

Correlations at STAR: interferometry and event structure

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1 STAR fluctuation/correlation physics in this talk

Equilibration: Arguably the central issue of RHIC hadronic physics. Is it taking place ? What is the mechanism ? And **what** is equilibrating ?

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- Azimuthal dependence in Bose-Einstein correlations
- Arguments for Blast Wave from HBT and p_t fluctuations

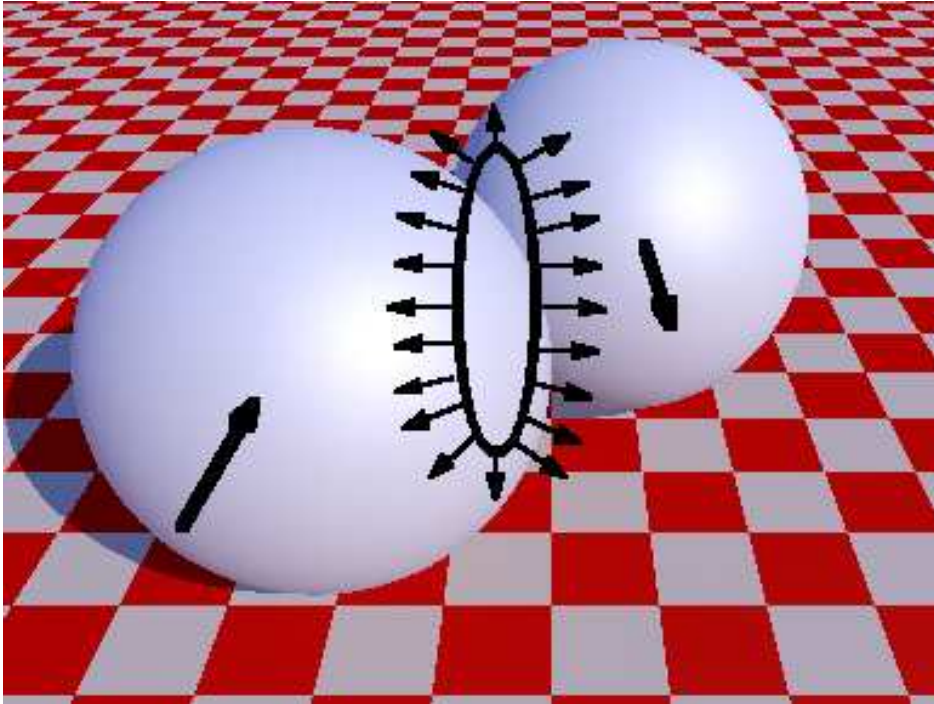


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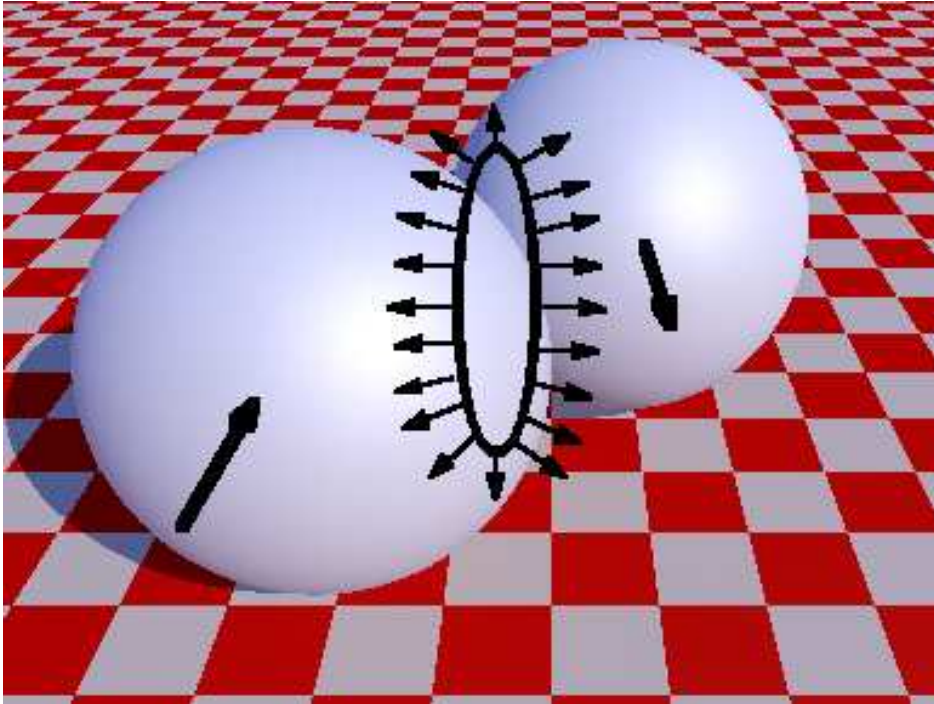
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- Novel techniques throughout...

2 Flow – directed and elliptic



(x, y) anisotropy \rightarrow
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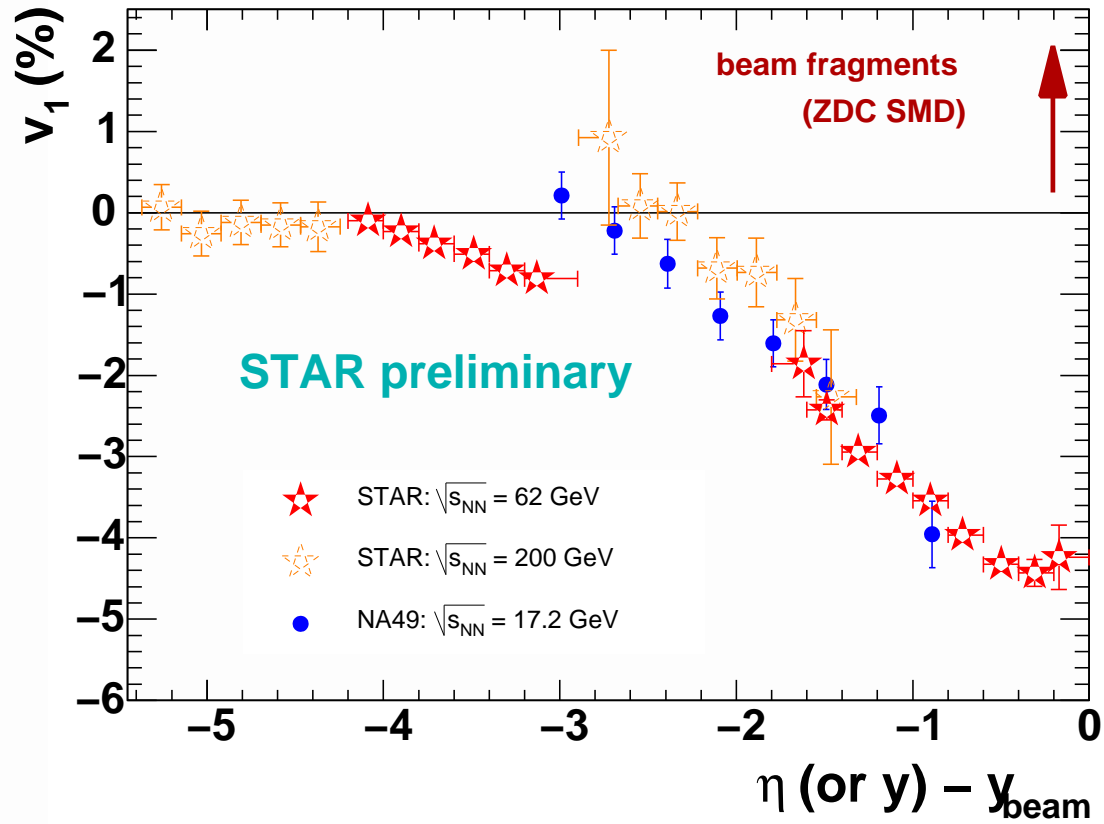


(x, y) anisotropy \rightarrow
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$$E \frac{d^3 N}{d^3 p} = \frac{1}{2\pi} \frac{d^2 N}{p_t dp_t dy} \left\{ 1 + \sum_{m=1}^{\infty} 2v_m \cos[m(\phi - \Psi_r)] \right\} \quad (1)$$

- flow starts early – perhaps before hydro is applicable (stopping stage)
- testifies to equilibration
- sensitive to pressure and density gradients
- flow is a multiparticle effect; there is “non-flow”

Directed flow

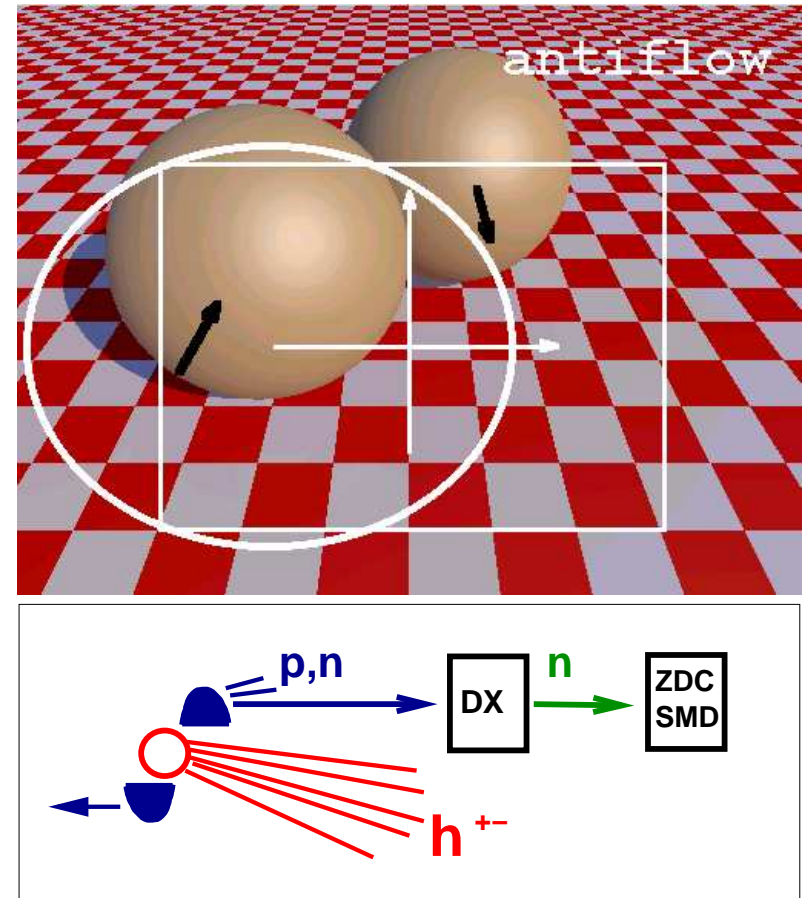
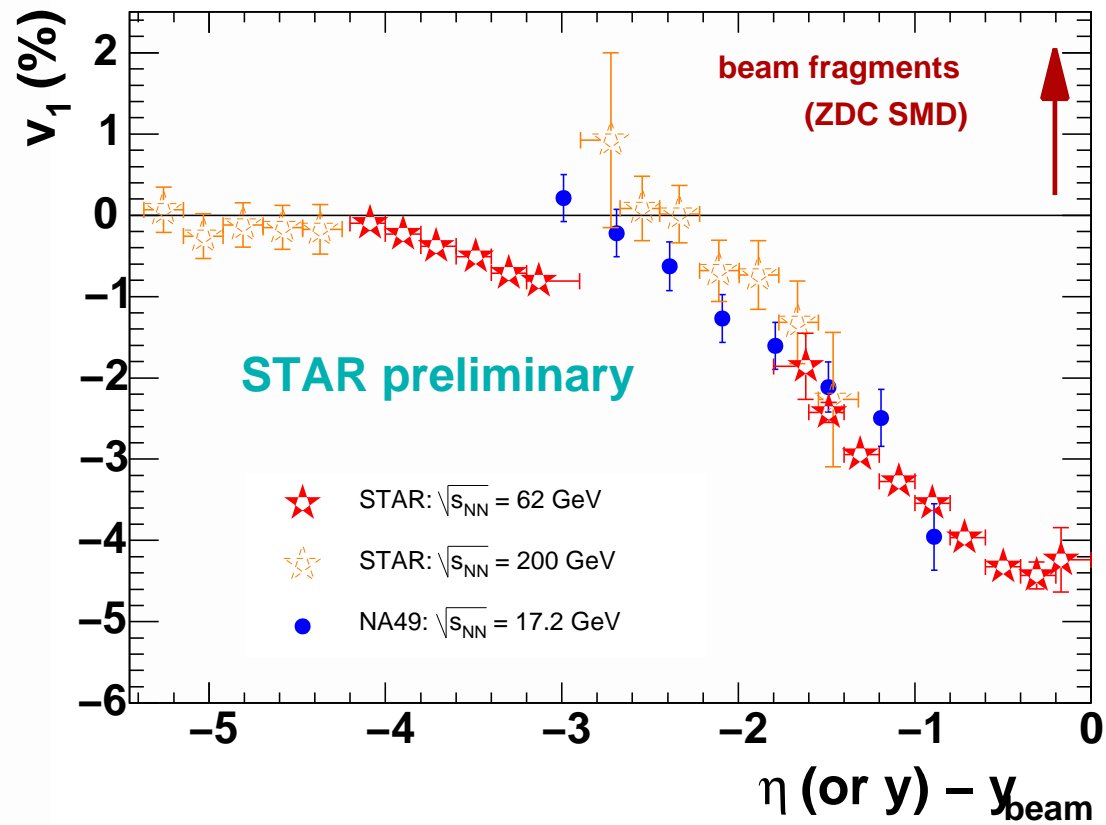


Charged-particles v_1 from

3-particle cumulants in the projectile frame.

- monotonic around midrapidity
- Supports limiting fragmentation
- Antiflow !

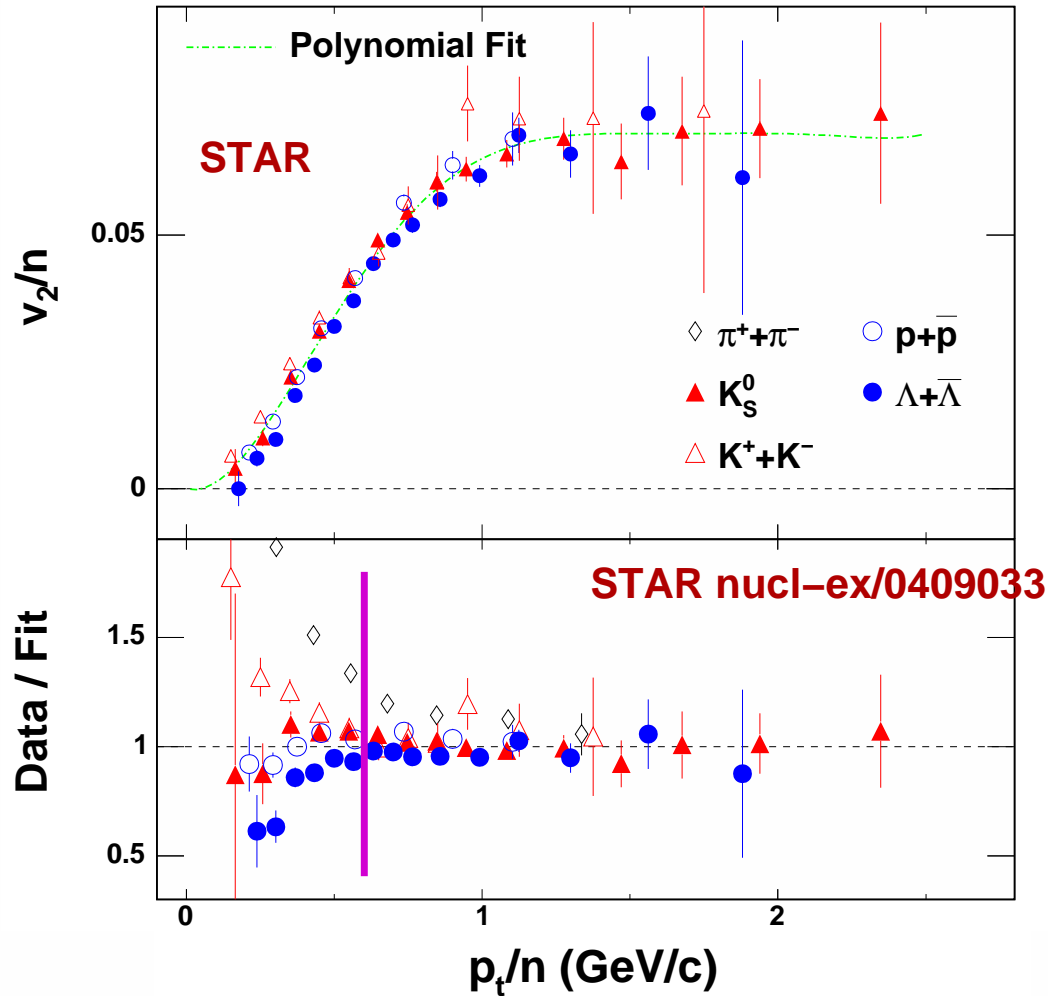
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4 Elliptic flow and quark coalescence



$$\frac{dN}{d\phi} \propto 1 + 2v_2 \cos(2\phi) \quad (2)$$

$$\frac{dN_{\text{clscnc},n}}{d\phi}(p_t) \propto \left(\frac{dN(\frac{p_t}{n})}{d\phi} \right)^n \quad (3)$$

$$(1 + 2v_2 \cos(2\phi))^n = 1 + 2v_2 n \cos(2\phi) + \mathcal{O}(v_2^2) \quad (4)$$

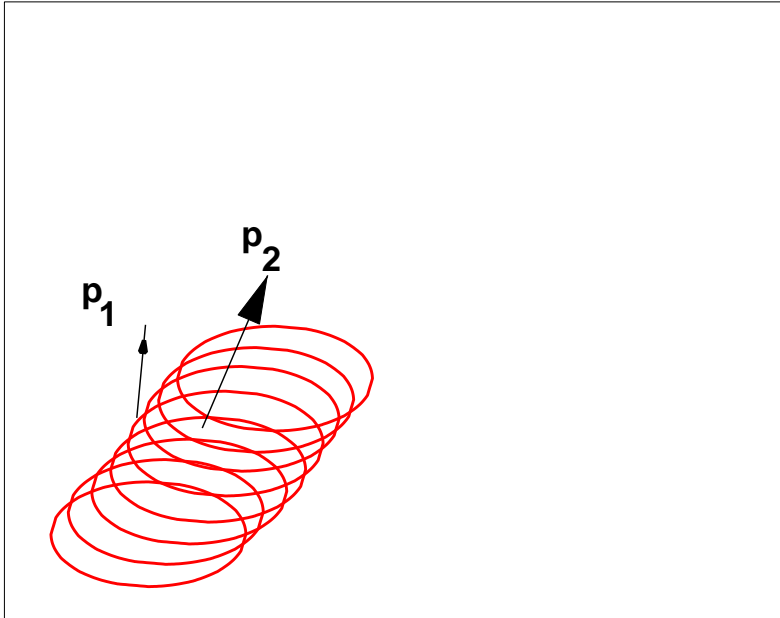
STAR AuAu 200 GeV minbas; n is number of constituent quarks. Expect universality if quark coalescence dominates hadronization after the universal flow sets in. Valid at $p_t/n > 0.6$ GeV/c for $K_S^0, K^\pm, p, \bar{p}, \Lambda, \bar{\Lambda}$.

