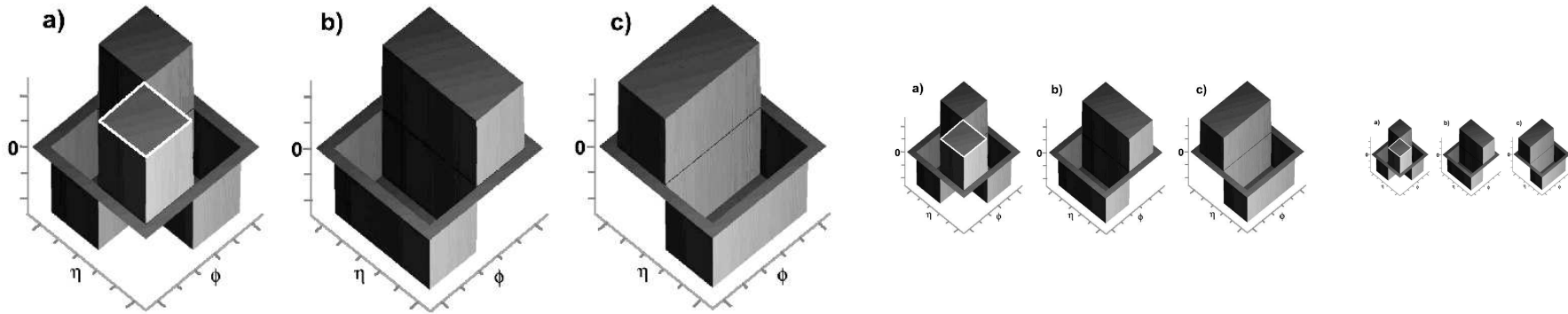
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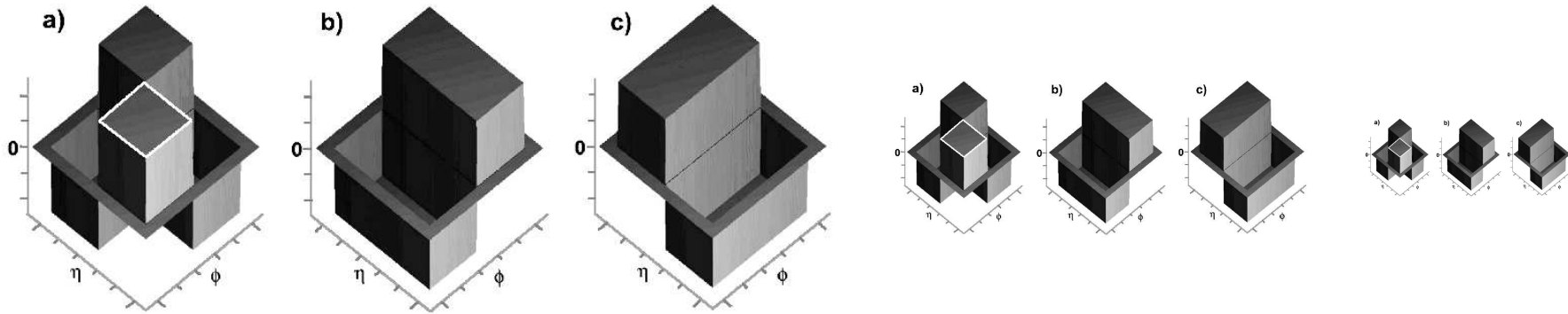


40 Local hadron density fluctuations and Discrete Wavelet Transform (DWT)



$F_{m,l,k}^\lambda(\phi, \eta)$ —Haar wavelet **orthonormal basis** in (ϕ, η) . scale fineness (m), directional modes of sensitivity (λ), track density $\rho(\eta, \phi, p_t)$, locations in 2D (l, k) . **DWT** is an expansion in this basis.

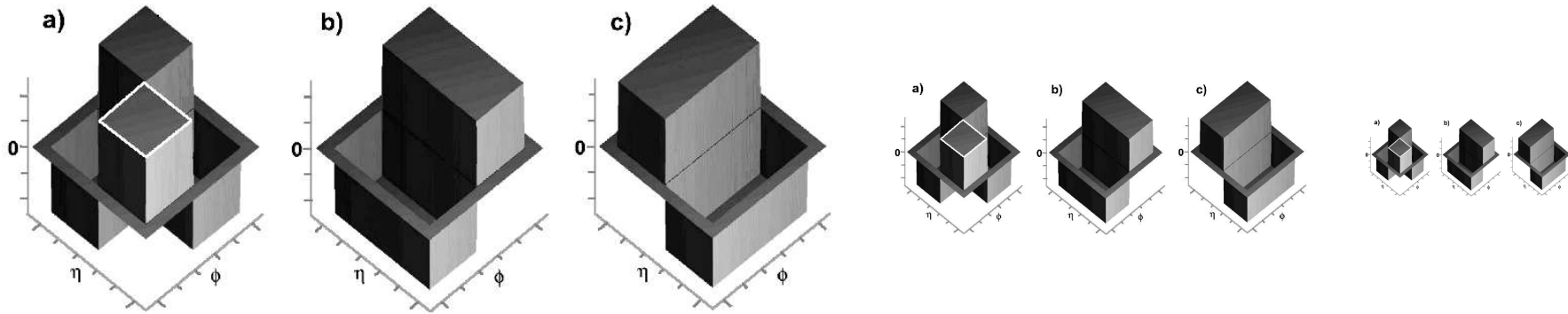
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Power of local fluctuations, mode λ :

$$P^\lambda(m) = 2^{-2m} \sum_{l,k} \langle \rho, F_{m,l,k}^\lambda \rangle^2 \quad (25)$$

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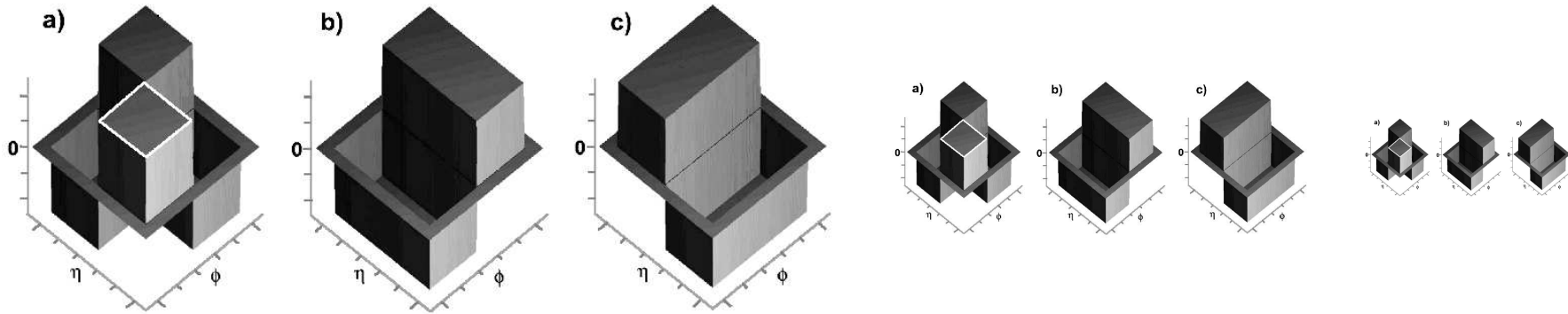
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“dynamic texture”:

$$P_{\text{dyn}}^\lambda(m) \equiv P_{\text{true}}^\lambda(m) - P_{\text{mix}}^\lambda(m) \quad (26)$$

