Correlation structure of STAR events

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STAR Detector



- **1** Content of the talk
 - Equilibration

Arguably the central issue of RHIC hadronic physics. Is it taking place ? What is the mechanism ? And what is equilibrating ?

• Methods

Initial state: a "calibrated" source of correlations. Watch their evolution into final state in time = system size.

- Direct construction of a correlation function.
- Inversion of scale-dependent variance
- Discrete Wavelet Transform
- Observations
- Conclusions

2 Autocorrelation



$$\left(\begin{array}{c} x_1 \\ x_2 \end{array}
ight)
ightarrow \left(\begin{array}{c} x_\Sigma \equiv x_1 + x_2 \\ x_\Delta \equiv x_1 - x_2 \end{array}
ight),$$

always a lossless transformation of data. **Autocorrelation** A is a projection of a two-point distribution onto difference variable(s) x_{Δ} , lossless for x_{Σ} -invariant (homogenous, stationary) problems.

$$\Delta R(x_1, x_2) = \frac{\rho(x_1, x_2)}{\rho_{\text{ref}}(x_1, x_2)} - 1$$

3 Uncorrelated event reference for DWT



mixed events: no pixel used twice; ≤ 1 pixel from any event in the same mixed event; no mixing of events with largely different multiplicity and vertex.



4 Local hadron density fluctuations and Discrete Wavelet Transform (DWT)



 $F_{m,l,k}^{\lambda}(\phi,\eta)$ -Haar wavelet **orthonormal basis** in (ϕ,η) . scale fineness (m), directional modes of sensitivity (λ) , track density $\rho(\eta,\phi,p_t)$, locations in 2D (l,k). **DWT is an expansion in this basis.** Power of local fluctuations, mode λ :

$$P^{\lambda}(m) = 2^{-2m} \sum_{l,k} \langle \rho, F_{m,l,k}^{\lambda} \rangle^2$$
(1)

"dynamic texture":

$$P_{\rm dyn}^{\lambda}(m) \equiv P_{\rm true}^{\lambda}(m) - P_{\rm mix}^{\lambda}(m)$$
⁽²⁾

Normalized:

$$P_{\rm dyn}^{\lambda}(m)/P_{\rm mix}^{\lambda}(m)/n(p_t)$$
(3)

5 A flow-like example



Elliptic flow-inspired example: x axis - an angle in "naturalunits" $(2\pi = 1)$, y axis multiplicity. The multiresolution theorem: **a4** = **a0+b0+b1+b2+b3**, can have better fineness.

6 Example of a DWT power spectrum



Power spectrum of that flow event as a function of "fineness" m. The dominant contrubution is m = 1 (the " v_2 " harmonic, **b1**). Statistical fluctuations also contribute.

 $P(m) = 2^{-m} \sum_{i} \langle \rho, F_{m,i} \rangle^2.$

Computational complexity O(N)!

7 "Dynamic texture" as a nonparametric measure of the correlation shape



8 "Dynamic texture" p_t dependence: peripheral events, $\sqrt{s_{NN}} = 200 \text{ GeV}$



Peripheral (60-84%) events: normalized dynamic texture for fineness scales m = 0, 1, 0from left to right panels, respectively, as a function of p_t . • – STAR data; solid line – standard HIJING, dash-dotted line – HIJING without jets.

Qualitative trends in peripheral data are as expected. What signal to expect in the central data, if correlation does not change ?

$$\left(\frac{P_{\text{true}}}{P_{\text{mix}}} - 1\right) \frac{1}{N} \bigg|_{\text{centr}} = \left(\frac{P_{\text{true}}}{P_{\text{mix}}} - 1\right) \bigg|_{\text{periph}} \frac{1}{N_{\text{centr}}}$$
(4)

9 Longitudinal minijet broadening: DWT data



Central (top 4%) events: normalized dynamic texture for fineness scales m = 0, 1, 0 from left to right panels, respectively, as a function of p_t . • STAR data; solid line – Hijing without jet quenching; dashed line – Hijing with quenching; \Box peripheral STAR data renormalized to compare. Minijet elongation \Rightarrow correlation broadening \Leftrightarrow reduced correlation gradient \Leftrightarrow reduced "texture"



10 "Dynamic texture" response

Dynamic texture response in various idealized situations (showing only one scale): (a) events of random (uncorrelated) particles (b) p_t -independent elliptic flow (c) Correlations at low Q_{inv} (Bose-Einstein correlations and Coulomb effect) (d) HIJING jets