5 HBT definitions for bosons

$$C(p_1, p_2) = \frac{\rho(p_1, p_2)}{\rho(p_1)\rho(p_2)} \rightarrow C_{\text{exp}}(p_1, p_2) = \frac{\rho(p_1, p_2)}{\rho_{\text{mix}}(p_1, p_2)} \quad (5)$$

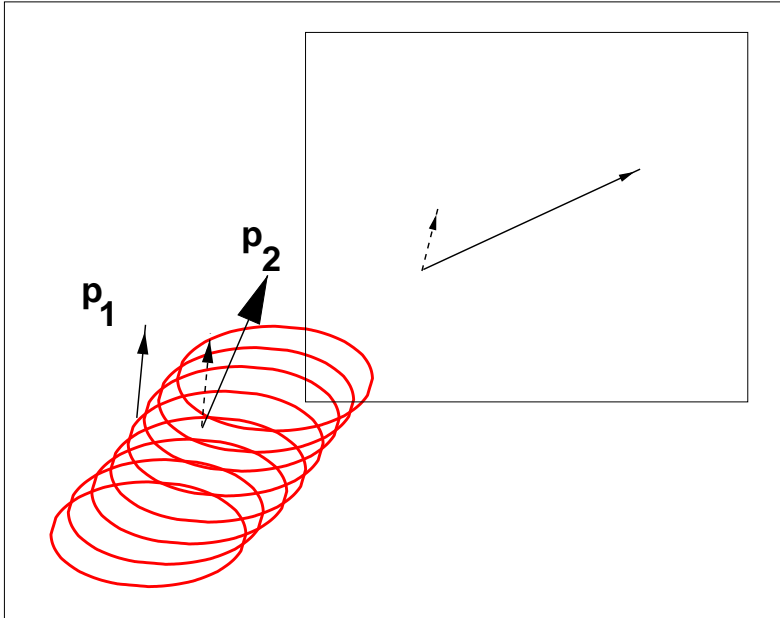
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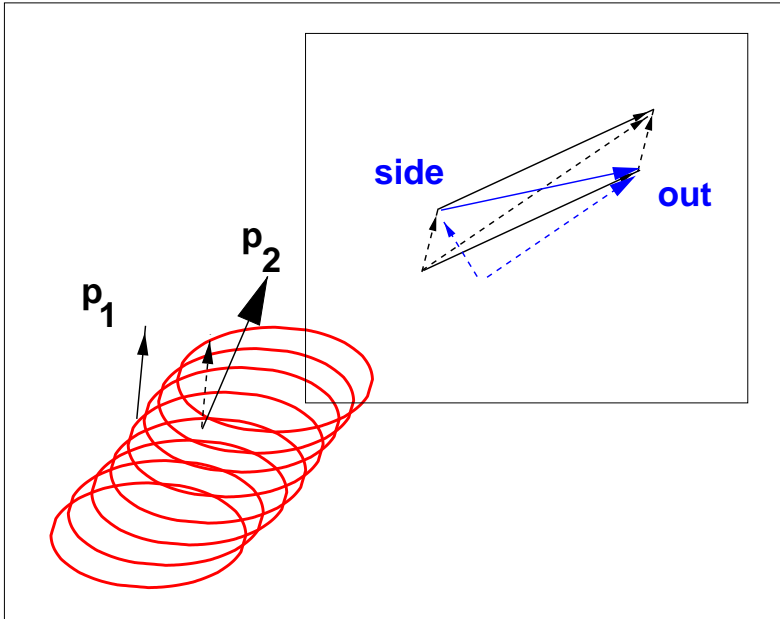
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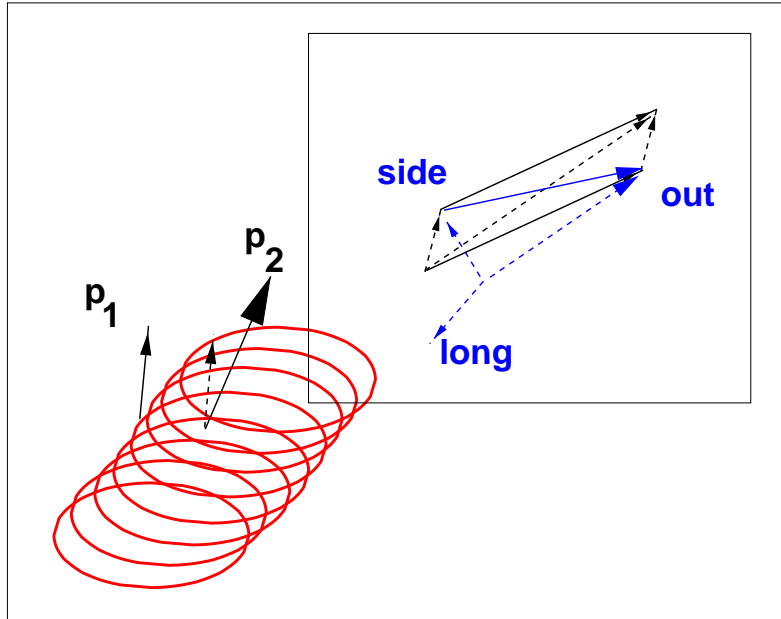
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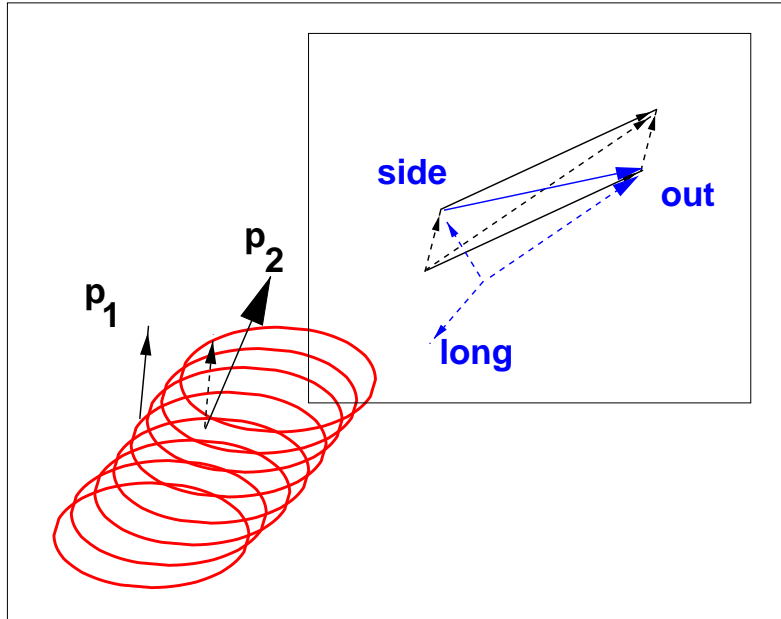
emission time $\rightarrow R_o^2$;
 transverse homogeneity $\rightarrow R_s^2$;
 longitudinal homogeneity $\rightarrow R_l^2$

Bertsch-Pratt parameterization: traditional

$$C_{\text{fit}}(\vec{q}) = 1 + \lambda \exp(-q_o^2 R_o^2 - q_s^2 R_s^2 - q_l^2 R_l^2 - 2q_o q_s R_{os}^2) \quad (6)$$

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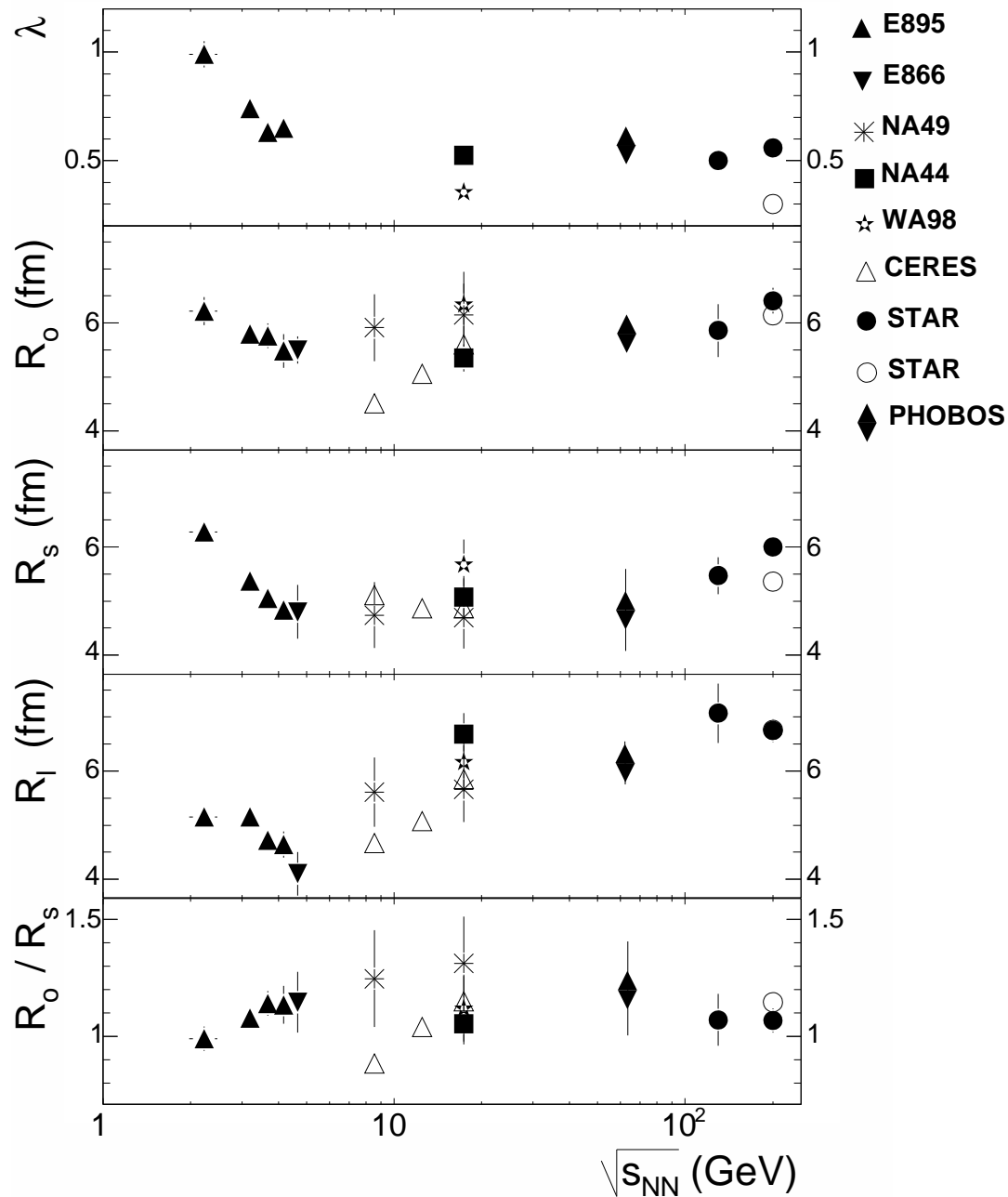
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including Coulomb effect with Bowler-Sinyukov method (implies complete chaoticity)

$$C_{\text{fit}}(\vec{q}) = (1 - \lambda) + \lambda K_{\text{Coulomb}} \exp(-q_o^2 R_o^2 - q_s^2 R_s^2 - q_l^2 R_l^2 - 2q_o q_s R_{os}^2) \quad (7)$$

6 The “HBT puzzle”



Open symbols:

Bowler-Sinyukov fits

- Causes of inhomogeneity ?
- $R_{out} \approx R_{side}$ – instantaneous emission ?
- R_{long} smaller than expected

7 Modified Blast Wave model (Retiere, Lisa)

- elliptic source (R_x, R_y) with diffuse edge (α_S)

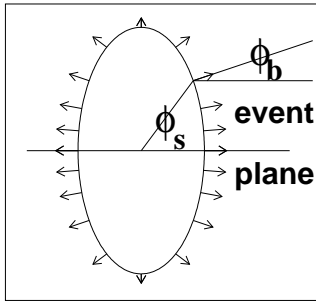
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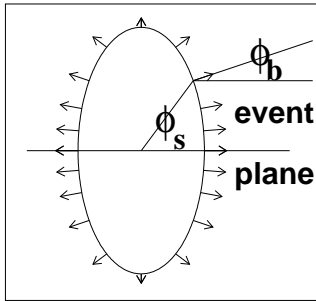
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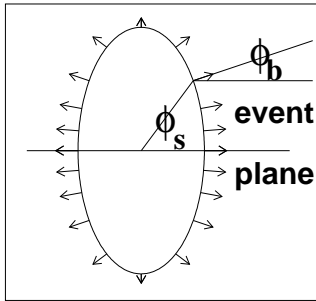
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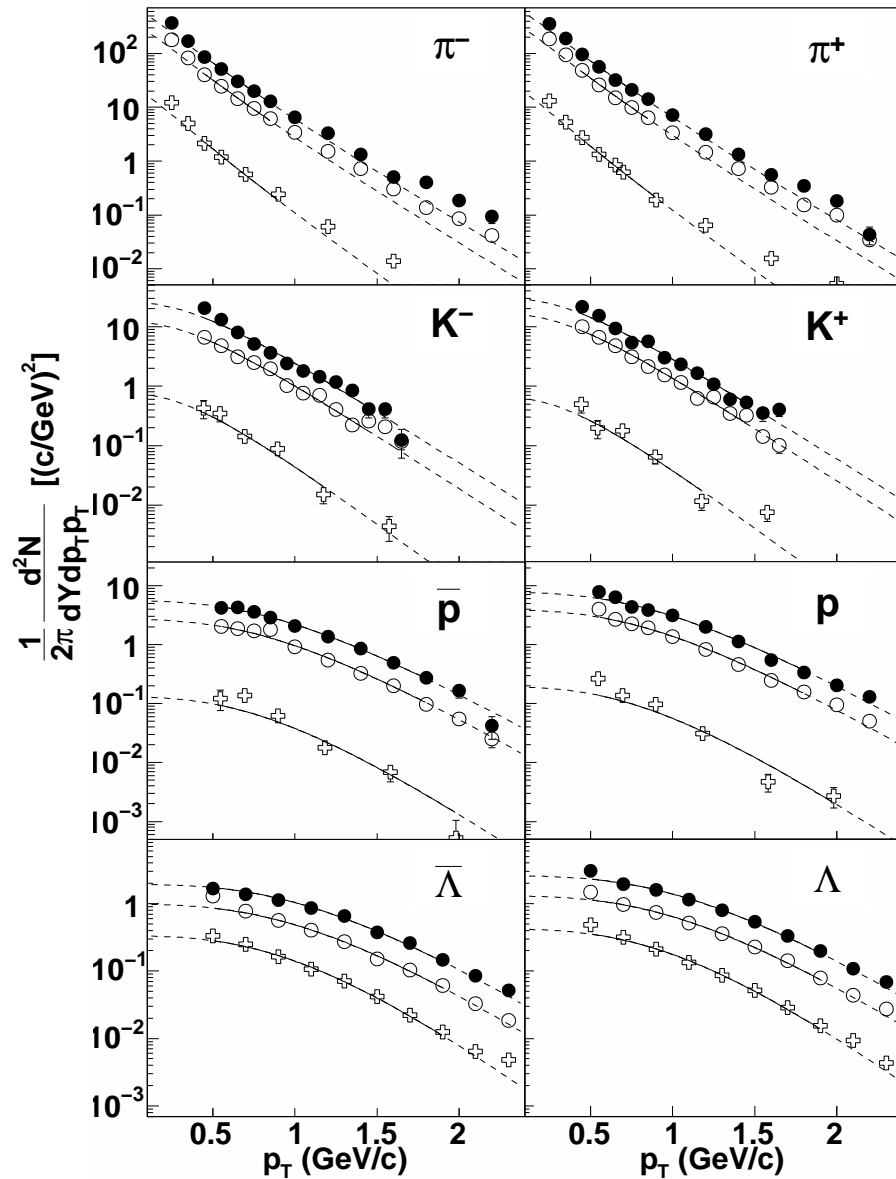
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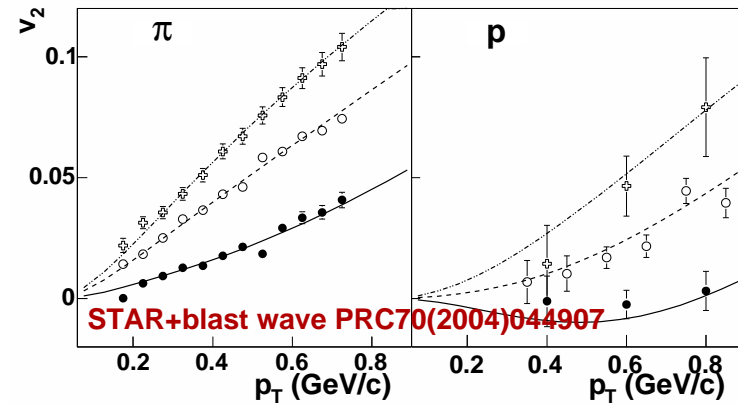
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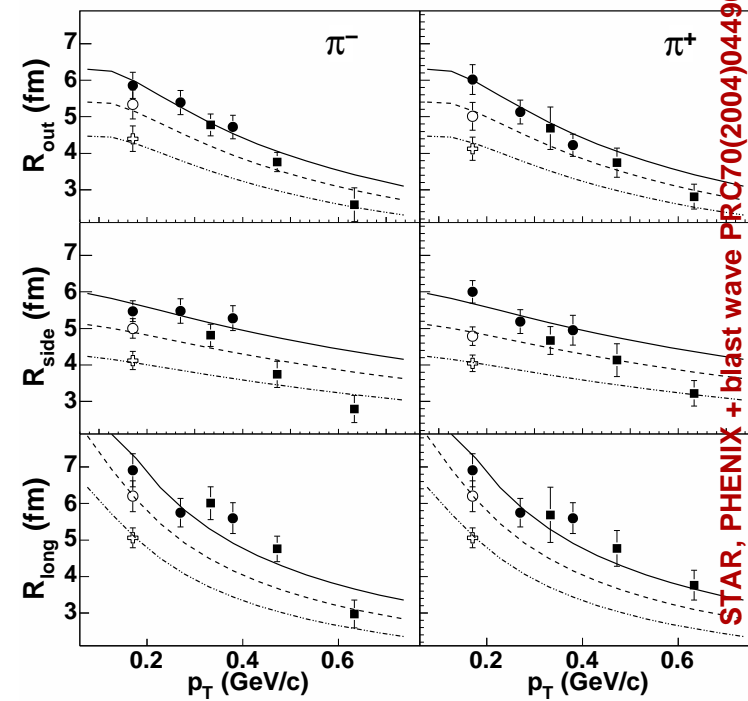
8 Combined fits with Blast Wave: spectra, v_2 , HBT



