

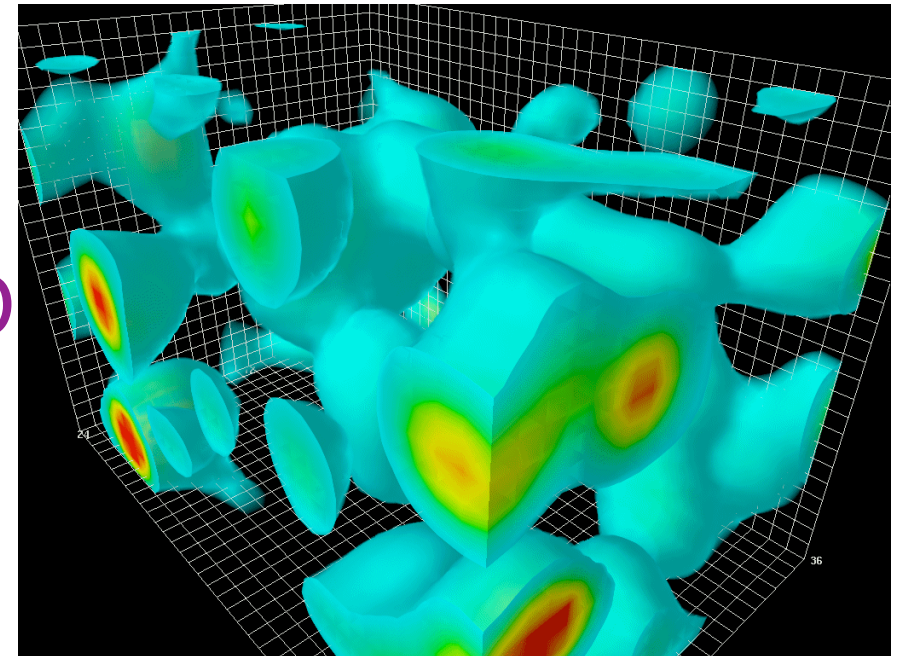
# $e+A$ physics at an Electron-Ion Collider - understanding the glue which binds us all

Matthew A. C. Lamont  
Brookhaven National Lab

# What do we know about gluons?

- **Gluons:**

- ➔ Mediators of the strong interaction
- ➔ Determine essential features of QCD
  - Asymptotic freedom from gluon loops
- ➔ Dominate structure of QCD vacuum

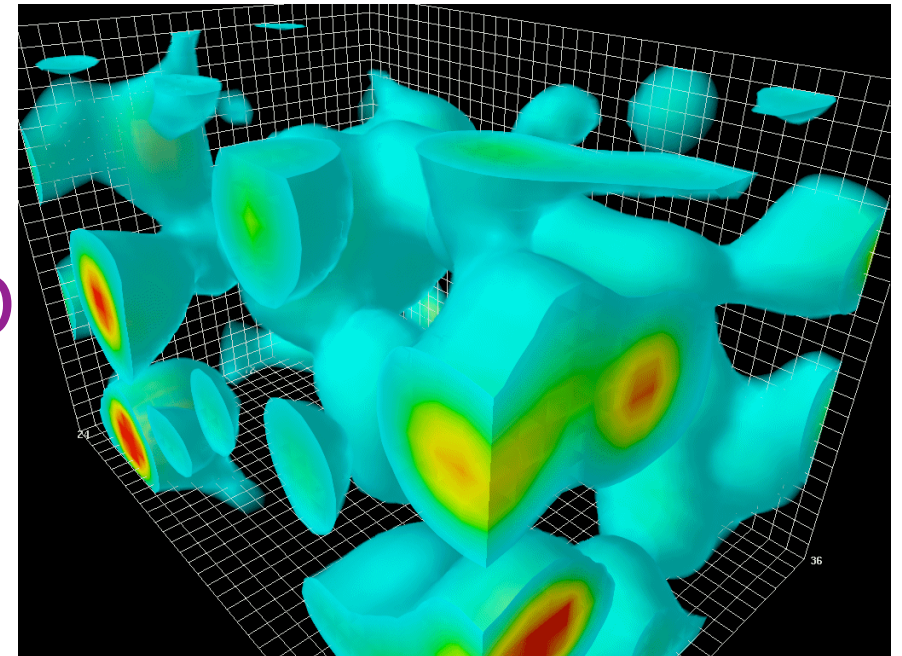


Action ( $\sim$ energy) density fluctuations of gluon-fields in QCD vacuum ( $2.4 \times 2.4 \times 3.6$  fm) (Derek Leinweber)

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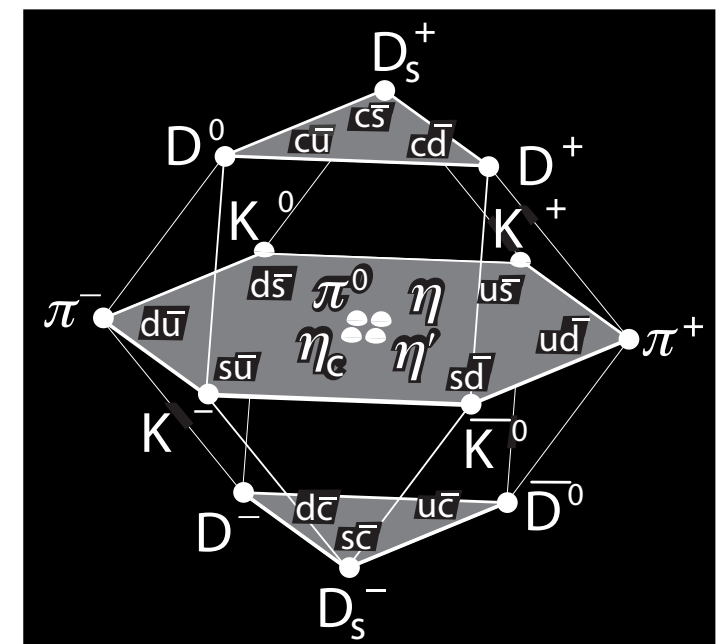


Action ( $\sim$ energy) density fluctuations of gluon-fields in QCD vacuum ( $2.4 \times 2.4 \times 3.6$  fm) (Derek Leinweber)

- **Hard to “see” glue in the low-energy world**

- ➔ Gluon degrees of freedom “missing” in hadronic spectrum
  - Constituent Quark Picture?
- ➔ From DIS:
  - Drive the structure of baryonic matter already at medium-x

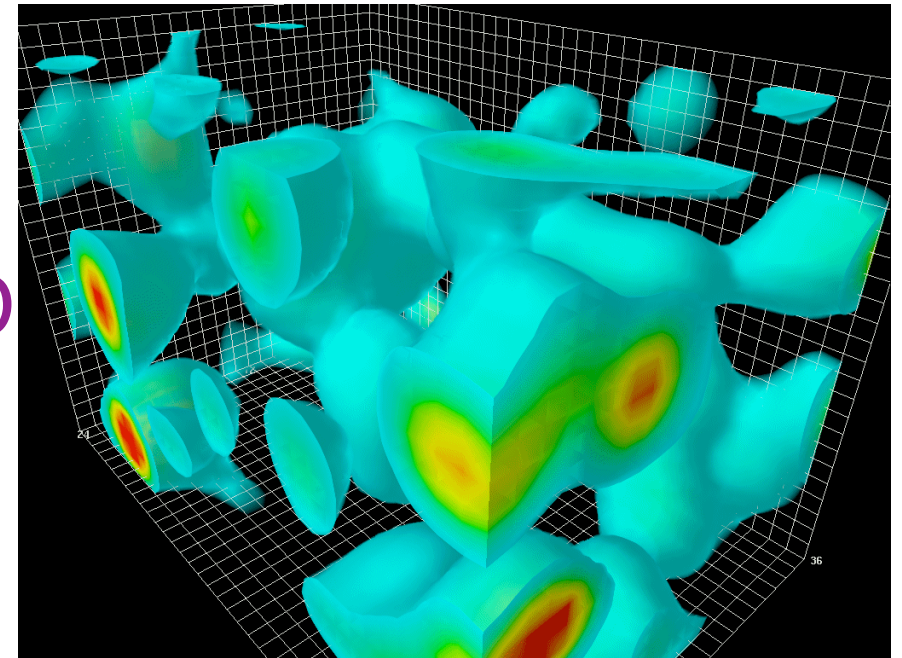
★ Crucial players at RHIC and LHC



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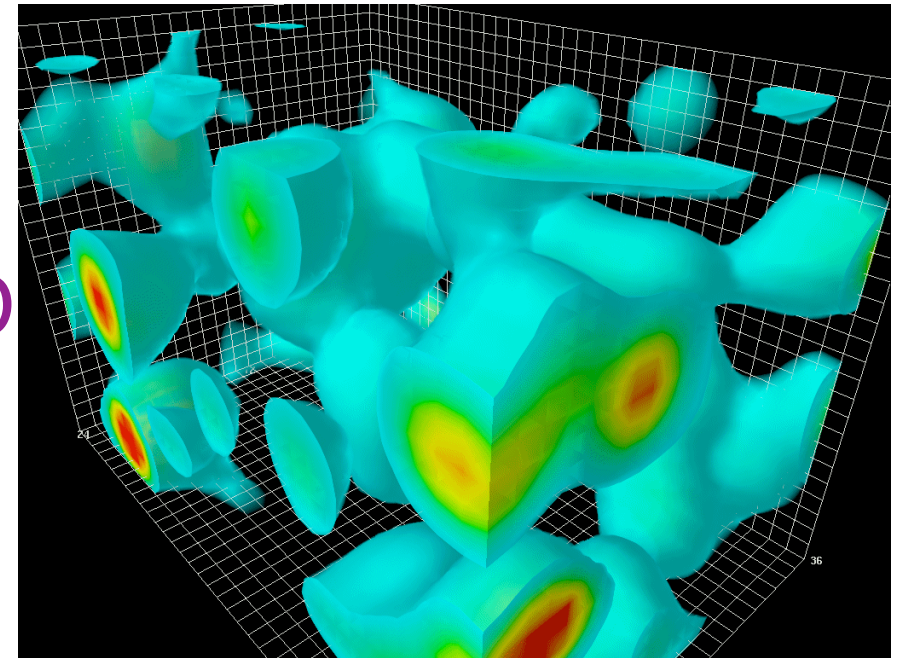
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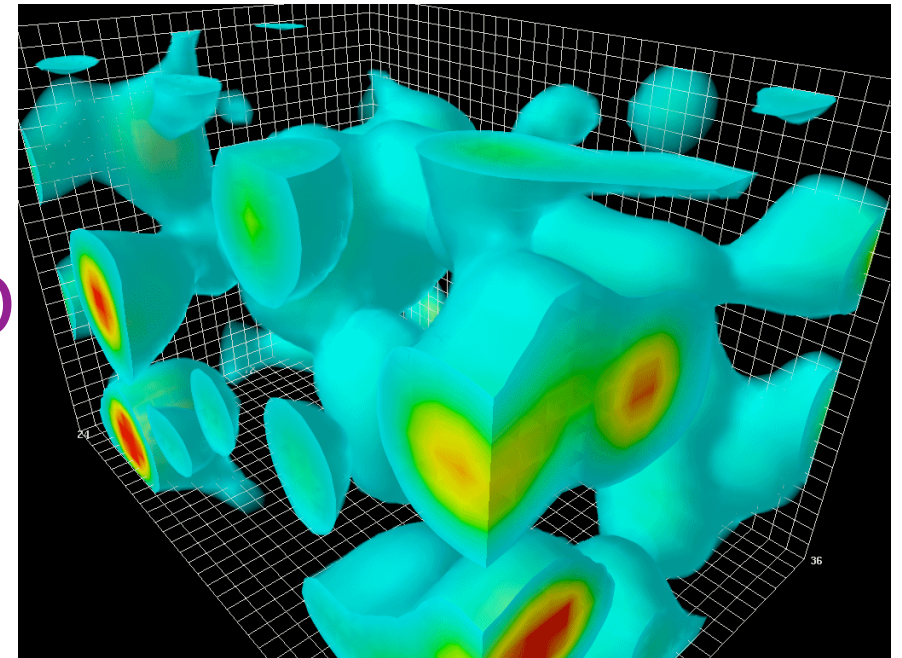
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- ➔ What is the **spatial** and **momentum** distribution of gluons in nuclei/nucleons?
- ➔ What are the **properties of high-density gluon matter**?
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- ➔ What role do the **gluons** play in the **spin structure of the nucleon**?

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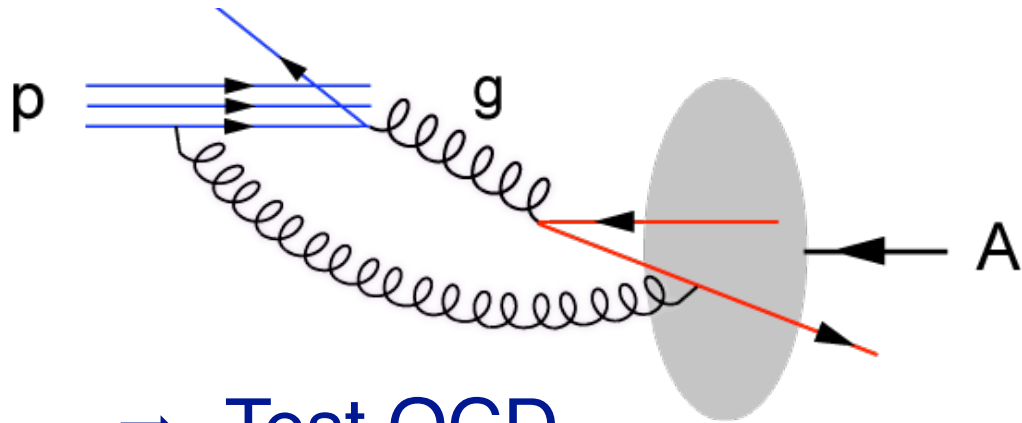
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How do we get to the answers?

# How do we measure Glue ?

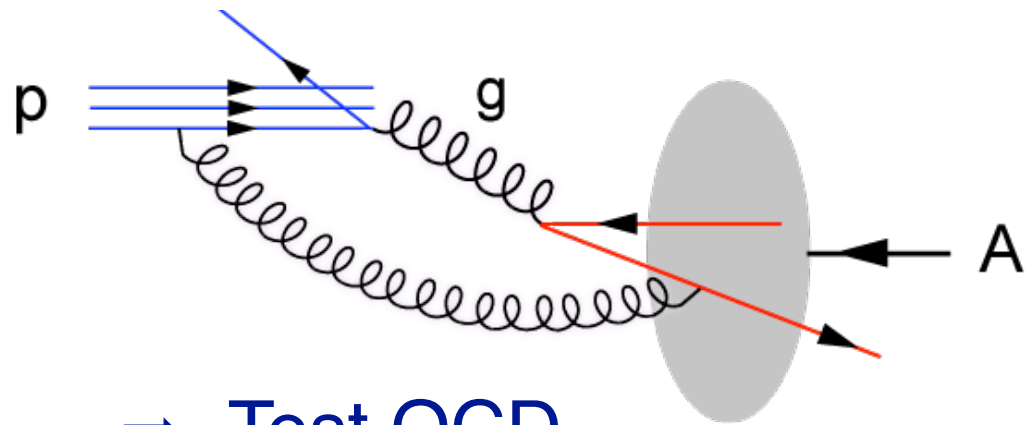
- Hadron-Hadron



- ➔ Test QCD
- ➔ Probe/Target interaction directly via gluons
- ➔ Lacks the direct access to collision kinematics
- ➔ Interactions with other partons modifies nuclear wave function

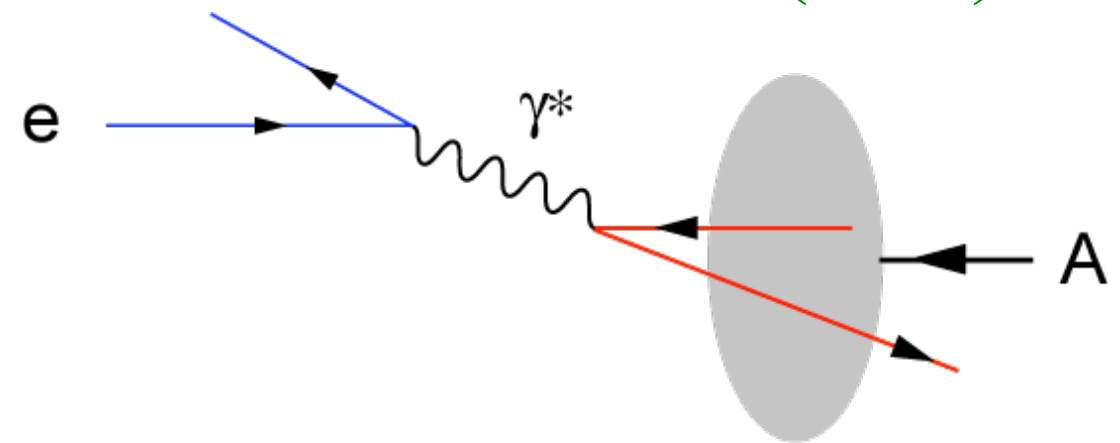
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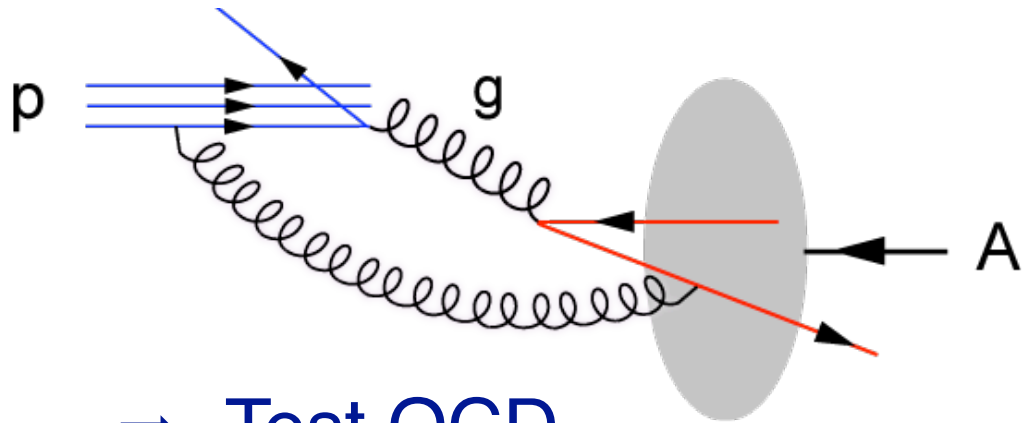
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- ★ Explore QCD & Hadron Structure
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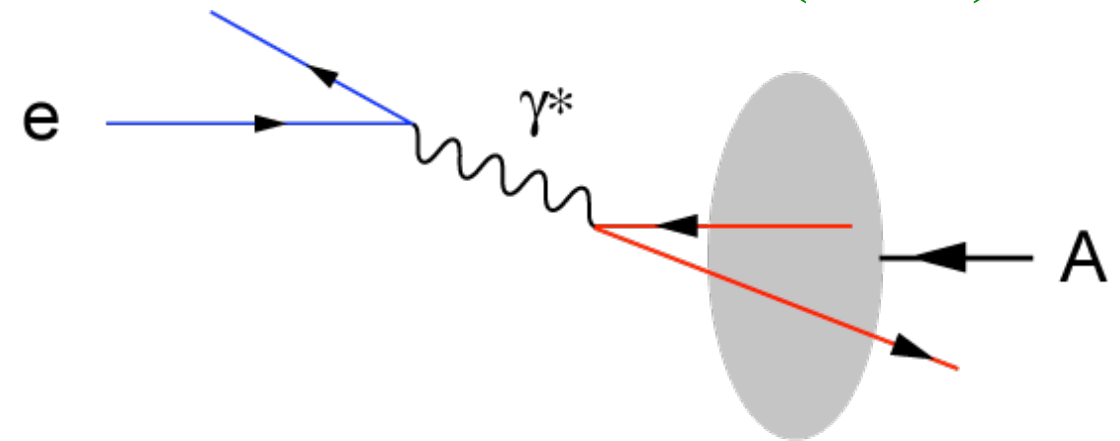
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- ★ Explore QCD & Hadron Structure
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Both are **complementary** and provide excellent information on properties of gluons in the nuclear wave functions

Precision measurements  $\Rightarrow$  DIS

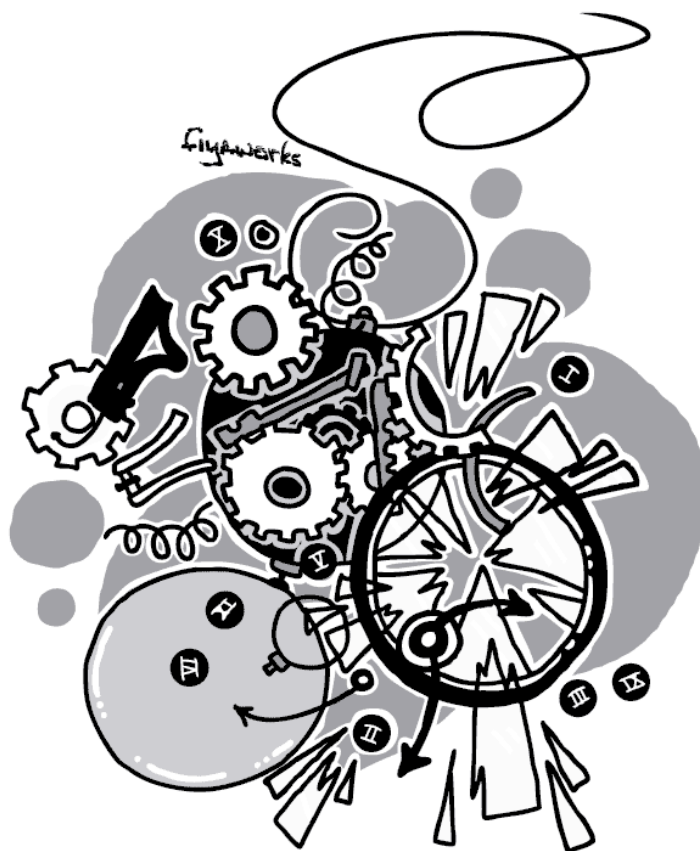
[macl@bnl.gov](mailto:macl@bnl.gov)



# How do we measure Glue ?

*Scattering of hadrons on hadrons  
is like colliding Swiss watches to find out how  
they are build.*

R. Feynman



Precision measurements  $\Rightarrow$  DIS

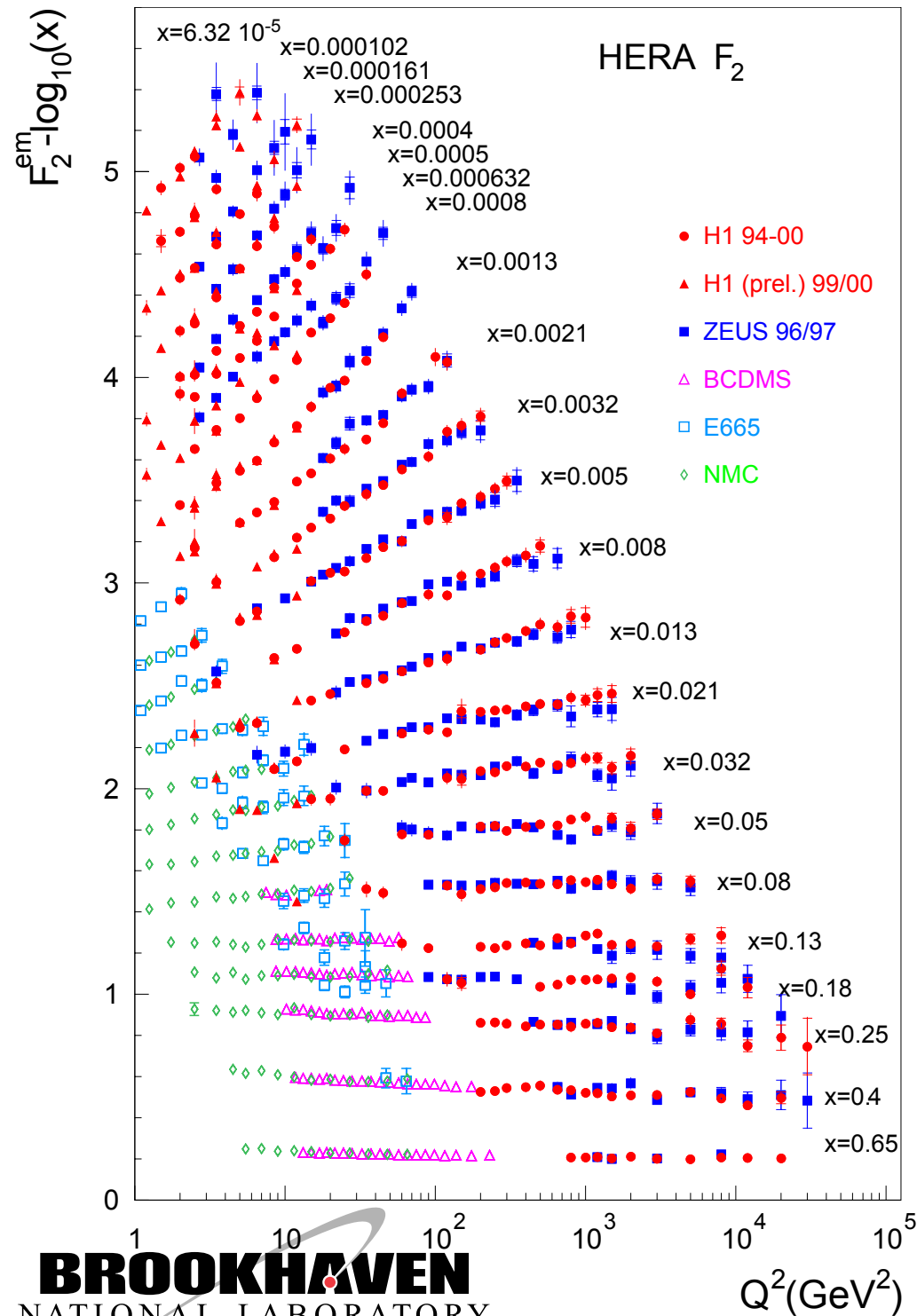
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# How to measure the glue?