3 2.5 **STAR Prelimir STAR Preli** 1.5 0.5 -0.5-2 2 -2-2-2 0 2 0 0 2 A 2 \mathbf{X}_{Δ} \mathbf{X}_{Δ} **(d)** $\mathbf{X}_{\!\Delta}$ $\mathbf{X}_{\!\Delta}$ (c) **(b) (a)** Projections of $\bar{N}[\rho(\eta_{\Delta},\phi_{\Delta})/\rho(\eta_{\Delta},\phi_{\Delta})_{\text{ref}}-1]|_{CI}$ on x_{Δ} which is ϕ_{Δ} (Δ) or η_{Δ} (•). v_1 and v_2 are subtracted.

11 Longitudinal minijet broadening: correlation data

12 Longitudinal minijet broadening: centrality dependence



13 Charge-dependent correlations = Like sign - Unlike sign



The driving physics: charge conservation in hadronization. Suppress short range correlations – BEC and conversion e^+e^- – by a kinematic pair cut. The $\bar{N} \times$ is good when number of correlation sources $\propto N$.

14 Modified hadronization geometry ?



Projections of $\overline{N}(\rho(\eta_{\Delta}, \phi_{\Delta})/\rho(\eta_{\Delta}, \phi_{\Delta})_{ref} - 1)|_{CD}$ on x_{Δ} which is ϕ_{Δ} (Δ) or η_{Δ} (\bullet). $\eta - \phi$ width disparity (d, peripheral) is gone in (a) \Rightarrow transition from (string) 1D to bulk (>2D) fragmentation symmetrizes η and ϕ .





15 Number and p_t correlations



This p_t field may have elliptic flow (number effect). Abounds at RHIC.

Also elliptic... flow $(p_t \text{ effect})$! Pro: blast wave fits. Is there a **direct** measurement ?

16 Towards p_t correlation/fluctuation analysis

Problem: need to tell apart $p_{t,i}$ and number contributions to the $p_t \equiv \sum_{i \in (\eta,\phi) \text{bin}} p_{t,i} \Rightarrow$ can extract the p_t correlation alone.

Solution: use $p_t - n\hat{p_t}$

Q:When is the *n*-contribution into $Var[p_t - n\hat{p_t}]$ canceled ?

$$\sigma^2(p_t:n) \equiv \operatorname{Var}[p_t - n\hat{p_t}] = \operatorname{Var}[p_t] + \hat{p_t}^2 \operatorname{Var}[n] - 2\hat{p_t} \operatorname{Cov}[n, p_t]$$
(5)

$$\operatorname{Var}[p_{t}] = \operatorname{Var}[\sum_{i}^{n} p_{t,i}] = \operatorname{Var}[\sum_{i}^{n} (\hat{p_{t}} + u_{i})] = \hat{p_{t}}^{2} \operatorname{Var}[n] + \operatorname{Var}[u] + 2\hat{p_{t}} \operatorname{Cov}[n, u]$$
(6)

 $\operatorname{Cov}[n, p_t] = \overline{np_t} - \overline{n}\overline{p_t} = \hat{p}_t \operatorname{Var}[n]$ A:For independent p_t and n production, when $\operatorname{Cov}[n, u] \equiv \overline{nu} = 0$, where $u \equiv \sum_{i}^{n} u_i, \ u_i = p_{t,i} - \hat{p_t}.$ (7)



17 Get correlations from fluctuations

Extract correlation structure of random field X from the scale dependence of variance (van Marcke "Random Fields" MIT 1983; Trainor,Porter,Prindle hep-ph/0410180)



$$\operatorname{Var}[X;\delta\eta,\delta\phi] = \int_{-\delta\eta/2}^{\delta\eta/2} d\eta_1 \int_{-\delta\phi/2}^{\delta\phi/2} d\phi_1 \int_{-\delta\eta/2}^{\delta\eta/2} d\eta_2 \int_{-\delta\phi/2}^{\delta\phi/2} d\phi_2 \qquad (8)$$
$$\times [\overline{X(\eta_1,\phi_1)X(\eta_2,\phi_2)} - \overline{X(\eta_1,\phi_1)} \times \overline{X(\eta_2,\phi_2)}]$$

Compare with uncorrelated reference; recognize autocorrelation $\rho(X, t_{\Delta}) \equiv \overline{X(t)X(t + t_{\Delta})}$ (*t*-average).