STAR, PHENIX+blast wave PRC70(2004)044907



● central; ○ mid-central; + peripheral



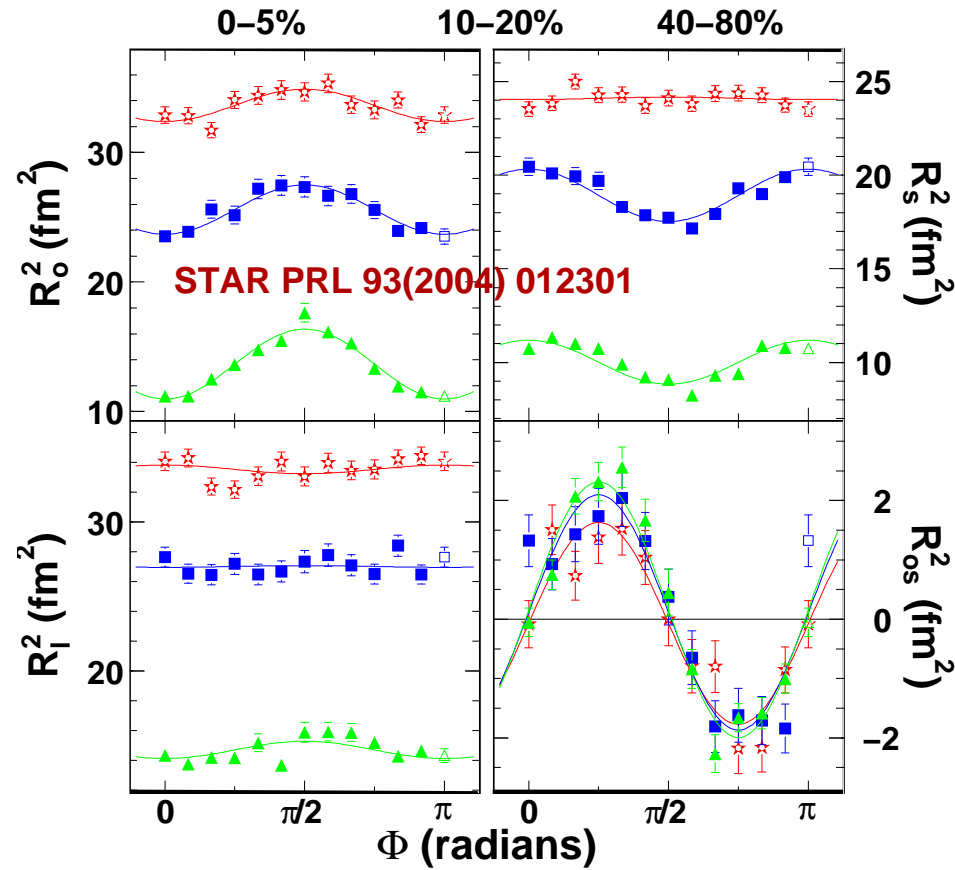
STAR ● central; ○ mid-central;
+ peripheral; ■ PHENIX 30% central

9 Blast Wave fit parameters at RHIC

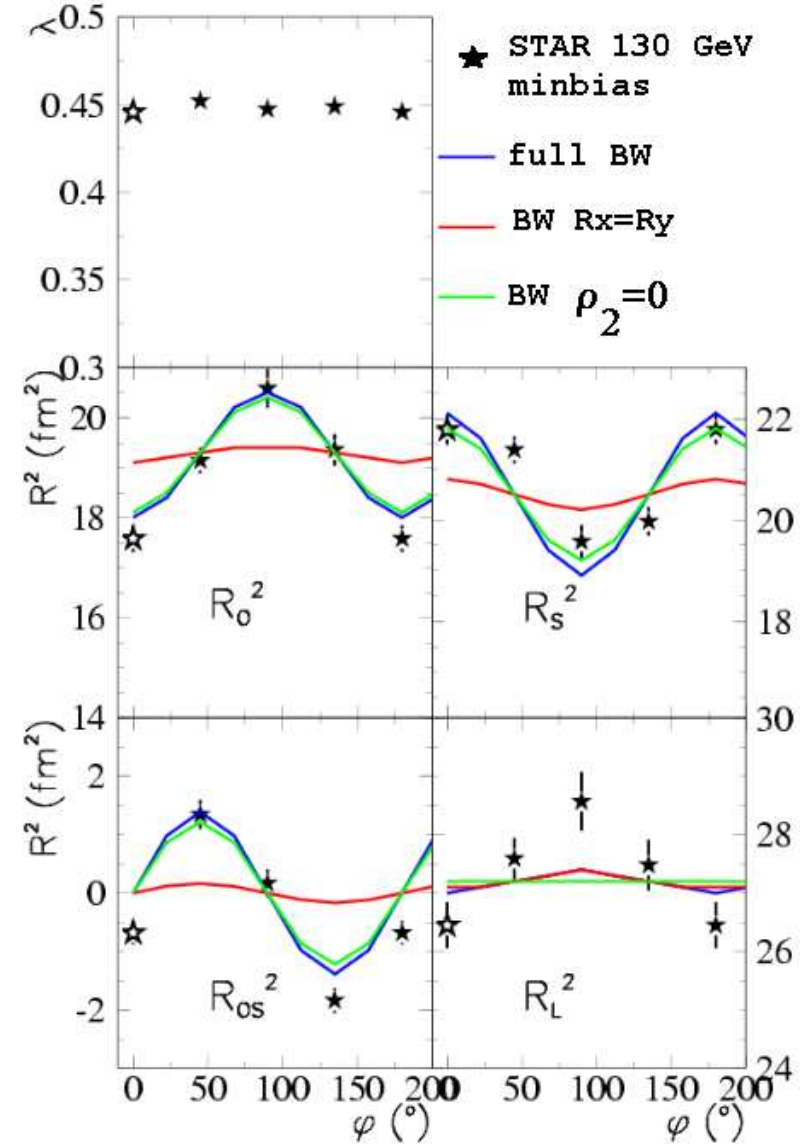
	centr	mid-c	periph
$T(\text{MeV})$	106	107	100
ρ_0	0.89	0.85	0.79
ρ_2	$(6.0 \pm 0.8)10^{-2}$	$(5.8 \pm 0.5)10^{-2}$	$(5 \pm 1)10^{-2}$
$R_x(\text{fm})$	13.2	10.4	8.0
$R_y(\text{fm})$	13.0	11.8	10.1
$\tau(\text{fm}/c)$	9.2	7.7	6.5
$\Delta\tau (\text{fm}/c)$	0.003 ± 1.3	0.06 ± 1.3	0.6 ± 1.8

Emission duration consistent with 0...

10 Blast Wave and the azimuthally-dependent HBT results

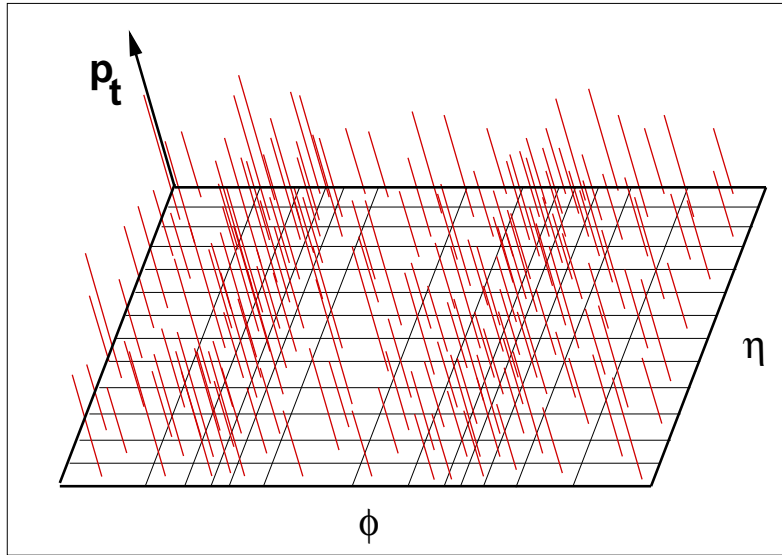


Left: STAR AuAu 200 GeV. Right: calculations for non-round source $10.8 \times 11.4 \text{fm}$, $\tau=8.3 \text{fm}/c$, $\Delta\tau=0$ compare with STAR AuAu 130 GeV data



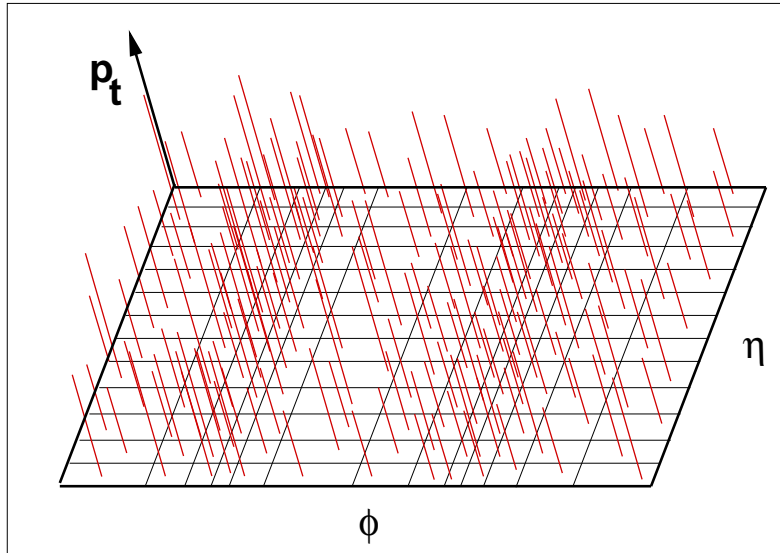
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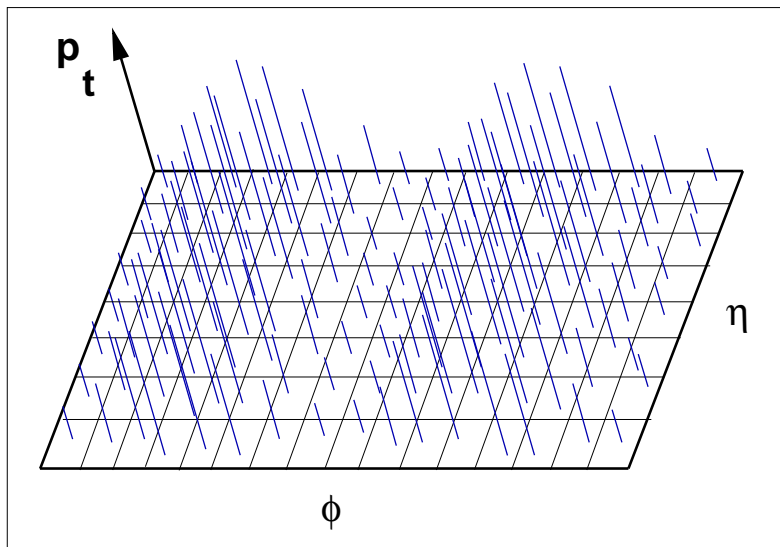


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Also elliptic... flow (p_t effect) !
Pro: blast wave fits. Is there a **direct** measurement ?

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Problem: need to tell apart $p_{t,i}$ and number contributions to the

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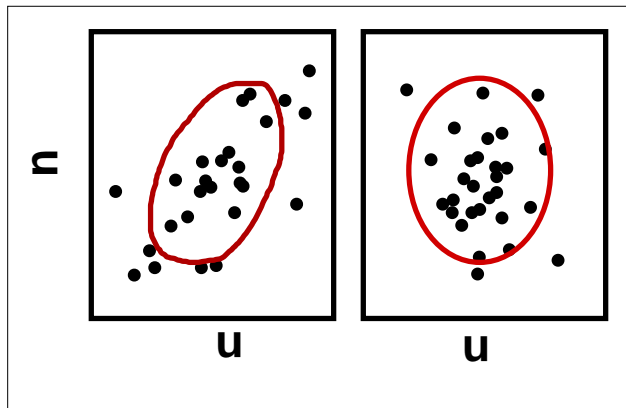
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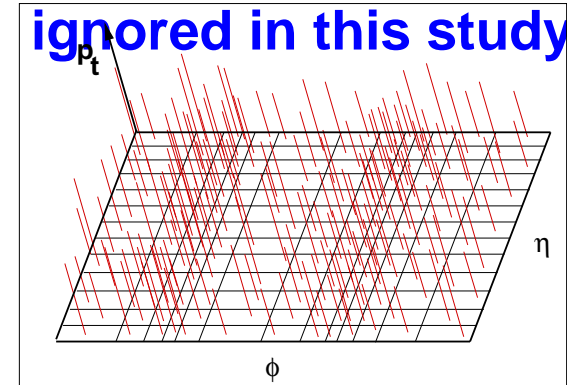
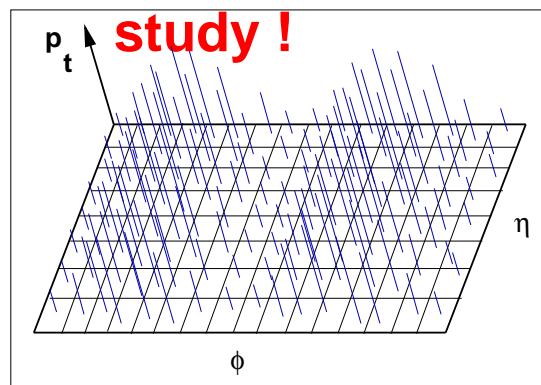
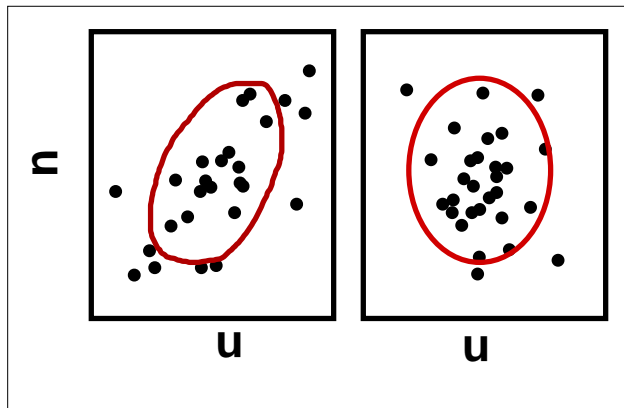
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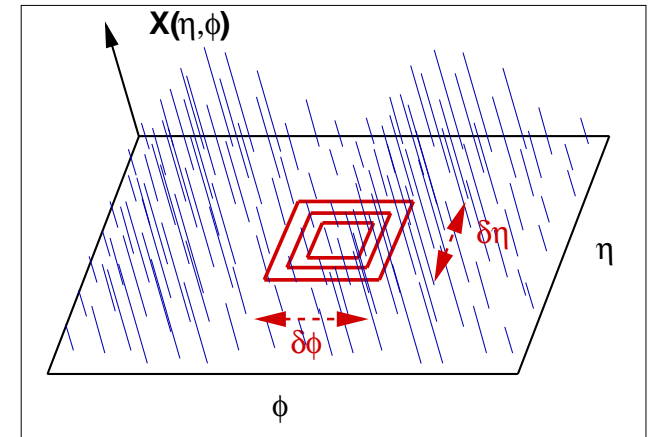
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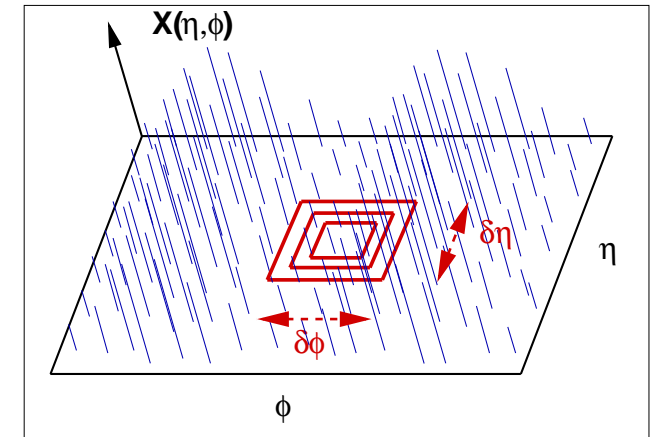
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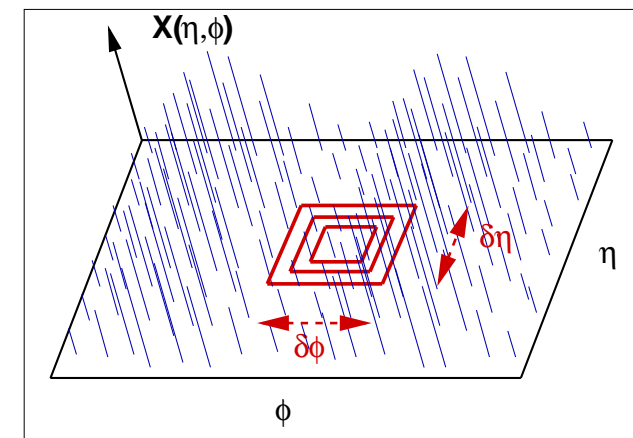
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$$\times [\overline{X(\eta_1, \phi_1)X(\eta_2, \phi_2)} - \overline{X(\eta_1, \phi_1)} \times \overline{X(\eta_2, \phi_2)}]$$

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Extract correlation structure of random field X from the scale dependence of variance (van Marcke "Random Fields" MIT 1983; Trainor, Porter, Prindle hep-ph/0410180)



$$\text{Var}[X; \delta\eta, \delta\phi] = \int_{-\delta\eta/2}^{\delta\eta/2} d\eta_1 \int_{-\delta\phi/2}^{\delta\phi/2} d\phi_1 \int_{-\delta\eta/2}^{\delta\eta/2} d\eta_2 \int_{-\delta\phi/2}^{\delta\phi/2} d\phi_2 \quad (12)$$

$$\times [\overline{X(\eta_1, \phi_1)X(\eta_2, \phi_2)} - \overline{X(\eta_1, \phi_1)} \times \overline{X(\eta_2, \phi_2)}]$$

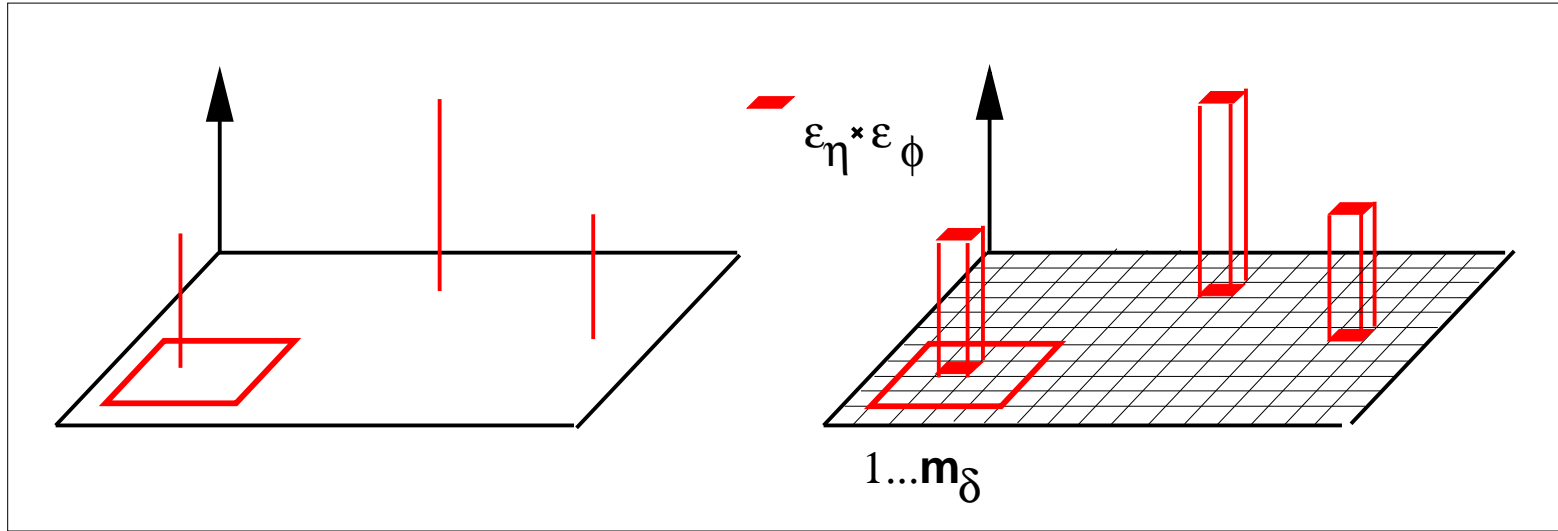
Compare with uncorrelated reference; recognize autocorrelation $\rho(X, t_\Delta) \equiv \overline{X(t)X(t + t_\Delta)}$ (t -average).