40 Local hadron density fluctuations and Discrete Wavelet Transform (DWT)



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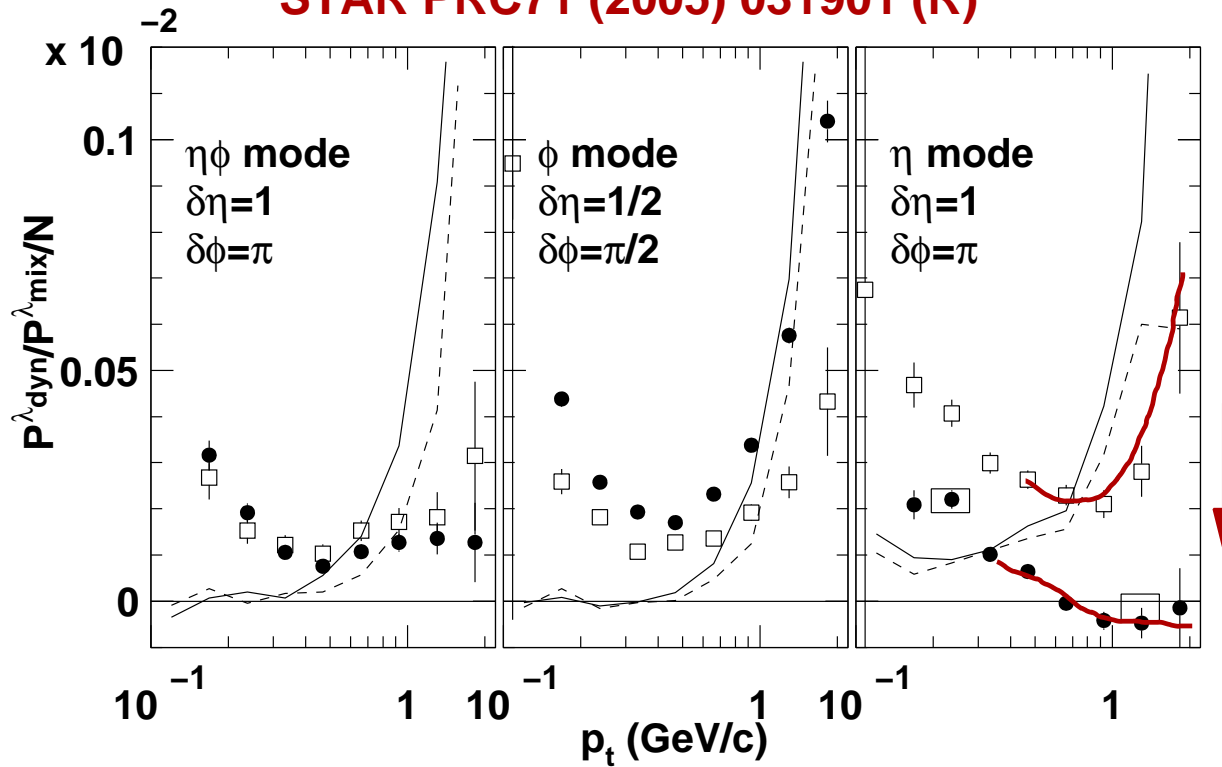
$$P_{\text{dyn}}^\lambda(m) \equiv P_{\text{true}}^\lambda(m) - P_{\text{mix}}^\lambda(m) \quad (26)$$

Normalized:

$$P_{\text{dyn}}^\lambda(m) / P_{\text{mix}}^\lambda(m) / n(p_t) \quad (27)$$

41 Longitudinal minijet broadening – wavelet-based technique

STAR PRC71 (2005) 031901 (R)



Central events: normalized dynamic texture for fineness scales $m = 0, 1, 0$ from left to right panels, respectively, as a function of p_t . ● STAR data; solid line – Hijing without jet quenching; dashed line – Hijing with quenching; □ peripheral STAR data renormalized to compare.

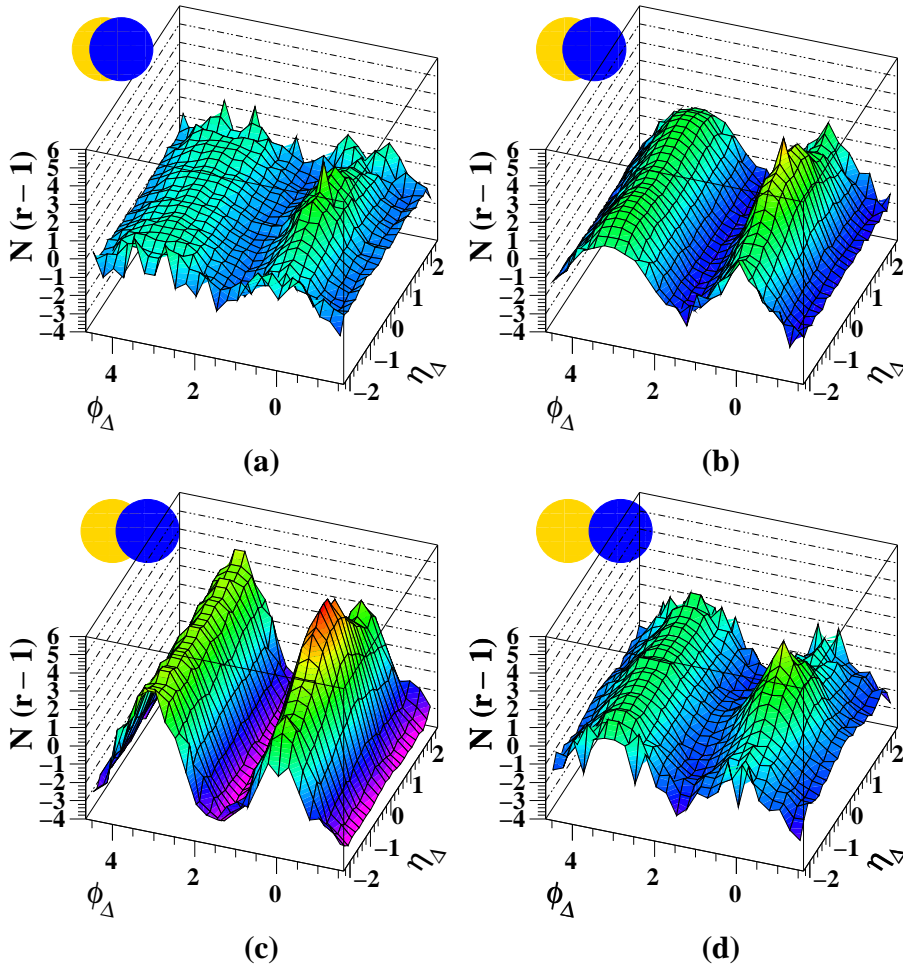
42 Longitudinal minijet broadening – traditional technique

$$r(\vec{p}_1, \vec{p}_2) \equiv \rho_{sib}(\vec{p}_1, \vec{p}_2) / \rho_{mix}(\vec{p}_1, \vec{p}_2). \quad (28)$$

42 Longitudinal minijet broadening – traditional technique

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STAR PRC73 (2006) 064907

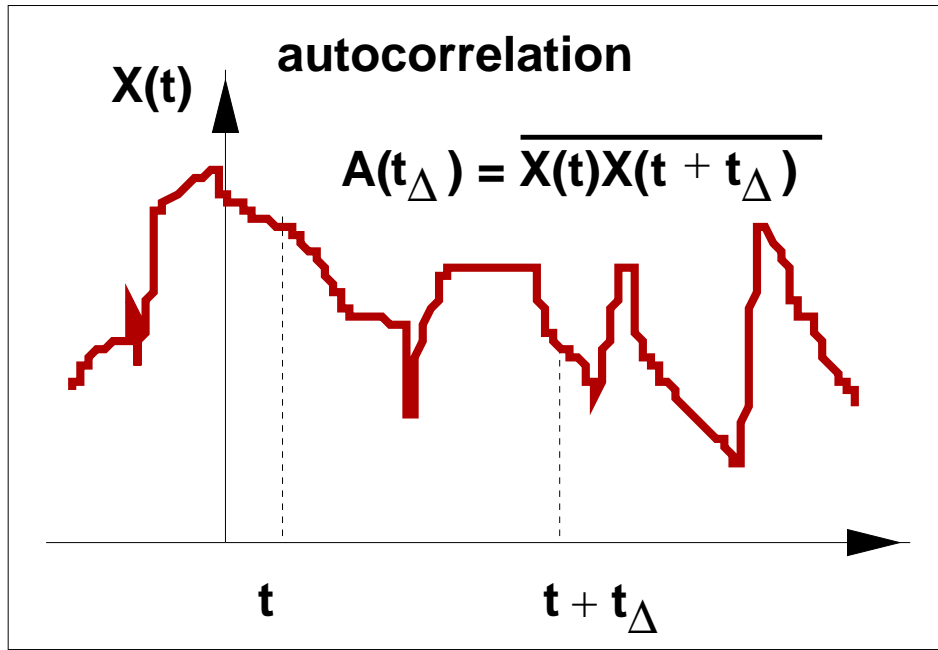


Two-particle charge-independent joint correlations $\bar{N}(\hat{r} - 1)$ on $(\eta_\Delta \equiv \eta_1 - \eta_2, \phi_\Delta \equiv \phi_1 - \phi_2)$ for central (a) to peripheral (d) collisions.