$$\frac{d^2\sigma^{ep\rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[ \left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$

**quark+anti-quark  
momentum distributions**

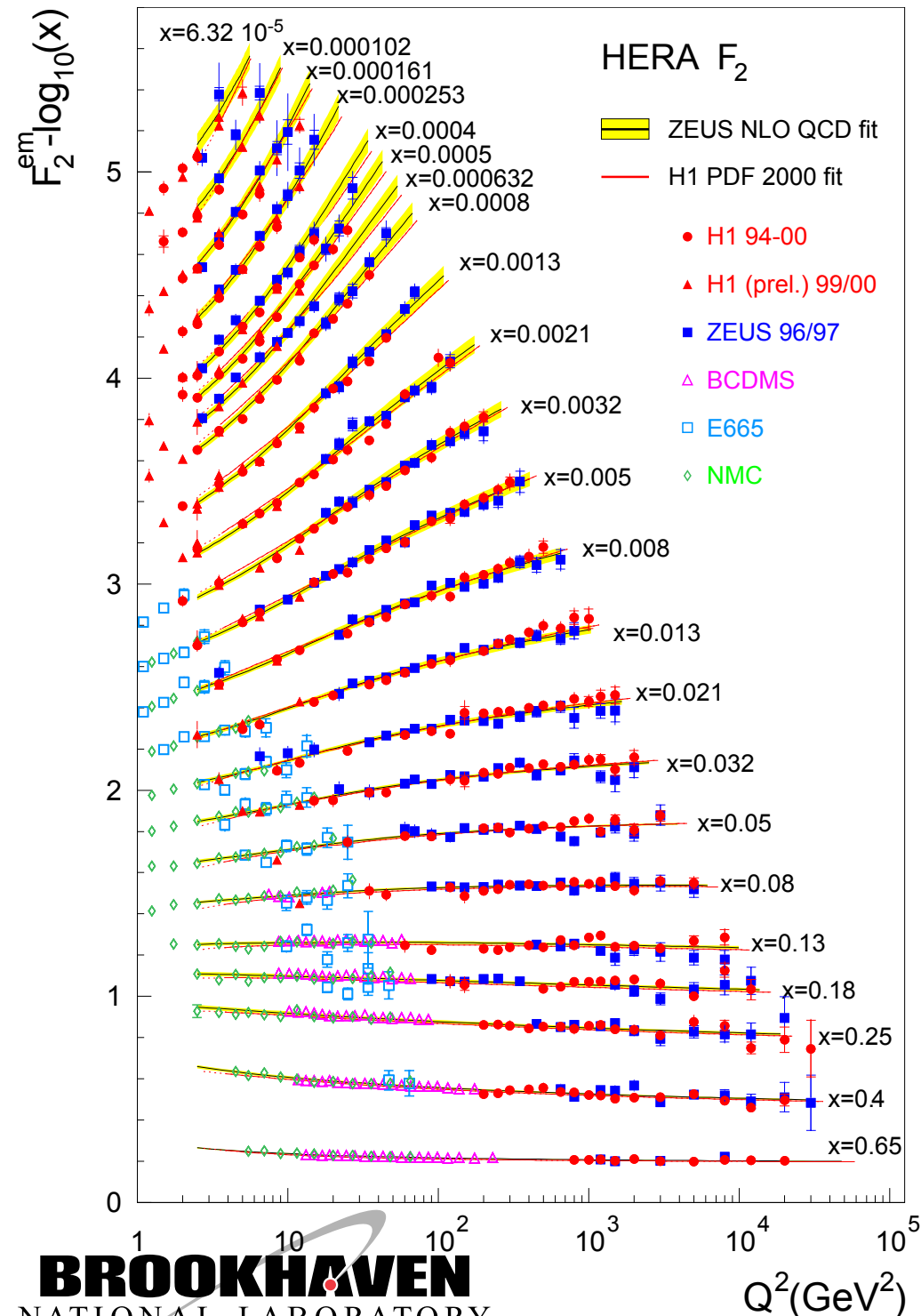
**gluon momentum  
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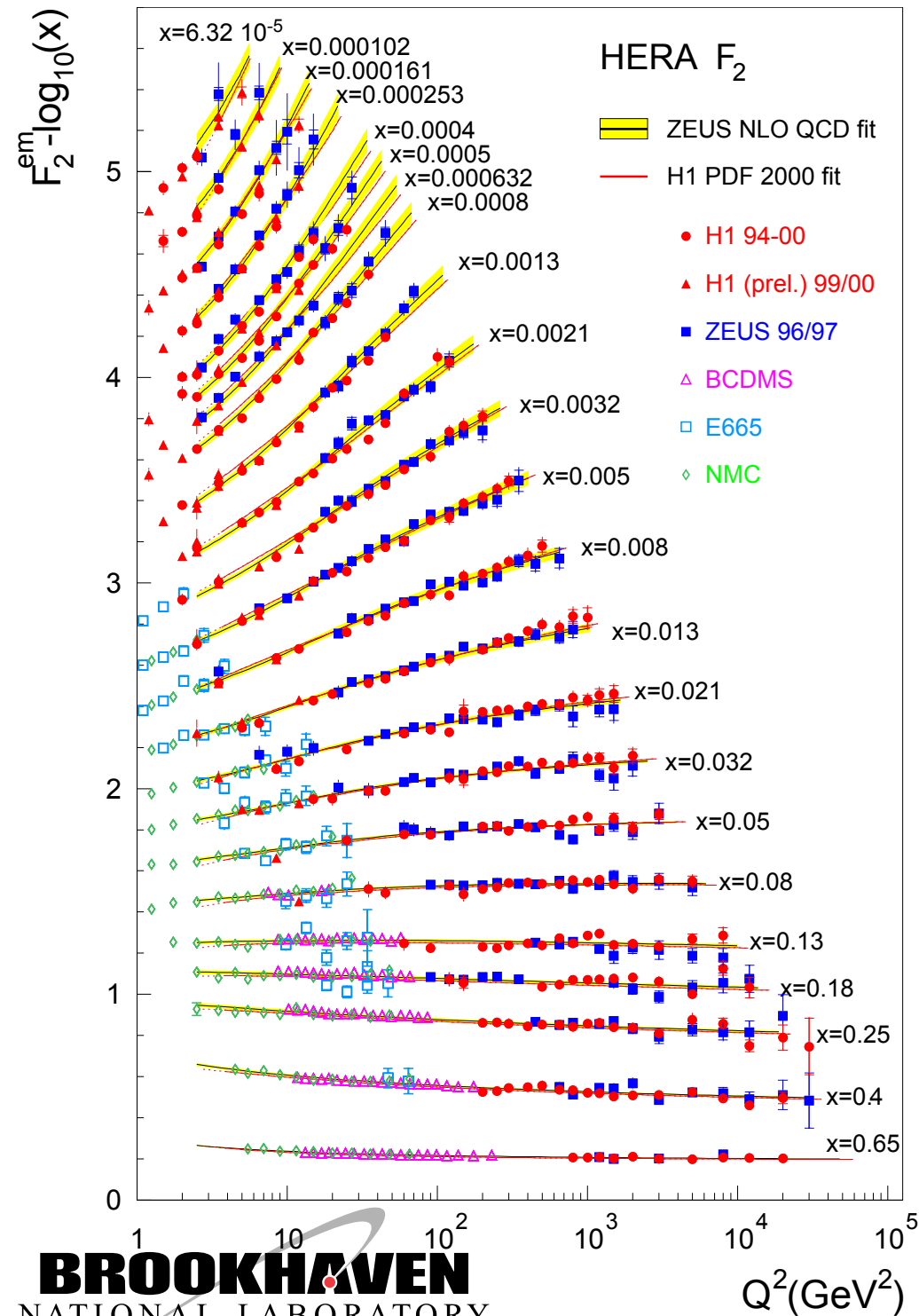
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Scaling violation:  $dF_2/d\ln Q^2$  and linear DGLAP  
Evolution  $\Rightarrow G(x, Q^2)$

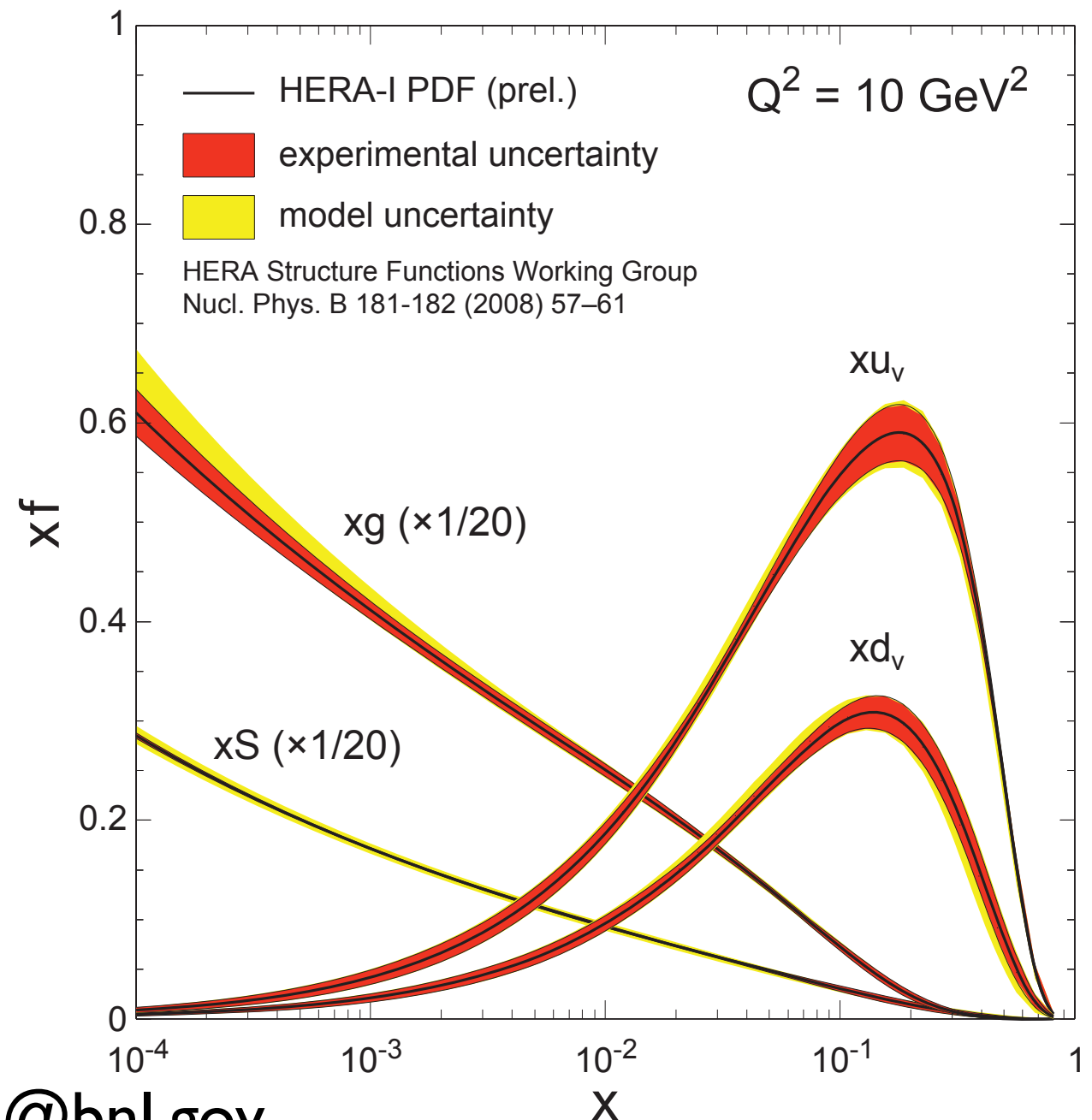


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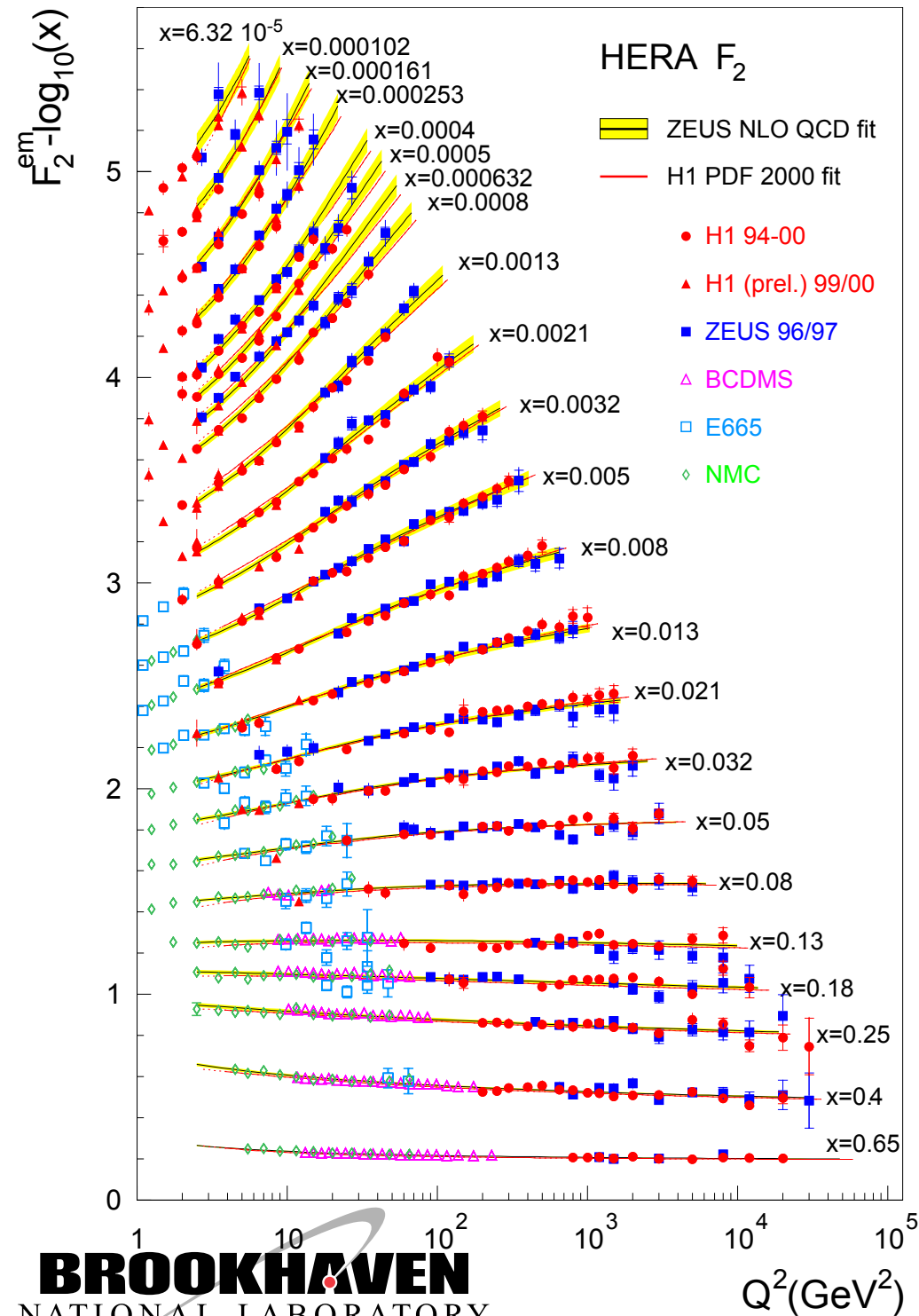
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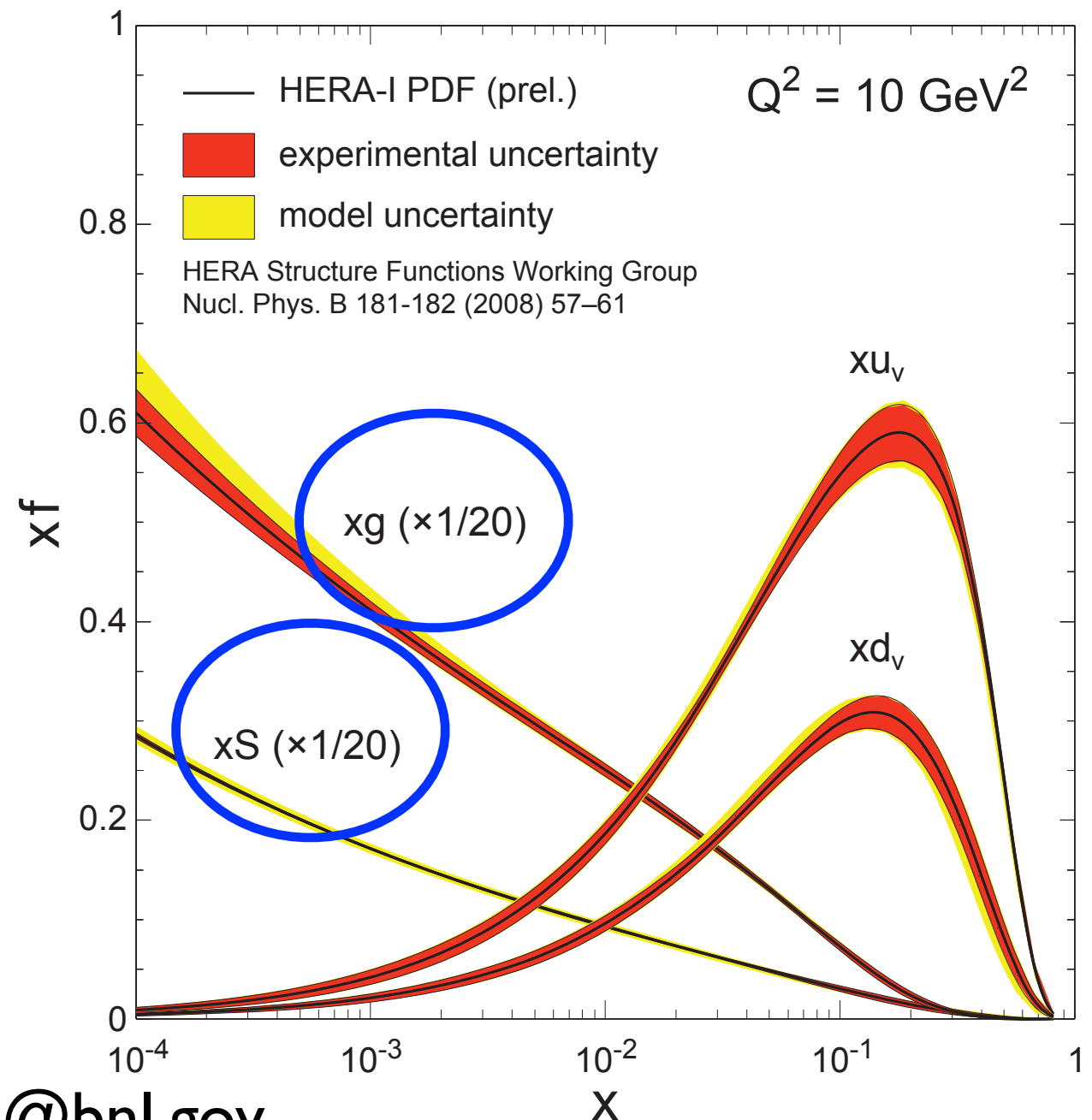
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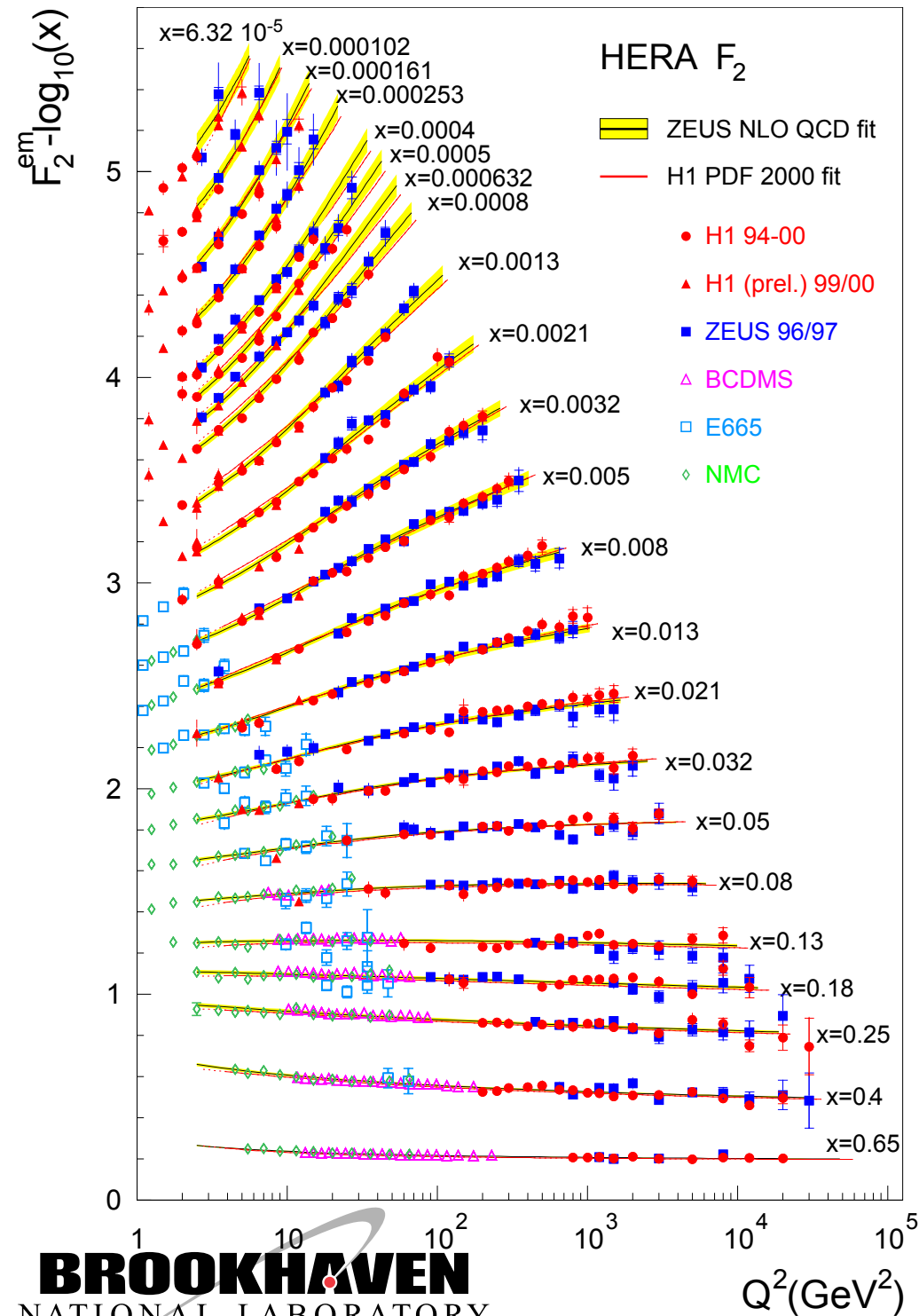