$$\Delta\sigma^2(X, \delta\eta, \delta\phi) = \quad (13)$$

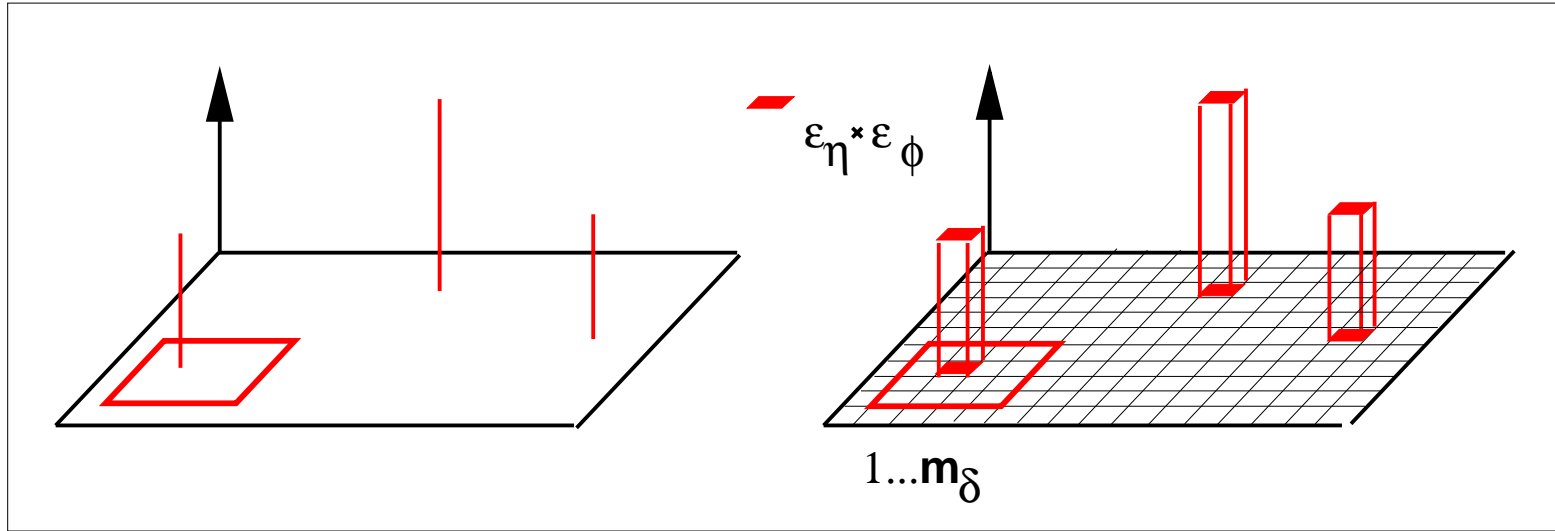
$$\int_{-\delta\eta/2}^{\delta\eta/2} d\eta_1 \int_{-\delta\phi/2}^{\delta\phi/2} d\phi_1 \int_{-\delta\eta/2}^{\delta\eta/2} d\eta_2 \int_{-\delta\phi/2}^{\delta\phi/2} d\phi_2 \Delta\rho(X, \eta_1 - \eta_2, \phi_1 - \phi_2) \quad (14)$$

$$= 2 \int_0^{\delta\eta} d\eta_\Delta 2 \int_0^{\delta\phi} d\phi_\Delta (\delta\eta - \eta_\Delta)(\delta\phi - \phi_\Delta) \Delta\rho(X, \eta_\Delta, \phi_\Delta) \quad (15)$$

14 The actual analysis is discrete: $\int \rightarrow \Sigma$



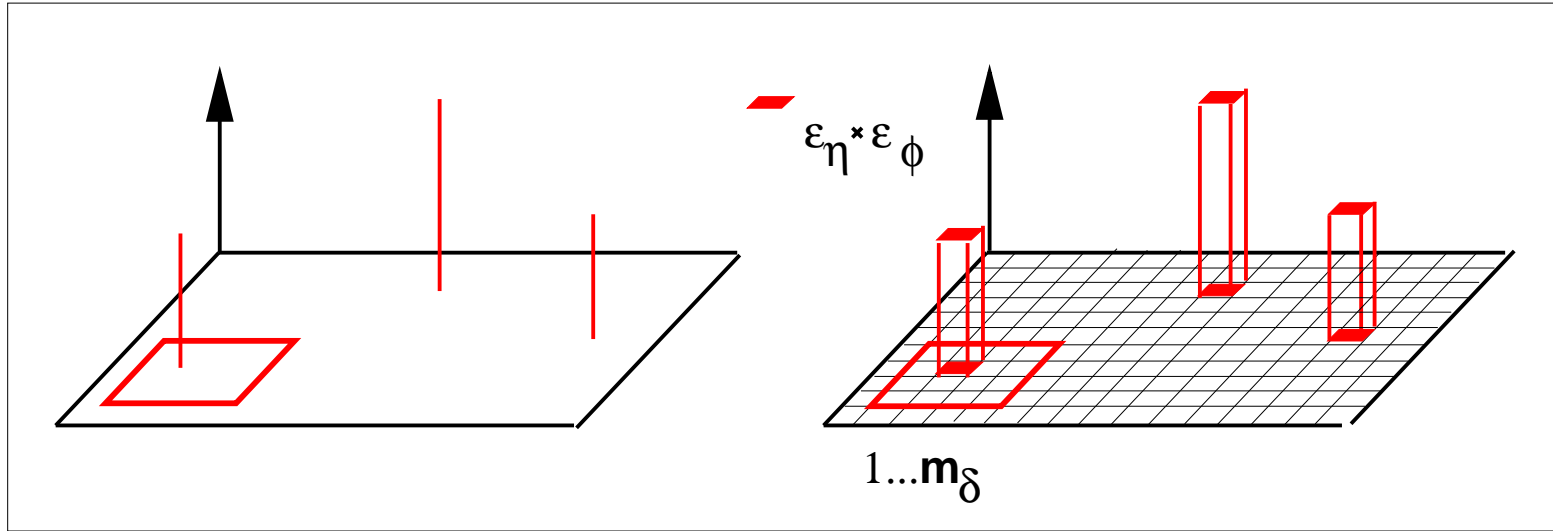
14 The actual analysis is discrete: $\int \rightarrow \Sigma$



kernel K :

$$(\delta\eta - \eta_\Delta)(\delta\phi - \phi_\Delta) \rightarrow \varepsilon_\eta \varepsilon_\phi K_{m_\delta n_\delta:kl} \equiv \varepsilon_\eta \varepsilon_\phi (m_\delta - k + \frac{1}{2})(n_\delta - l + \frac{1}{2}) \quad (16)$$

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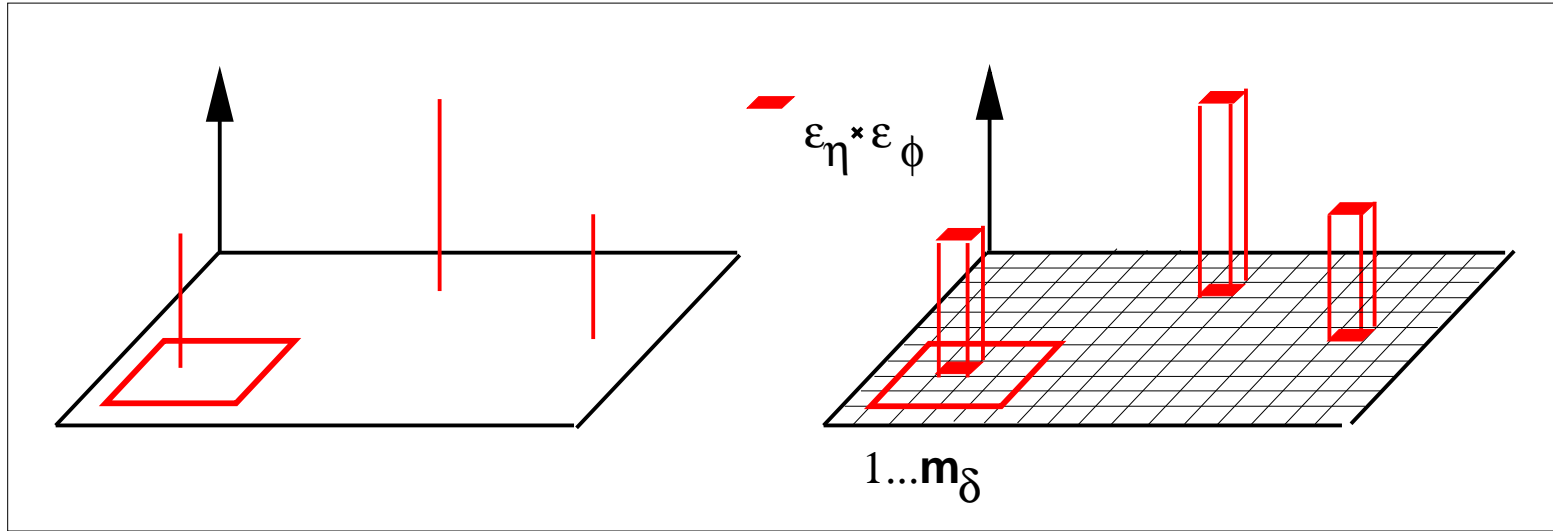
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reference density ρ_{ref} makes a per-particle measure:

$$\rho_{\text{ref}} \propto \bar{n}^2 \Rightarrow \frac{1}{\sqrt{\rho_{\text{ref}}}} \propto \frac{1}{\bar{n}} \quad (17)$$

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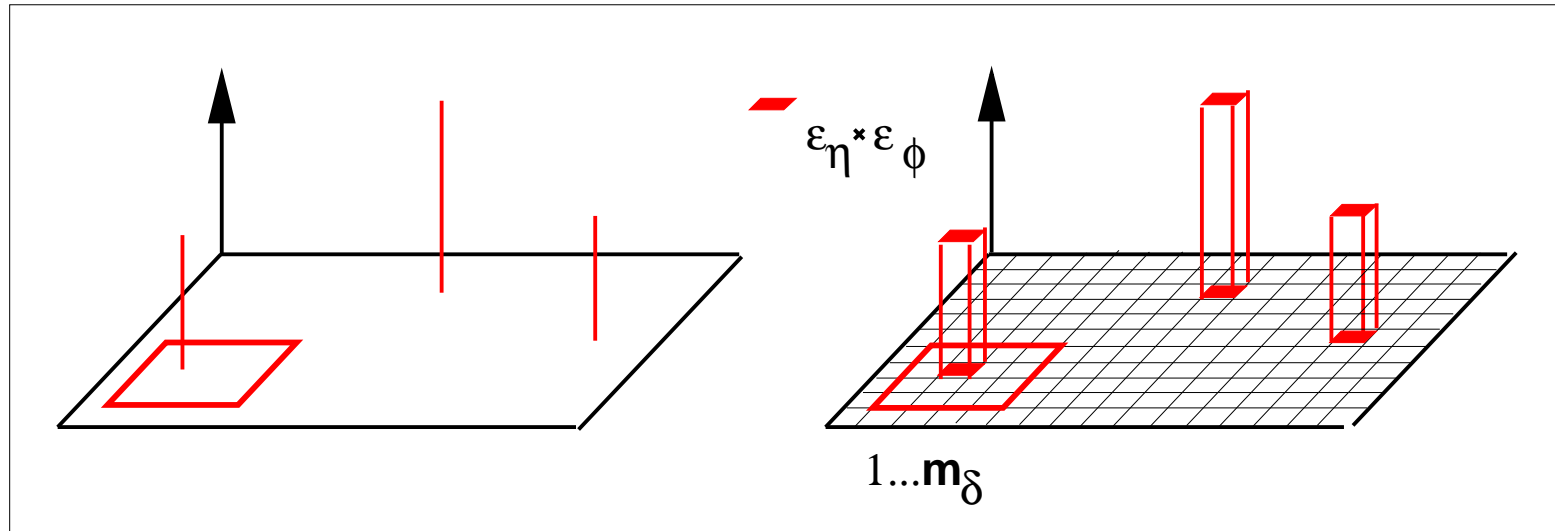
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$$\Delta\sigma_{p_t:n}^2(m_\delta \varepsilon_\eta, n_\delta \varepsilon_\phi) = 4 \sum_{k,l=1}^{m_\delta, n_\delta} \varepsilon_\eta \varepsilon_\phi K_{m_\delta n_\delta:kl} \frac{\Delta\rho(p_t : n; k\varepsilon_\eta, l\varepsilon_\phi)}{\sqrt{\rho_{\text{ref}}(n; k\varepsilon_\eta, l\varepsilon_\phi)}} \quad (18)$$

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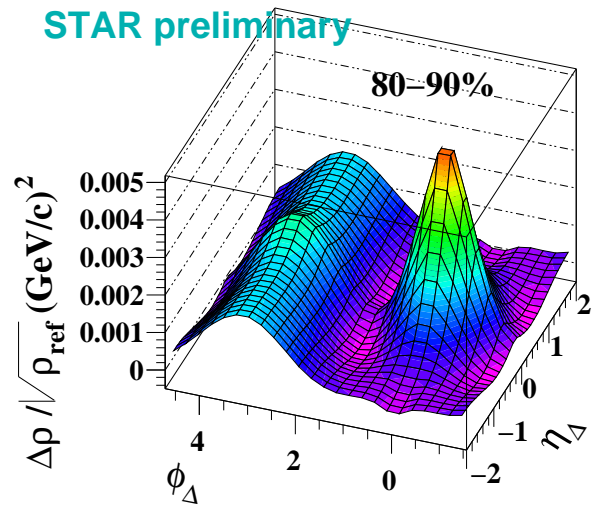
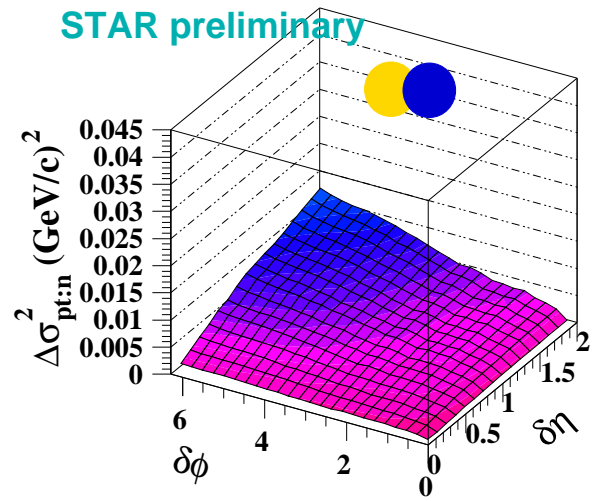
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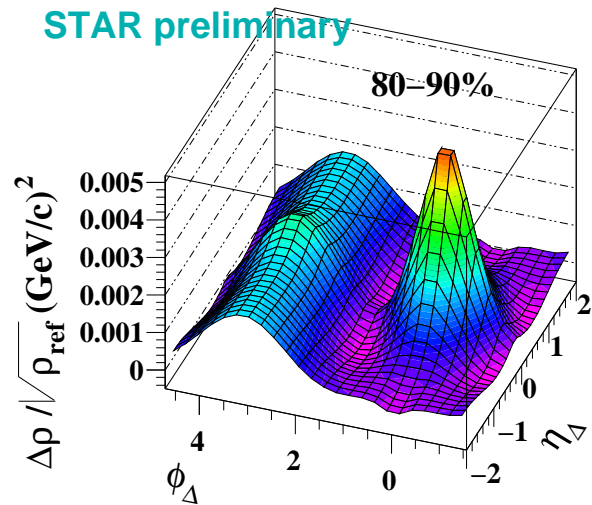
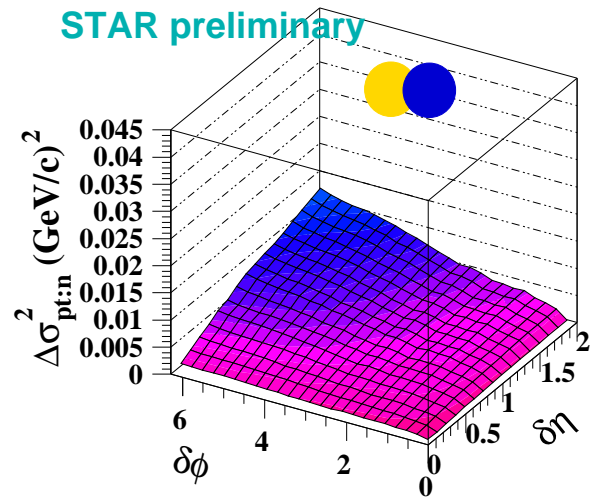
Inverse problem: knowing $\Delta\sigma^2$, solve for $\Delta\rho/\sqrt{\rho_{\text{ref}}} \Rightarrow$ save $O(N)$ in CPU time !