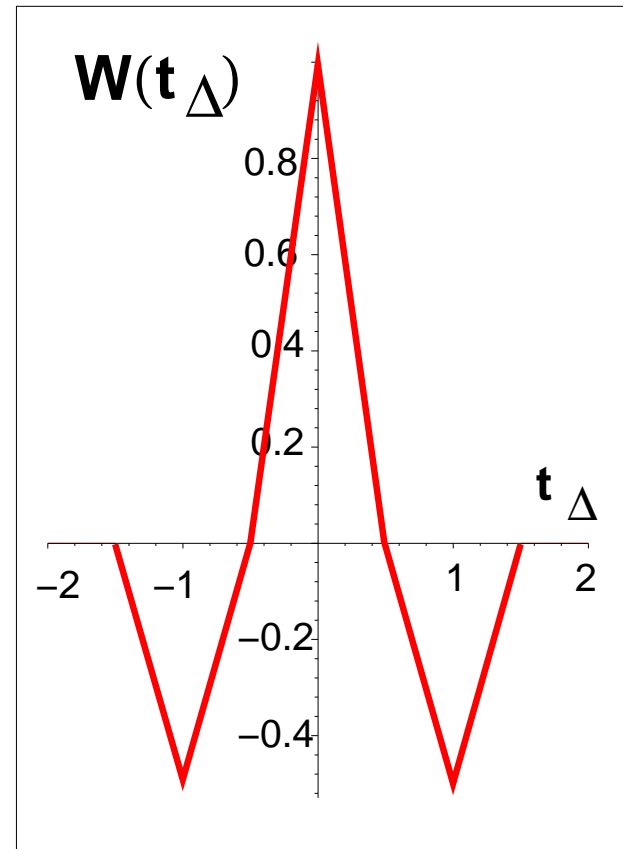
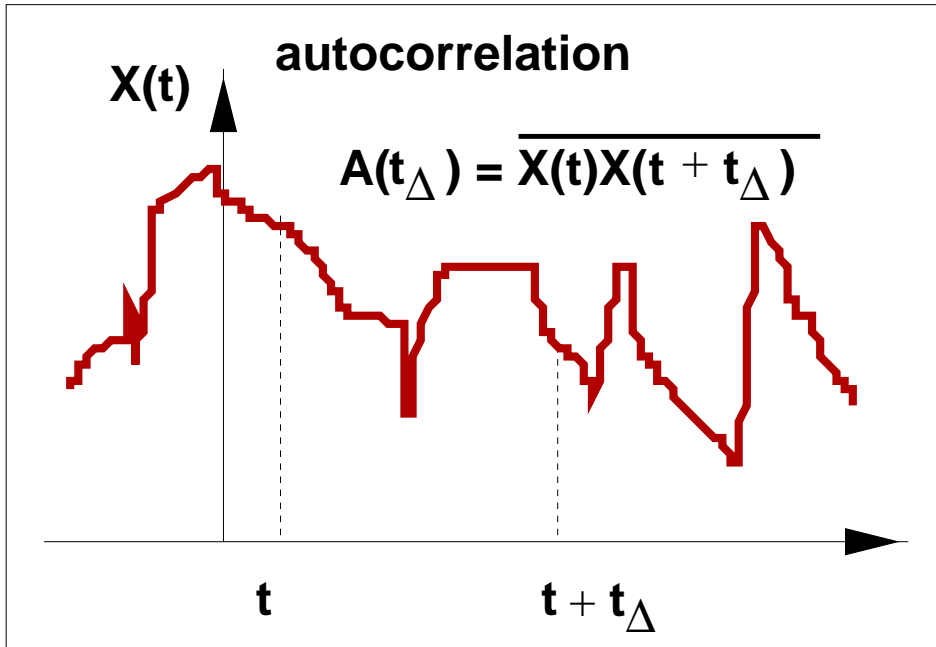
43 Connecting DWT and two-point correlation measures

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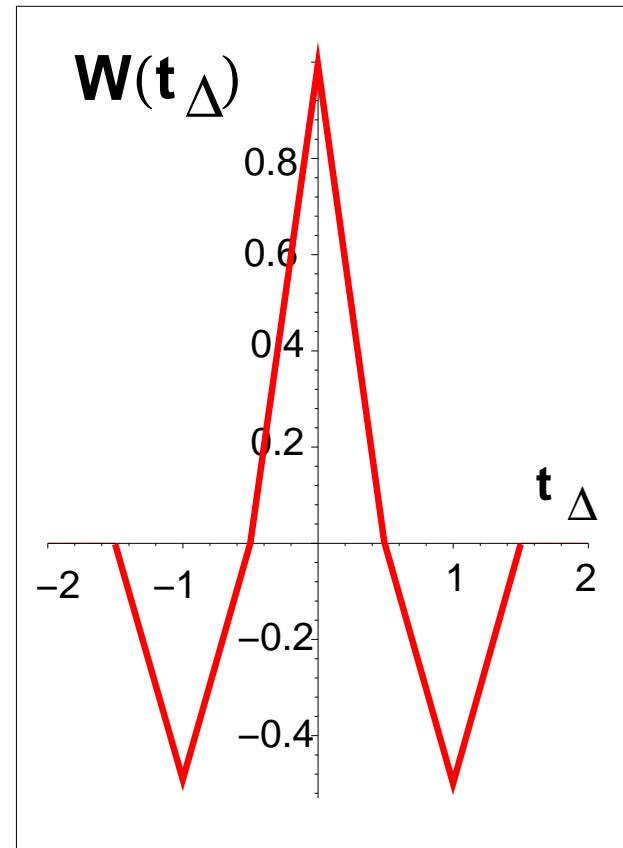
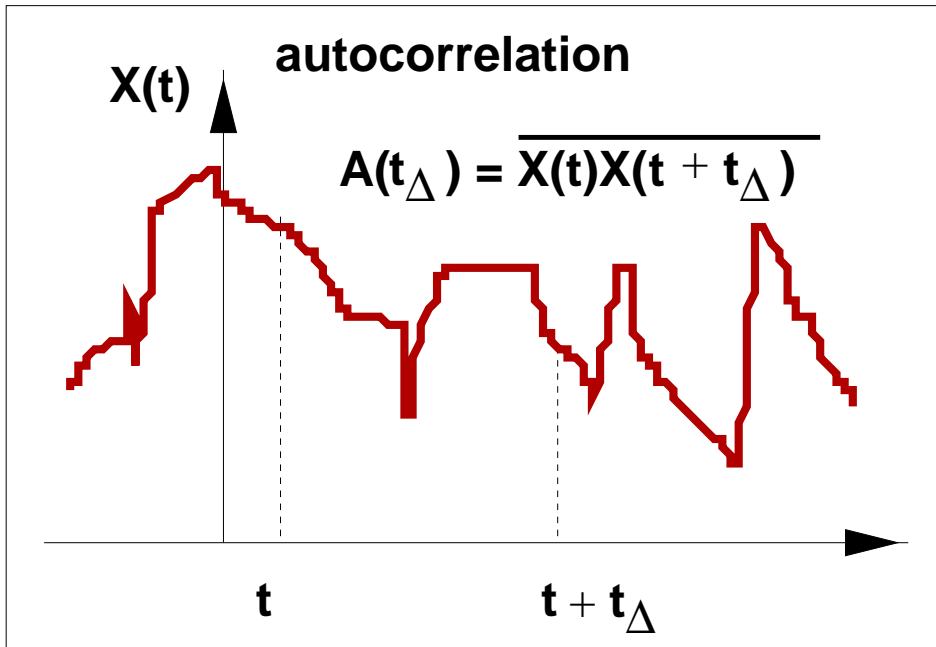
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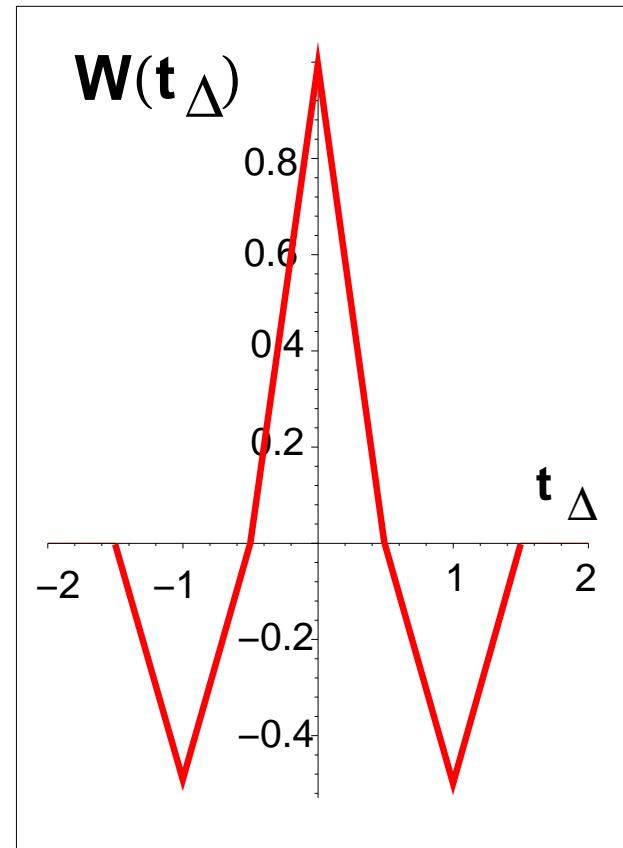
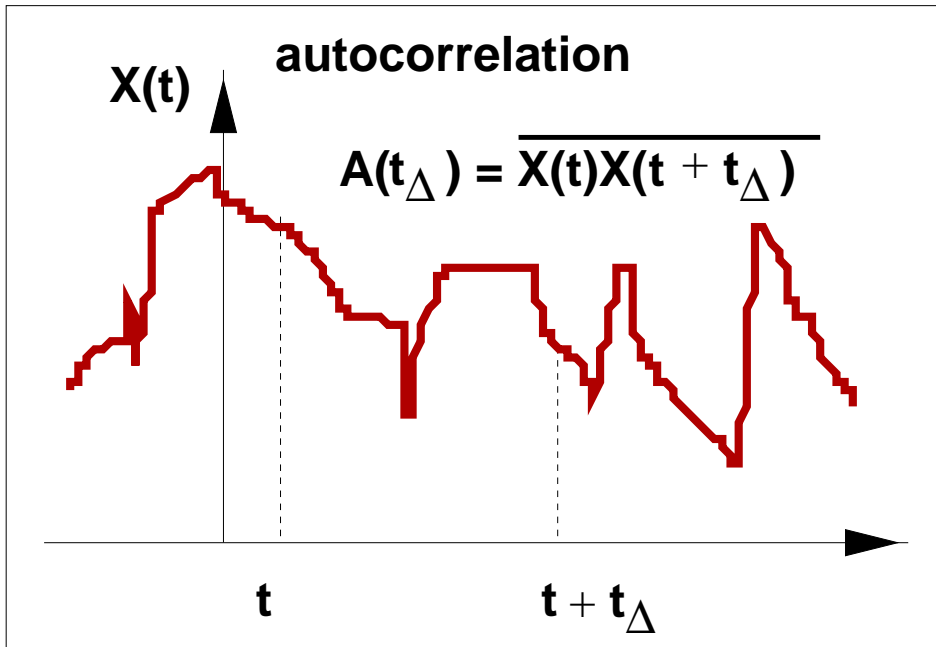