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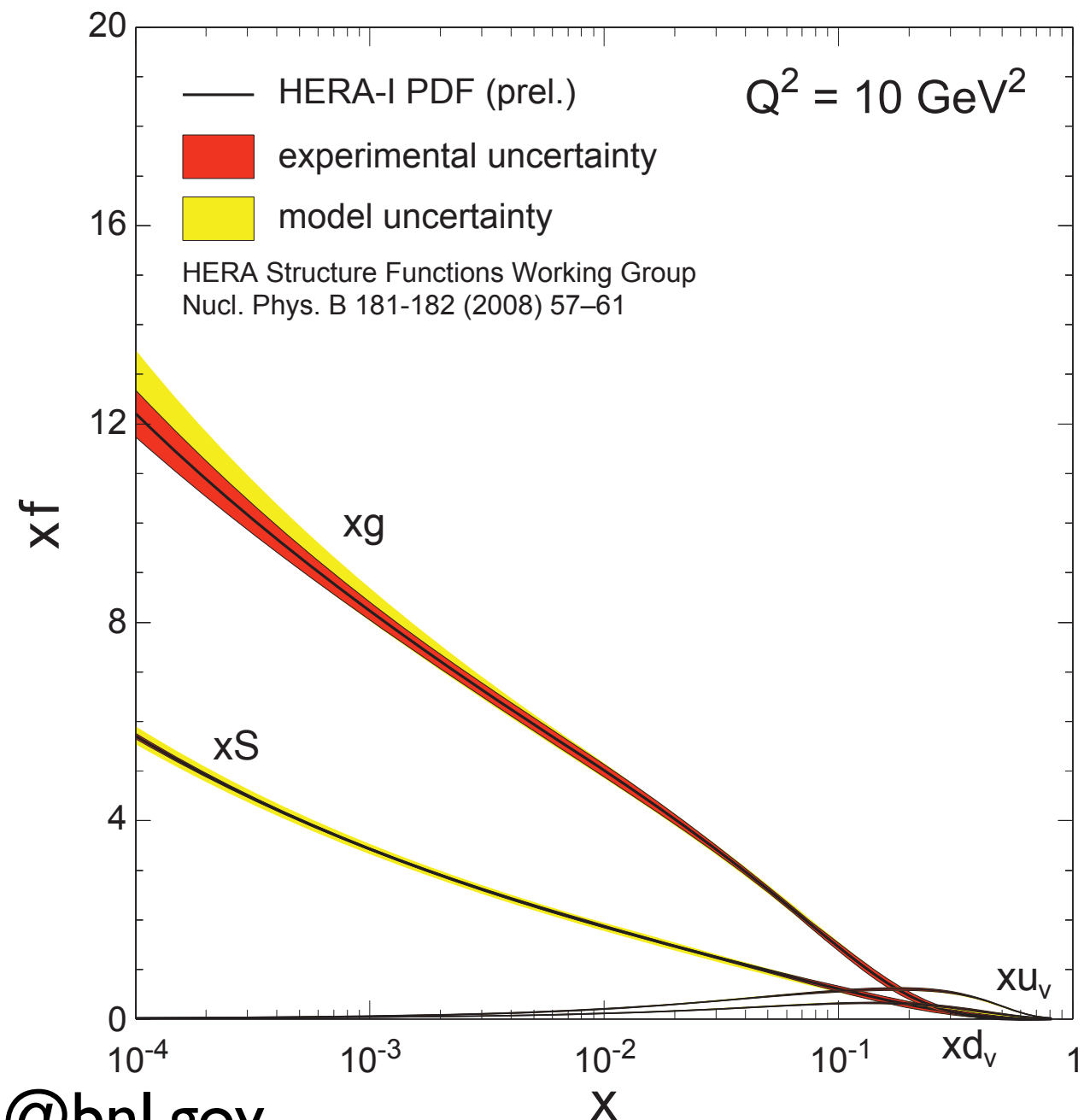


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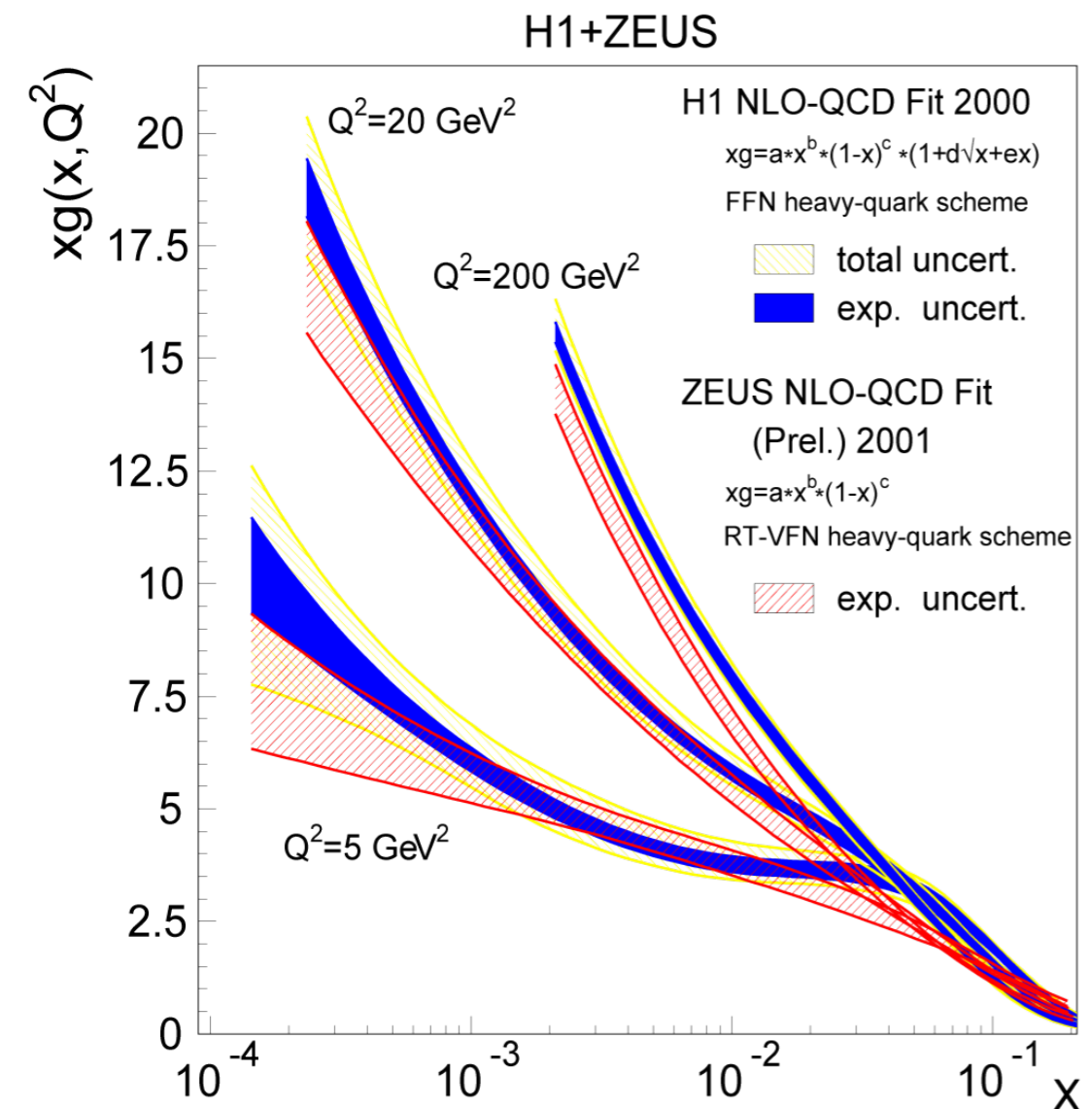
# The problem with our current understanding

- Using the Linear DGLAP evolution model:

- Linear evolution has a built-in high-energy “catastrophe”
- $xG$  has rapid rise with decreasing  $x$  (and increasing  $Q^2$ )  $\Rightarrow$  violation of Froissart unitarity bound

$$\sigma_{tot} = \frac{\pi}{m_{\pi}^2} (\ln s)^2$$

- Must have saturation to tame the growth



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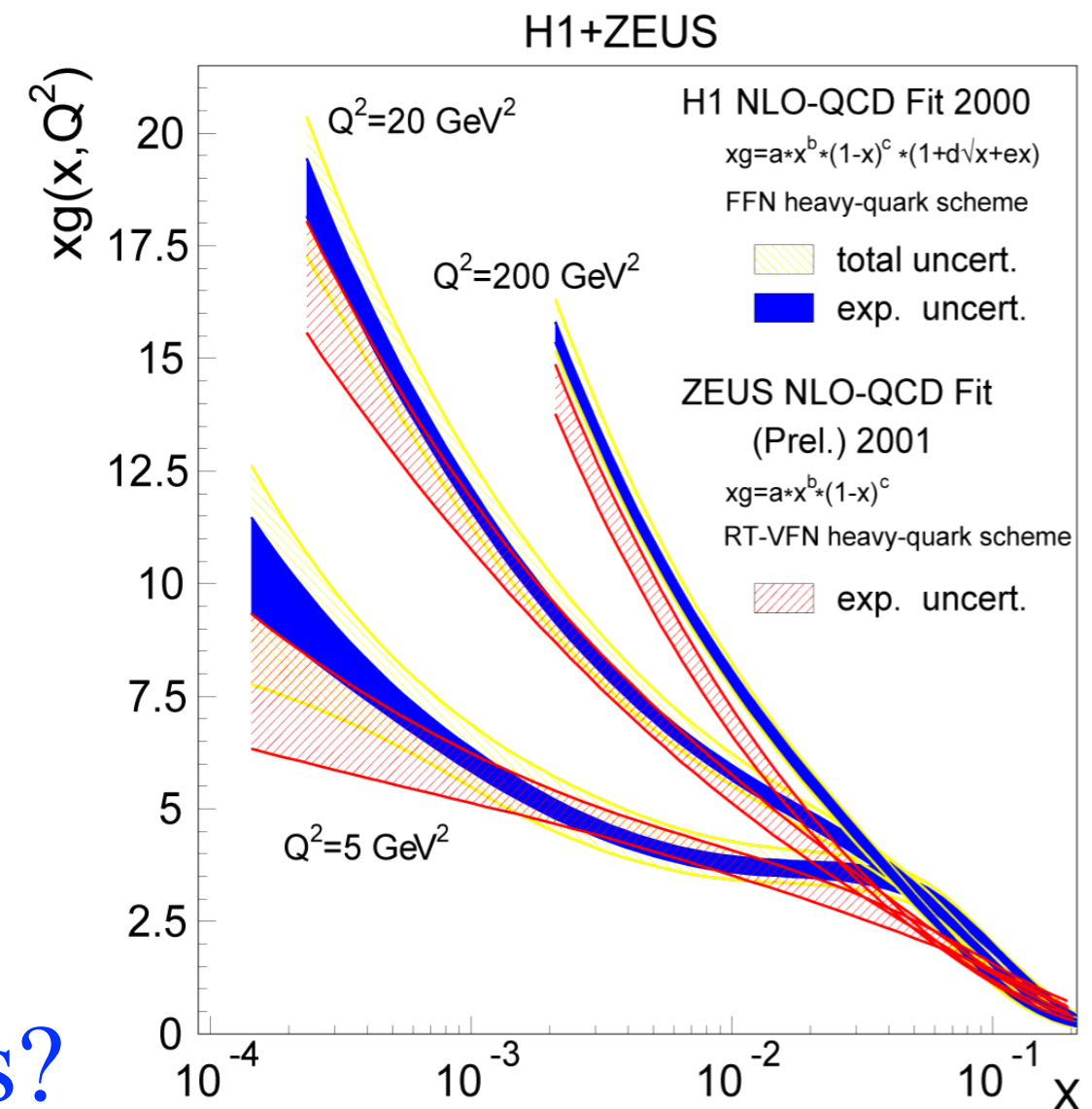
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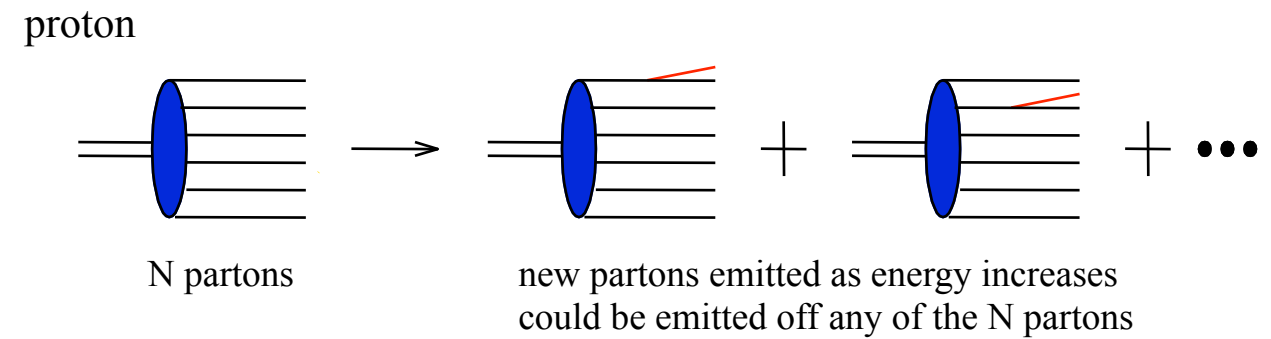
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What's the underlying dynamics?



# Non-linear QCD - Saturation

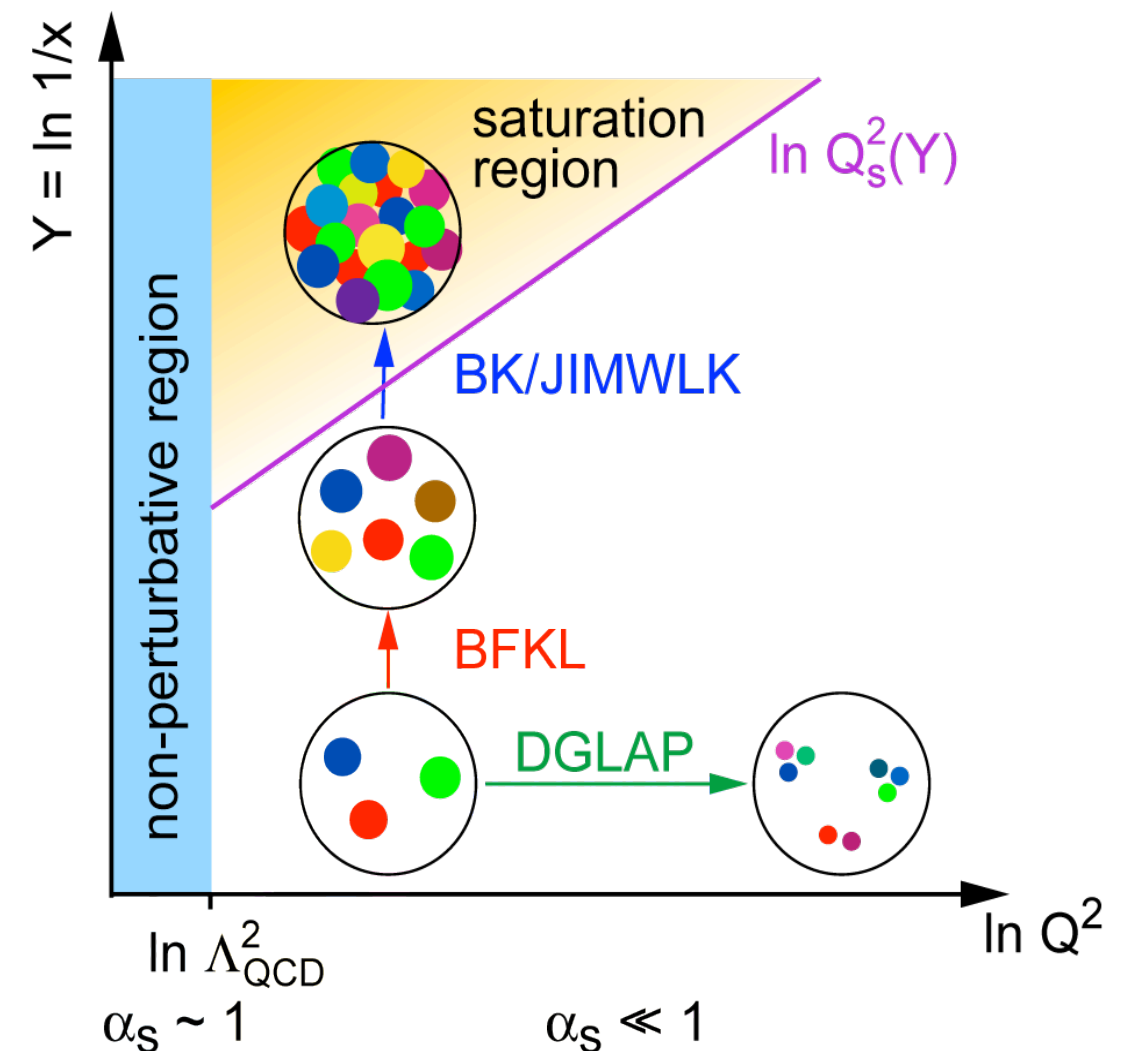
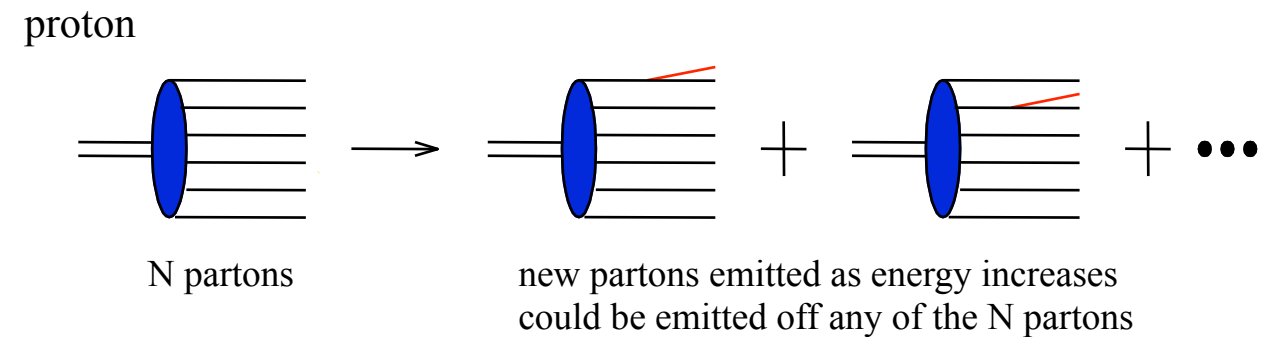


# Non-linear QCD - Saturation

- **BFKL**: evolution in  $x$

→ linear

- ▶ explosion in colour field at low- $x$





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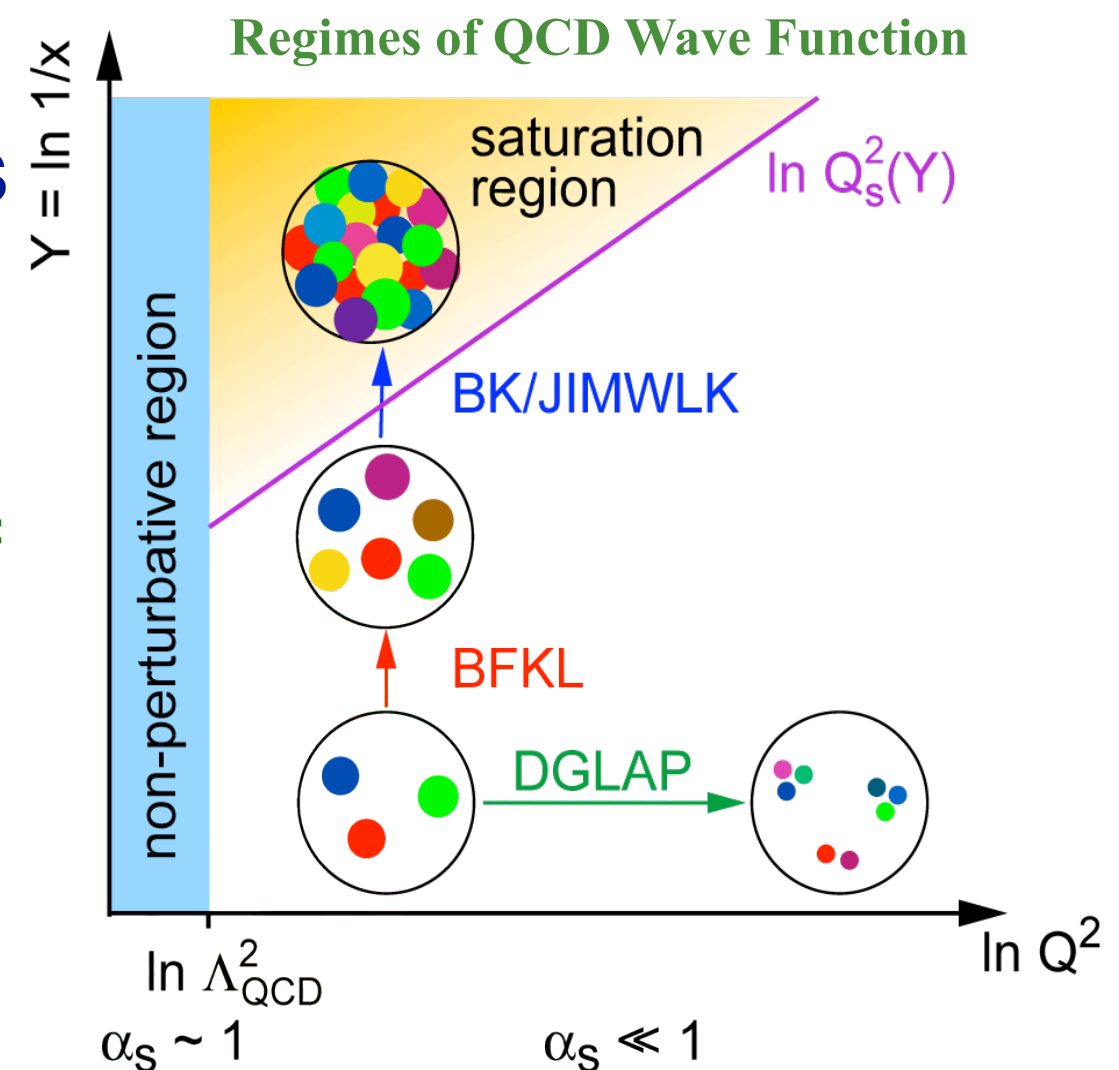
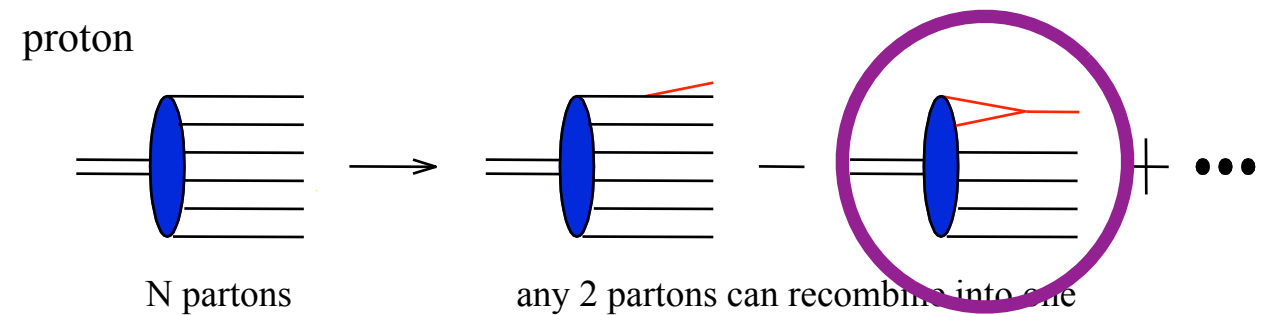
- explosion in colour field at low- $x$