15 p_t correlations from the inversion

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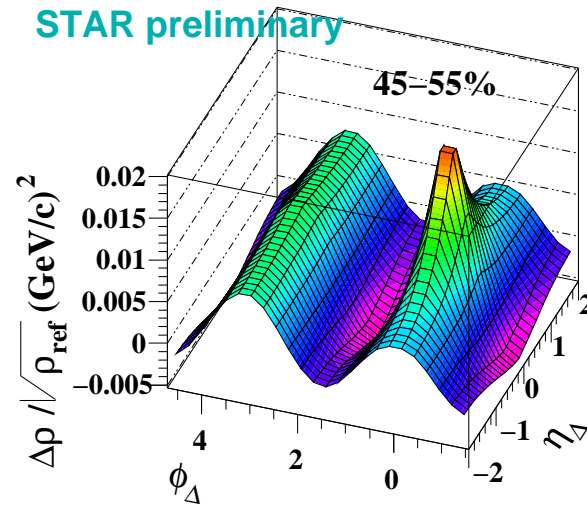
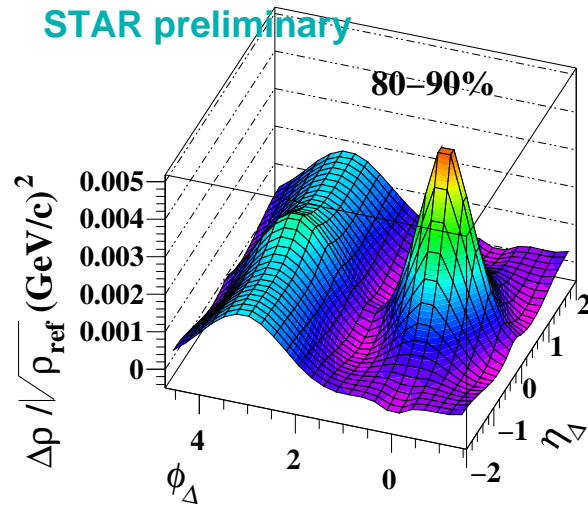
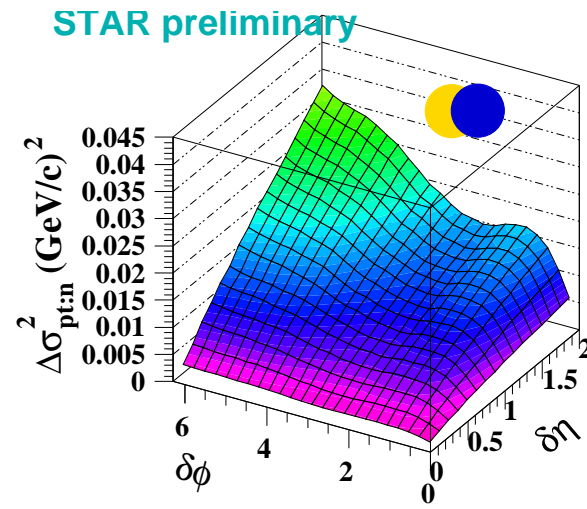
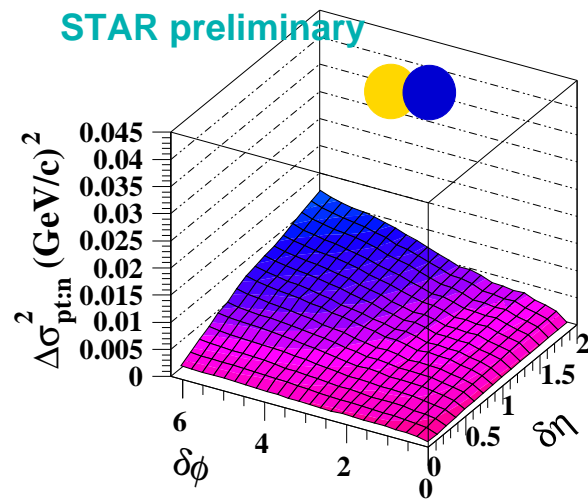


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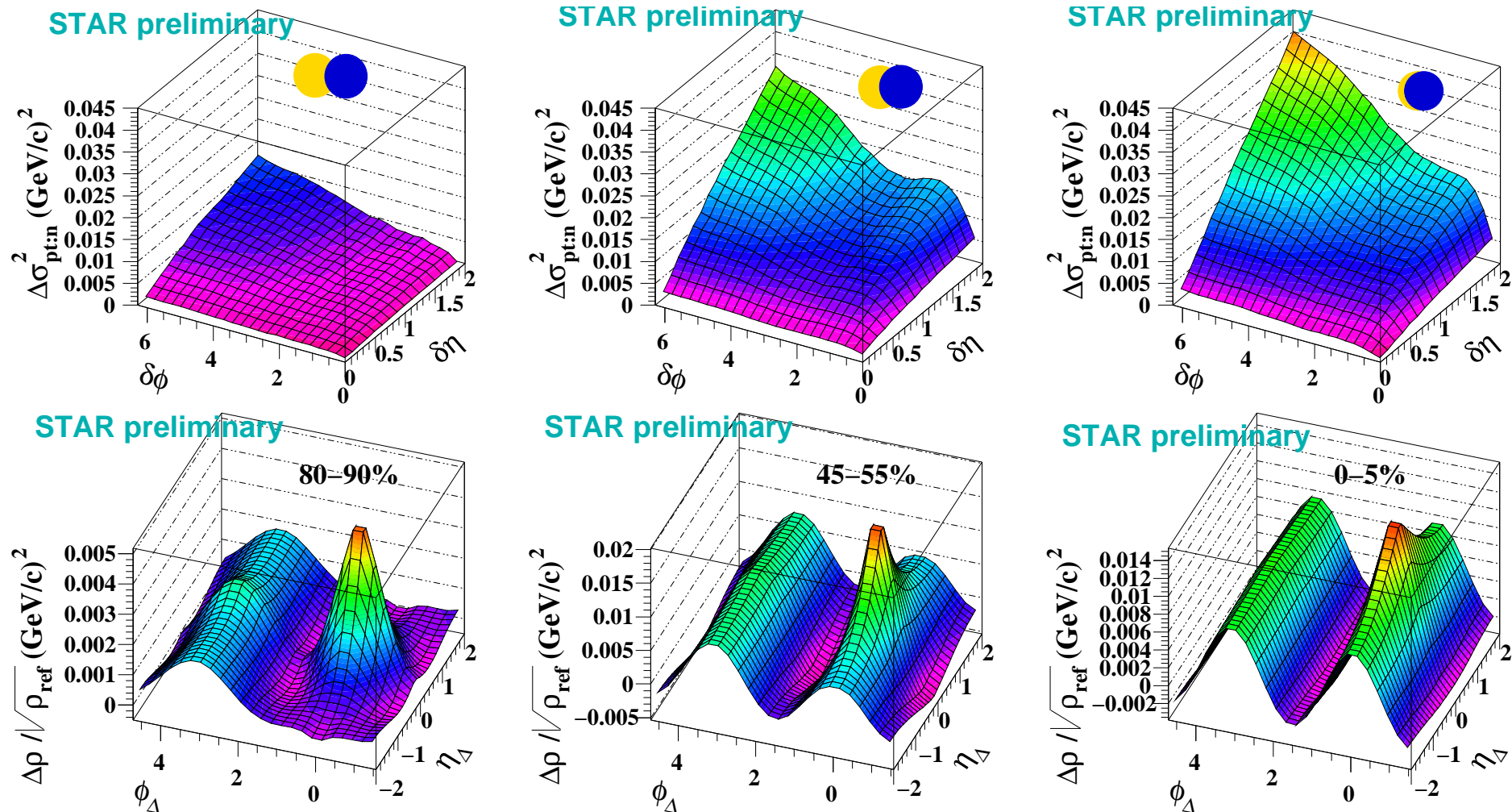
Top: AuAu 200 GeV,
scale dependence of
the “pure” p_t variance.
Bottom:
corresponding
autocorrelation

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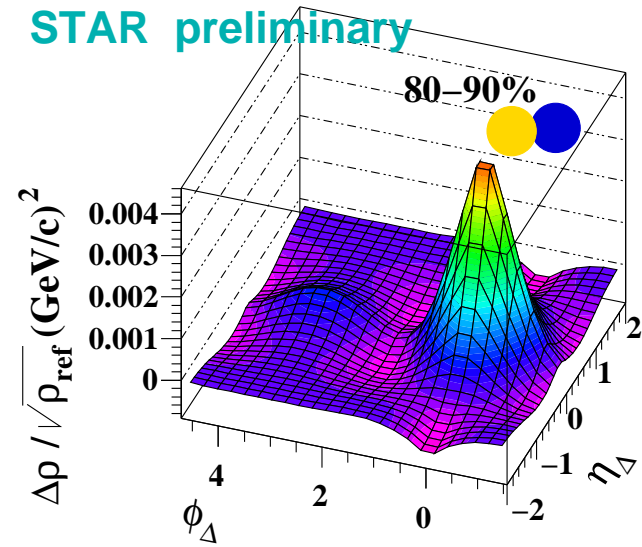
Bottom: corresponding autocorrelation

First direct evidence of elliptic flow as a velocity phenomenon at RHIC. Next, subtract the flow contribution to look at minijets.

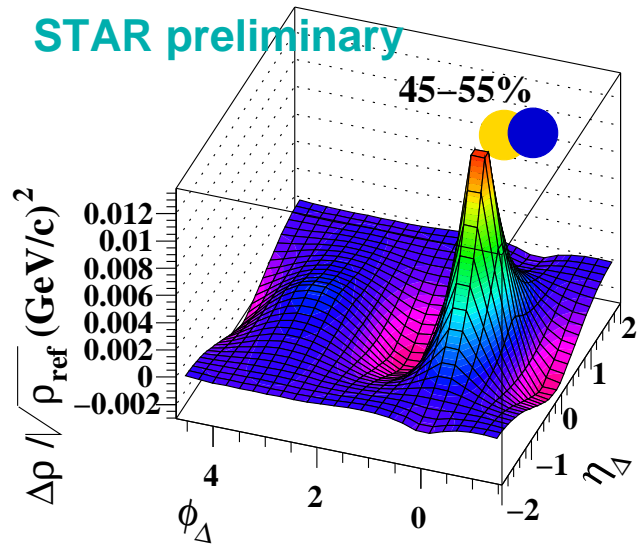
16 Localized p_t correlations: minijets

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STAR preliminary

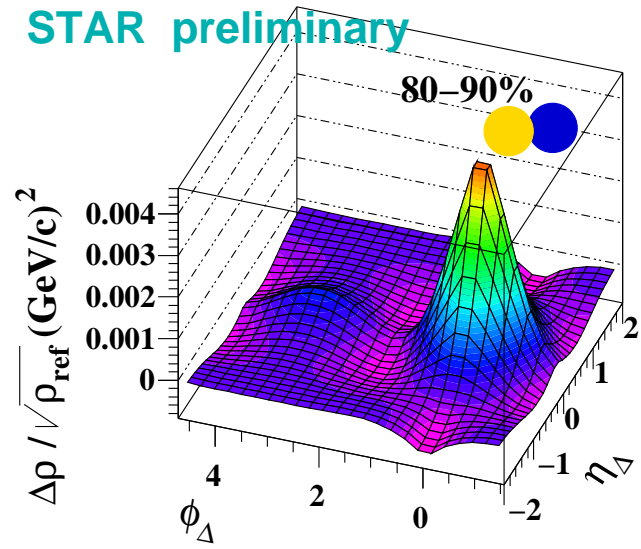


STAR preliminary

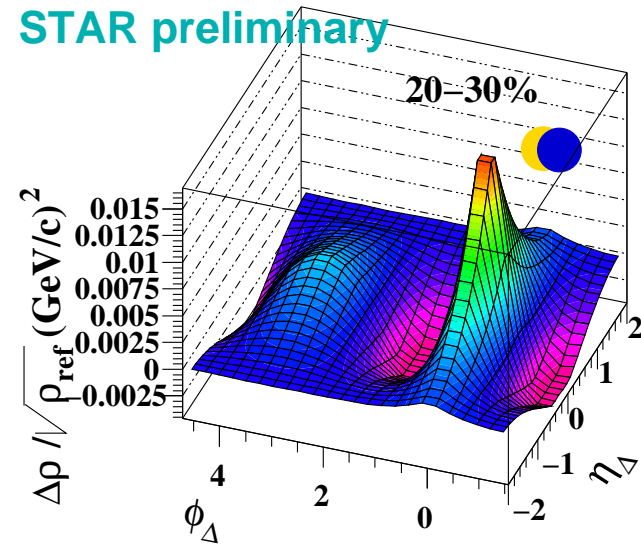


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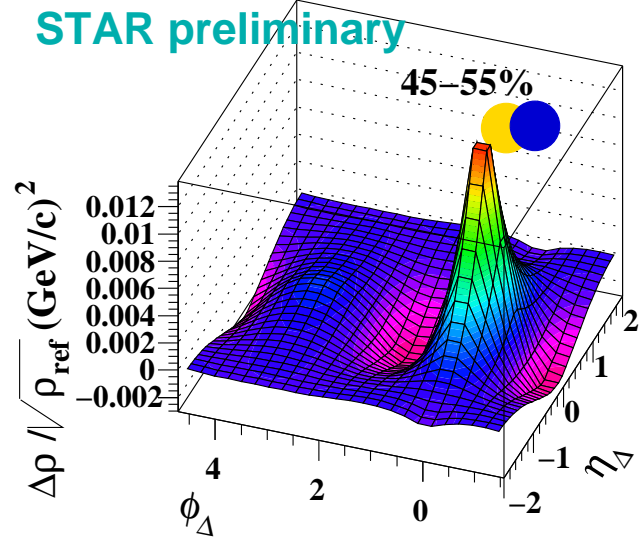
STAR preliminary