43 Connecting DWT and two-point correlation measures



$$P(m) = \int_{-\infty}^{\infty} X(t_{\Delta}/2)X(-t_{\Delta}/2)W(t_{\Delta}, m) dt_{\Delta}, \quad (29)$$

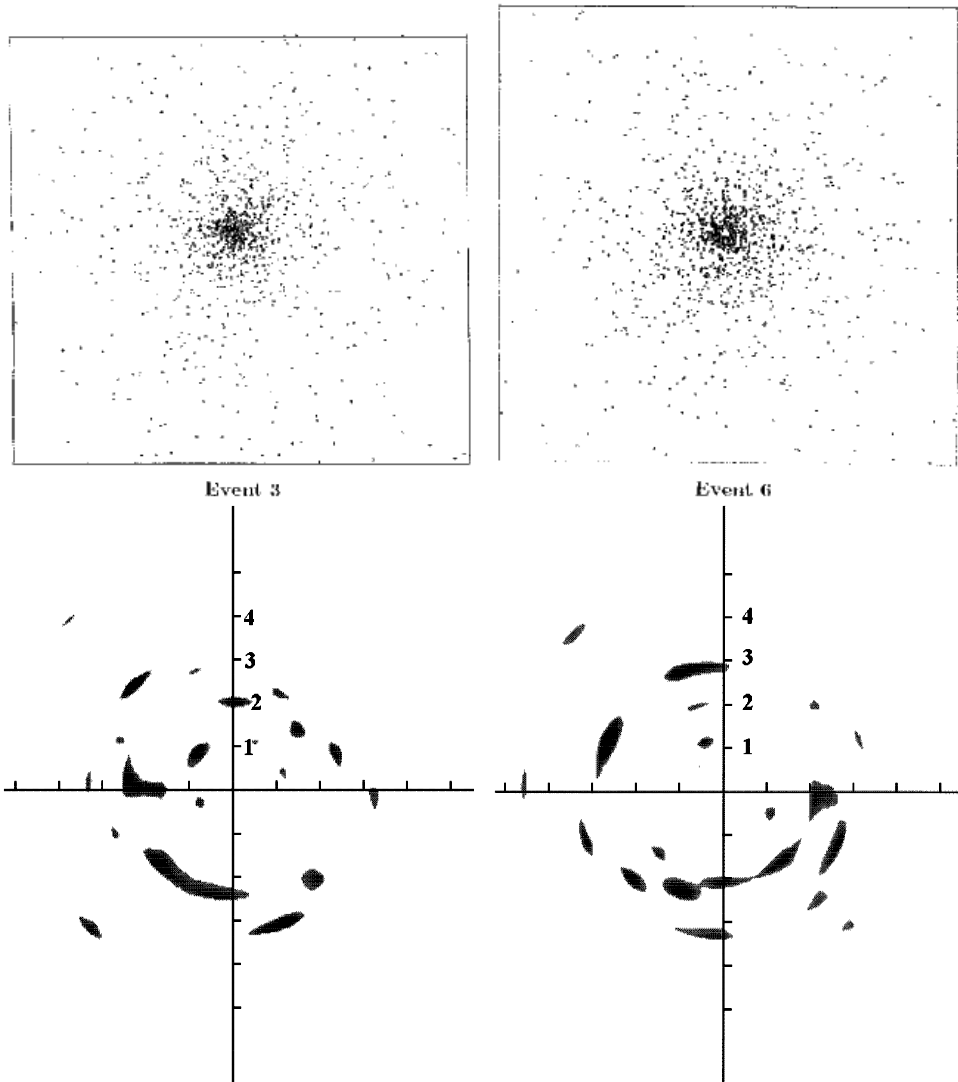
43 Connecting DWT and two-point correlation measures



$P(m)$ differentiates correlation on scale m . Minijet elongation \Rightarrow correlation broadening \Leftrightarrow reduced correlation gradient \Leftrightarrow reduced “texture”

$$P(m) = \overline{\int_{-\infty}^{\infty} X(t_{\Delta}/2)X(-t_{\Delta}/2)W(t_{\Delta}, m) dt_{\Delta}}, \quad (29)$$

44 Future of wavelet correlations: Kent — P.N.Lebedev — MEPhI project



Dremin et al.,
Phys.Lett.B499:97-103,2001.
Emulsion plates exposed at
SPS.

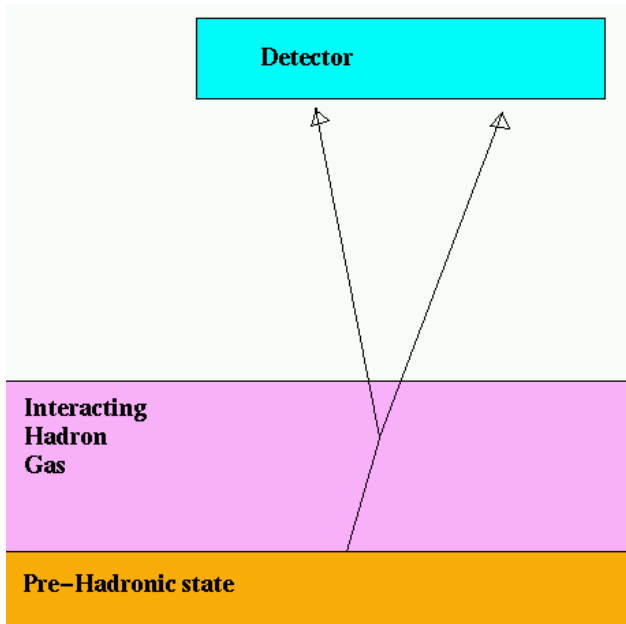
Top image → DWT →
suppress certain scales →
inverse DWT → bottom
image.

Rings of Cherenkov gluons ?
Could determine "dielectric
permeability" of QCD matter
at RHIC.