- Non-linear **BK/JIMWLK** equations

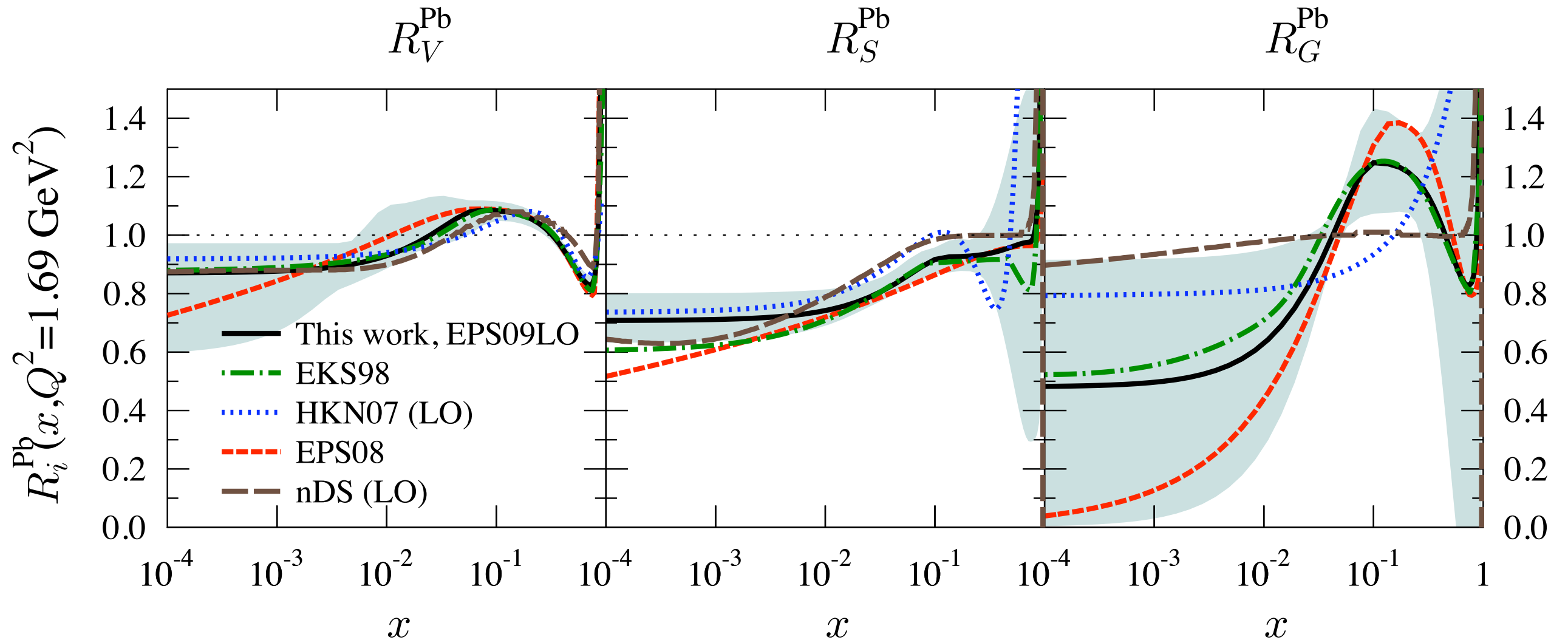
➔ non-linearity  $\Rightarrow$  saturation

- Allows for the recombination of gluons in a dense gluonic medium

➔ characterised by the saturation scale,  $Q_s(x,A)$

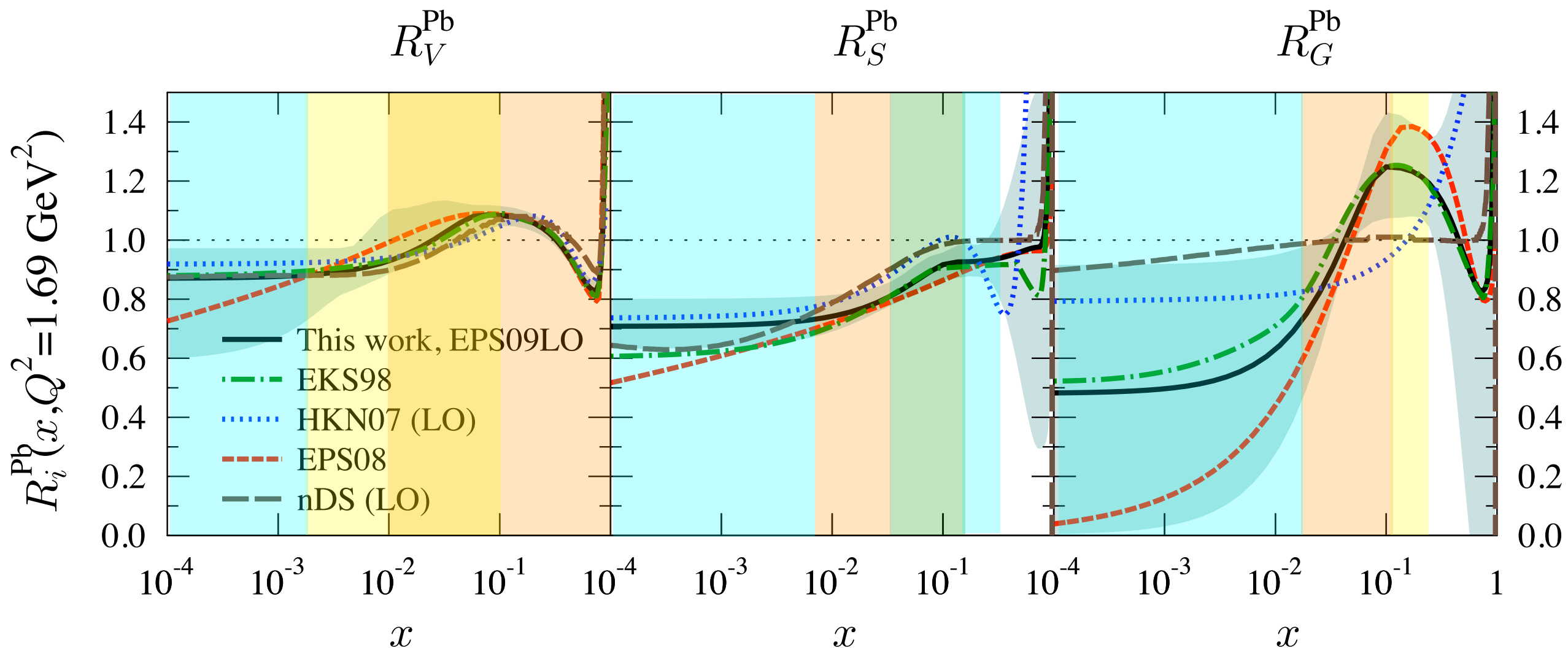


# What about gluons in nuclei?



The distribution of valence and sea quarks  
are relatively well known in nuclei -  
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# What about gluons in nuclei?



Constrained by DIS

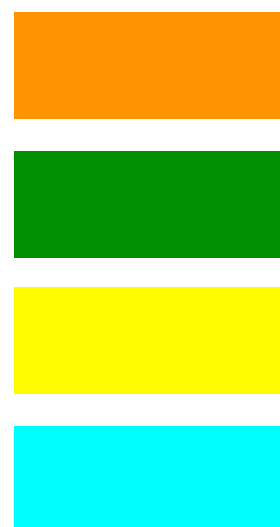
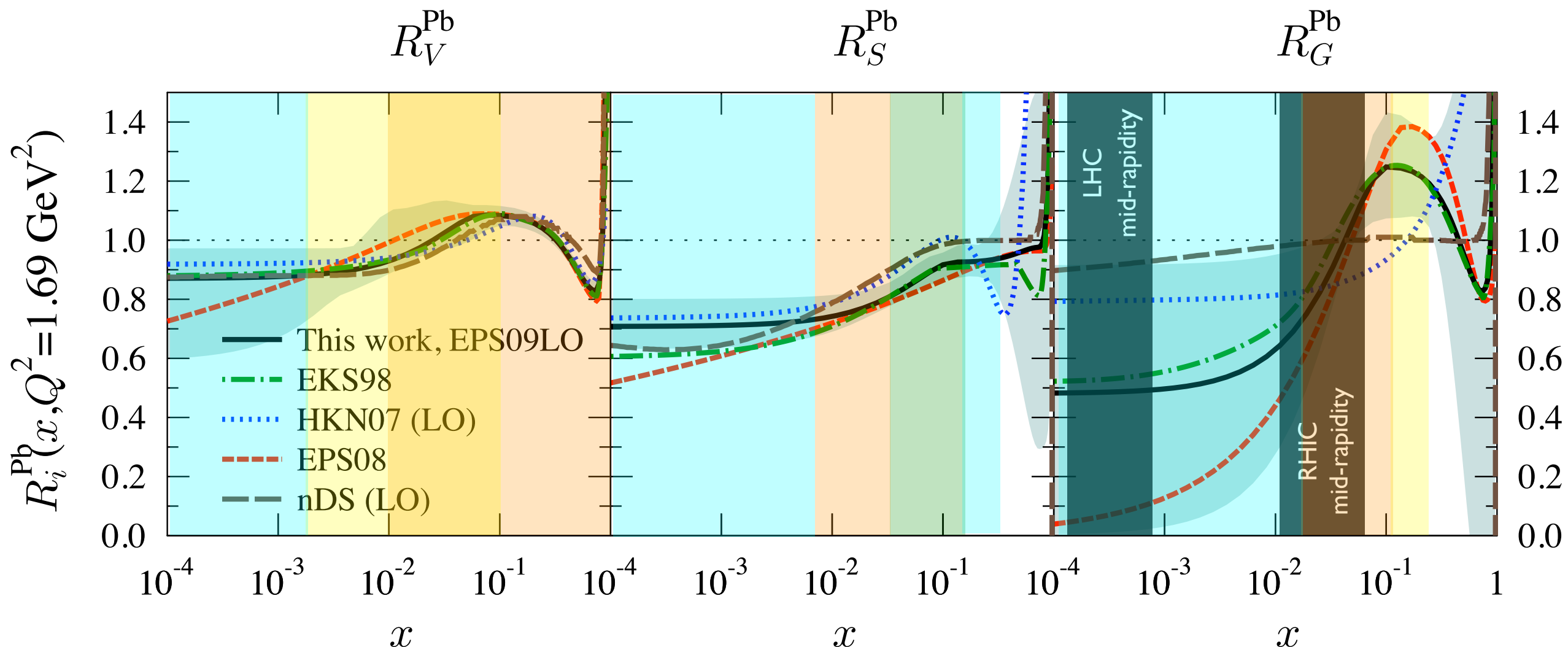
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Assumptions

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 Assumptions

The distribution of valence and sea quarks are relatively well known in nuclei - theories agree well

Large discrepancies exist in the gluon distributions from models for mid-rapidity LHC and forward RHIC rapidities !!

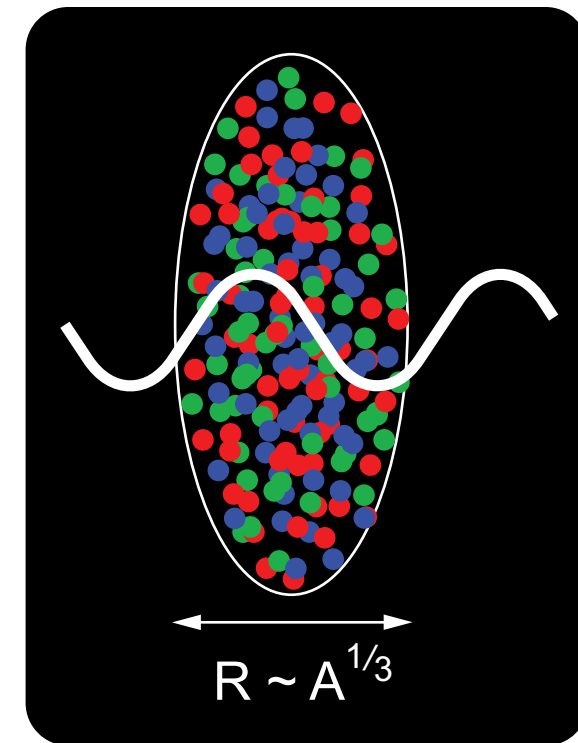
# The Nuclear “Oomph Factor”

- Enhancing Saturation effects:

➡ Probes interact over distances  $L \sim (2m_n\lambda)^{-1}$

For probes where  $L > 2R_A$  ( $\sim A^{1/3}$ ) cannot distinguish between nucleons in front or back of the nucleus.

▶ Probe interacts coherently with all nucleons.





# The Nuclear “Oomph Factor”

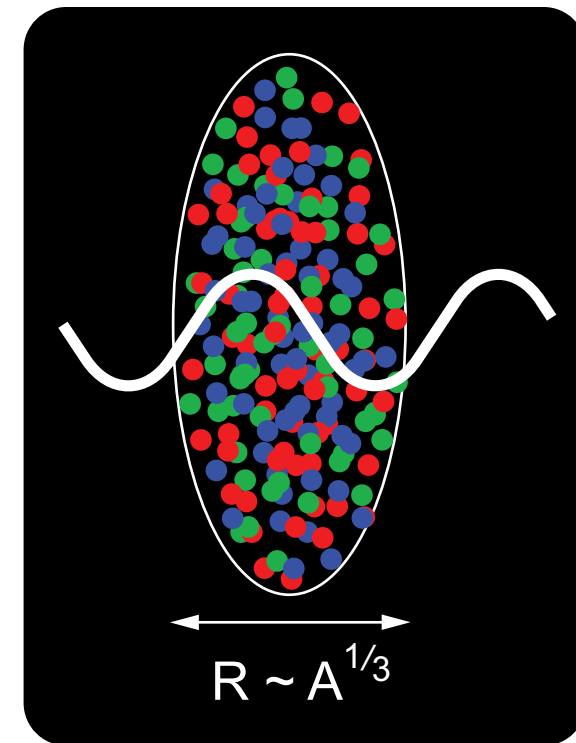
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Simple geometric considerations lead to:



$$Q_s^2 \propto \frac{\alpha_s x G(x, Q_s^2)}{\pi R_A^2} \quad \text{HERA: } xG \propto \frac{1}{x^{1/3}} \quad \text{A dependence: } xG_A \propto A$$

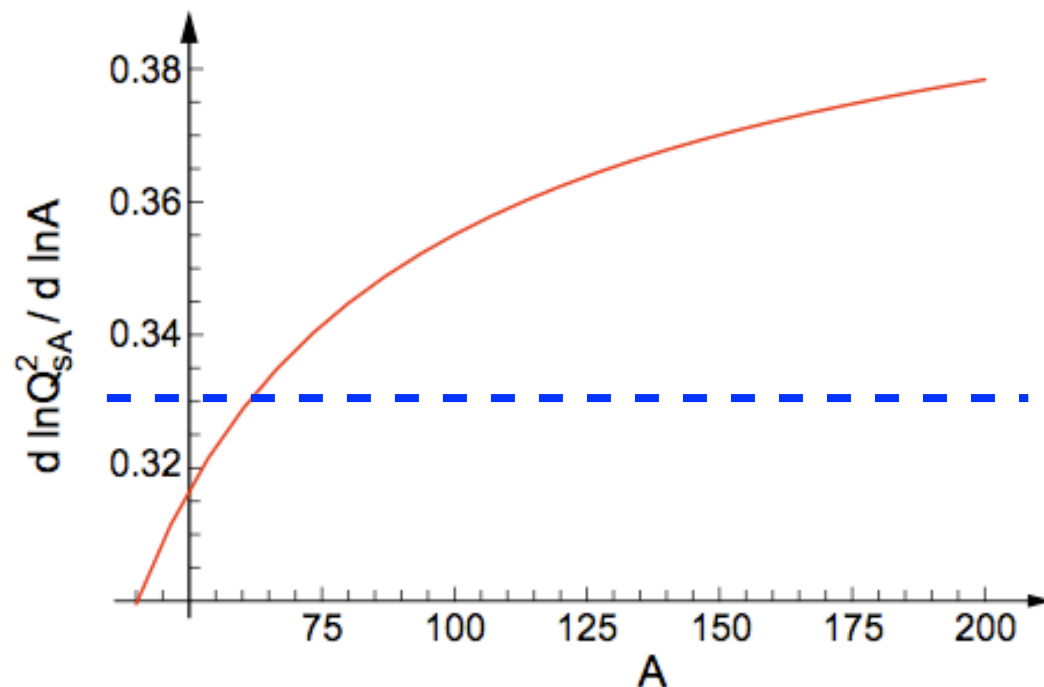
Nuclear “Oomph” Factor:  $(Q_s^A)^2 \approx c Q_0^2 \left( \frac{A}{x} \right)^{1/3}$

*Enhancement* of  $Q_s$  with  $A$ :  $\Rightarrow$  non-linear QCD regime reached at significantly lower energy in  $e+A$  than in  $e+p$

# The Nuclear “Oomph Factor”

More sophisticated analyses  
⇒ confirm (exceed) pocket  
formula for high A

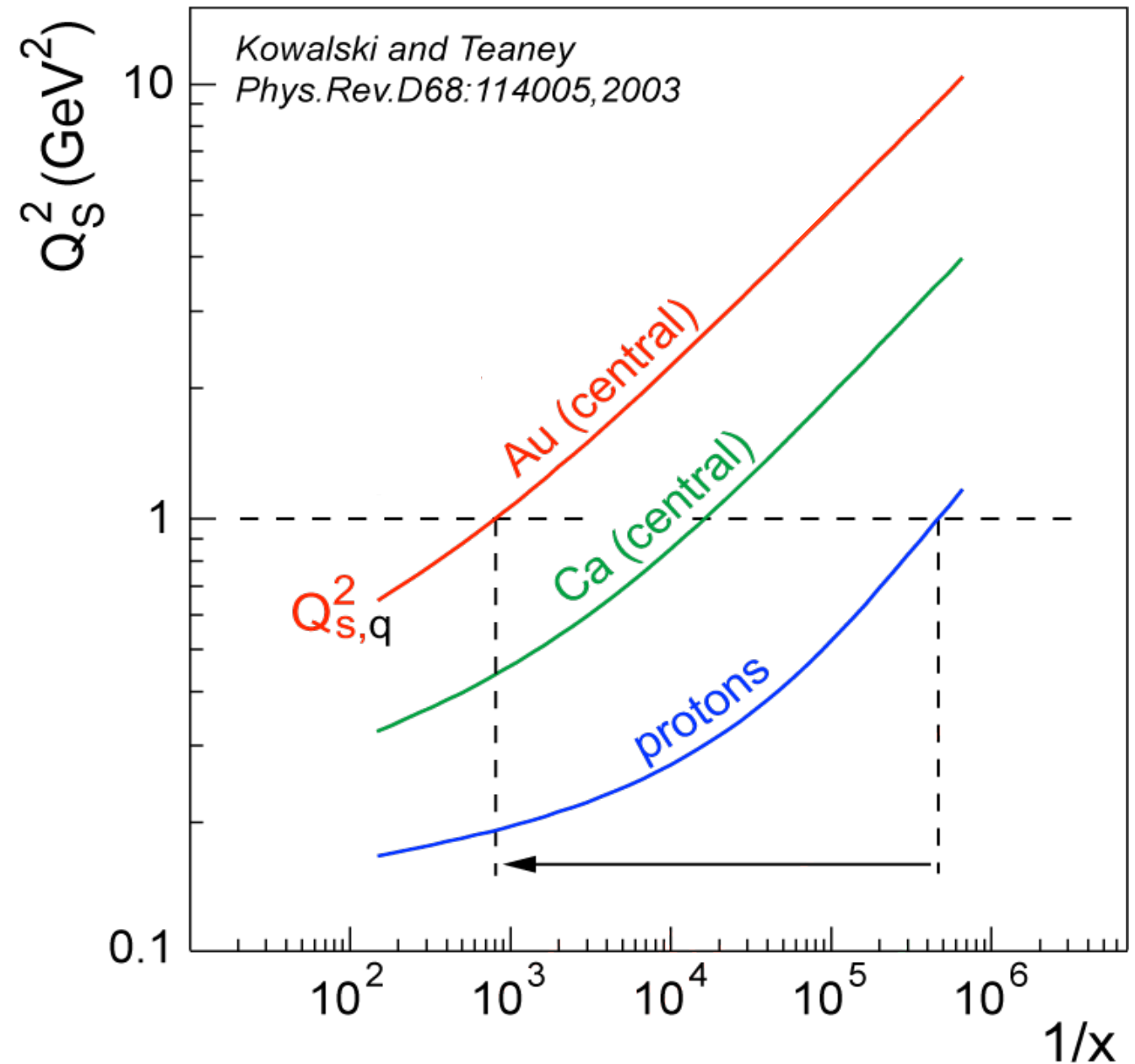
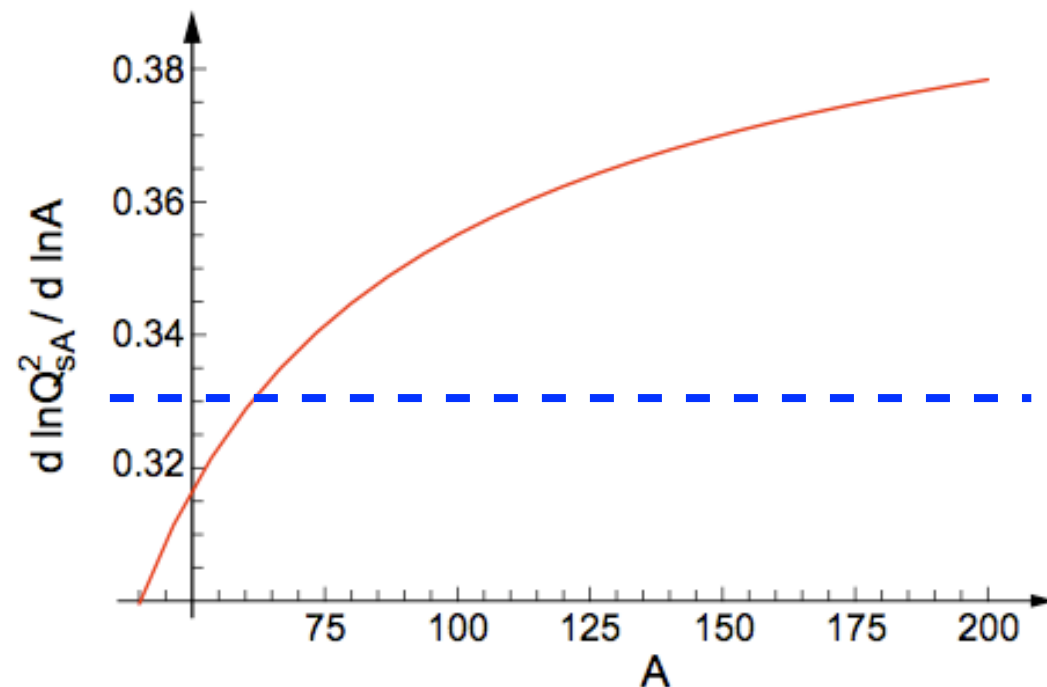
e.g. Kowalski, Lappi and Venugopalan,  
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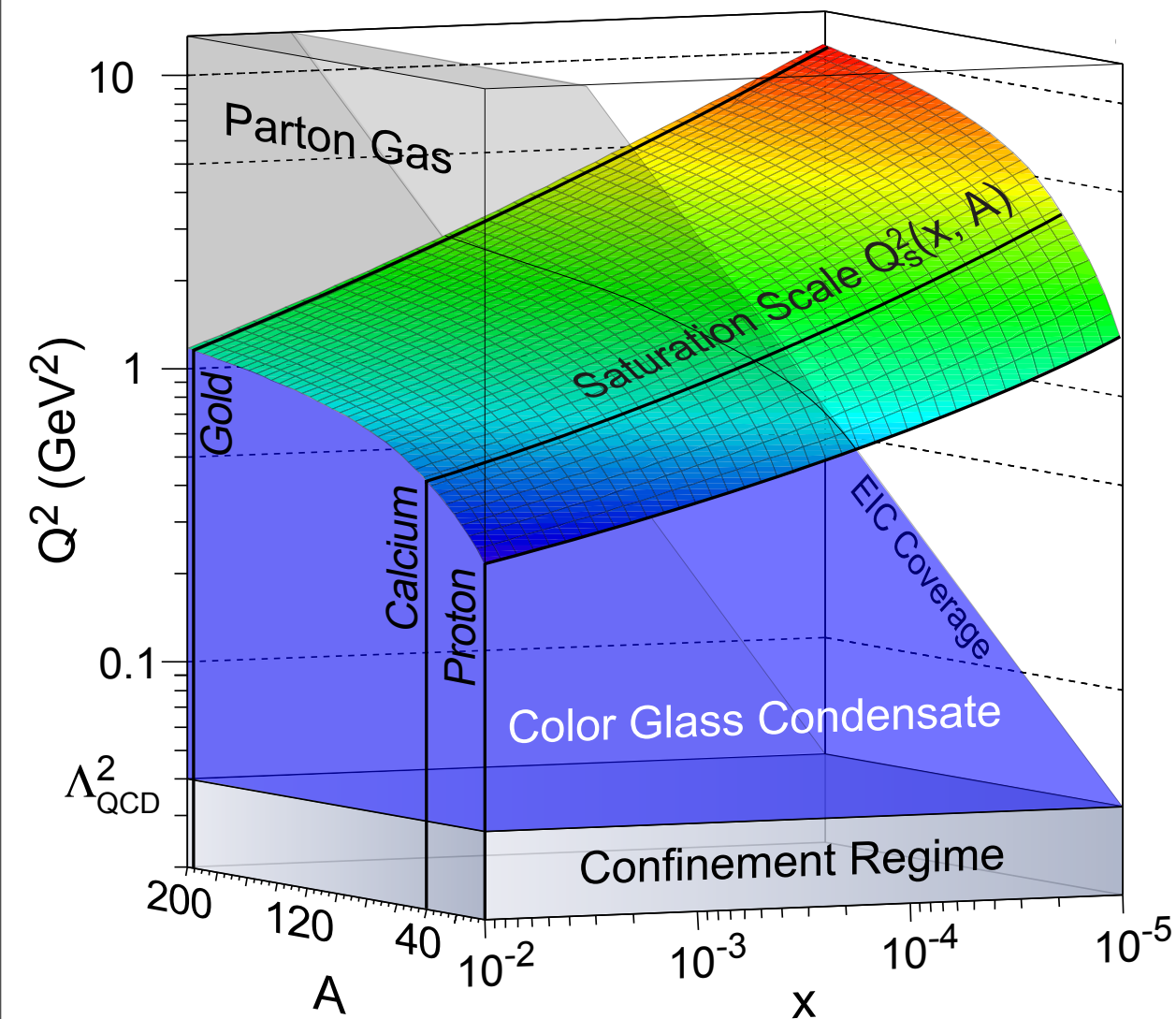
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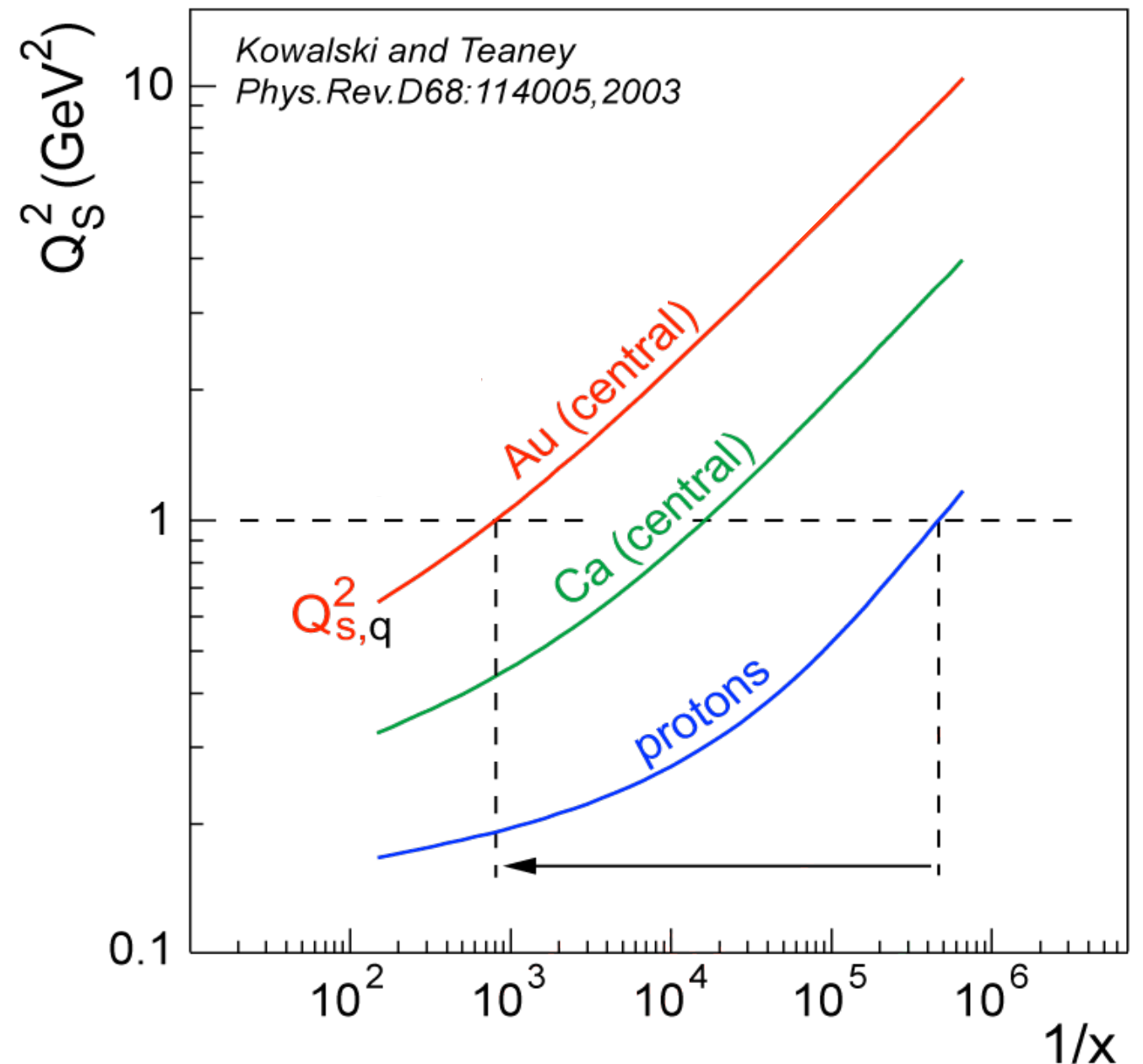


