## Correlation structure of STAR events

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## STAR Detector



## 1 Content of the talk

- Equilibration

Arguably the central issue of RHIC hadronic physics. Is it taking place ? What is the mechanism ? And what is equilibrating ?

- Methods

Initial state: a "calibrated" source of correlations. Watch their evolution into final state in time $=$ system size.

- Direct construction of a correlation function.
- Inversion of scale-dependent variance
- Discrete Wavelet Transform
- Observations
- Conclusions


## 2 Autocorrelation



$$
\binom{x_{1}}{x_{2}} \rightarrow\binom{x_{\Sigma} \equiv x_{1}+x_{2}}{x_{\Delta} \equiv x_{1}-x_{2}}
$$

always a lossless transformation of data. Autocorrelation $A$ is a projection of a two-point distribution onto difference variable(s) $x_{\Delta}$, lossless for $x_{\Sigma}$-invariant (homogenous, stationary) problems.

$$
\Delta R\left(x_{1}, x_{2}\right)=\frac{\rho\left(x_{1}, x_{2}\right)}{\rho_{\mathrm{ref}}\left(x_{1}, x_{2}\right)}-1
$$

3 Uncorrelated event reference for DWT

mixed events: no pixel used twice; $\leq 1$ pixel from any event in the same mixed event; no mixing of events with largely different multiplicity and vertex.

true

mixed

## 4 Local hadron density fluctuations and Discrete Wavelet Transform (DWT)


$F_{m, l, k}^{\lambda}(\phi, \eta)$-Haar wavelet orthonormal basis in $(\phi, \eta)$. scale fineness $(m)$, directional modes of sensitivity ( $\lambda$ ), track density $\rho\left(\eta, \phi, p_{t}\right)$, locations in 2D $(l, k)$. DWT is an expansion in this basis.
Power of local fluctuations, mode $\lambda$ :

$$
\begin{equation*}
P^{\lambda}(m)=2^{-2 m} \sum_{l, k}\left\langle\rho, F_{m, l, k}^{\lambda}\right\rangle^{2} \tag{1}
\end{equation*}
$$

"dynamic texture":

Normalized:

$$
\begin{equation*}
P_{\mathrm{dyn}}^{\lambda}(m) \equiv P_{\mathrm{true}}^{\lambda}(m)-P_{\mathrm{mix}}^{\lambda}(m) \tag{2}
\end{equation*}
$$

$$
\begin{equation*}
P_{\mathrm{dyn}}^{\lambda}(m) / P_{\text {mix }}^{\lambda}(m) / n\left(p_{t}\right) \tag{3}
\end{equation*}
$$

## 5 A flow-like example



Elliptic flow-inspired example: $x$ axis - an angle in "natural units" $(2 \pi=1), y$ axis multiplicity. The multiresolution theorem: a4
$=\mathrm{a} 0+\mathrm{b} 0+\mathrm{b} 1+\mathrm{b} 2+\mathrm{b} 3$, can have better fineness.

## 6 Example of a DWT power spectrum



Power spectrum of that flow event as a function of "fineness" $m$. The dominant contrubution is $m=1$ (the " $v_{2}$ " harmonic, b1). Statistical fluctuations also contribute.
$P(m)=2^{-m} \sum_{i}\left\langle\rho, F_{m, i}\right\rangle^{2}$.
Computational complexity $O(N)$ !

## 7 "Dynamic texture" as a nonparametric measure of the correlation shape



8 "Dynamic texture" $p_{t}$ dependence: peripheral events, $\sqrt{s_{N N}}=200 \mathbf{G e V}$


Peripheral (60-84\%) events: normalized dynamic texture for fineness scales $m=0,1,0$ from left to right panels, respectively, as a function of $p_{t}$. - - STAR data; solid line standard HIJING, dash-dotted line - HIJING without jets.
Qualitative trends in peripheral data are as expected. What signal to expect in the central data, if correlation does not change ?

$$
\begin{equation*}
\left.\left(\frac{P_{\text {true }}}{P_{\text {mix }}}-1\right) \frac{1}{N}\right|_{\text {centr }}=\left.\left(\frac{P_{\text {true }}}{P_{\text {mix }}}-1\right)\right|_{\text {periph }} \frac{1}{N_{\text {centr }}} \tag{4}
\end{equation*}
$$

9 Longitudinal minijet broadening: DWT data


Central (top 4\%) events: normalized dynamic texture for fineness scales $m=0,1,0$ from left to right panels, respectively, as a function of $p_{t}$.