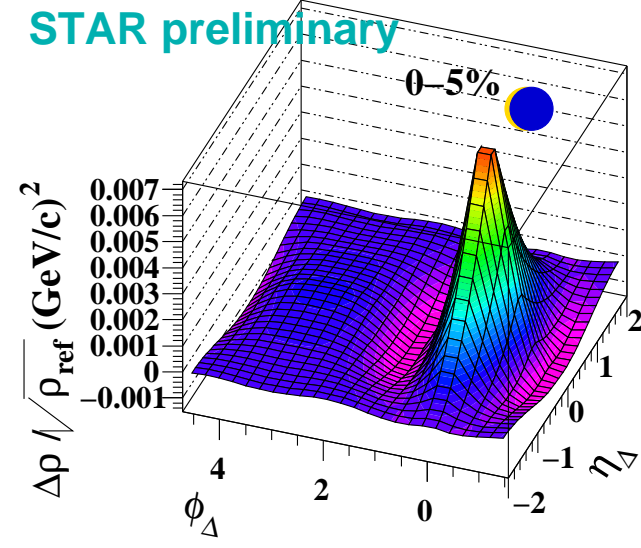
STAR preliminary



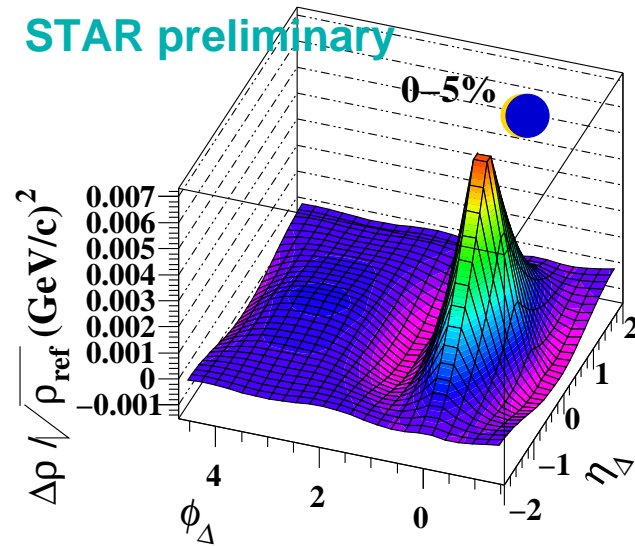
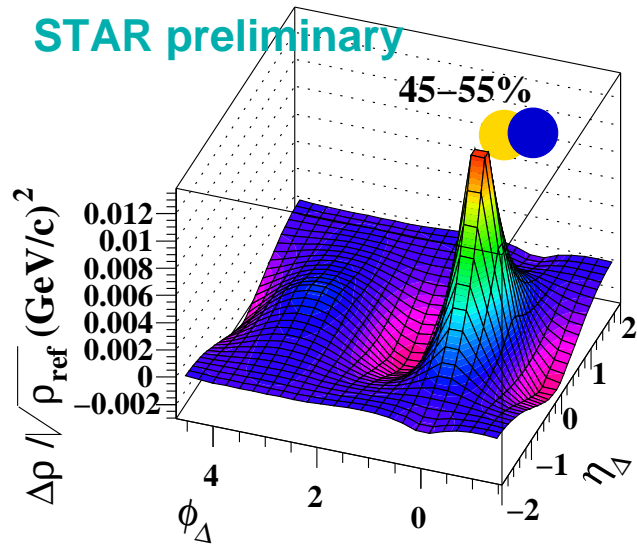
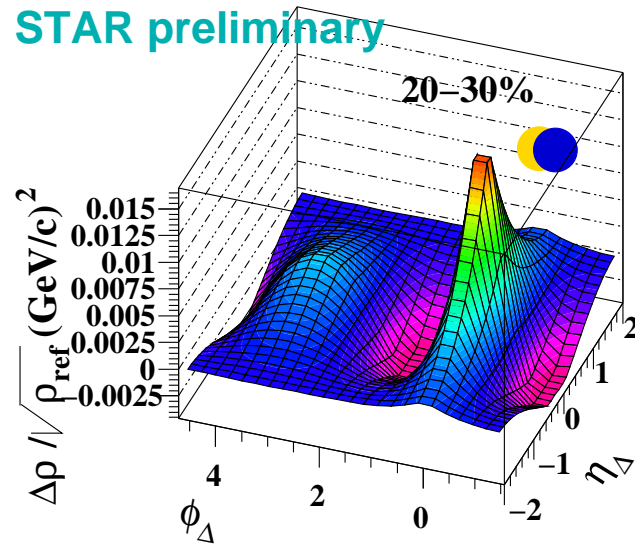
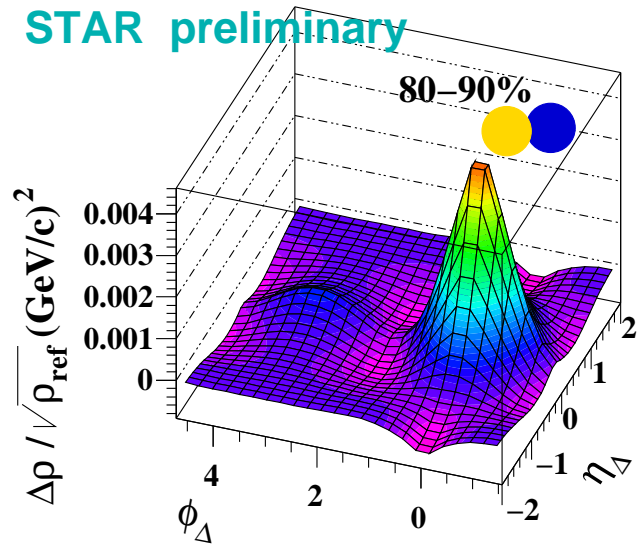
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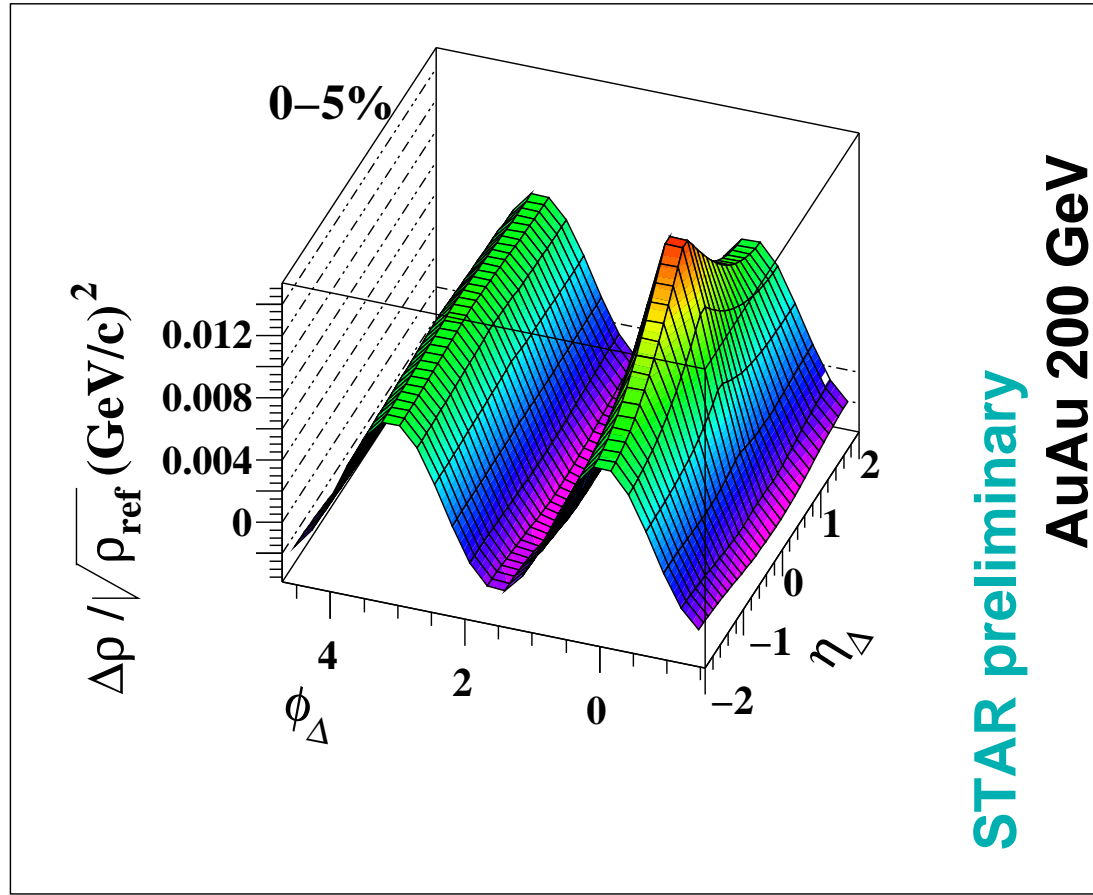
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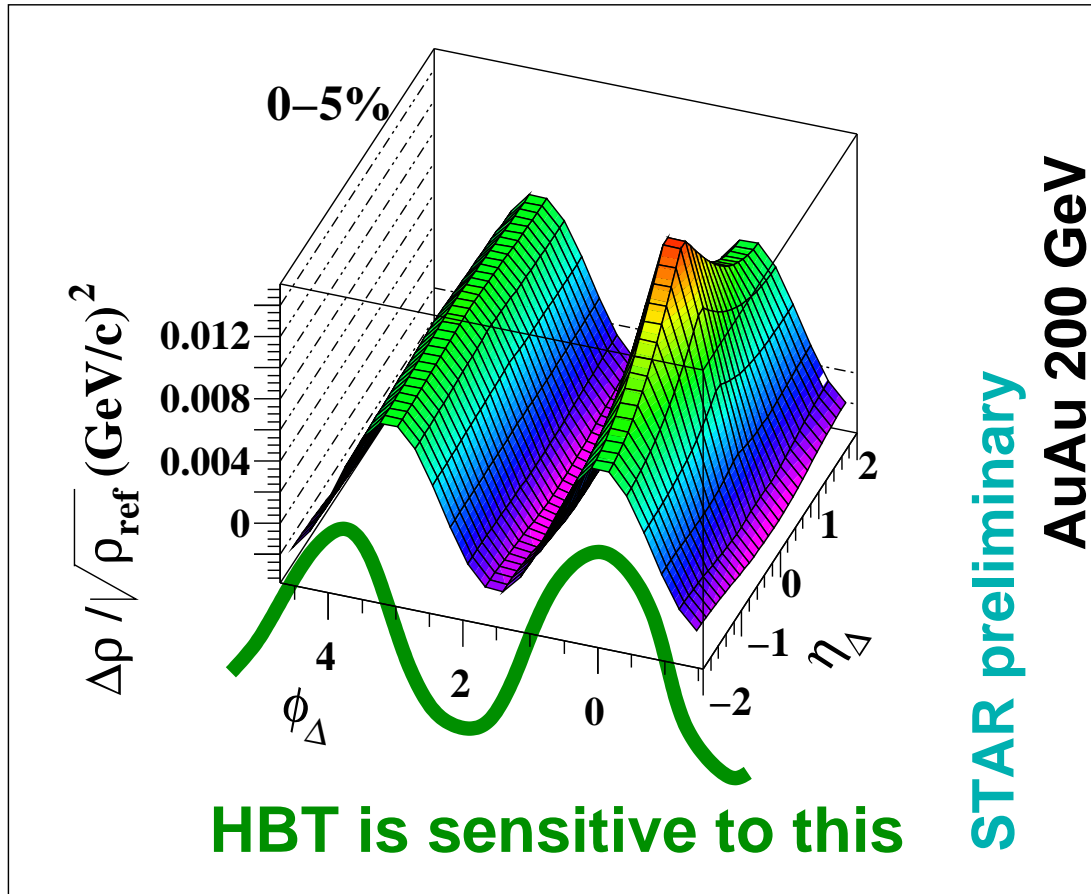
AuAu 200 GeV. In η , correlation broadens with centrality; in ϕ the trend is opposite. The surrounding background seems to recoil.

17 Minijet implications for HBT

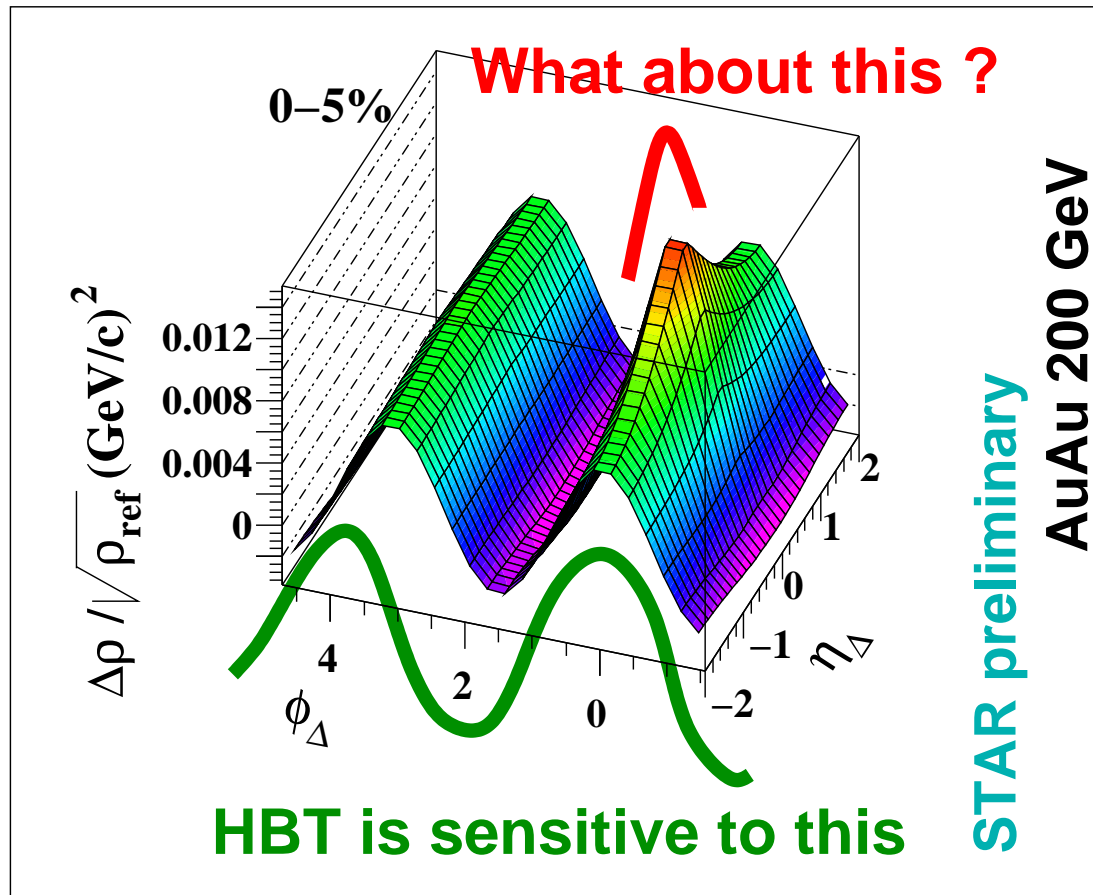
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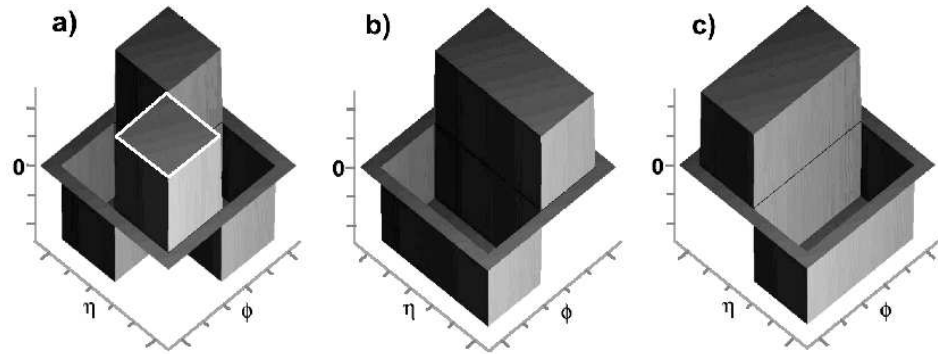
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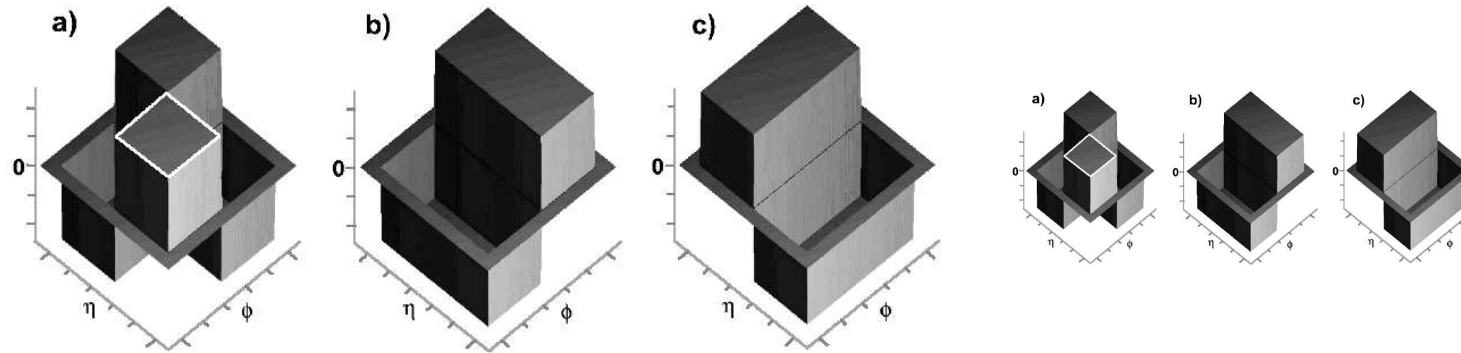
Minijet contribution at soft p_t has been hitherto ignored in the HBT studies. It is likely to contribute to the HBT puzzle by reducing homogeneity lengths/two current correlation length, as compared to the fully equilibrated case. Source of space-momentum correlation.

18 Local hadron density fluctuations and Discrete Wavelet Transform (DWT)

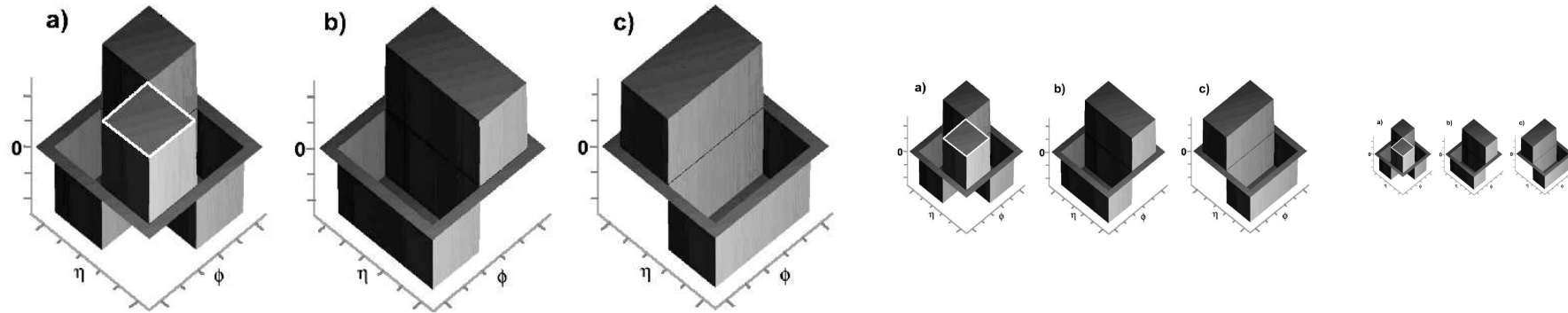
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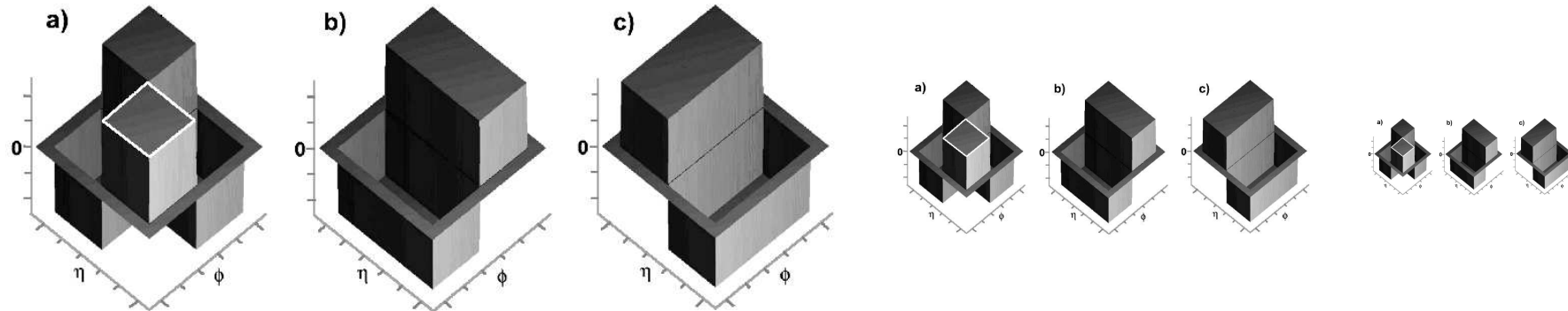
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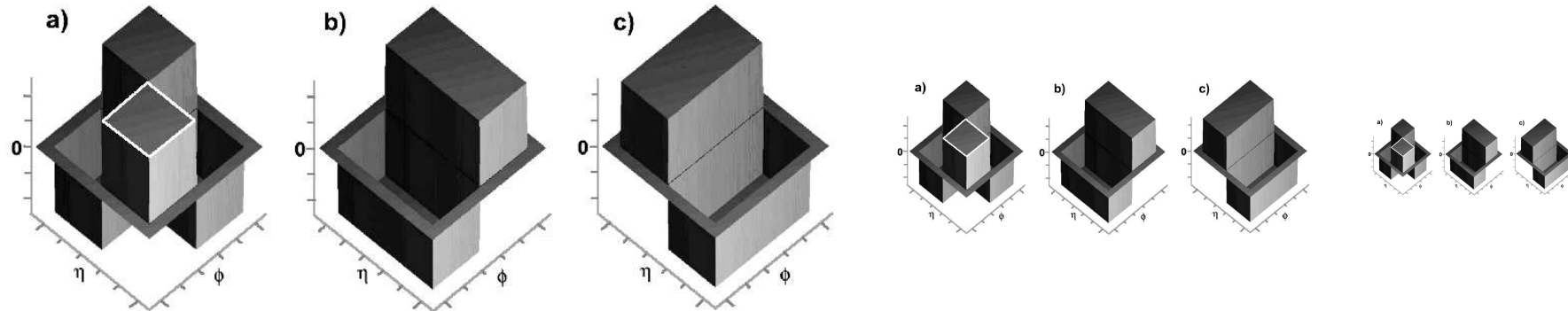


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$F_{m,l,k}^\lambda(\phi, \eta)$ —Haar wavelet **orthonormal basis** in (ϕ, η) . scale fineness (m), directional modes of sensitivity (λ), track density $\rho(\eta, \phi, p_t)$, locations in 2D (l, k) . **DWT is an expansion in this basis.**

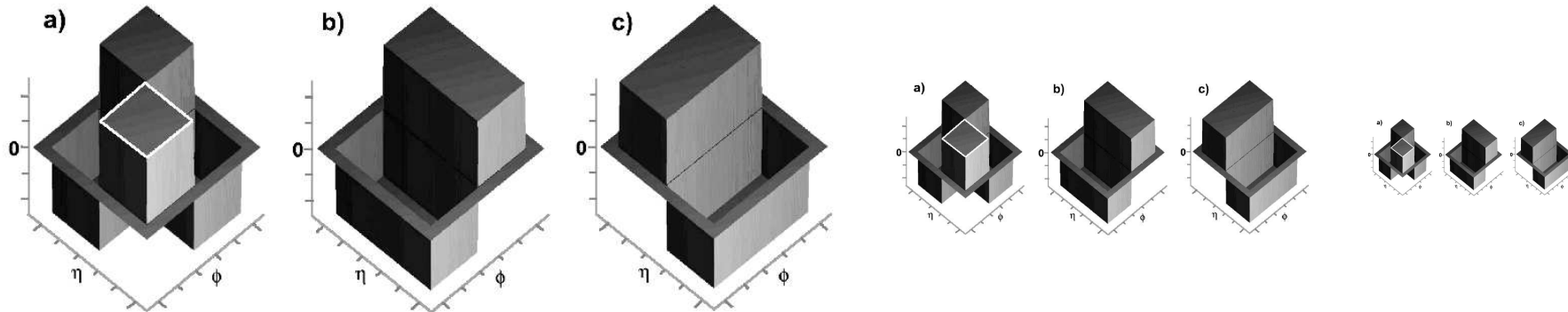
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Power of local fluctuations, mode λ :

$$P^\lambda(m) = 2^{-2m} \sum_{l,k} \langle \rho, F_{m,l,k}^\lambda \rangle^2 \quad (19)$$