One would require an energy in e+p  
~ 10-100 x e+A to get to same  $Q_s^2$

# The Nuclear “Oomph Factor”



Plot by T. Ullrich



One would require an energy in e+p  
 $\sim 10\text{-}100 \times e+A$  to get to same  $Q_s^2$



# Do EIC energies match the requirements?

**eRHIC** = RHIC +  
Energy-Recovery Linac



**ELIC** = CEBAF +  
Hadron Ring



Both  
designs in  
2 stages

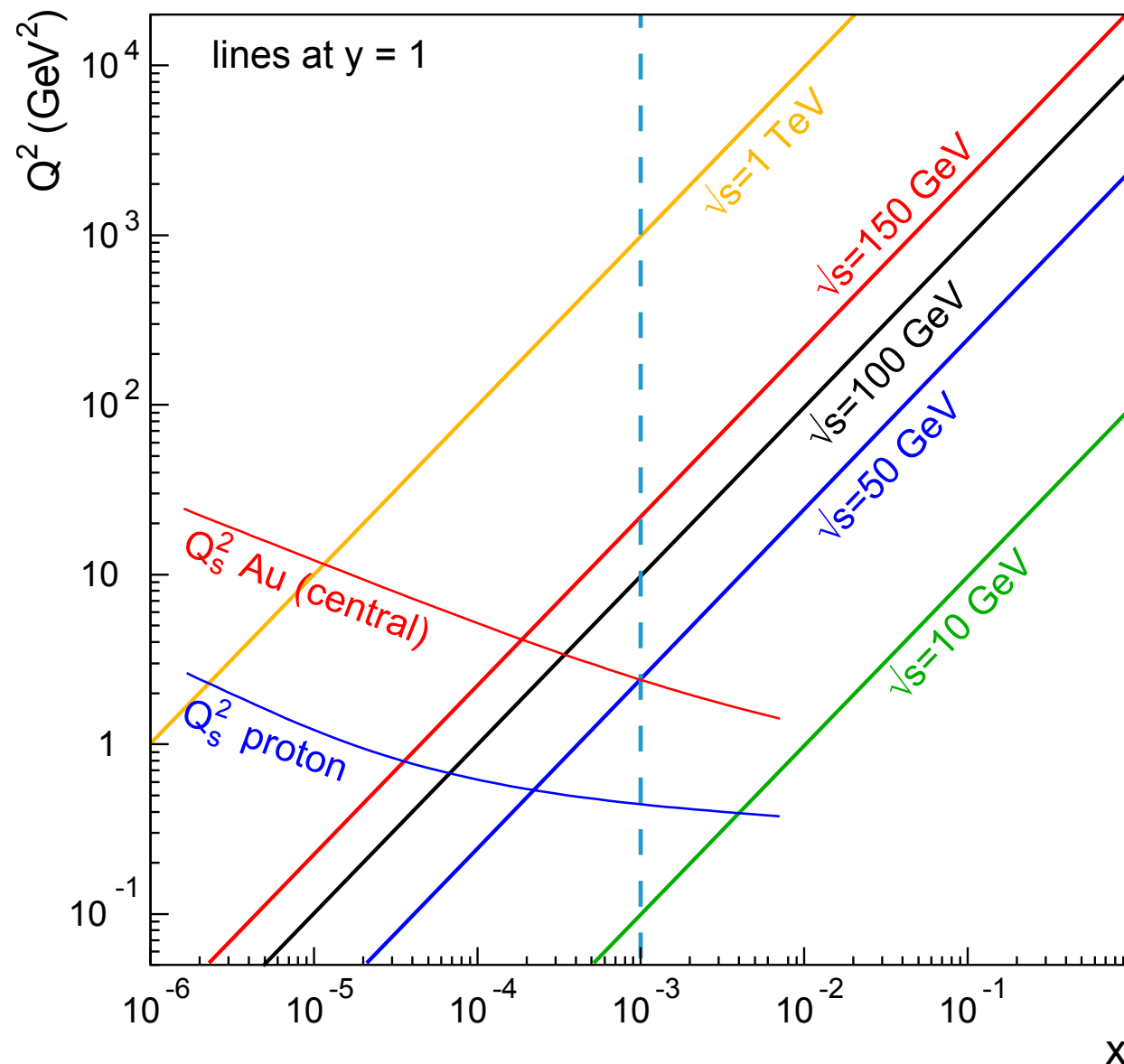
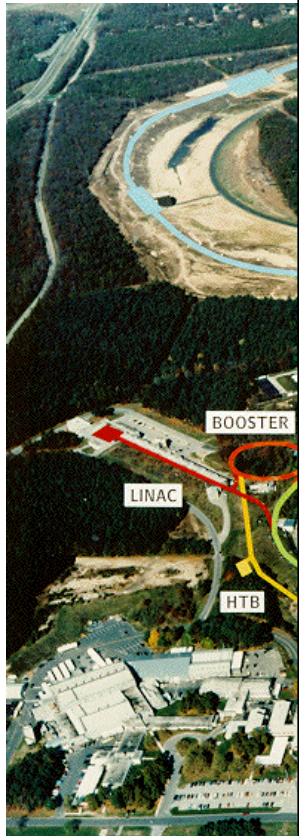
1. stage: 5+100 GeV/n e+Au ( $\sqrt{s}=45$  GeV/n)
2. stage: 30+130 GeV/n e+Au ( $\sqrt{s}=125$  GeV/n)

1. stage: 11+40 GeV/n e+Au ( $\sqrt{s}=42$  GeV/n)
2. stage: 20+100 GeV/n e+Au ( $\sqrt{s}=89$  GeV/n)



# Do EIC energies match the requirements?

eRH  
Energy



F +  
g



1. stage:

- In both cases 1<sup>st</sup> stage is ~OK but offers little  $Q^2$  lever arm

2. sta

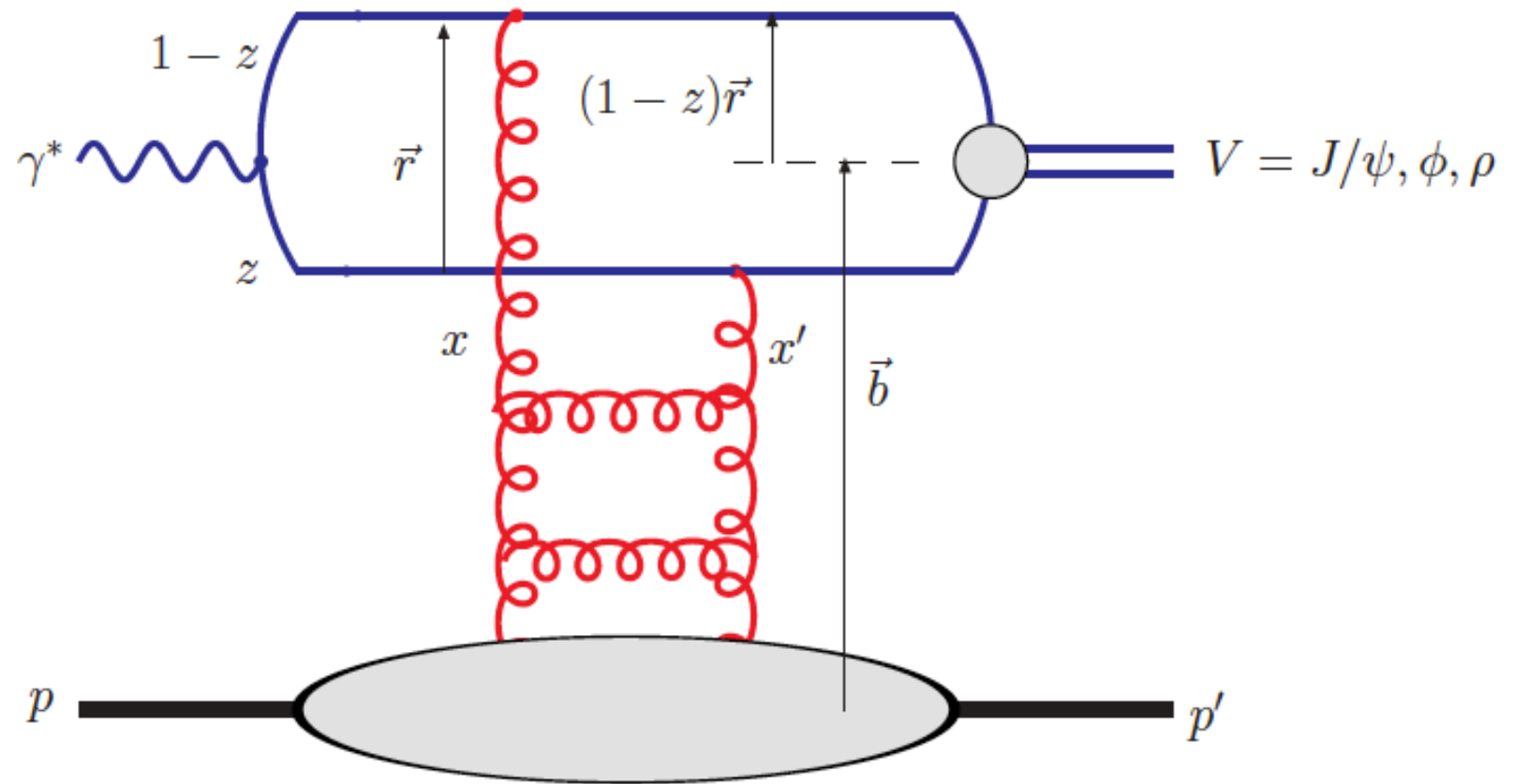
- 2<sup>nd</sup> stage will match requirements fully

/n e+Au  
(/n)  
(/n e+Au  
(/n)



# Getting a “Feel” for Non-Linear QCD

## Dipole Model:



$$\frac{d\sigma_{q\bar{q}}}{d^2b} = 2\mathcal{N}(x, r, b)$$

$$\mathcal{N}(x, r, b) = 2 \left[ 1 - \exp \left( -r^2 \frac{\pi^2}{2N_c} \alpha_s(\mu^2) x G(x, \mu^2) T(b) \right) \right]$$

$\mathcal{N}$  = Dipole Scattering Amplitude

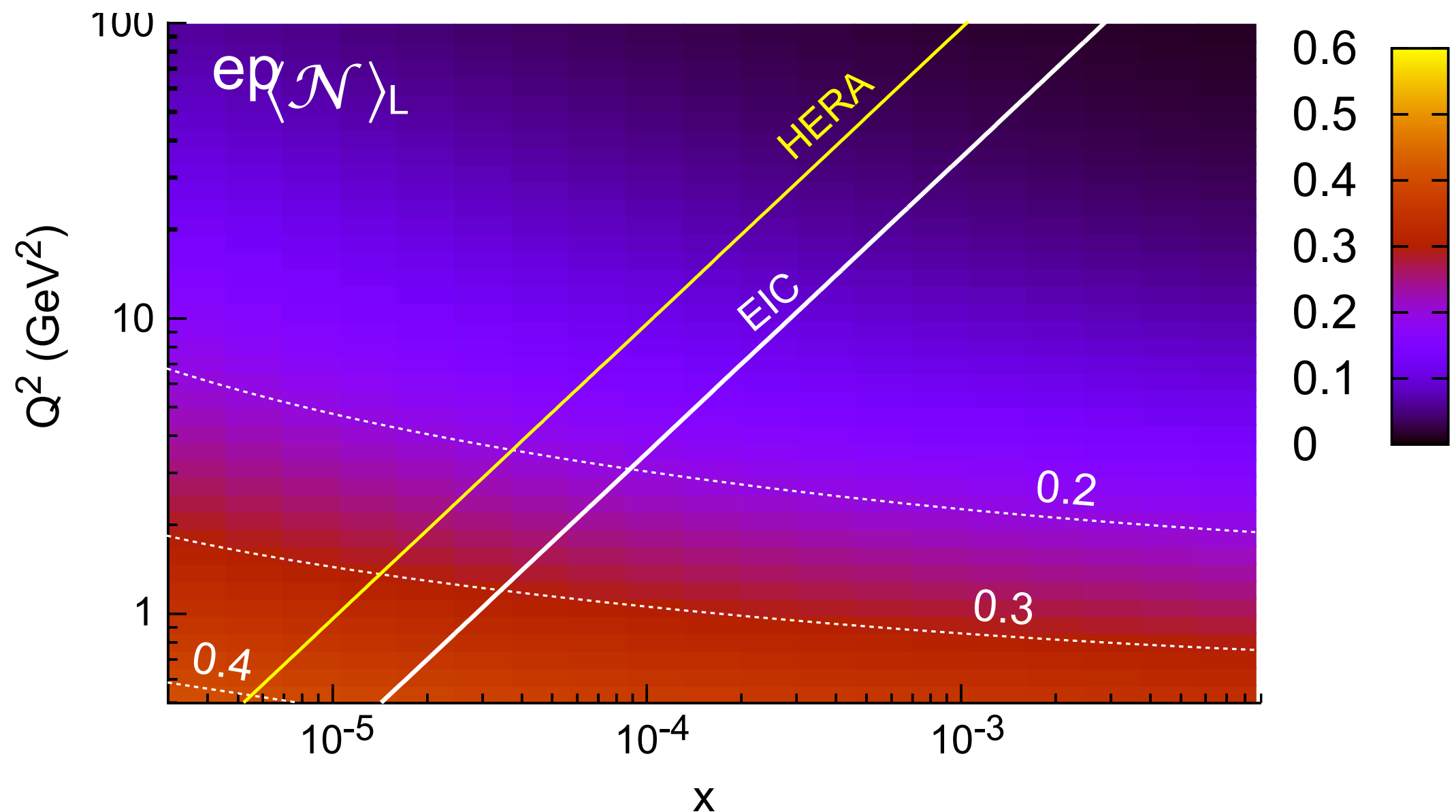
0 dilute system, linear QCD

1 saturated, non-linear regime

# Getting a “Feel” for Non-Linear QCD

To assess typical values of  $\mathcal{N}$  calculate average:

$$\langle \mathcal{N} \rangle_{2,L} = \frac{\int d^2b d^2r dz [\psi^* \psi]_{2,L} \mathcal{N}^2}{\int d^2b d^2r dz [\psi^* \psi]_{2,L} \mathcal{N}}$$

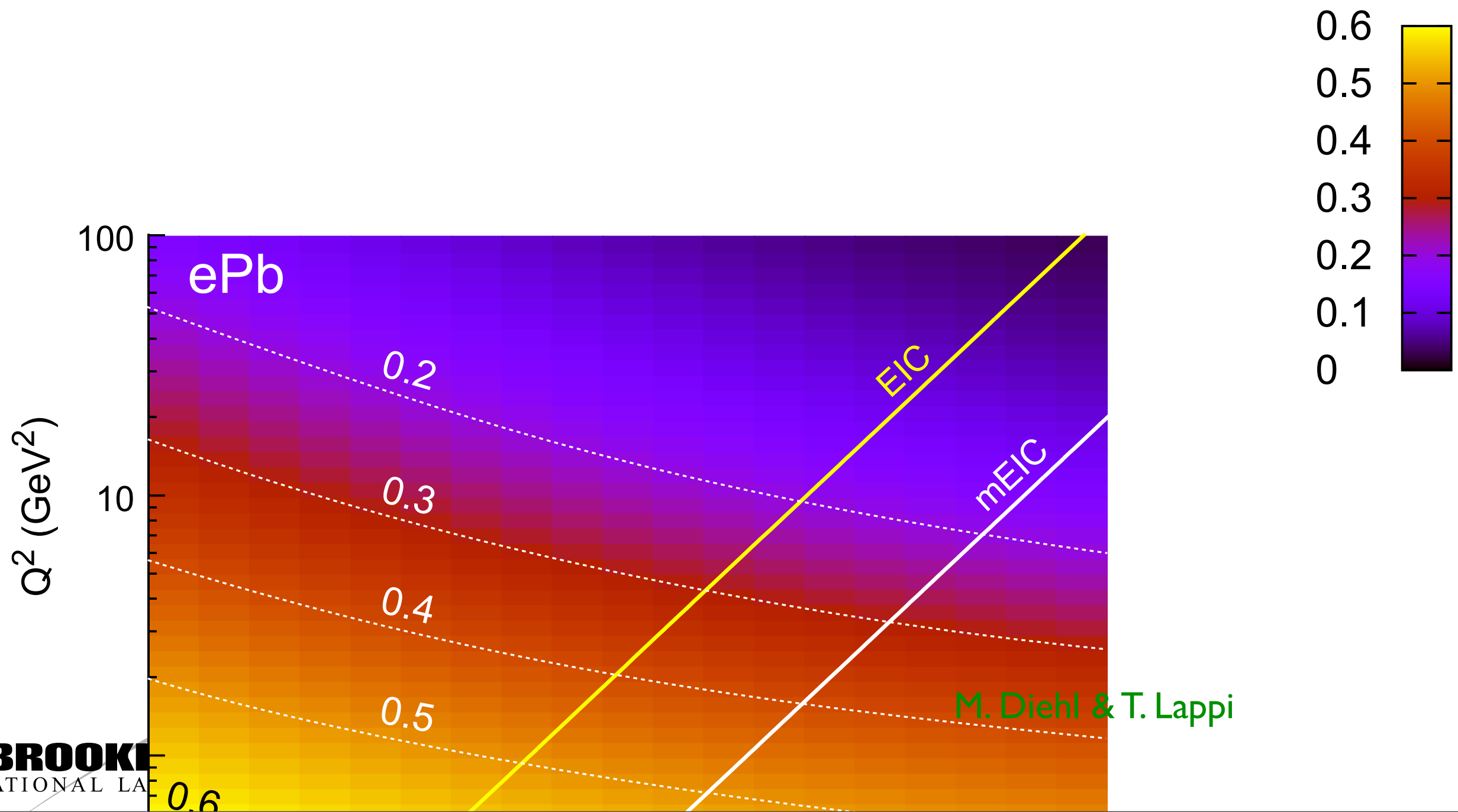


M. Diehl & T. Lappi

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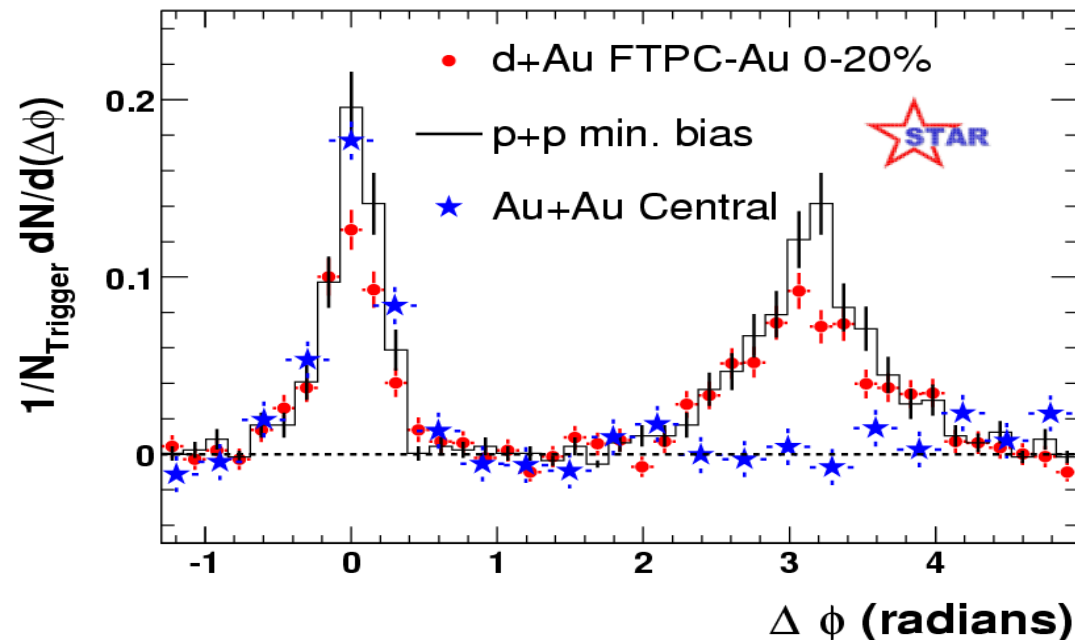
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# di-hadron angular correlations in d+A