- STAR data; solid line - Hijing without jet quenching; dashed line Hijing with quenching; $\square$ peripheral STAR data renormalized to compare. Minijet elongation $\Rightarrow$ correlation broadening $\Leftrightarrow$ reduced correlation gradient $\Leftrightarrow$ reduced "texture"


## 10 "Dynamic texture" response



Dynamic texture response in various idealized situations (showing only one scale):
(a) events of random (uncorrelated) particles
(b) $p_{t}$-independent elliptic flow
(c) Correlations at Iow $Q_{\text {inv }}$ (Bose-Einstein correlations and Coulomb effect)
(d) HIJING jets

## 11 Longitudinal minijet broadening: correlation data




 Projections of $\left.\bar{N}\left[\rho\left(\eta_{\Delta}, \phi_{\Delta}\right) / \rho\left(\eta_{\Delta}, \phi_{\Delta}\right)_{\text {ref }}-1\right]\right|_{C I}$ on $x_{\Delta}$ which is $\phi_{\Delta}$ ( $\Delta$ ) or $\eta_{\Delta}$ ( ) . $v_{1}$ and $v_{2}$ are subtracted.

12 Longitudinal minijet broadening: centrality dependence


## 13 Charge-dependent correlations $=$ Like sign - Unlike sign



The driving physics: charge conservation in hadronization. Suppress short range correlations - BEC and conversion $e^{+} e^{-}$- by a kinematic pair cut. The $\bar{N} \times$ is good when number of correlation sources $\propto N$.

14 Modified hadronization geometry ?

(d)

(c)

(b)


Projections of $\left.\bar{N}\left(\rho\left(\eta_{\Delta}, \phi_{\Delta}\right) / \rho\left(\eta_{\Delta}, \phi_{\Delta}\right)_{\text {ref }}-1\right)\right|_{C D}$ on $x_{\Delta}$ which is $\phi_{\Delta}(\Delta)$ or $\eta_{\Delta}(\bullet) . \eta-\phi$ width disparity (d, peripheral) is gone in (a) $\Rightarrow$ transition from (string) 1D to bulk (>2D) fragmentation symmetrizes $\eta$ and $\phi$.


## 15 Number and $p_{t}$ correlations



This $p_{t}$ field may have elliptic flow (number effect). Abounds at RHIC.

Also elliptic... flow ( $p_{t}$ effect) ! Pro: blast wave fits. Is there a direct measurement?

## 16 Towards $p_{t}$ correlation/fluctuation analysis

Problem: need to tell apart $p_{t, i}$ and number contributions to the $p_{t} \equiv \sum_{i \in(\eta, \phi) \text { bin }} p_{t, i} \Rightarrow$ can extract the $p_{t}$ correlation alone.

Solution: use $p_{t}-n \hat{p_{t}}$
$\mathrm{Q}:$ When is the $n$-contribution into $\operatorname{Var}\left[p_{t}-n \hat{p_{t}}\right]$ canceled ?

$$
\begin{equation*}
\sigma^{2}\left(p_{t}: n\right) \equiv \operatorname{Var}\left[p_{t}-n \hat{p_{t}}\right]=\operatorname{Var}\left[p_{t}\right]+{\hat{p_{t}}}^{2} \operatorname{Var}[n]-2 \hat{p_{t}} \operatorname{Cov}\left[n, p_{t}\right] \tag{5}
\end{equation*}
$$

$$
\begin{equation*}
\operatorname{Var}\left[p_{t}\right]=\operatorname{Var}\left[\sum_{i}^{n} p_{t, i}\right]=\operatorname{Var}\left[\sum_{i}^{n}\left(\hat{p_{t}}+u_{i}\right)\right]={\hat{p_{t}}}^{2} \operatorname{Var}[n]+\operatorname{Var}[u]+2 \hat{p_{t}} \operatorname{Cov}[n, u] \tag{6}
\end{equation*}
$$

$$
\begin{equation*}
\operatorname{Cov}\left[n, p_{t}\right]=\overline{n p_{t}}-\bar{n} \overline{p_{t}}=\hat{p_{t}} \operatorname{Var}[n] \tag{7}
\end{equation*}
$$

A:For independent $p_{t}$ and $n$ production, when $\operatorname{Cov}[n, u] \equiv \overline{n u}=0$, where $u \equiv \sum_{i}^{n} u_{i}, u_{i}=p_{t, i}-\hat{p_{t}}$.


17 Get correlations from fluctuations

Extract correlation structure of random field $X$ from the scale dependence of variance (van Marcke "Random Fields" MIT 1983;
Trainor,Porter,Prindle hep-ph/0410180)


$$
\begin{align*}
\operatorname{Var}[X ; \delta \eta, \delta \phi]= & \int_{-\delta \eta / 2}^{\delta \eta / 2} d \eta_{1} \int_{-\delta \phi / 2}^{\delta \phi / 2} d \phi_{1} \int_{-\delta \eta / 2}^{\delta \eta / 2} d \eta_{2} \int_{-\delta \phi / 2}^{\delta \phi / 2} d \phi_{2}  \tag{8}\\
& \times\left[\overline{X\left(\eta_{1}, \phi_{1}\right) X\left(\eta_{2}, \phi_{2}\right)}-\overline{X\left(\eta_{1}, \phi_{1}\right)} \times \overline{X\left(\eta_{2}, \phi_{2}\right)}\right]
\end{align*}
$$