$$\Delta\sigma^2(X,\delta\eta,\delta\phi) =$$
 (9)

$$\int_{-\delta\eta/2}^{\delta\eta/2} d\eta_1 \int_{-\delta\phi/2}^{\delta\phi/2} d\phi_1 \int_{-\delta\eta/2}^{\delta\eta/2} d\eta_2 \int_{-\delta\phi/2}^{\delta\phi/2} d\phi_2 \Delta\rho(X, \eta_1 - \eta_2, \phi_1 - \phi_2)$$
(10)
= $2 \int_{-\delta\eta/2}^{\delta\eta} d\eta_2 2 \int_{-\delta\phi/2}^{\delta\phi} d\phi_2 (\delta\eta - \eta_2) (\delta\phi - \phi_2) \Delta\rho(X, \eta_2, \phi_2)$ (11)

$$= 2 \int_{0} d\eta_{\Delta} 2 \int_{0} d\phi_{\Delta} (\delta \eta - \eta_{\Delta}) (\delta \phi - \phi_{\Delta}) \Delta \rho(X, \eta_{\Delta}, \phi_{\Delta})$$

18 The actual analysis is discrete: $\int \rightarrow \sum$



kernel K:

$$(\delta\eta - \eta_{\Delta})(\delta\phi - \phi_{\Delta}) \to \varepsilon_{\eta}\varepsilon_{\phi}K_{m_{\delta}n_{\delta}:kl} \equiv \varepsilon_{\eta}\varepsilon_{\phi}(m_{\delta} - k + \frac{1}{2})(n_{\delta} - l + \frac{1}{2})$$
(12)

reference density $ho_{
m ref}$ makes a per-particle measure:

$$\rho_{\rm ref} \propto \bar{n}^2 \Rightarrow \frac{1}{\sqrt{\rho_{\rm ref}}} \propto \frac{1}{\bar{n}}$$
(13)

$$\Delta \sigma_{p_t:n}^2(m_{\delta}\varepsilon_{\eta}, n_{\delta}\varepsilon_{\phi}) = 4 \sum_{k,l=1}^{m_{\delta}, n_{\delta}} \varepsilon_{\eta}\varepsilon_{\phi} K_{m_{\delta}n_{\delta}:kl} \frac{\Delta \rho(p_t:n; k\varepsilon_{\eta}, l\varepsilon_{\phi})}{\sqrt{\rho_{\text{ref}}(n; k\varepsilon_{\eta}, l\varepsilon_{\phi})}}$$
(14)

Inverse problem: knowing $\Delta \sigma^2$, solve for $\Delta \rho / \sqrt{\rho_{\rm ref}} \Rightarrow$ save O(N) in CPU time ! 20

19 p_t correlations from the inversion



minijets.

scale dependence of the "pure" p_t variance. Bottom:

corresponding autocorrelation

20 Localized p_t correlations: minijets



AuAu 200 GeV. In η , correlation broadens with centrality; in ϕ the trend is opposite. The surrounding background seems to recoil.

21 Conclusions

- Semi-hard scattering leaves a trace in the soft p_t domain new at RHIC !
- First direct measurements of p_t correlation structure reveal azimuthal anisotropy of p_t field \Rightarrow elliptic flow is a velocity phenomenon
- The minijet correlation structure is modified with centrality; the effect appears to "turn on" around $\nu = (N_{\rm part}/2)^{1/3} \approx 3$. Broadening of the correlation in η and weakening of $P_{\rm dyn}^{\eta}$ on the coarse scale are consistent descriptions of the effect. How exactly does the coupling between longitudinal flow and mini-jets work ? What do we learn about the expanding fluid ?
- Increased symmetry of the charge-dependent correlation on (η, ϕ) in the central collisions may point to a change in the hadronization geometry in the medium

22 Extra slides

23 DWT of a photographic image



Reproduced from textbook: I.Daubechies, "Ten lectures on wavelets". The original caption: "A real image, and its wavelet decomposition into three multiresolution layers. On the wavelet components one clearly sees that the $d^{j,v}, d^{j,h}, d^{j,d}$ emphasize, respectively, vertical, horizontal, and diagonal edges. In this figure, the bottom picture has been overexposed to make details in the $d^{j,\lambda}$ more apparent. I would like to thank M.Barlaud for providing this figure." The colored marks are mine. 25