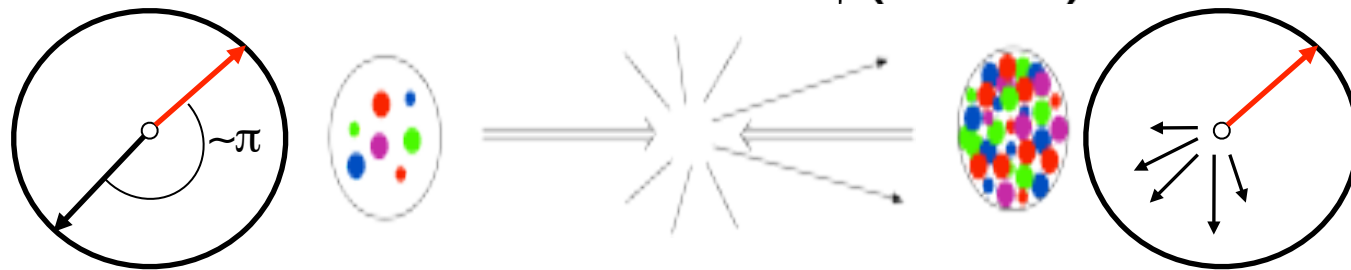
comparisons between d+Au  $\rightarrow h_1 h_2 X$  (or p+Au  $\rightarrow h_1 h_2 X$ ) and p+p  $\rightarrow h_1 h_2 X$



- At  $y=0$ , suppression of away-side jet is observed in A+A collisions
- No suppression in p+p or d+A

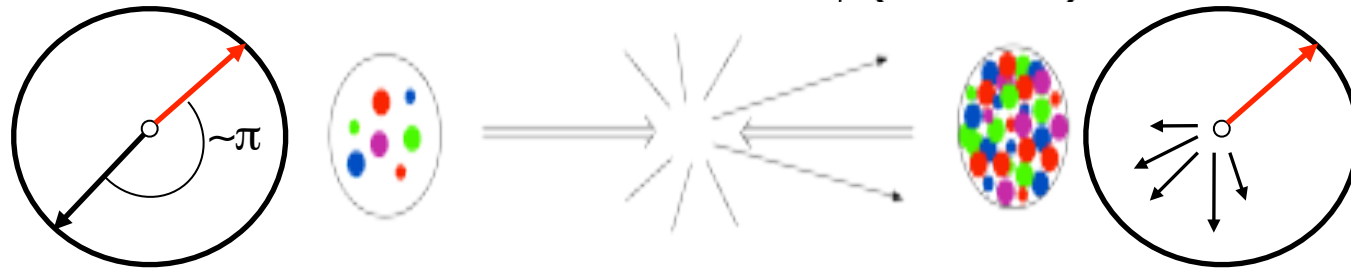
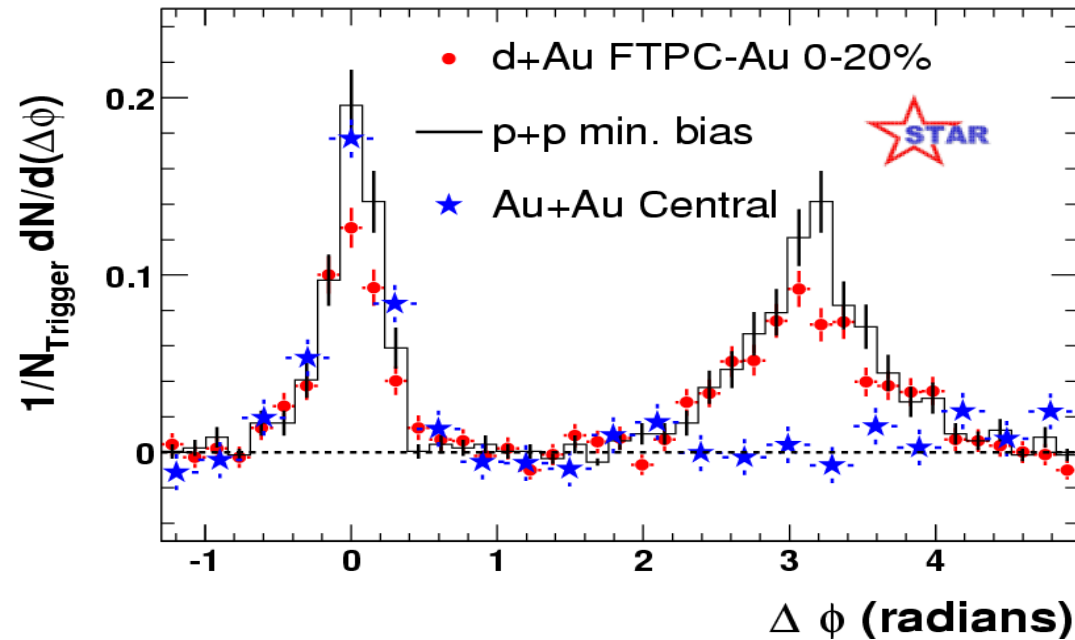
→  $x \sim 10^{-2}$

$$x_A = \frac{k_1 e^{-y_1} + k_2 e^{-y_2}}{\sqrt{s}} \ll 1$$



# di-hadron angular correlations in d+A

comparisons between d+Au  $\rightarrow h_1 h_2 X$  (or p+Au  $\rightarrow h_1 h_2 X$ ) and p+p  $\rightarrow h_1 h_2 X$



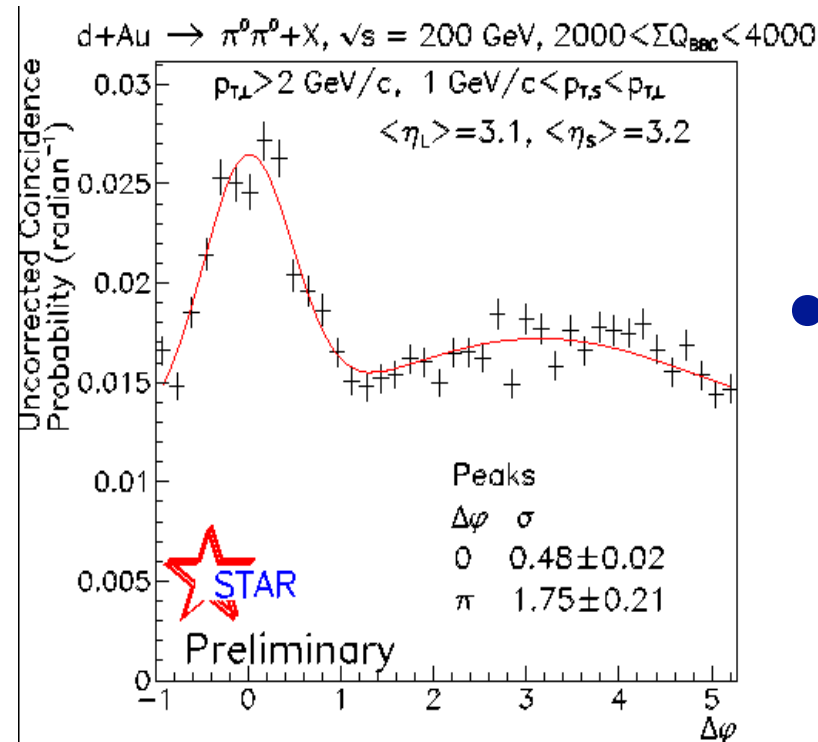
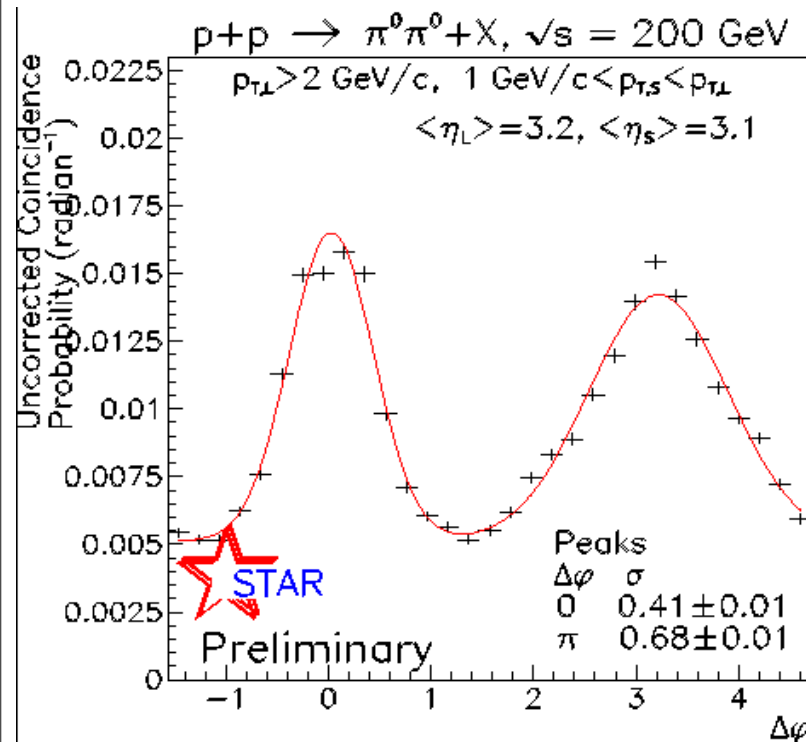
- At  $y=0$ , suppression of away-side jet is observed in A+A collisions
- No suppression in p+p or d+A

→  $x \sim 10^{-2}$

$$x_A = \frac{k_1 e^{-y_1} + k_2 e^{-y_2}}{\sqrt{s}} \ll 1$$

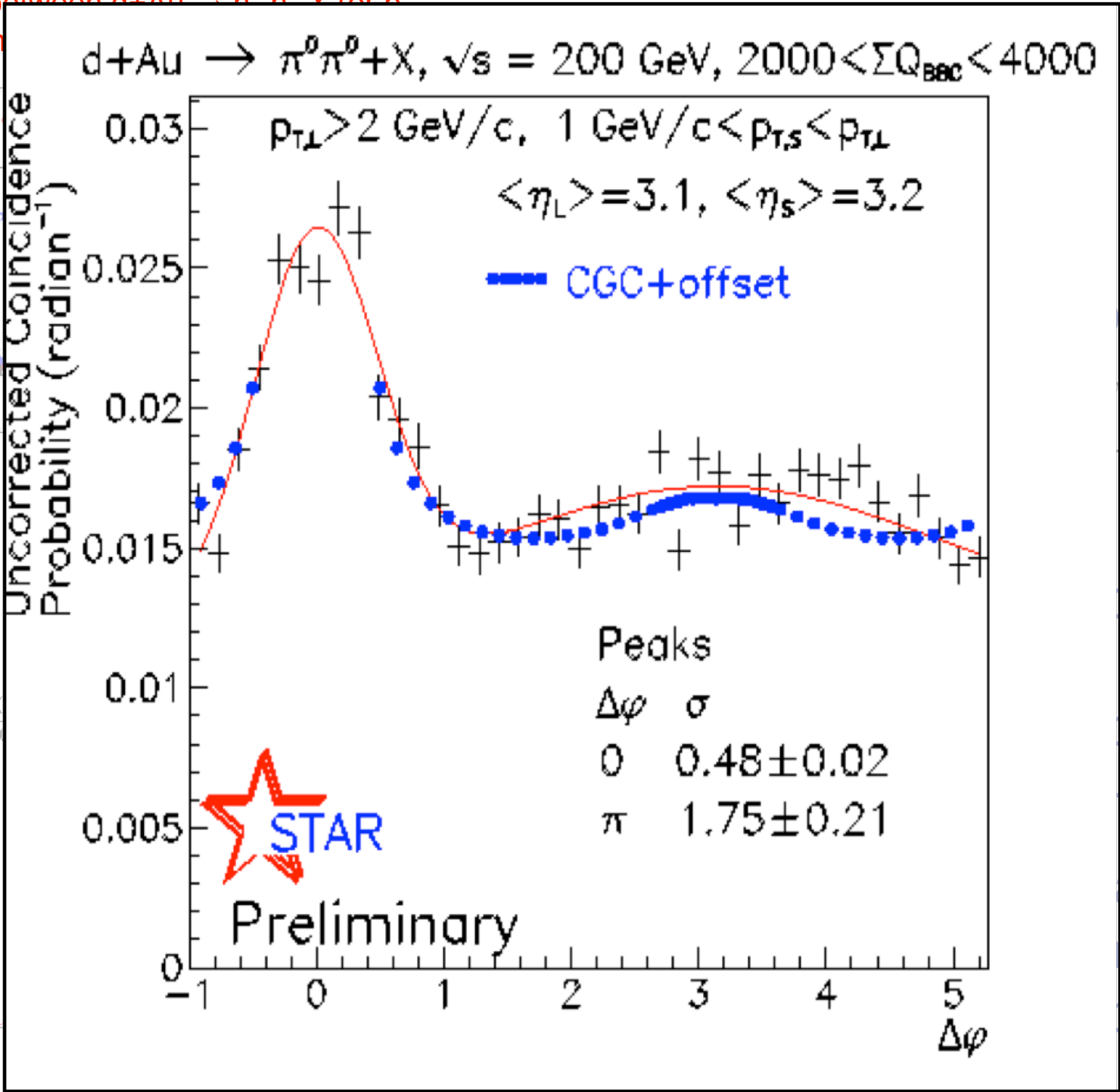
- However, at forward rapidities ( $y \sim 3.1$ ), an away-side suppression is observed in d+Au
- Away-side peak also much wider in d+Au compared to p+p

→  $x \sim 10^{-3}$



# di-hadron angular correlations in d+Au

comparisons between d+Au  $\rightarrow h_1 h_2 + X$  (or p+p  $\rightarrow h_1 h_2 + X$ )



of away-  
in A+A

+p or d+A

$$\frac{1 + k_2 e^{-y_2}}{\sqrt{s}} \ll 1$$

orward  
(3.1), an  
oppression is  
+Au

ak also  
d+Au

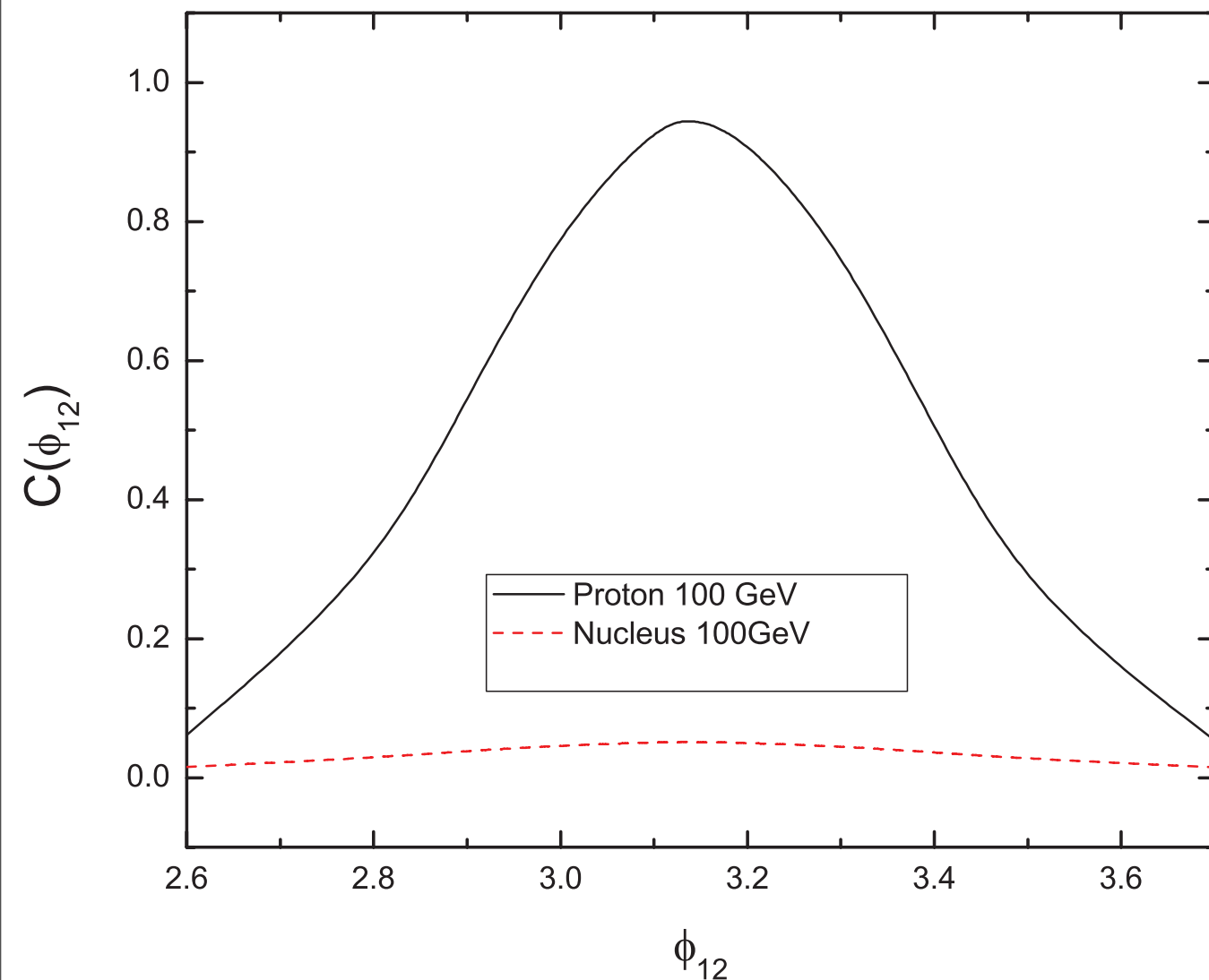
p+p

$\Rightarrow X \sim 10^{-3}$



# di-hadron correlations in e+A

Never been measured - we expect to see the same effect in e+A as in d+A



Dominguez, Xiao and Yuan (2010)

- At small-x, multi-gluon distributions are as important as single-gluon distributions and they contribute to di-hadron correlations

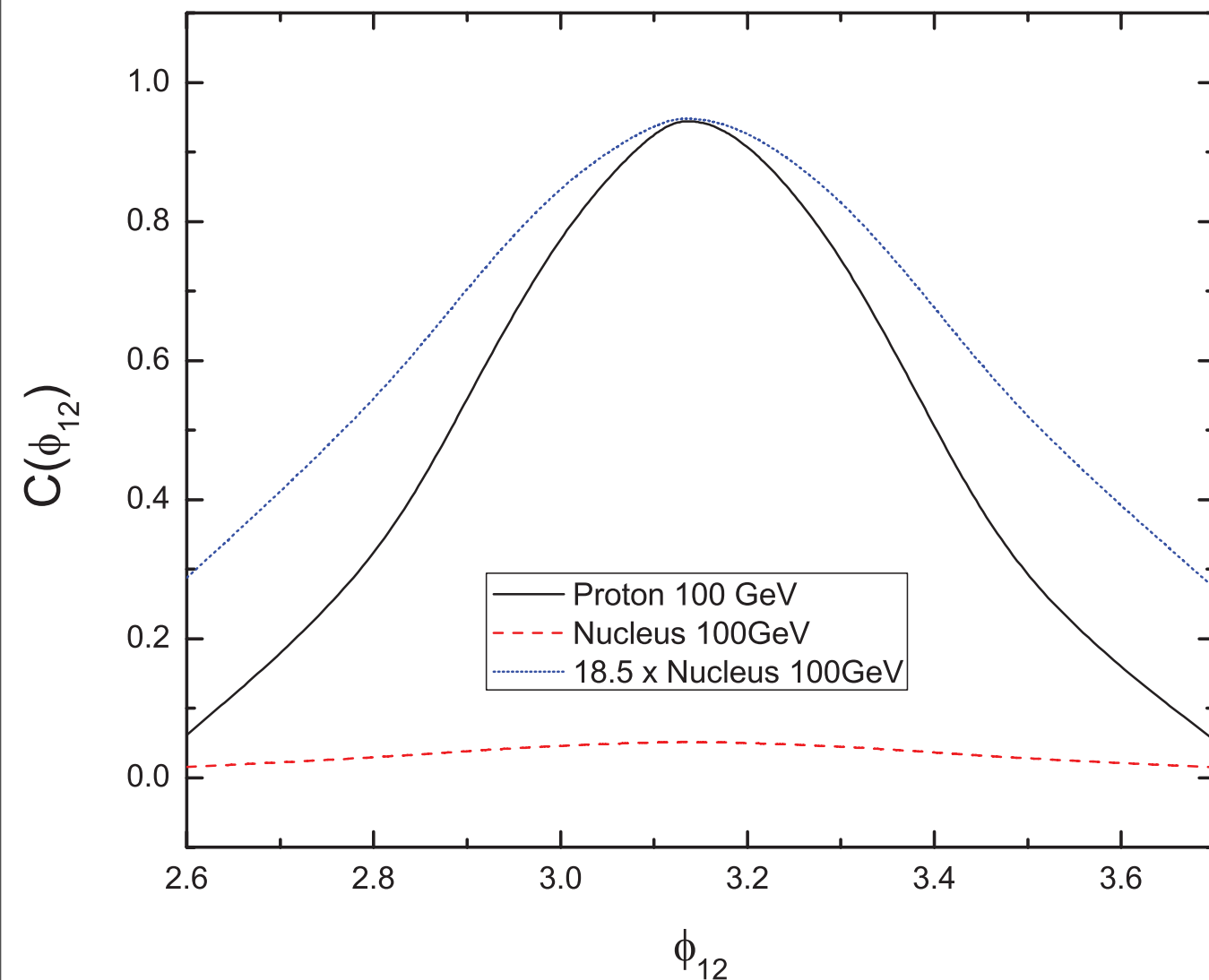
➔ The non-linear evolution of multi-gluon distributions is different from that of single-gluon distributions and it is **equally important** that we understand it