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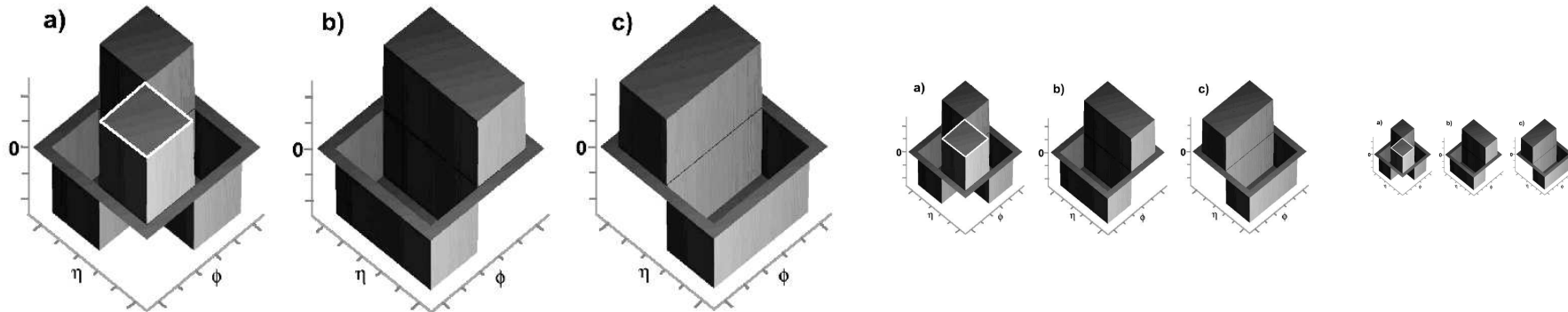
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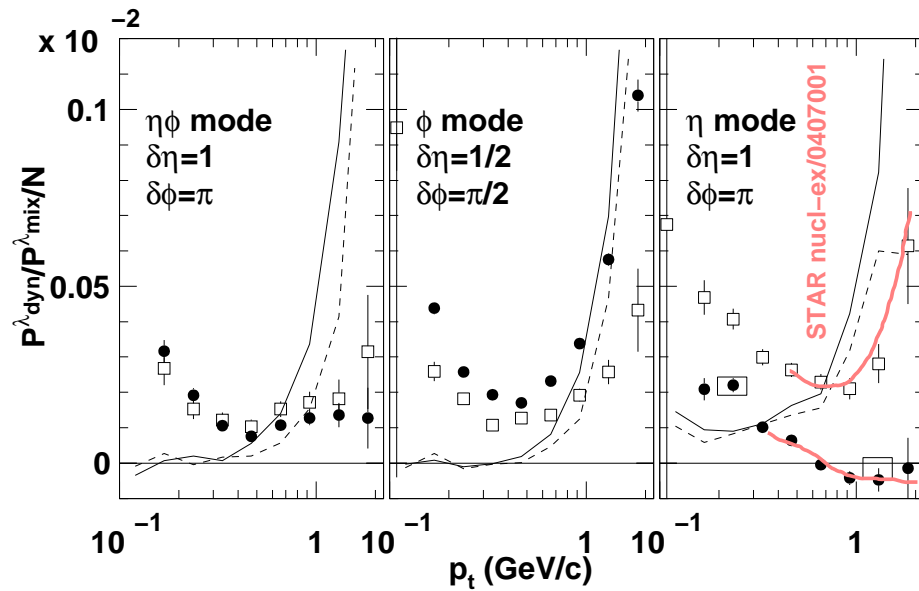
$$P_{\text{dyn}}^\lambda(m) \equiv P_{\text{true}}^\lambda(m) - P_{\text{mix}}^\lambda(m) \quad (20)$$

Normalized:

$$P_{\text{dyn}}^\lambda(m) / P_{\text{mix}}^\lambda(m) / n(p_t) \quad (21)$$

19 Longitudinal minijet broadening

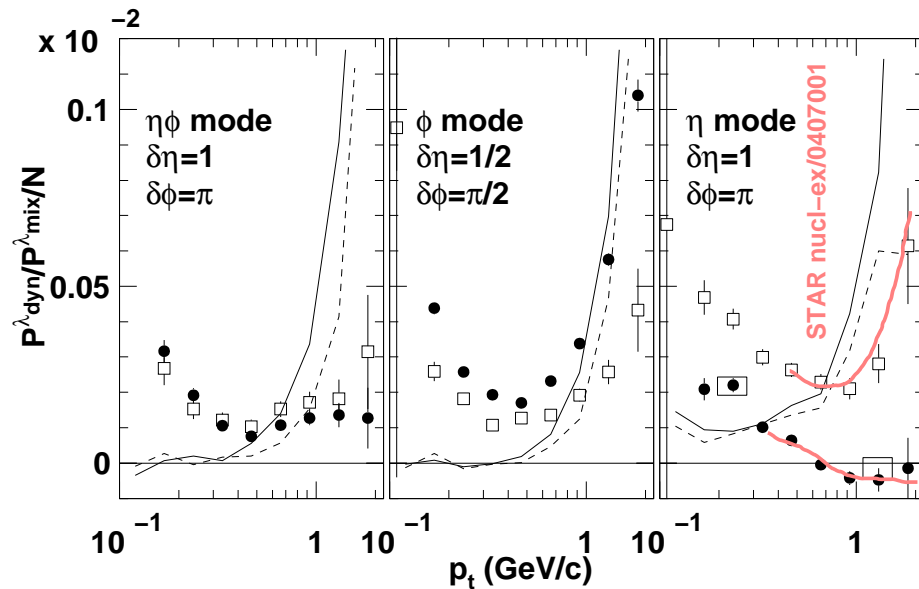
19 Longitudinal minijet broadening



Central events: normalized dynamic texture for fineness scales $m = 0, 1, 0$ from left to right panels, respectively, as a function of p_t .

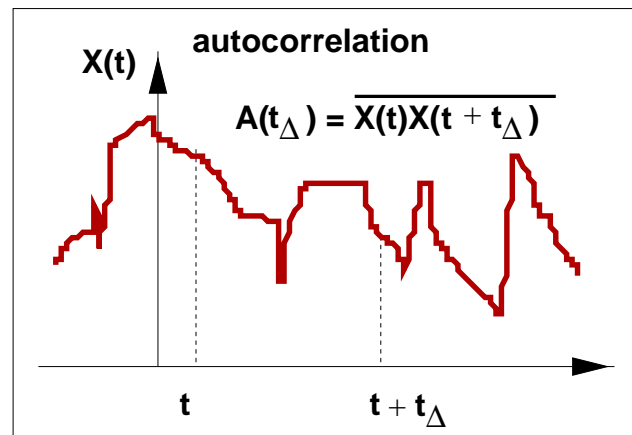
● STAR data; solid line – Hijing without jet quenching; dashed line – Hijing with quenching; □ peripheral STAR data renormalized to compare.

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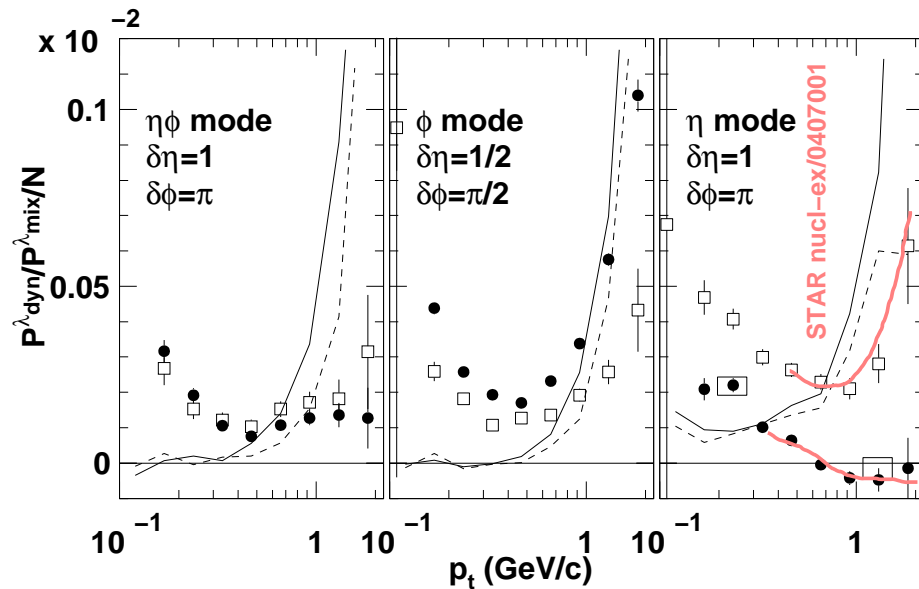


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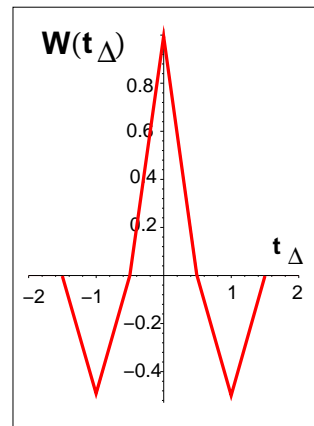
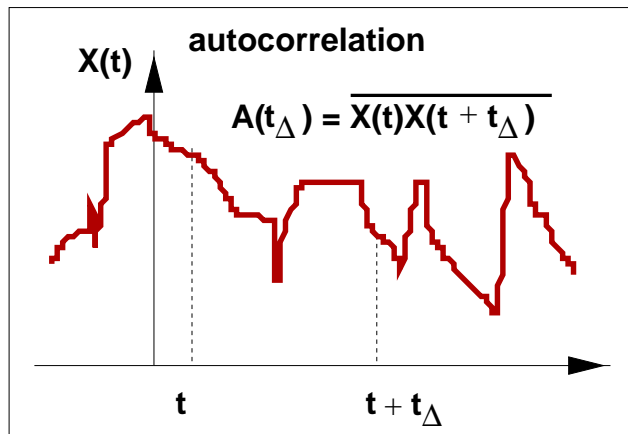


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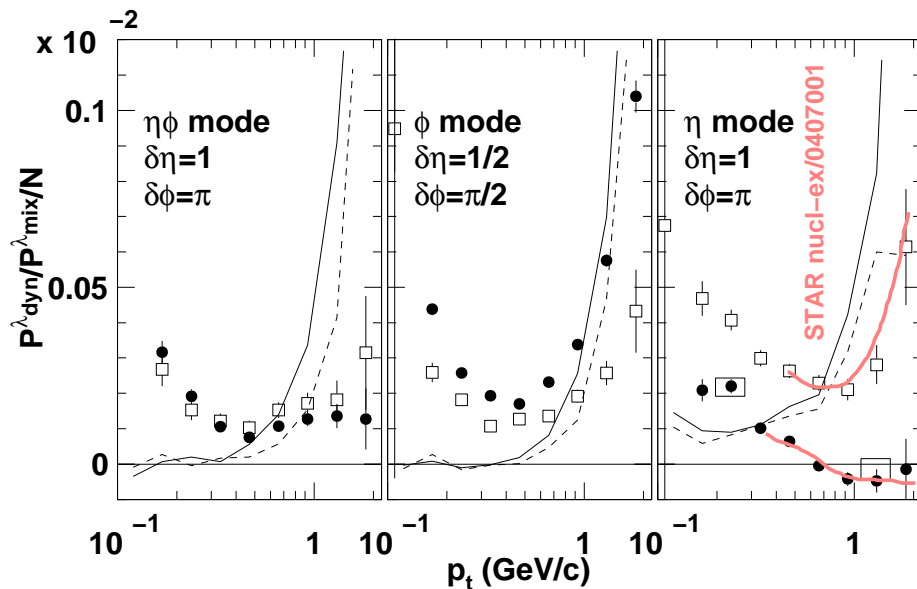


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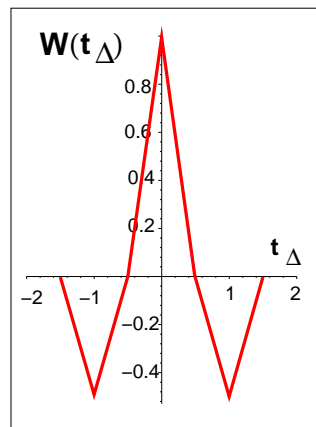
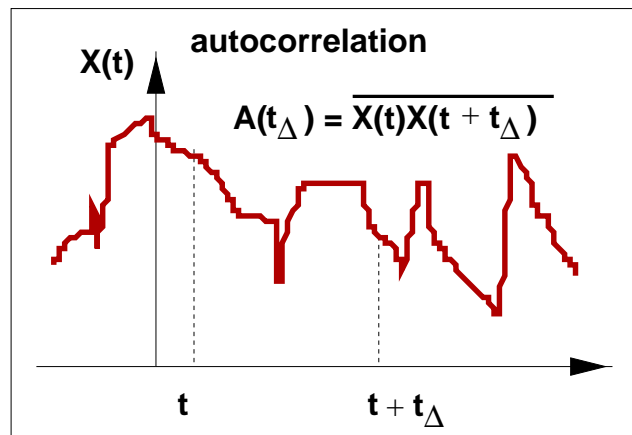


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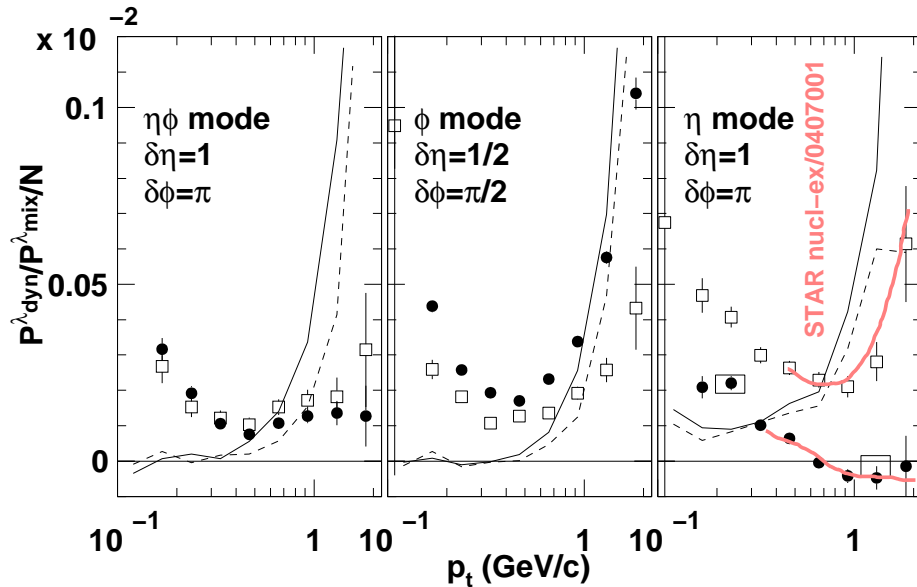
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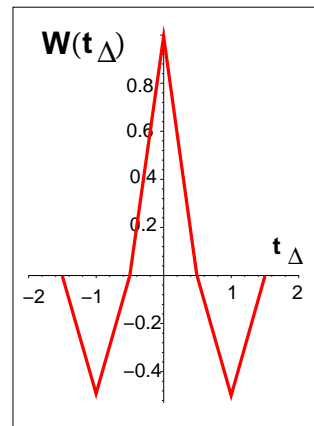
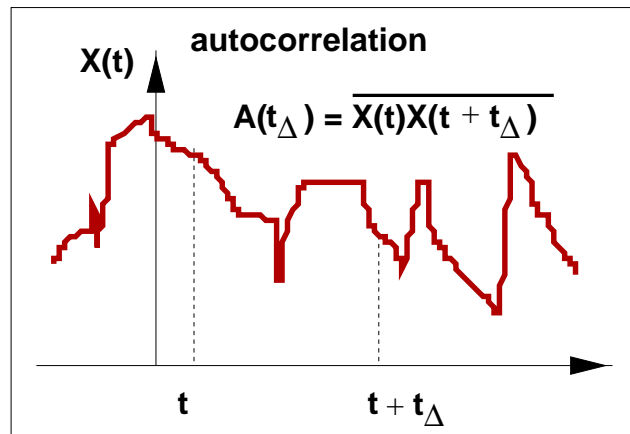
$$P(m) = \int_{-\infty}^{\infty} X(t_{\Delta}/2)X(-t_{\Delta}/2)W(t_{\Delta}, m) dt_{\Delta}, \quad (22)$$

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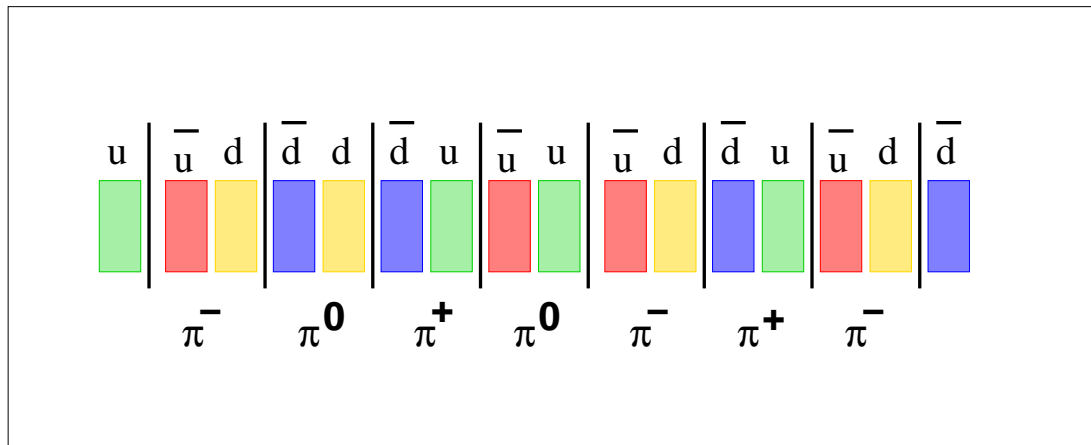


$P(m)$ differentiates correlation on scale m . Minijet elongation \Rightarrow correlation broadening \Leftrightarrow reduced correlation gradient \Leftrightarrow reduced “texture”