45 Resonances in hadronic matter

Markert, Torrieri, Rafelski, Campos do Jordao 2002

45 Resonances in hadronic matter



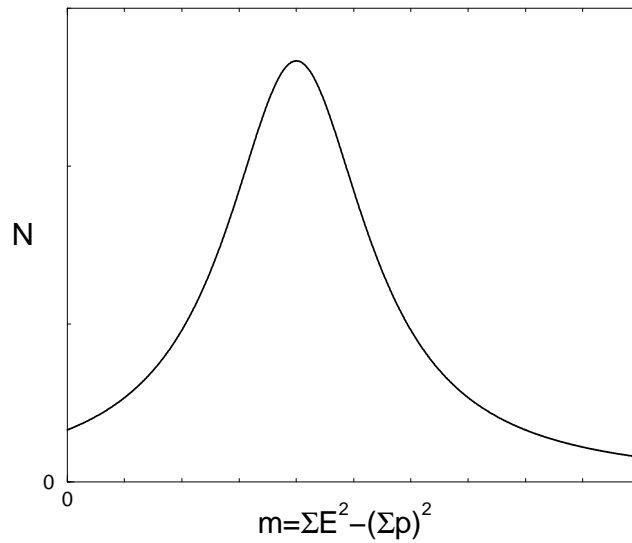
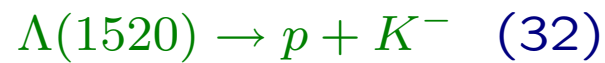
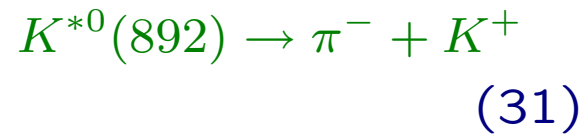
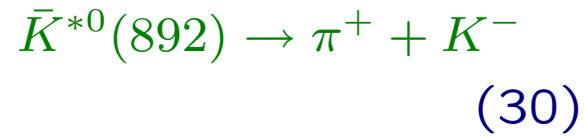
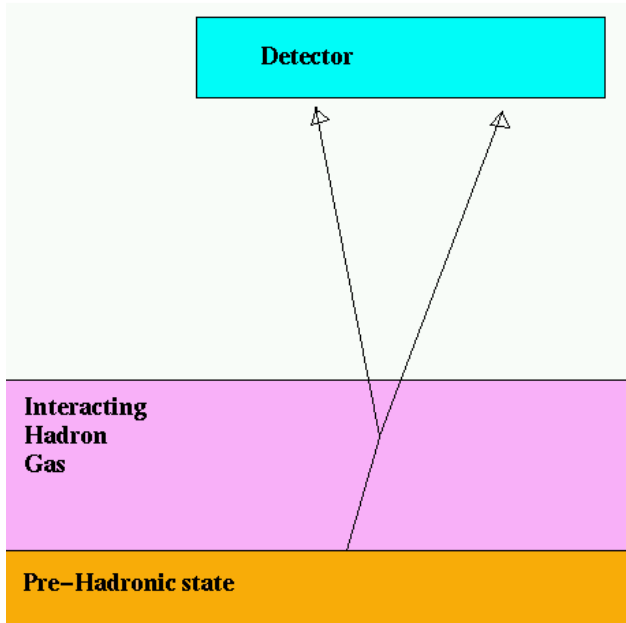
$$\bar{K}^{*0}(892) \rightarrow \pi^+ + K^- \quad (30)$$

$$K^{*0}(892) \rightarrow \pi^- + K^+ \quad (31)$$

$$\Lambda(1520) \rightarrow p + K^- \quad (32)$$

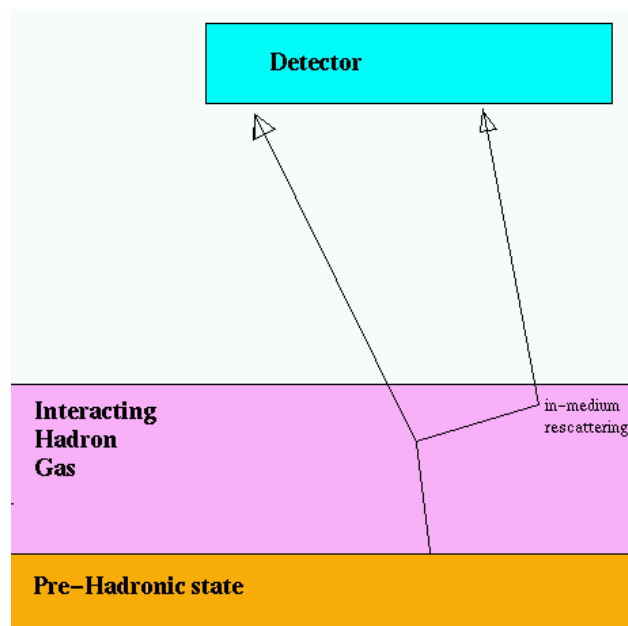
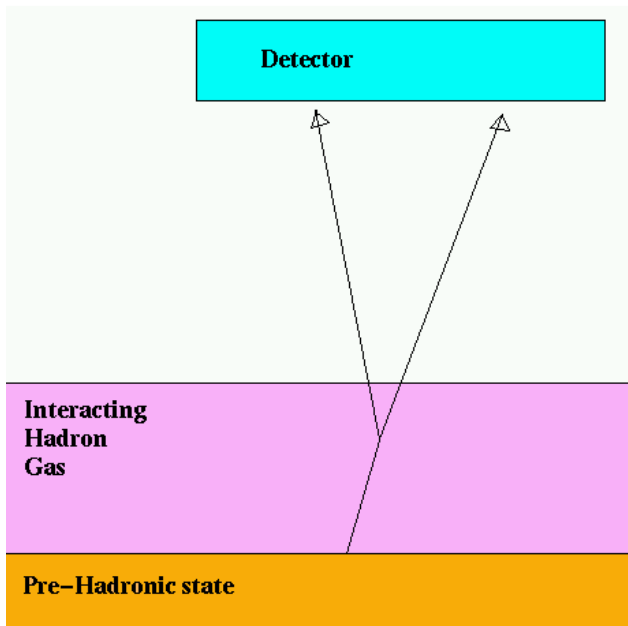
Markert, Torrieri, Rafelski, Campos do Jordao 2002

45 Resonances in hadronic matter



Markert, Torrieri, Rafelski, Campos do Jordao 2002

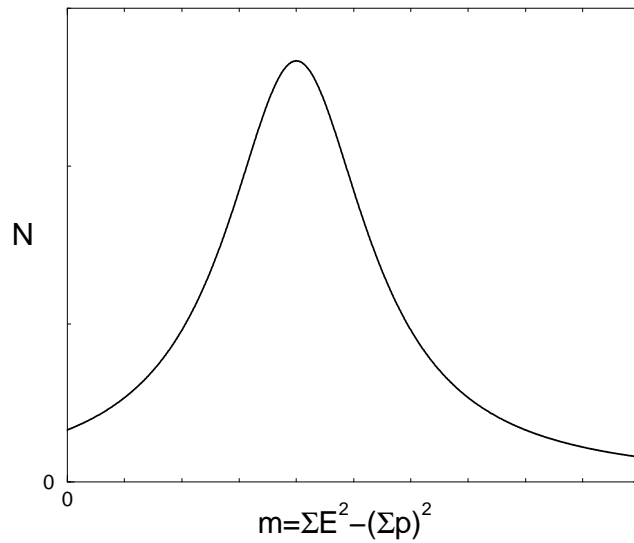
45 Resonances in hadronic matter



$$\bar{K}^{*0}(892) \rightarrow \pi^+ + K^- \quad (30)$$

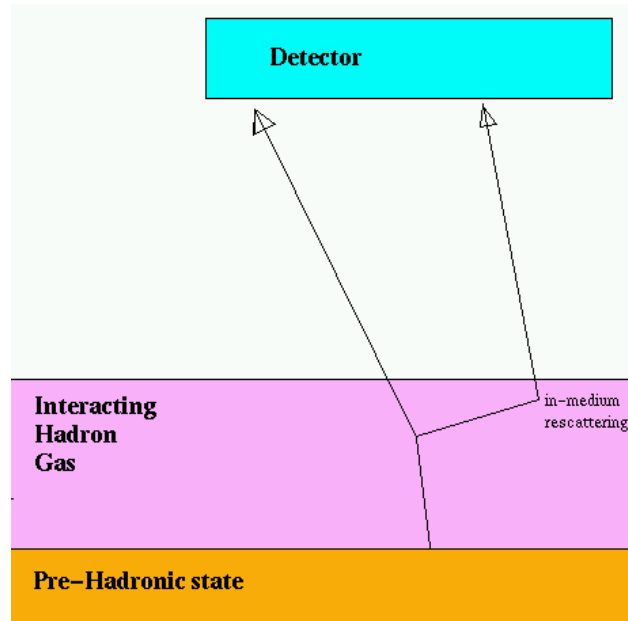
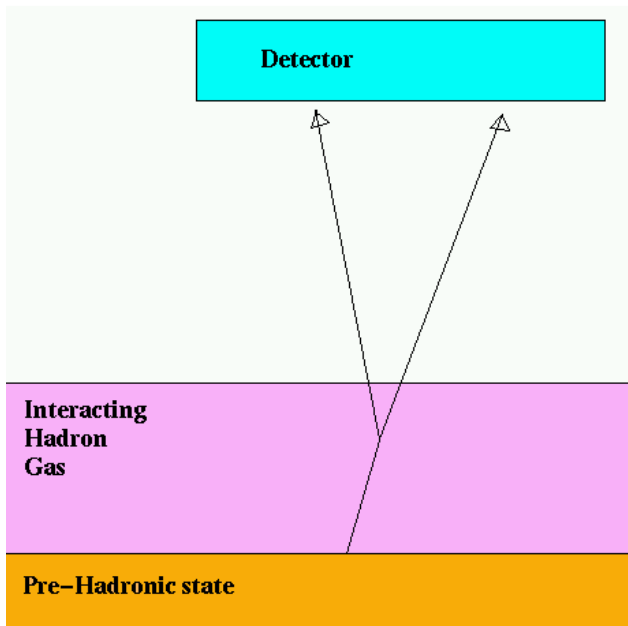
$$K^{*0}(892) \rightarrow \pi^- + K^+ \quad (31)$$