- The d+Au RHIC data is therefore subject to many uncertainties

➔ these correlations in e+A can help to constrain them better

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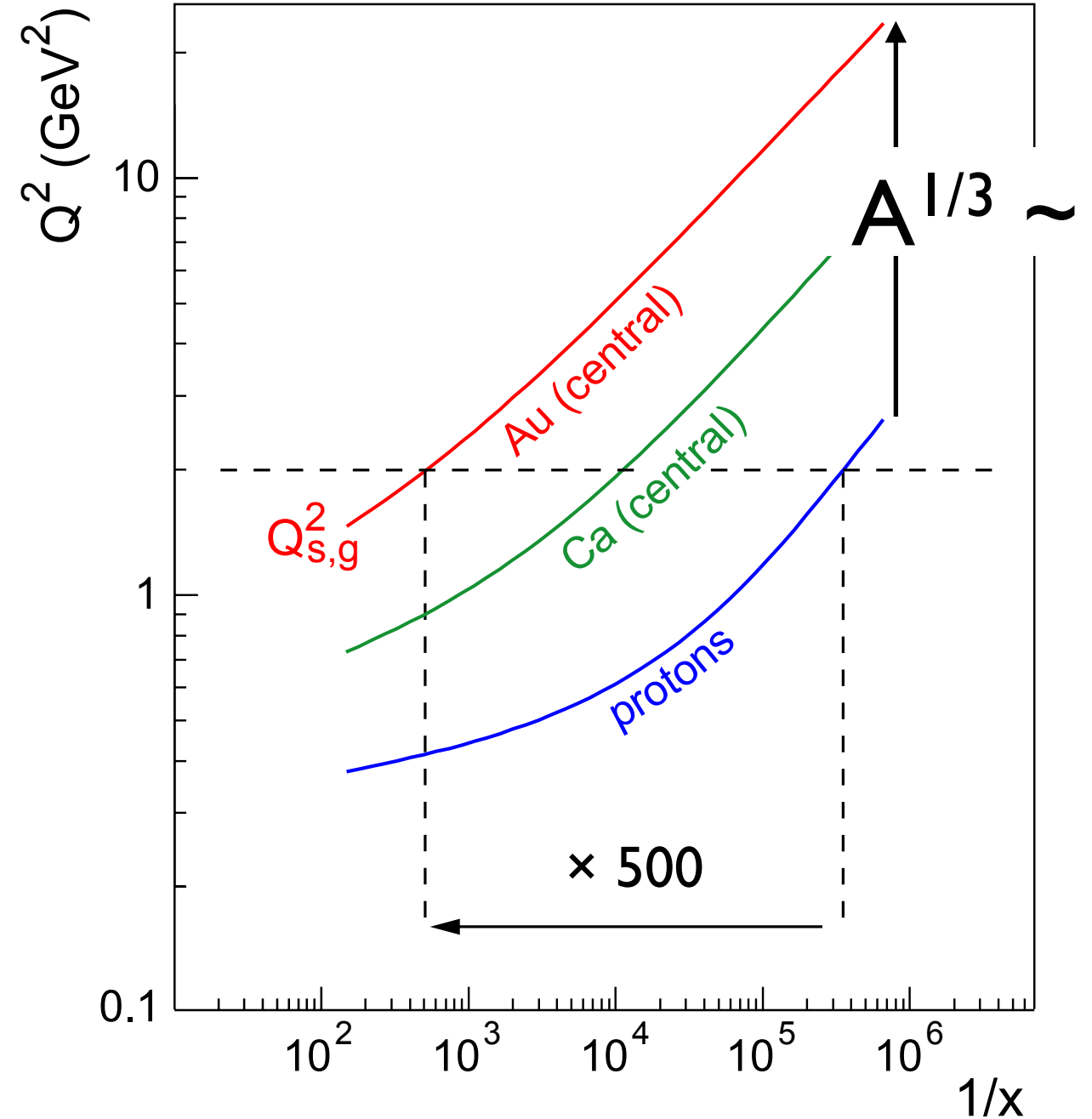
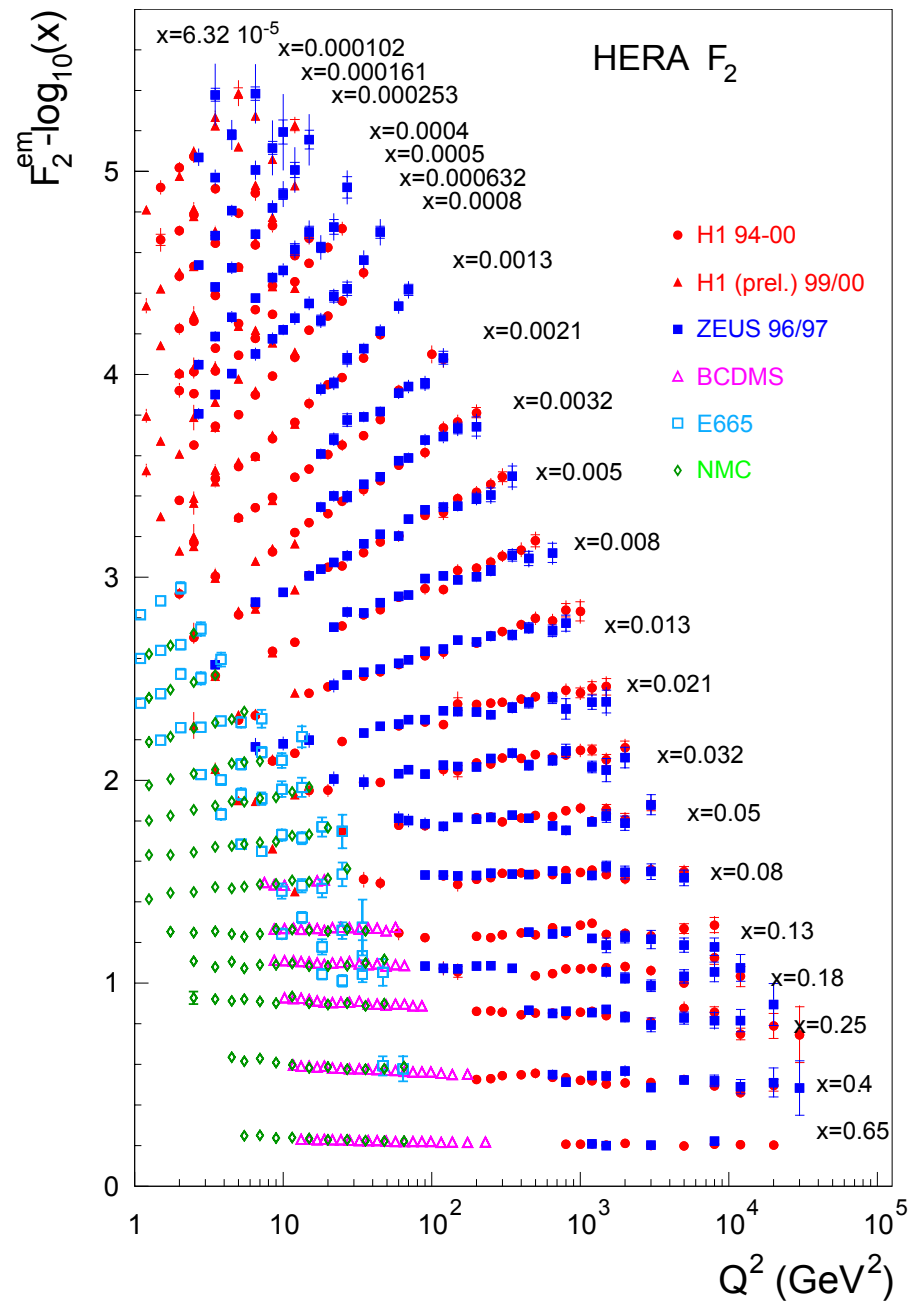
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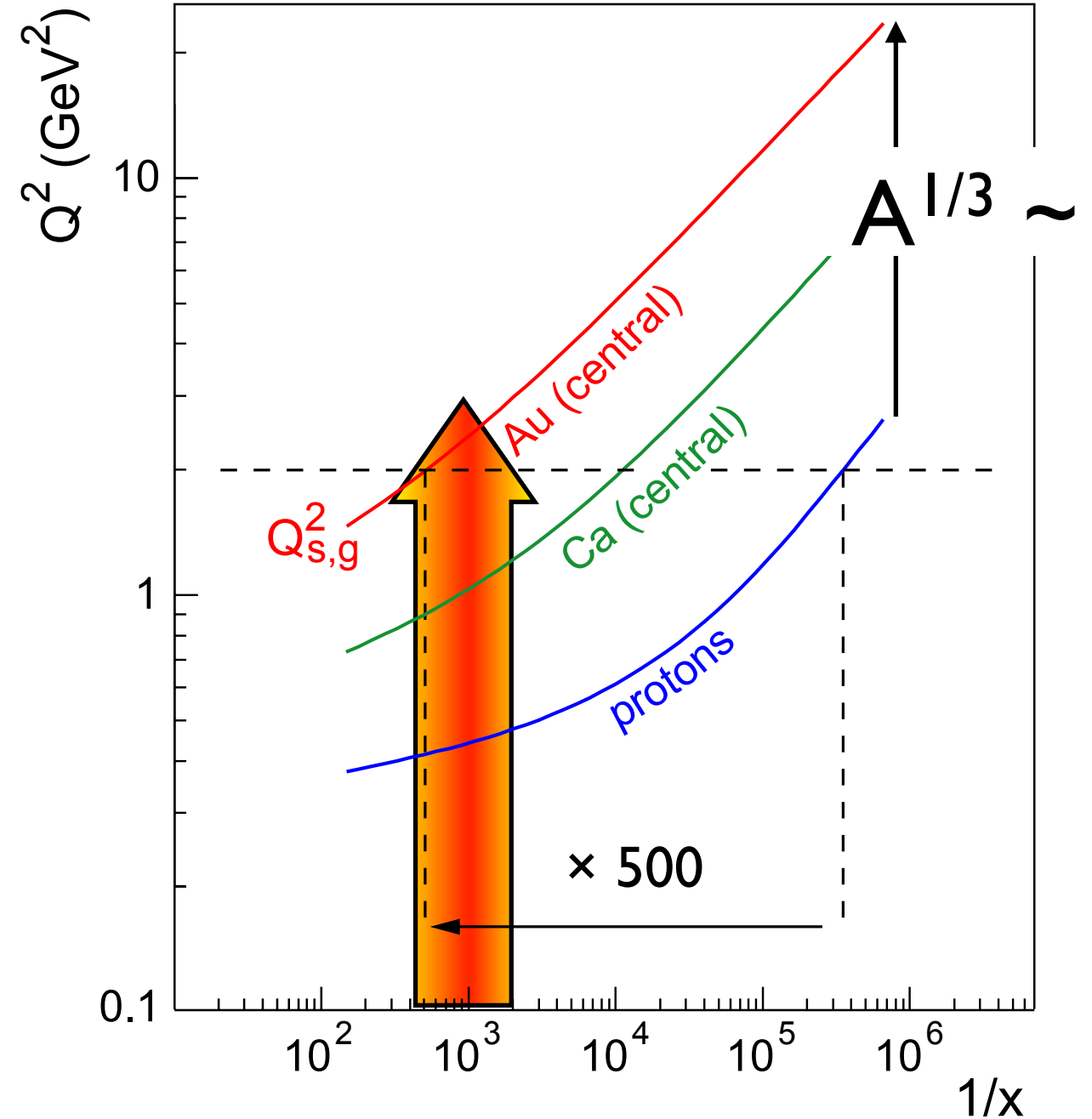
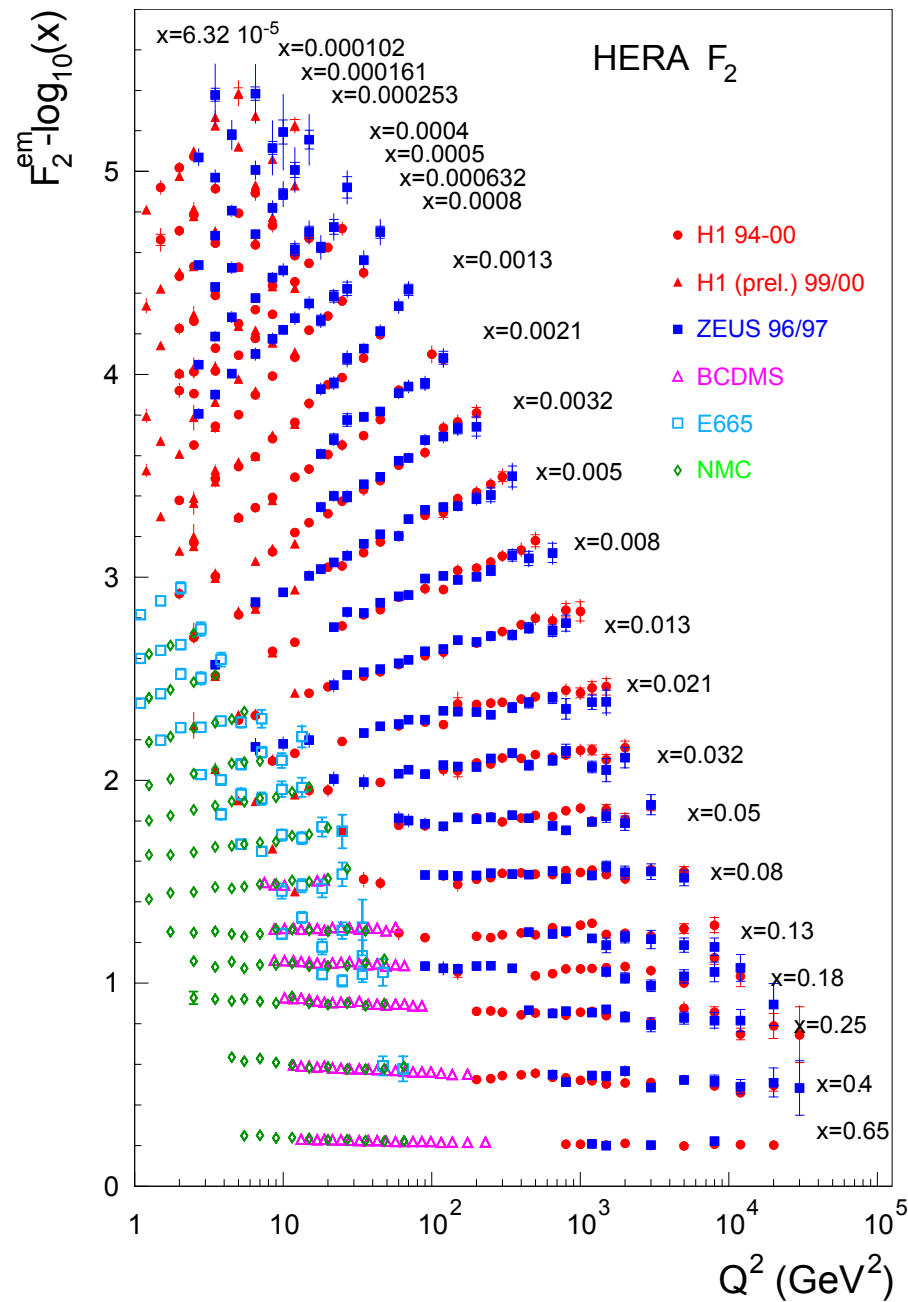
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# Are RHIC & HERA Results consistent?



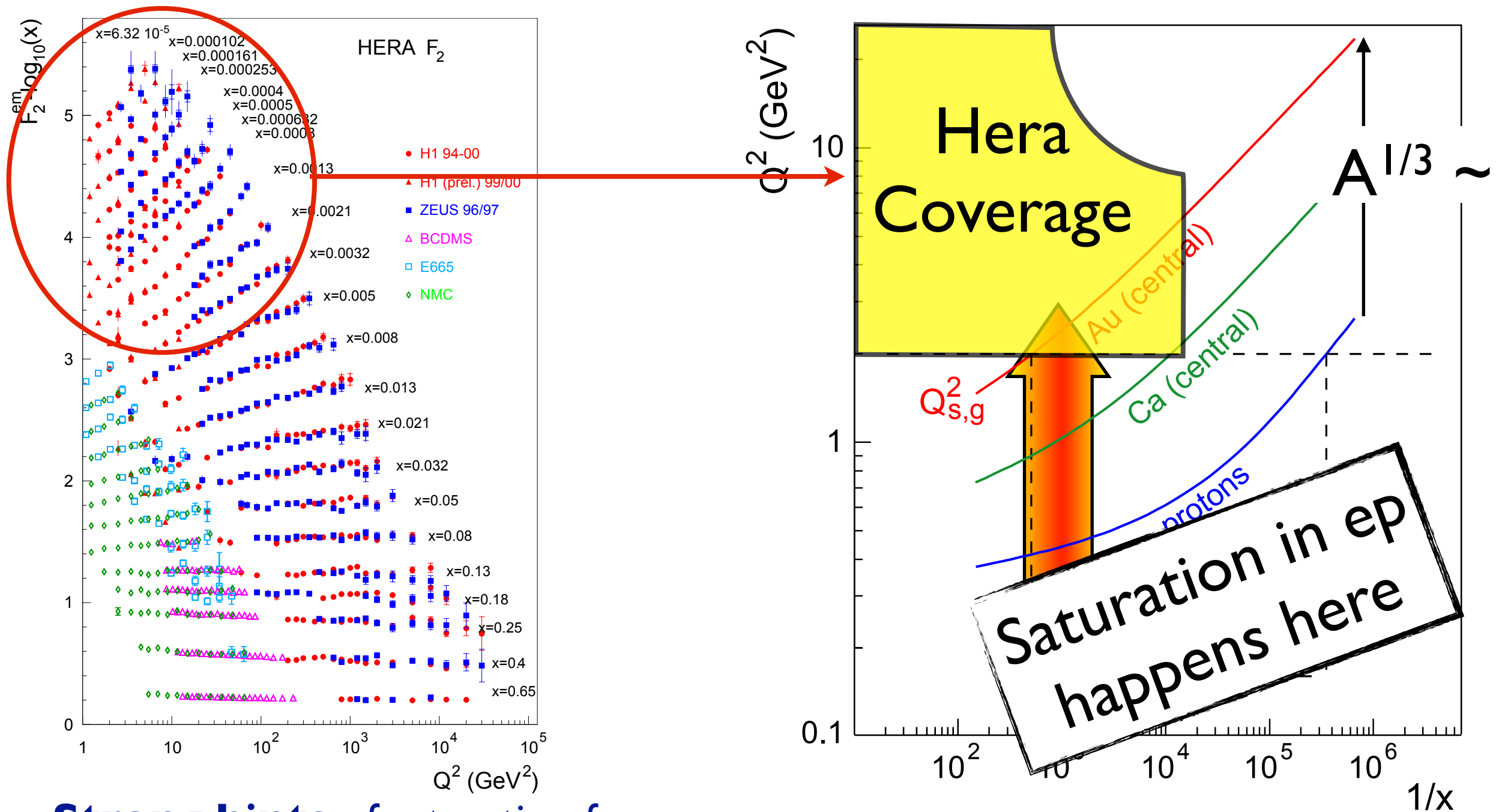
- **Strong hints** of saturation from RHIC:  $x \sim 10^{-3}$  in Au
- ep: **No/weak hints** in DIS at HERA up to  $x=6.32 \cdot 10^{-5}$ ,  $Q^2=1-5 \text{ GeV}^2$

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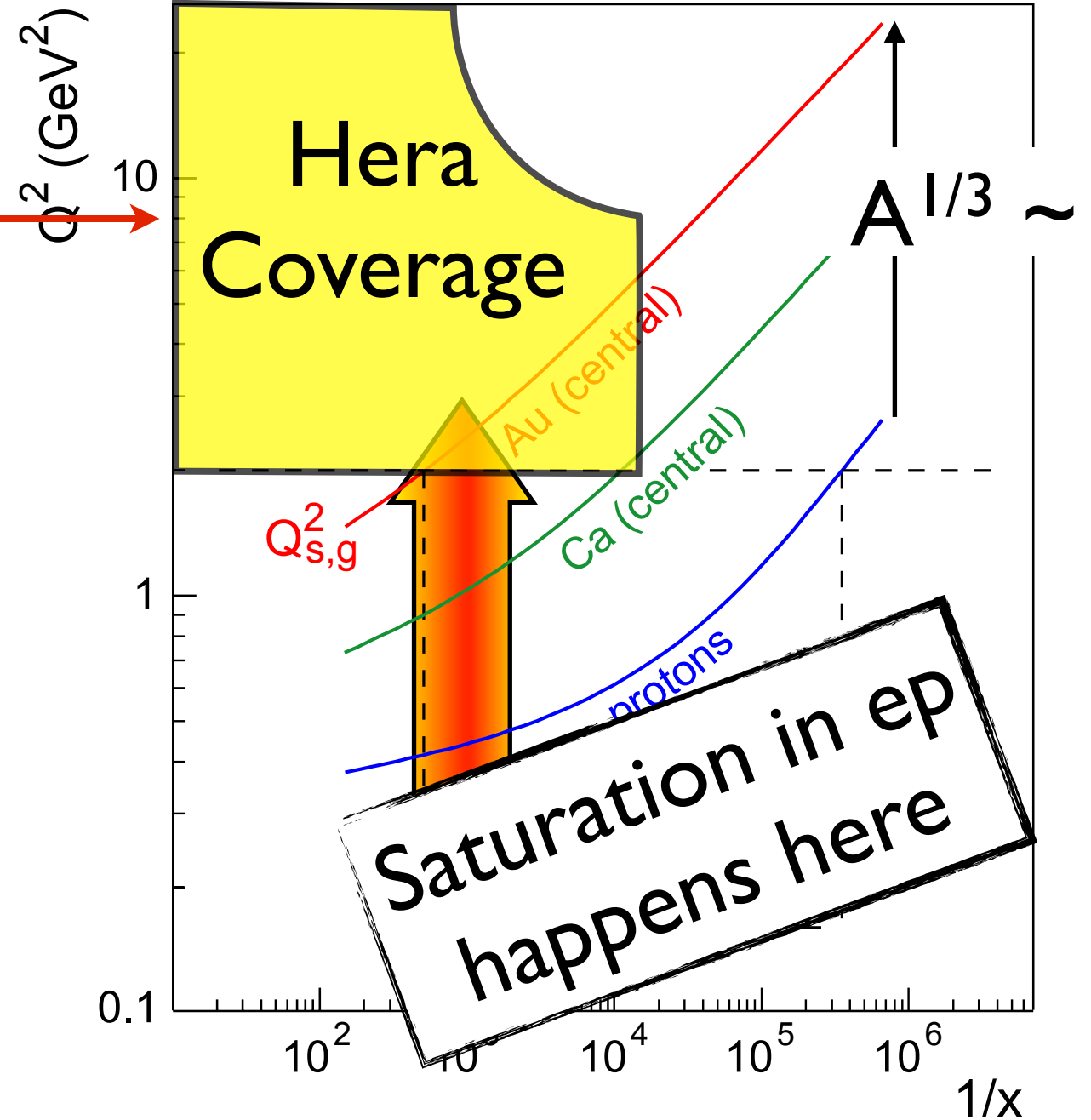
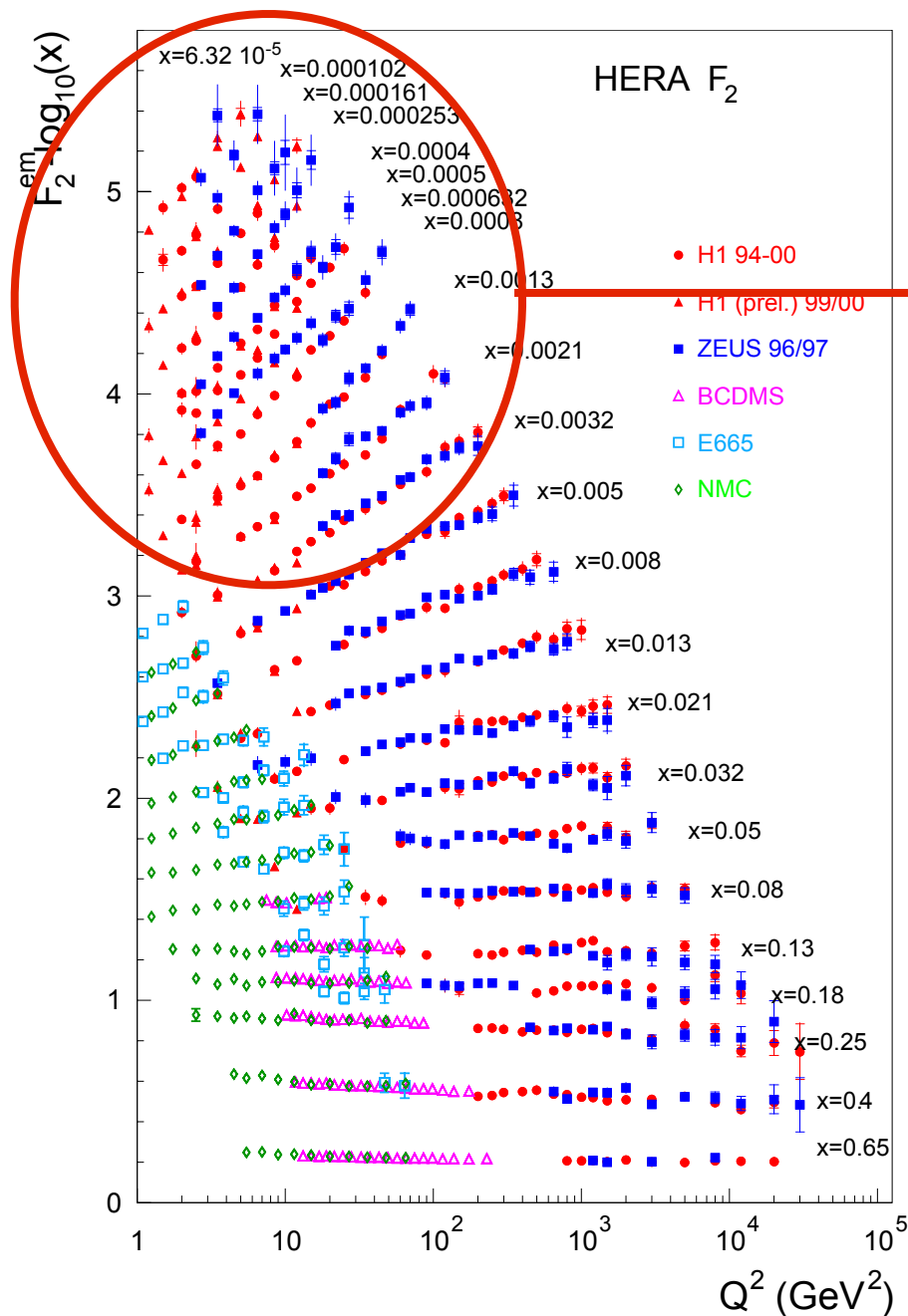
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- ep: **No/weak hints** in DIS at HERA up to  $x=6.32 \cdot 10^{-5}$ ,  $Q^2=1-5$  GeV<sup>2</sup>
- Finding RHIC and HERA &  $Q_s$  scalings consistent
- At d+A in RHIC we see the Nuclear “Oomph”  $Q_s^2 \sim Q_0^2 (A/x)^{1/3}$