Compare with uncorrelated reference; recognize autocorrelation $\rho\left(X, t_{\Delta}\right) \equiv \overline{X(t) X\left(t+t_{\Delta}\right)}$ (t-average).

$$
\begin{equation*}
\Delta \sigma^{2}(X, \delta \eta, \delta \phi)= \tag{9}
\end{equation*}
$$

$$
\begin{array}{r}
\int_{-\delta \eta / 2}^{\delta \eta / 2} d \eta_{1} \int_{-\delta \phi / 2}^{\delta \phi / 2} d \phi_{1} \int_{-\delta \eta / 2}^{\delta \eta / 2} d \eta_{2} \int_{-\delta \phi / 2}^{\delta \phi / 2} d \phi_{2} \Delta \rho\left(X, \eta_{1}-\eta_{2}, \phi_{1}-\phi_{2}\right) \\
\quad=2 \int_{0}^{\delta \eta} d \eta_{\Delta} 2 \int_{0}^{\delta \phi} d \phi_{\Delta}\left(\delta \eta-\eta_{\Delta}\right)\left(\delta \phi-\phi_{\Delta}\right) \Delta \rho\left(X, \eta_{\Delta}, \phi_{\Delta}\right) \tag{11}
\end{array}
$$

18 The actual analysis is discrete: $\int \rightarrow \sum$

kernel $K$ :

$$
\begin{equation*}
\left(\delta \eta-\eta_{\Delta}\right)\left(\delta \phi-\phi_{\Delta}\right) \rightarrow \varepsilon_{\eta} \varepsilon_{\phi} K_{m_{\delta} n_{\delta}: k l} \equiv \varepsilon_{\eta} \varepsilon_{\phi}\left(m_{\delta}-k+\frac{1}{2}\right)\left(n_{\delta}-l+\frac{1}{2}\right) \tag{12}
\end{equation*}
$$

reference density $\rho_{\text {ref }}$ makes a per-particle measure:

$$
\begin{gather*}
\rho_{\mathrm{ref}} \propto \bar{n}^{2} \Rightarrow \frac{1}{\sqrt{\rho_{\mathrm{ref}}}} \propto \frac{1}{\bar{n}}  \tag{13}\\
\Delta \sigma_{p_{t}: n}^{2}\left(m_{\delta} \varepsilon_{\eta}, n_{\delta} \varepsilon_{\phi}\right)=4 \sum_{k, l=1}^{m_{\delta}, n_{\delta}} \varepsilon_{\eta} \varepsilon_{\phi} K_{m_{\delta} n_{\delta}: k l} \frac{\Delta \rho\left(p_{t}: n ; k \varepsilon_{\eta}, l \varepsilon_{\phi}\right)}{\sqrt{\rho_{\mathrm{ref}}\left(n ; k \varepsilon_{\eta}, l \varepsilon_{\phi}\right)}} \tag{14}
\end{gather*}
$$

Inverse problem: knowing $\Delta \sigma^{2}$, solve for $\Delta \rho / \sqrt{\rho_{\text {ref }}} \Rightarrow$ save $O(N)$ in CPU time !

## $19 p_{t}$ correlations from the inversion




Top:
scale dependence of the "pure" $p_{t}$ variance.


First direct evidence of elliptic flow as a $p_{t}$ blast. Next, subtract the flow contribution to look at minijets.

## Bottom:

corresponding
autocorrelation

## 20 Localized $p_{t}$ correlations: minijets



STAR preliminary






SSTAR
AuAu 200 GeV . In $\eta$, correlation broadens with centrality; in $\phi$ the trend is opposite. The surrounding background seems to recoil.

## 21 Conclusions

- Semi-hard scattering leaves a trace in the soft $p_{t}$ domain - new at RHIC !
- First direct measurements of $p_{t}$ correlation structure reveal azimuthal anisotropy of $p_{t}$ field $\Rightarrow$ elliptic flow is a velocity phenomenon
- The minijet correlation structure is modified with centrality; the effect appears to "turn on" around $\nu=\left(N_{\text {part }} / 2\right)^{1 / 3} \approx 3$. Broadening of the correlation in $\eta$ and weakening of $P_{\text {dyn }}^{\eta}$ on the coarse scale are consistent descriptions of the effect. How exactly does the coupling between longitudinal flow and mini-jets work ? What do we learn about the expanding fluid ?
- Increased symmetry of the charge-dependent correlation on $(\eta, \phi)$ in the central collisions may point to a change in the hadronization geometry in the medium

22 Extra slides

## 23 DWT of a photographic image



Reproduced from textbook: I.Daubechies, "Ten lectures on wavelets". The original caption: "A real image, and its wavelet decomposition into three multiresolution layers. On the wavelet components one clearly sees that the $d^{j, v}, d^{j, h}, d^{j, d}$ emphasize, respectively, vertical, horizontal, and diagonal edges. In this figure, the bottom picture has been overexposed to make details in the $d^{j, \lambda}$ more apparent. I would like to thank M.Barlaud for providing this figure." The colored marks are mine.