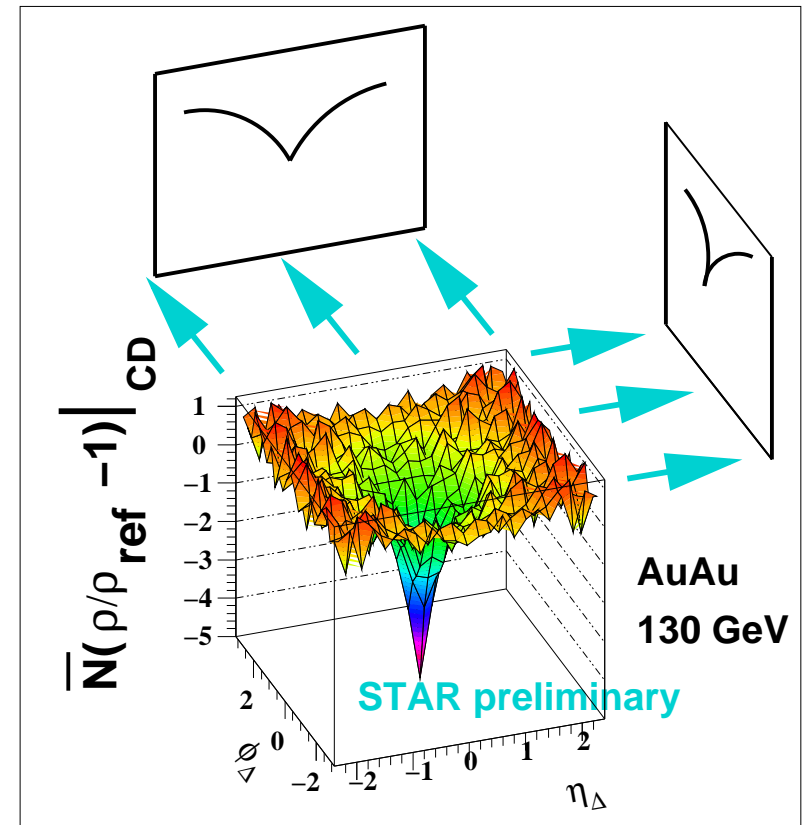
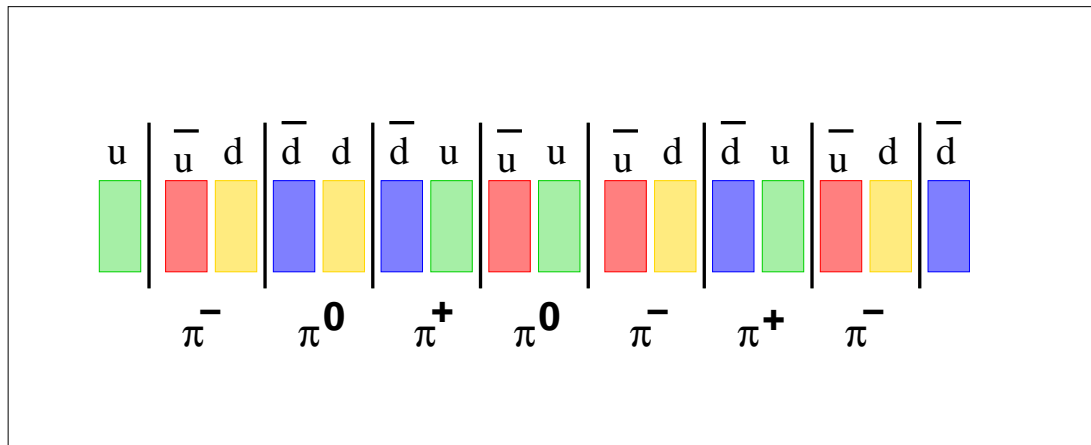
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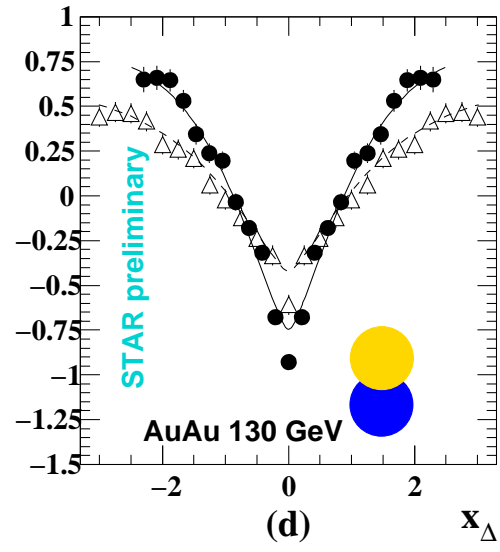
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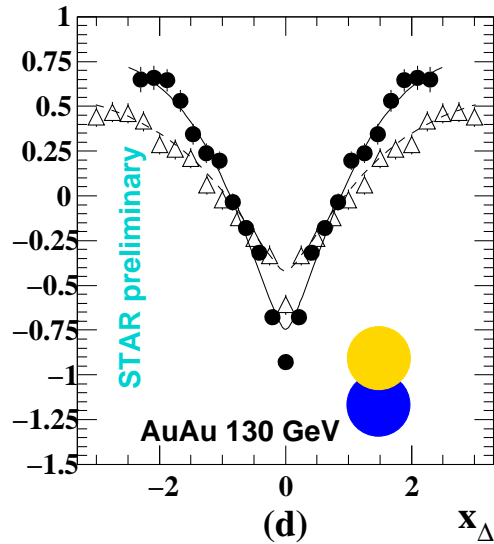
The driving physics: charge conservation in hadronization. Suppress short range correlations – BEC and conversion e^+e^- – by a kinematic pair cut. The \bar{N}_x is good when number of correlation sources $\propto N$.

21 Modified hadronization geometry ?

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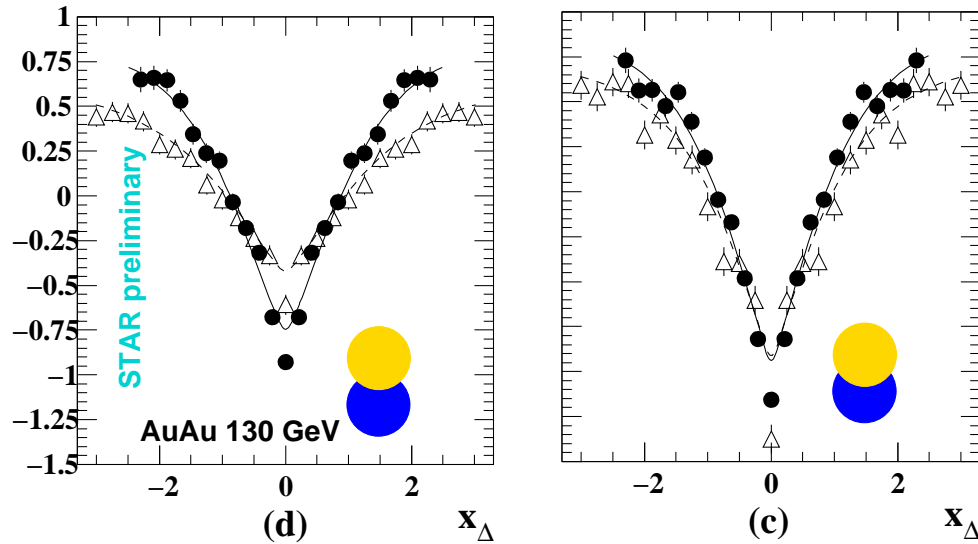


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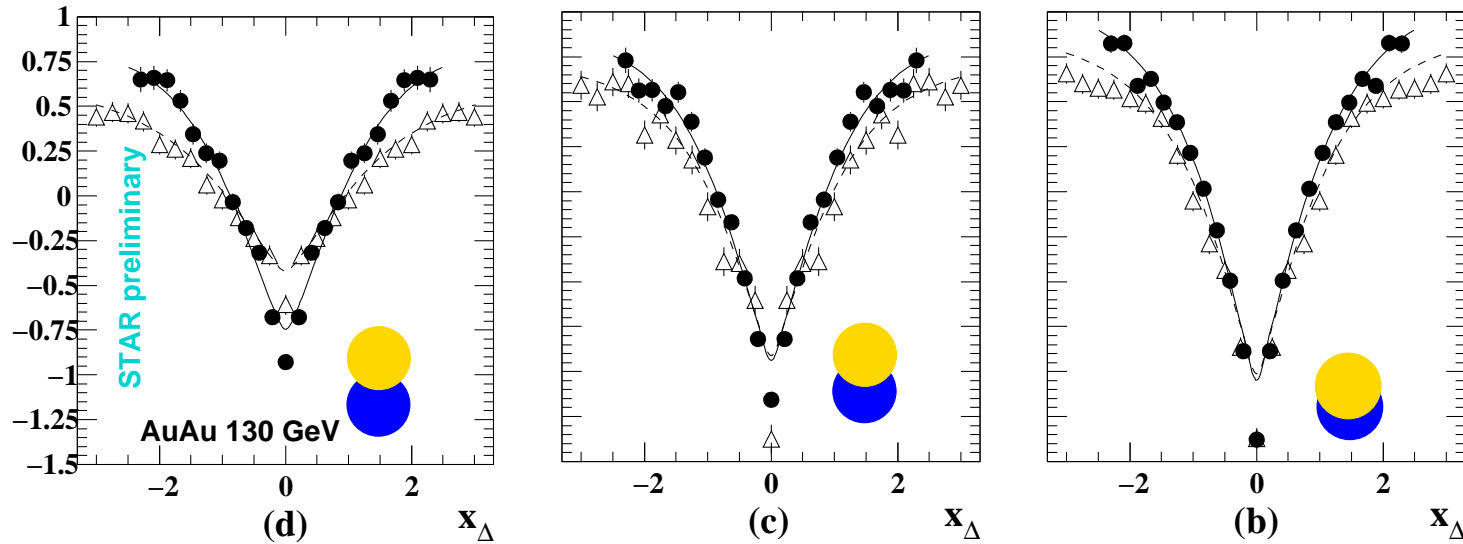
Projections of $\bar{N}[\rho(\eta_\Delta, \phi_\Delta)/\rho(\eta_\Delta, \phi_\Delta)_{\text{ref}} - 1]|_{CD}$ on x_Δ which is ϕ_Δ (Δ) or η_Δ (\bullet). $\eta - \phi$ width disparity (d, peripheral) is gone in (a) \Rightarrow transition from (string) 1D to bulk ($>2D$) fragmentation symmetrizes η and ϕ .

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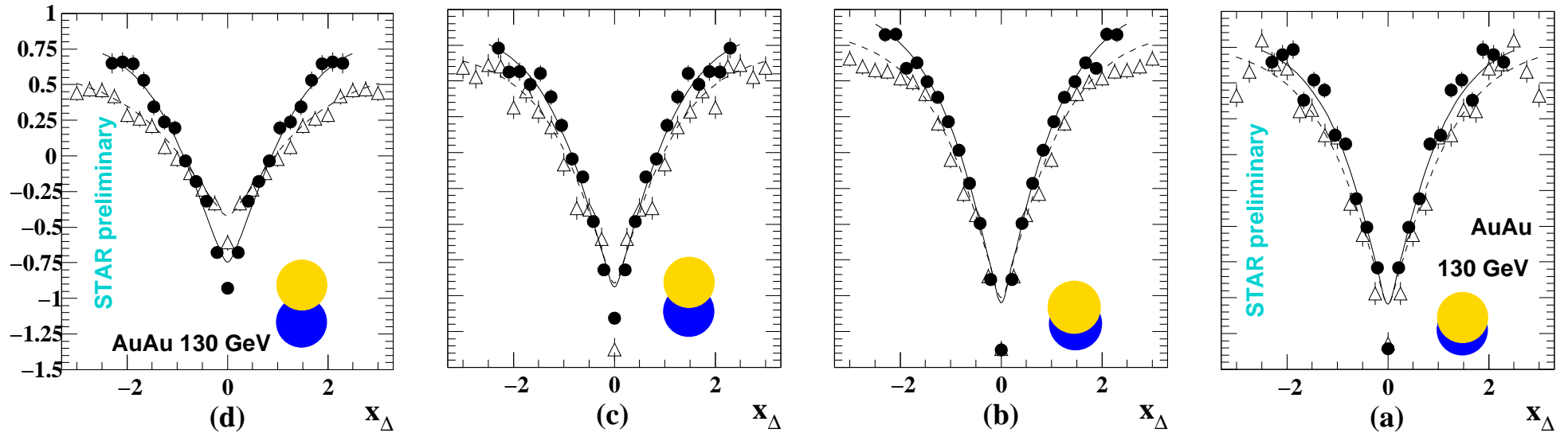
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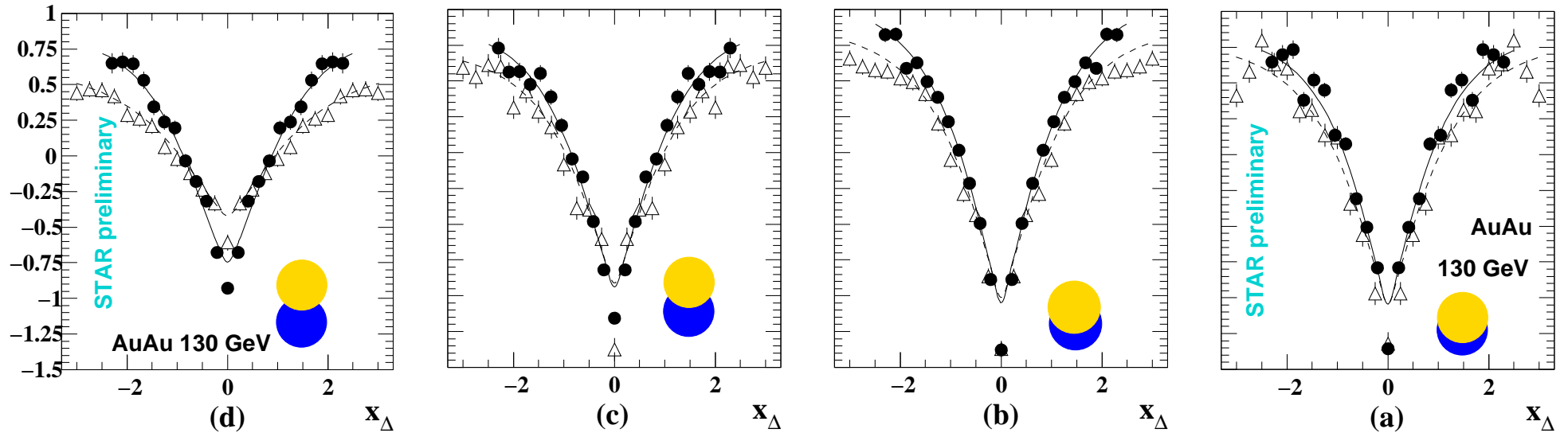
Projections of $\bar{N}[\rho(\eta_{\Delta}, \phi_{\Delta})/\rho(\eta_{\Delta}, \phi_{\Delta})_{\text{ref}} - 1]|_{CD}$ on x_{Δ} which is ϕ_{Δ} (Δ) or η_{Δ} (\bullet). $\eta - \phi$ width disparity (d, peripheral) is gone in (a) \Rightarrow transition from (string) 1D to bulk ($>2D$) fragmentation symmetrizes η and ϕ .

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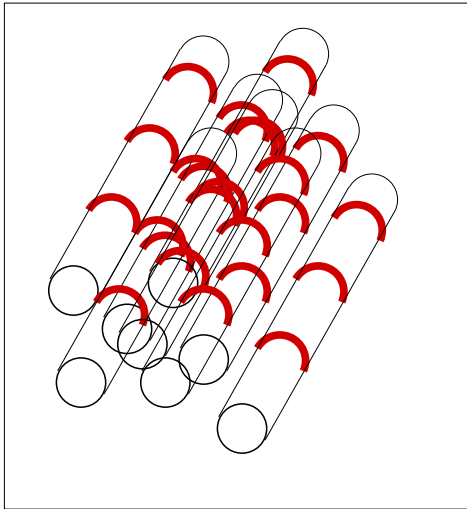


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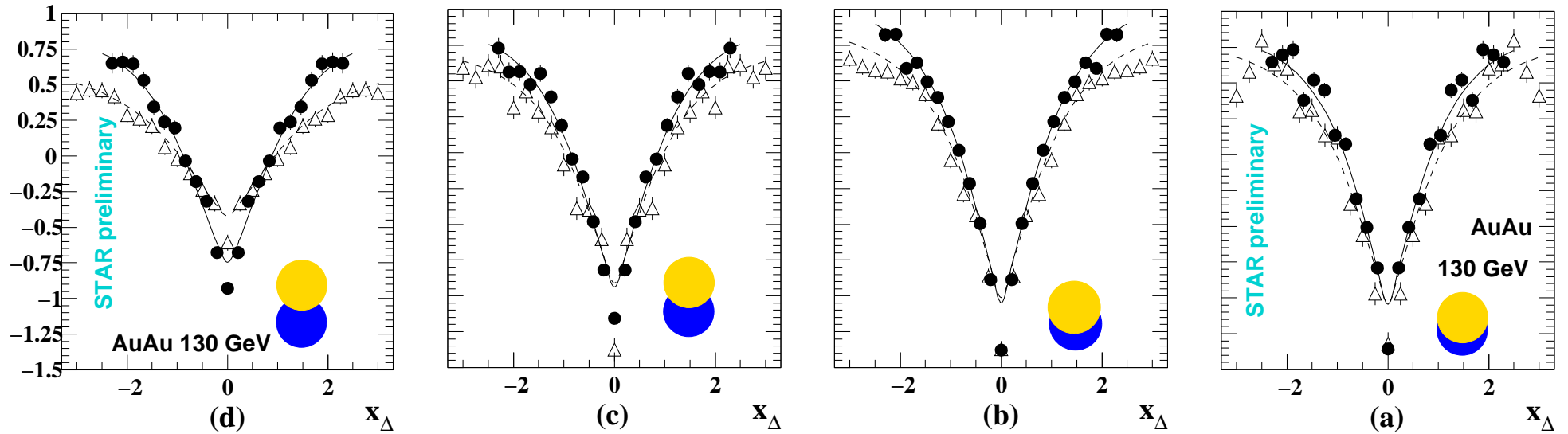
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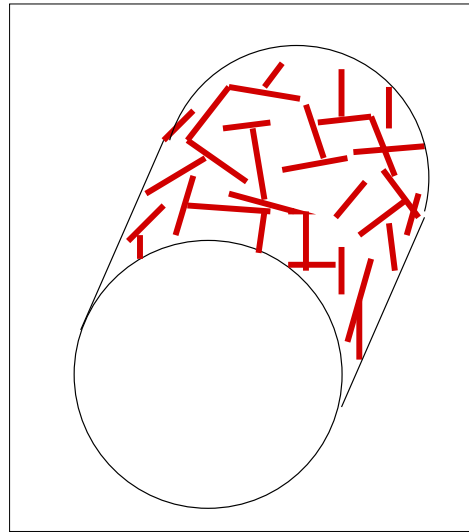
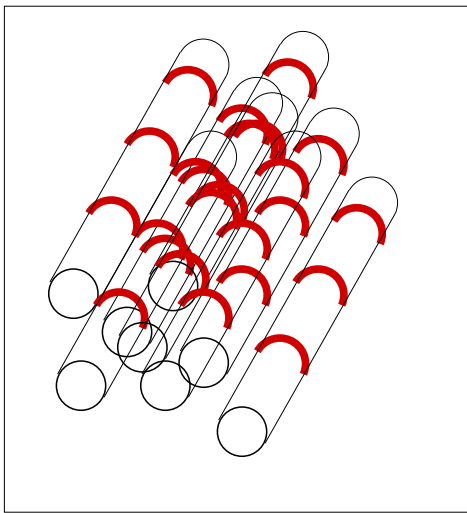
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- Increased symmetry of the charge-dependent correlation on (η, ϕ) in the central collisions may point to a change in the hadronization geometry in the medium

23 Extra slides

24 Notation and glossary

- HBT – Hanbury-Brown and Twiss technique (intensity interferometry)
- i – particle index
- N – total number of particles in an event
- $n, n(p_t)$ – number of particles within a kinematic cut (bin)
- $\overline{(\dots)}$ – average over events
- \hat{p}_t – inclusive mean p_t per particle
- x_Δ (variants: $t_\Delta, \eta_\Delta, \phi_\Delta \dots$) – difference variable $= x_i - x_{i'}$.
- δx – scale (range of local integration; e.g. see Fig.13)
- Δx – the upper limit on δx

25 Directed flow and antiflow