$$\Lambda(1520) \rightarrow p + K^- \quad (32)$$



Markert, Torrieri, Rafelski, Campos do Jordao 2002

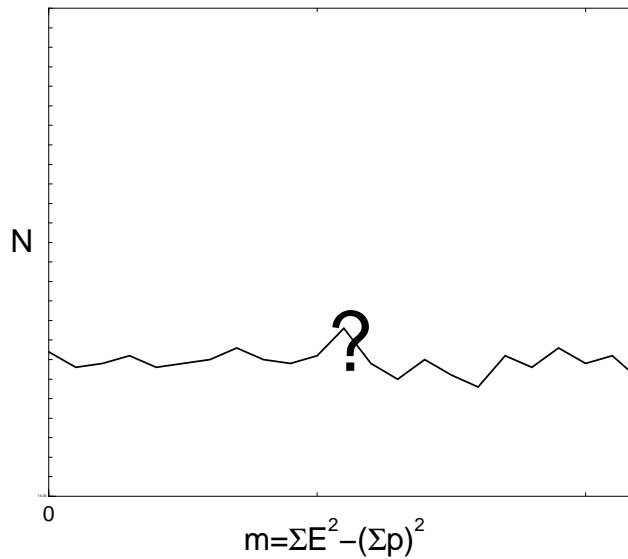
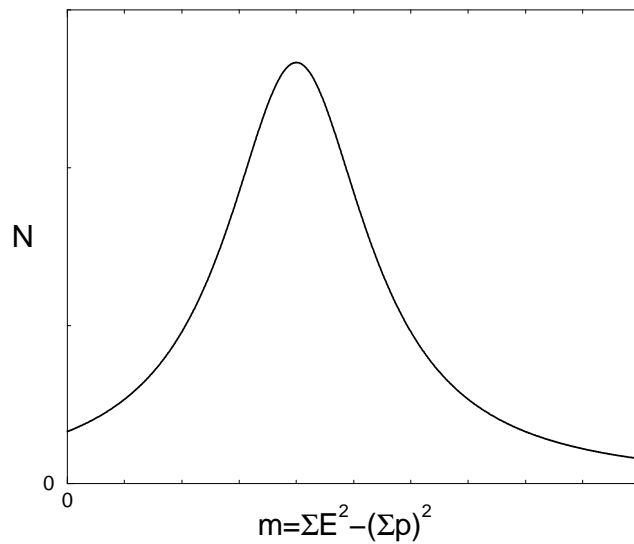
45 Resonances in hadronic matter



$$\bar{K}^{*0}(892) \rightarrow \pi^+ + K^- \quad (30)$$

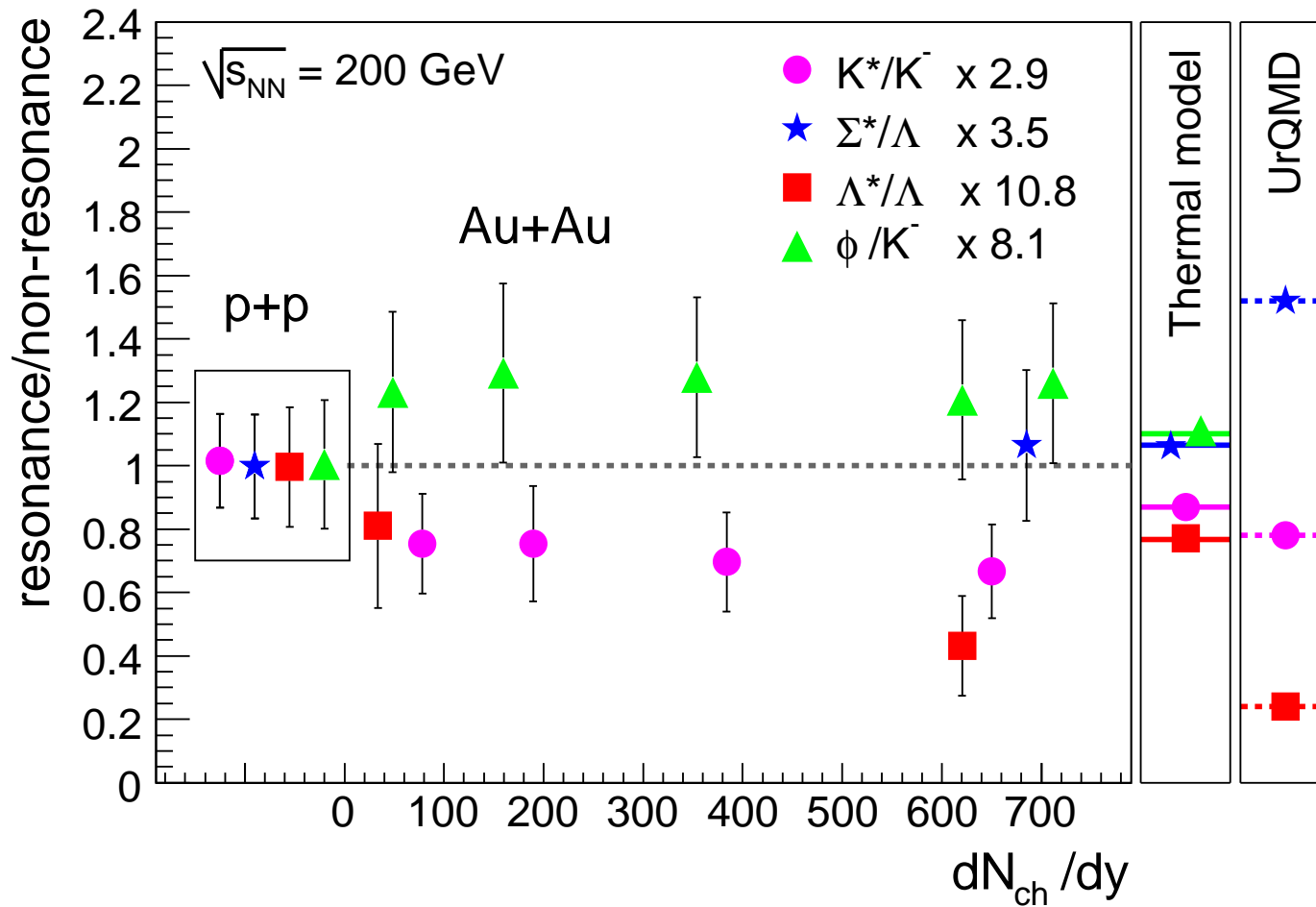
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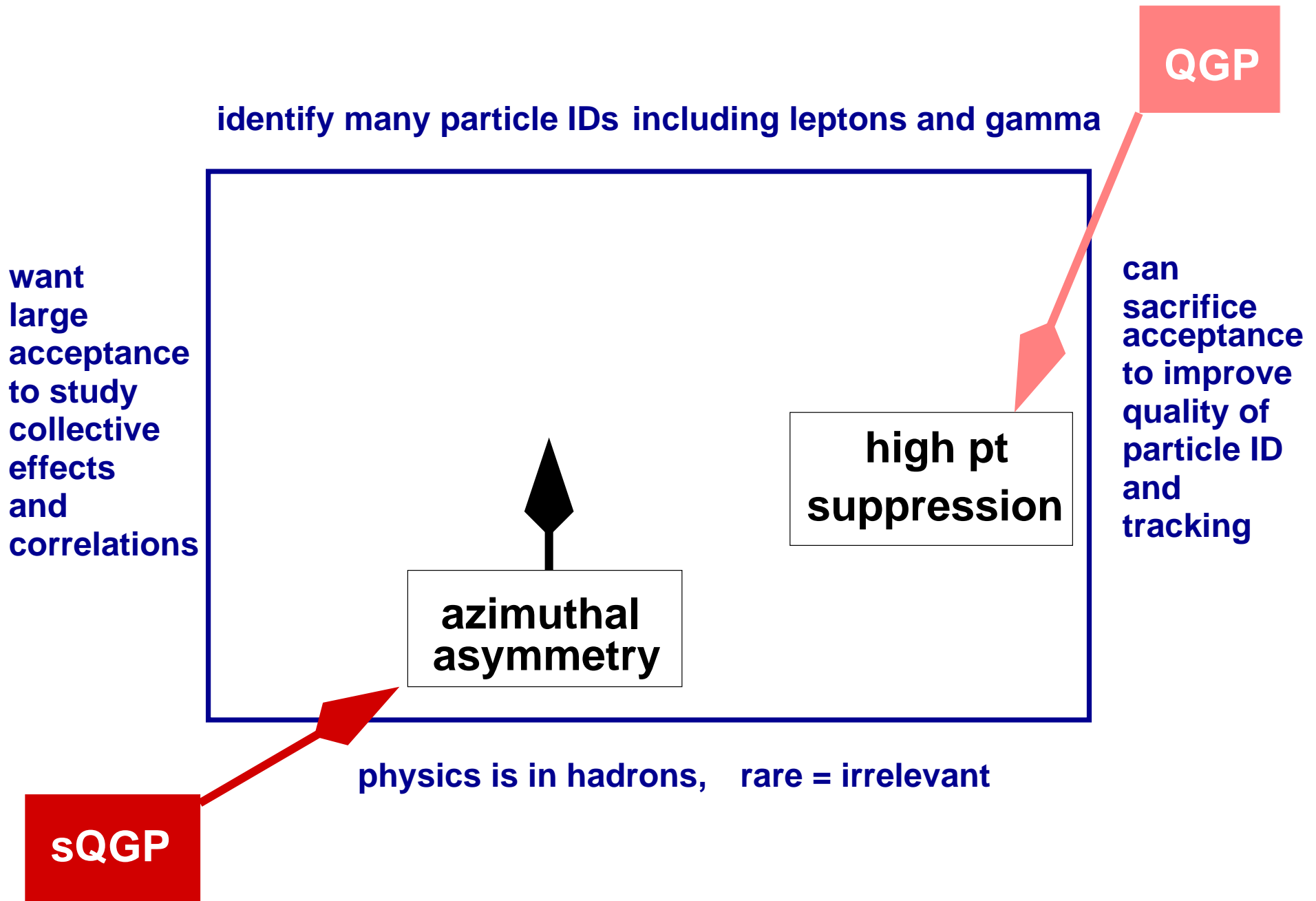
46 Resonances in hadronic matter



STAR PRL 97, 132301 (2006): a thermal model with rescattering. Can use models to put limits on the system evolution time.

47 QGP: physical reality or a justifiable ansatz?

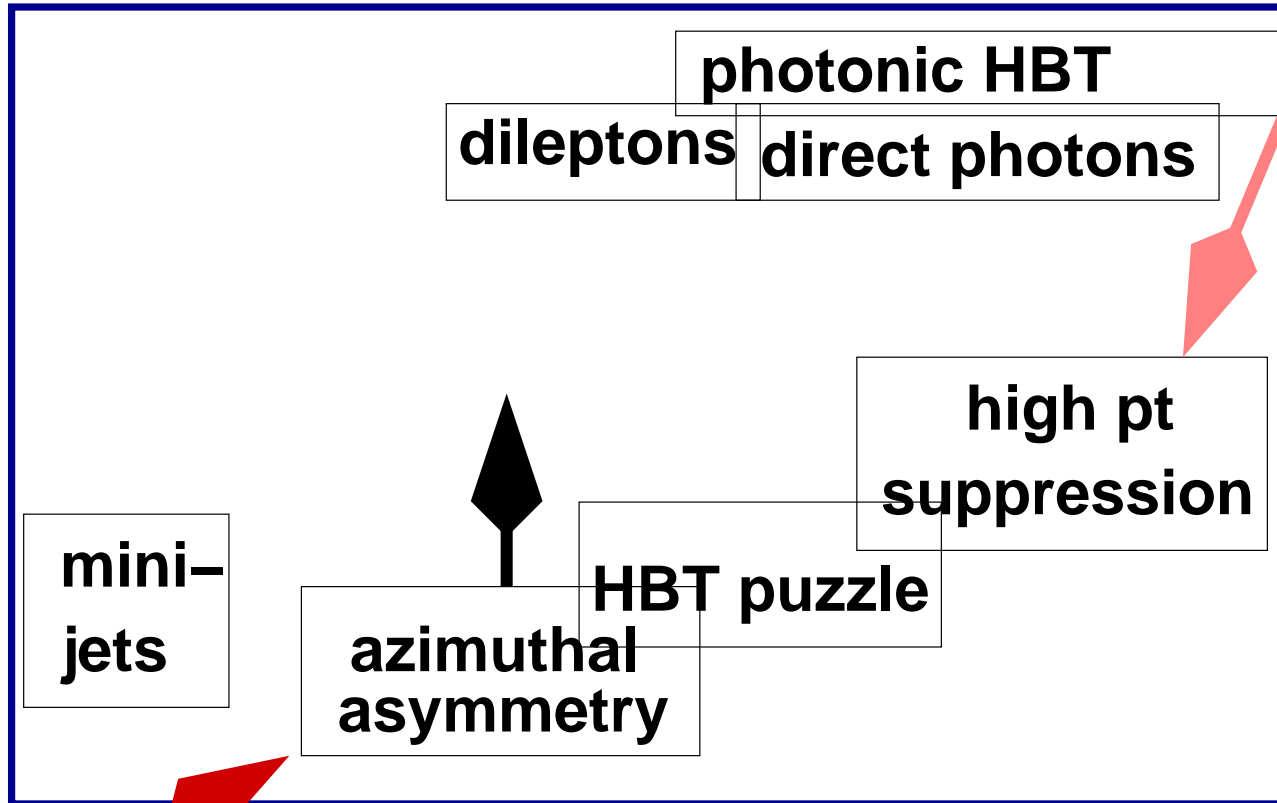
47 QGP: physical reality or a justifiable ansatz?



47 QGP: physical reality or a justifiable ansatz?

identify many particle IDs including leptons and gamma

want large acceptance to study collective effects and correlations



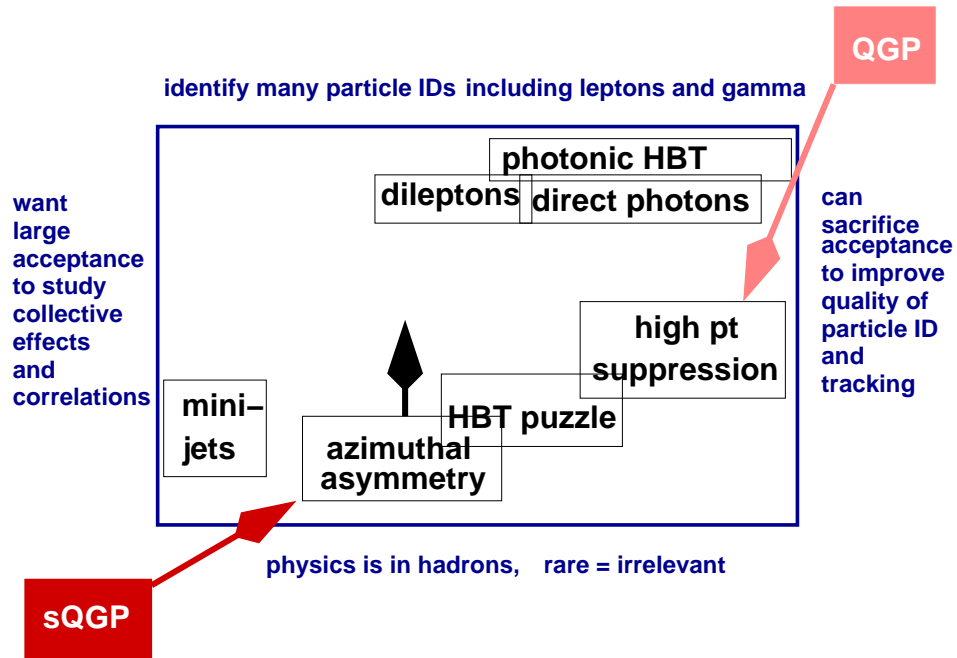
can sacrifice acceptance to improve quality of particle ID and tracking

physics is in hadrons, rare = irrelevant

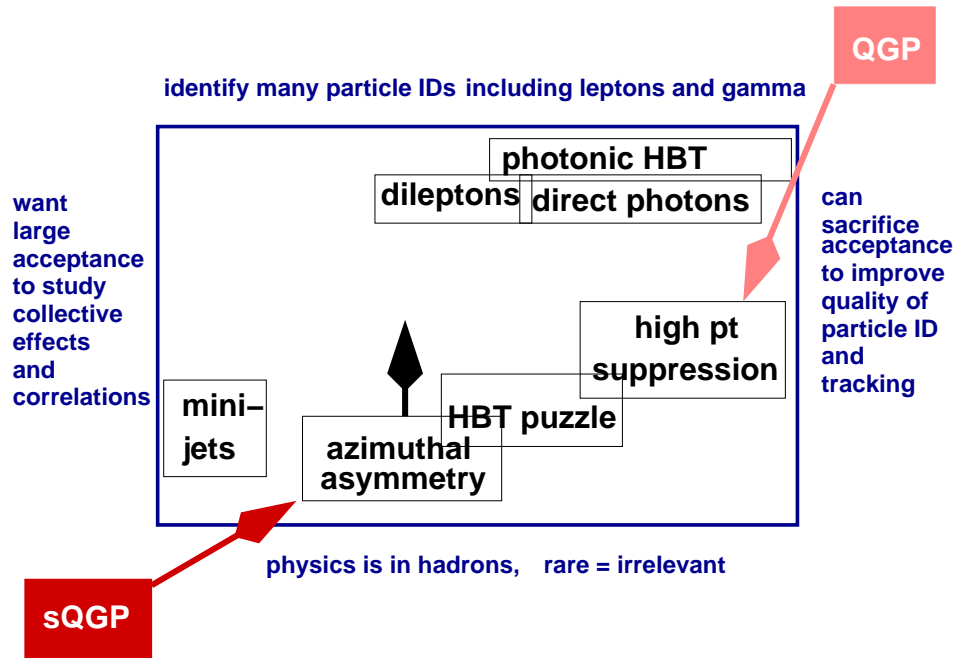
sQGP

QGP

48 So what happened at RHIC ? Conclusions so far:

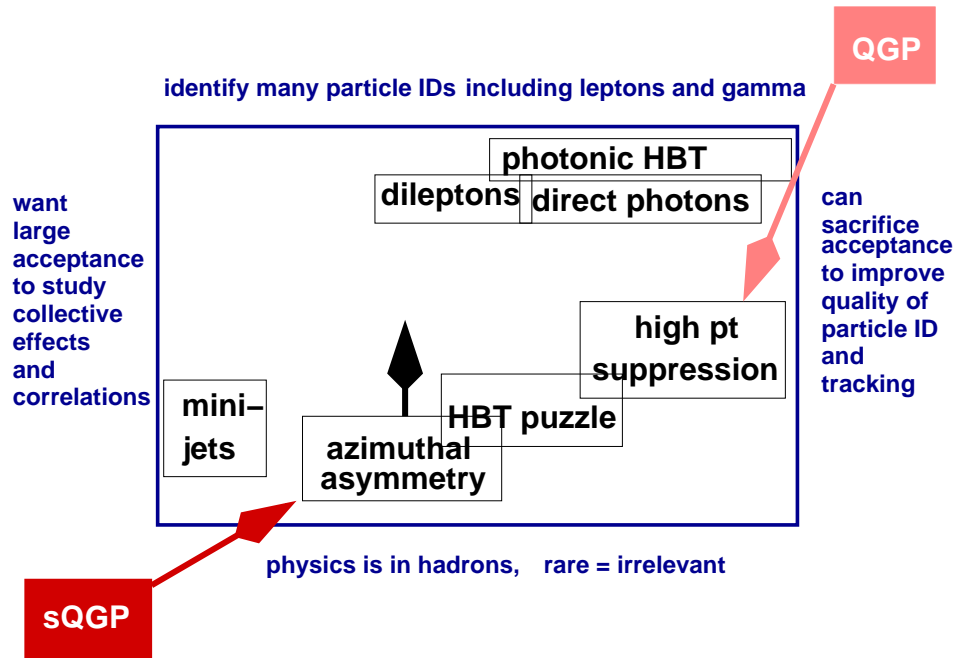


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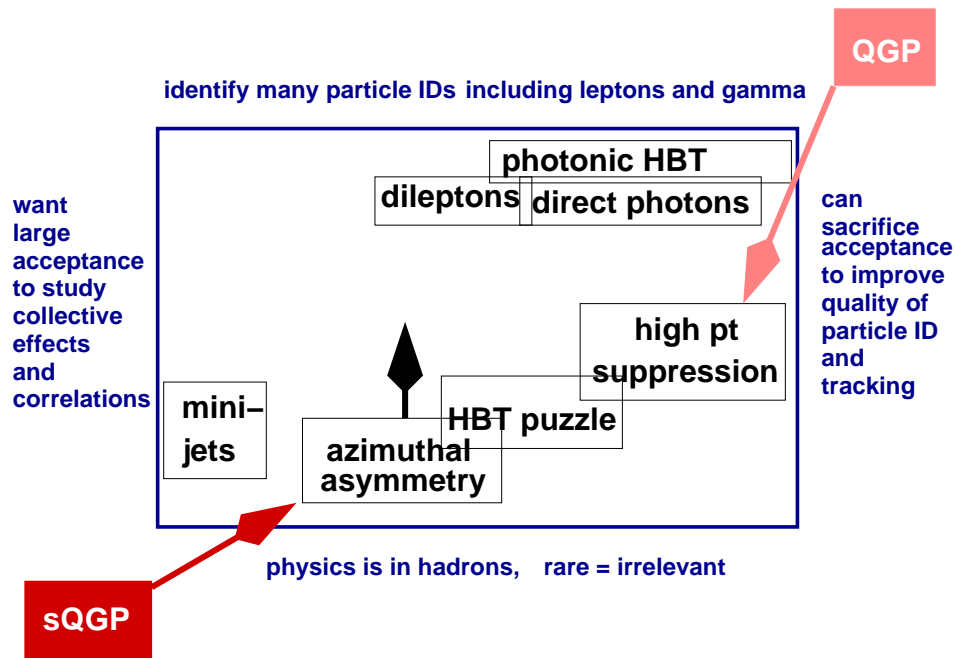
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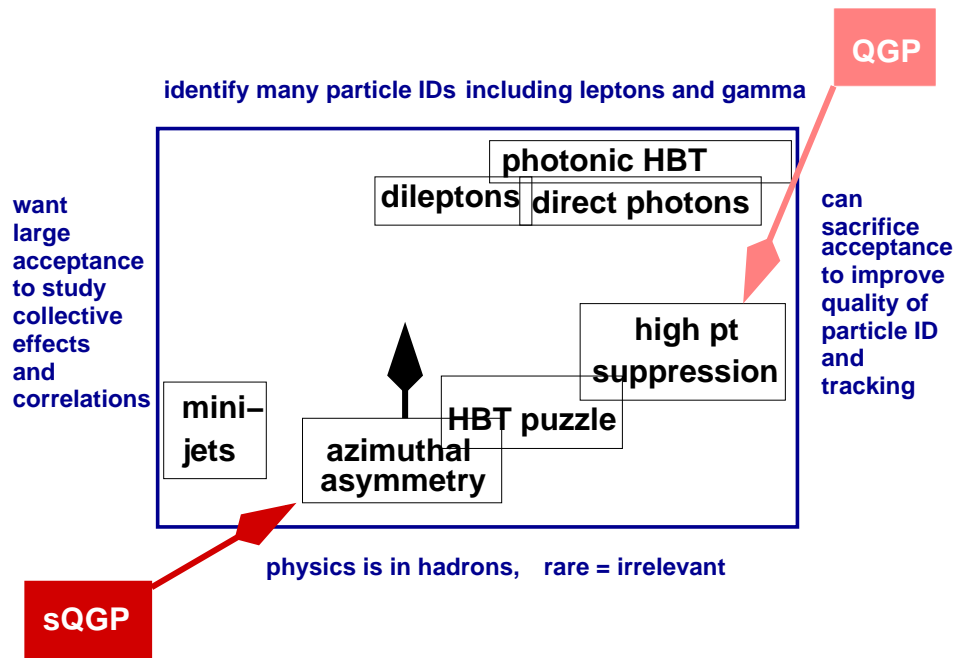
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48 So what happened at RHIC ? Conclusions so far:



- first phase of the campaign: unexpectedly, a lot of action is taking place on the South-Western front! (of minimum-bias hadronic correlations)
- QGP as a theory ansatz may have been **justified** by the data
- to elevate QGP to the status of a **discovered physical reality**, need to demonstrate uniqueness of the interpretations, embrace full gamut of the phenomena, avoid "confirmation bias".