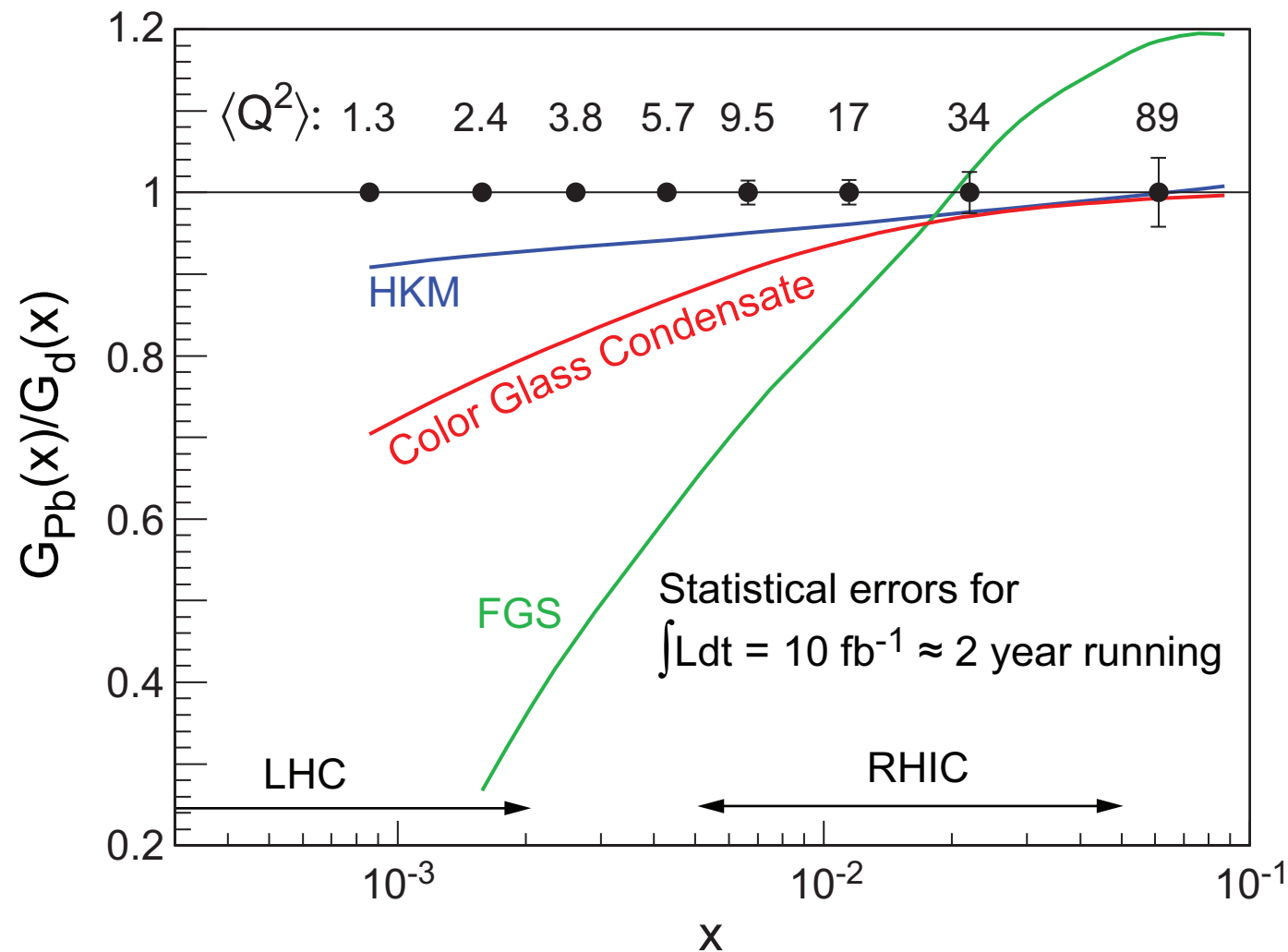


# Key Measurements in $e+A$

- **Momentum distribution of gluons  $xG(x, Q^2)$** 
  - ➔ Extract via scaling violation in  $F_2$ :  $\delta F_2 / \delta \ln Q^2$
  - ➔ Direct measurement:  $F_L \sim xG(x, Q^2)$  (requires  $\sqrt{s}$  scan)
  - ➔ 2+1 jet rates
  - ➔ Inelastic vector meson production (e.g.  $J/\psi$ )
  - ➔ Diffractive vector meson production  $\sim [xG(x, Q^2)]^2$

# Example of Key Measurements: $F_L$

$$\frac{d^2\sigma^{ep\rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[ \left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$



HKM and FGS are "standard" shadowing parameterizations that are evolved with DGLAP

$$F_L \sim \alpha_s xG(x, Q^2)$$

requires  $\sqrt{s}$  scan,  $Q^2/xs = y$

Here:

$$\begin{aligned} \int L dt &= 4/A \text{ fb}^{-1} \text{ (10+100) GeV} \\ &= 4/A \text{ fb}^{-1} \text{ (10+50) GeV} \\ &= 2/A \text{ fb}^{-1} \text{ (5+50) GeV} \end{aligned}$$

statistical error only

Syst. studies of  $F_L(A, x, Q^2)$ :

- $xG(x, Q^2)$  with great precision
- Distinguish between models

# Feasibility study: $\sigma_r = F_2(x, Q^2) - y^2/Y_+ \cdot F_L(x, Q^2)$

$$Y_+ = 1 + (1 - y)^2$$

## Strategies:

slope of  $y^2/Y_+$  for  
different  $s$  at fixed  $x$  &  
 $Q^2$

e+p: 1st stage

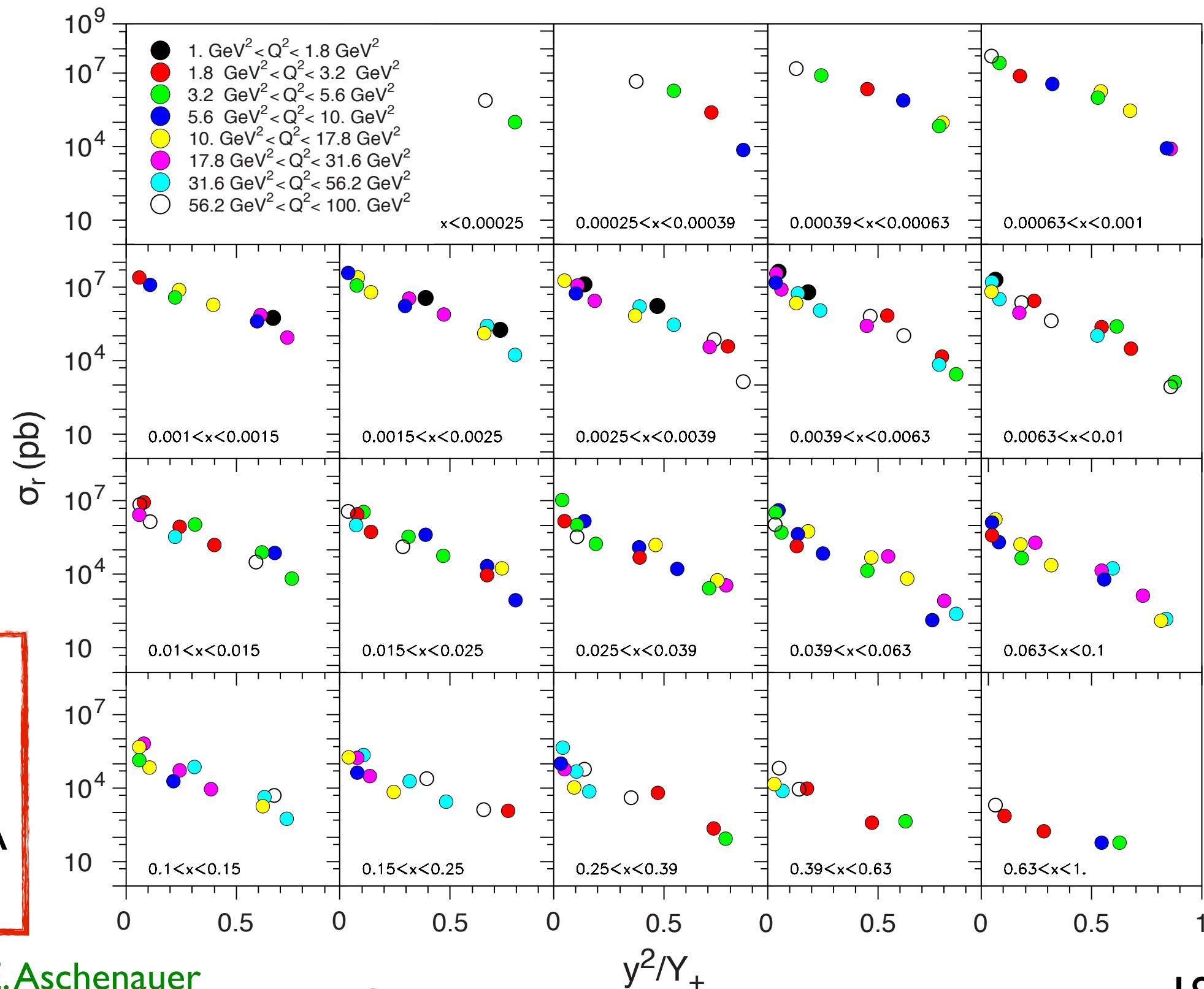
5x50 - 5x325

running combined  
4 weeks/each  
(50% eff)

stat. error shown  
and negligible

## To Do:

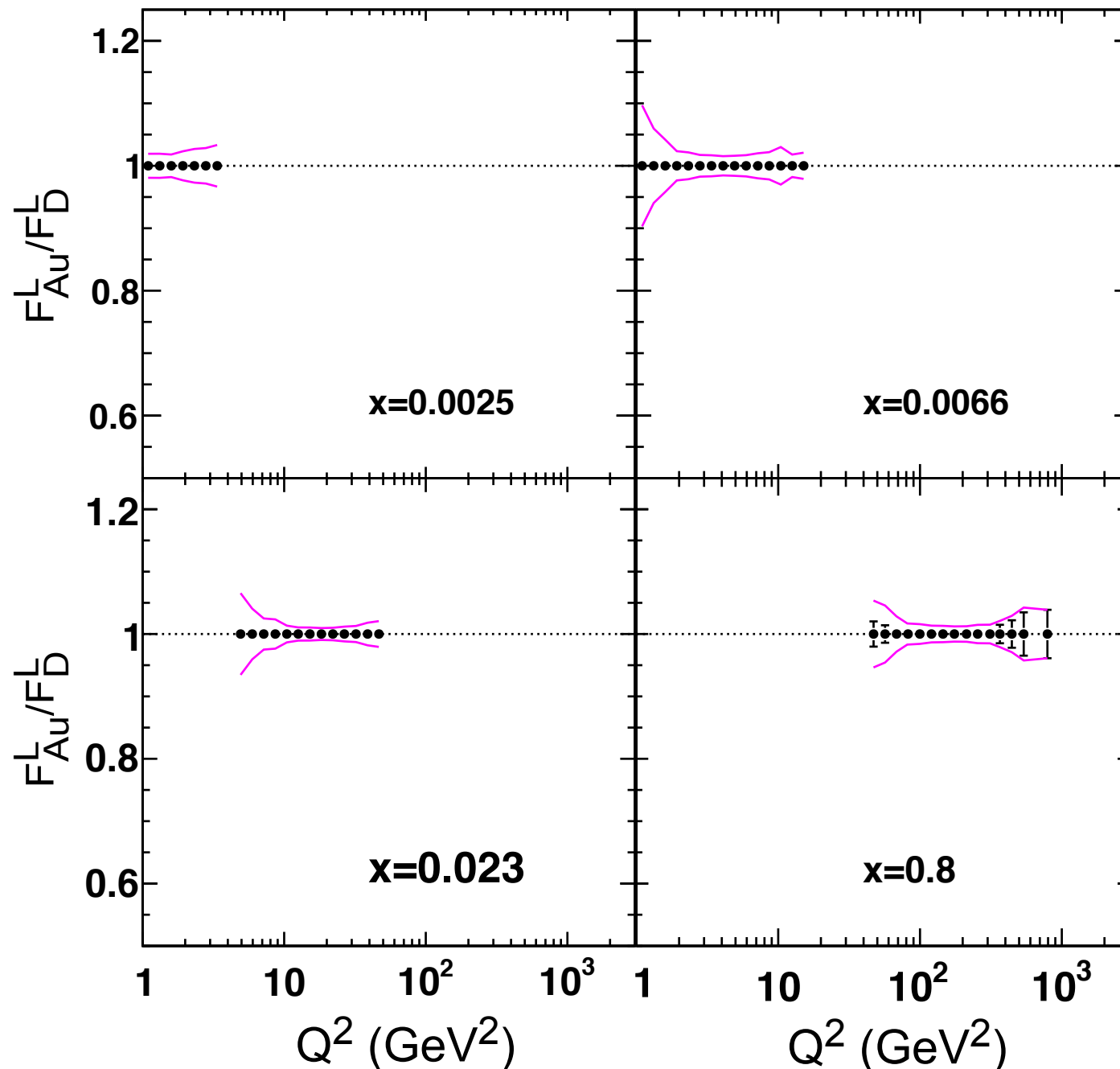
refine method &  
test how well we  
can extract  $F_L$  in e+A  
collisions



# Syst. Uncertainties in $F_L$ for staged EIC

- $F_L$  for electron energy fixed at 4 GeV and proton energies: 50, 70, 100, 250 GeV (4fb<sup>-1</sup> each)

Plot from J. Dunlop



The **magenta** curves show the statistical and systematic errors (1% uncertainty in normalization) added in quadrature.

Again, the extraction of  $F_L$  is **dominated by systematic uncertainties**

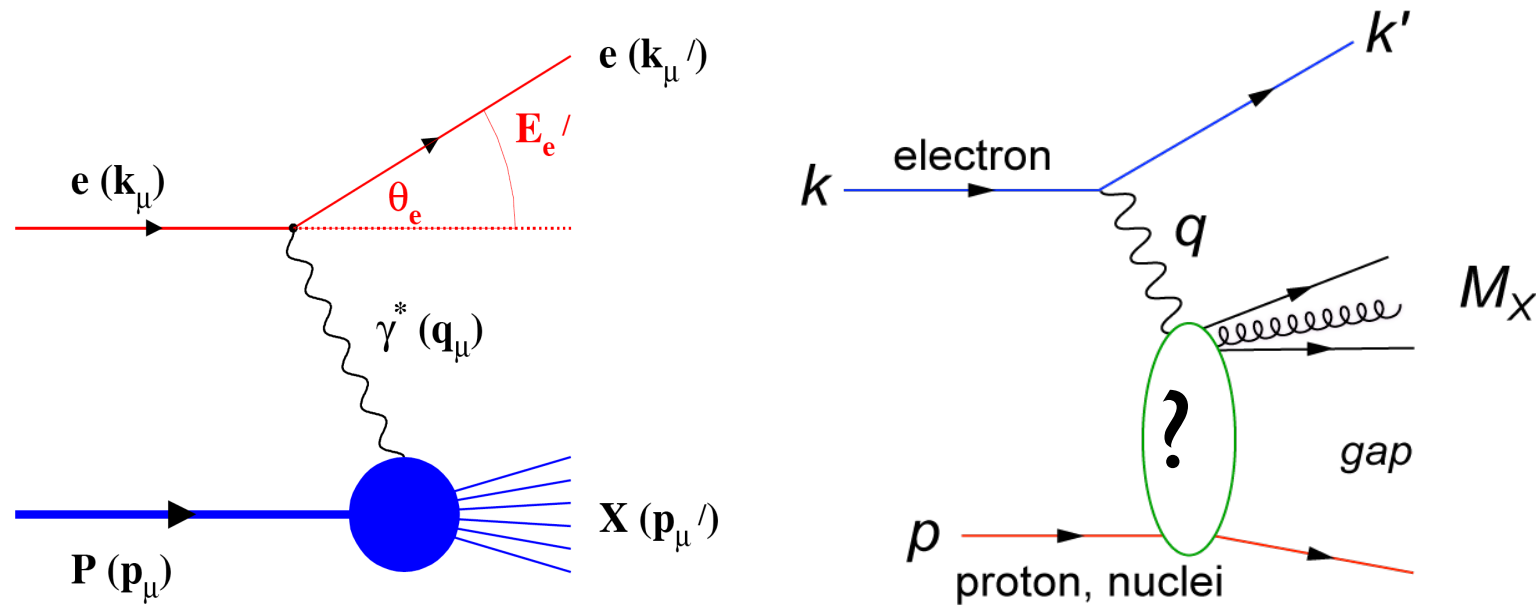
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  - ➔ Diffractive vector meson production  $\sim [xG(x, Q^2)]^2$
- **Space-time distributions of gluons in matter**
  - ➔ Exclusive final states (e.g. vector meson production  $\rho, J/\psi$ )
  - ➔ Deep Virtual Compton Scattering (DVCS) -  $\sigma \sim A^{4/3}$
  - ➔  $F_2, F_L$  for various  $A$  and impact parameter dependence
- **Interaction of fast probes with *gluonic* medium?**
  - ➔ Hadronization, Fragmentation
  - ➔ Energy loss (charm!)
- **Role of colour neutral excitations (Pomerons)**
  - ➔ Diffractive cross-section  $\sigma_{diff}/\sigma_{tot}$  (HERA/ $ep$ : 10% , EIC/ $eA$ : 30%?)
  - ➔ Diffractive structure functions and vector meson production
  - ➔ Abundance and distribution of rapidity gaps

See next talk by A. Accardi

# Role of colour-neutral (Pomeron) excitations

$$\frac{d\sigma}{dt} \Big|_{t=0} (\gamma^* \rightarrow M_X A) \propto \alpha^2 [G_A(x, Q^2)]^2$$

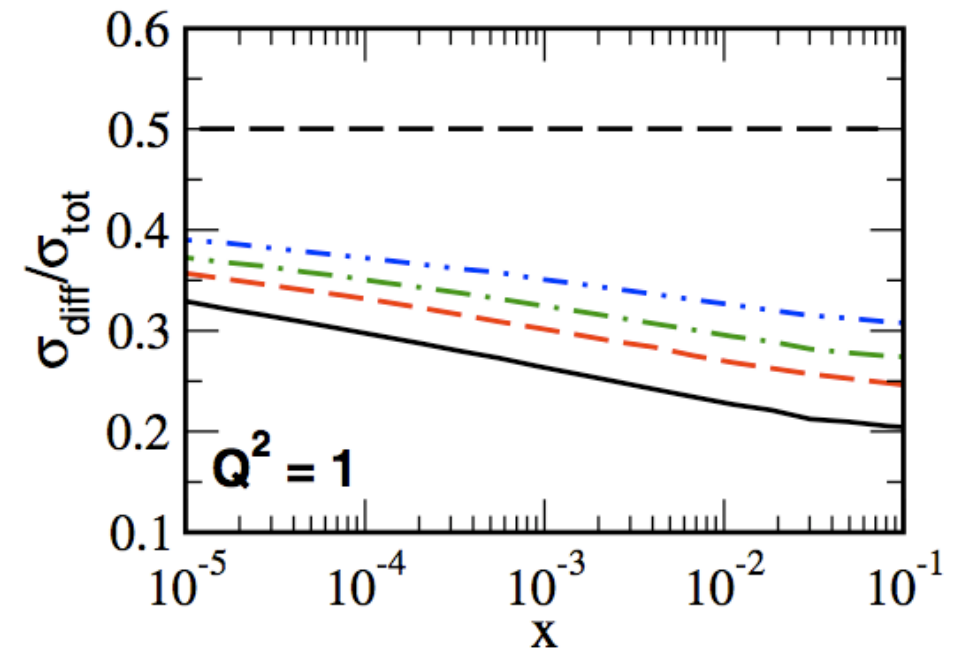
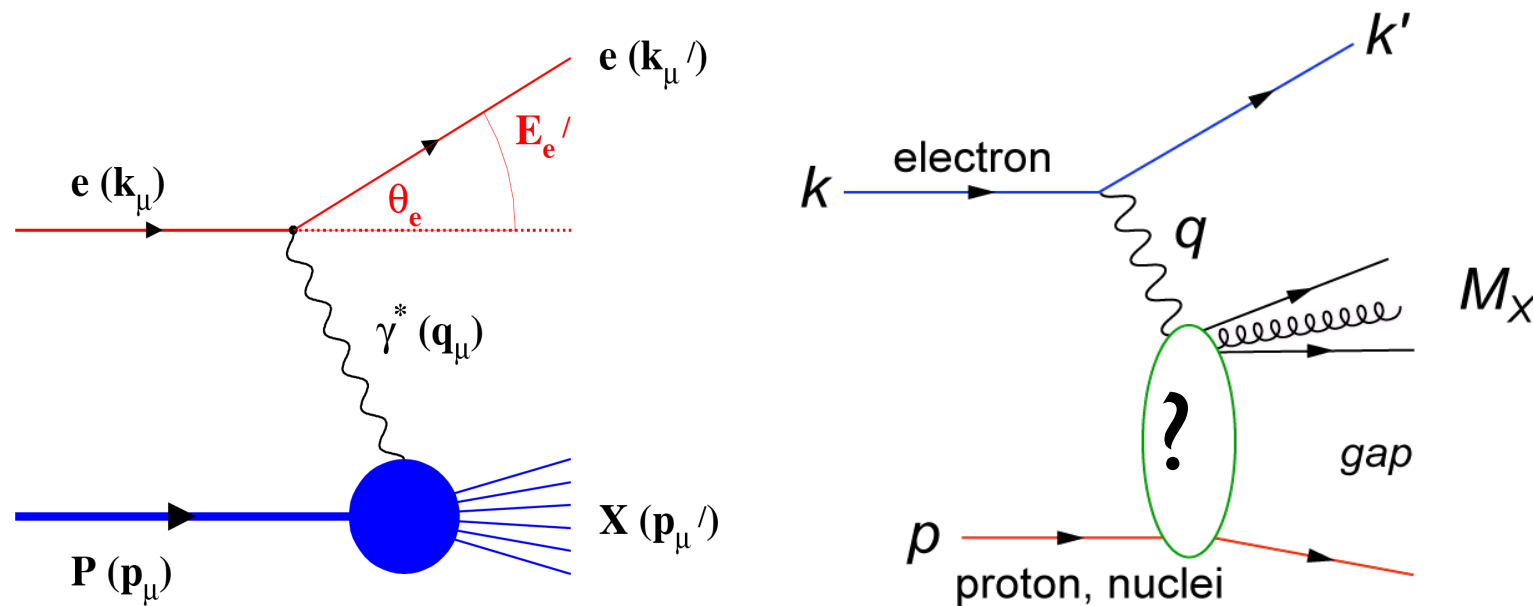




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Curves: Kugeratski, Goncalves,  
Navarra, EPJ C46, 413



HERA/ep: 15% of all events are hard diffractive

Diffractive cross-section  $\sigma_{\text{diff}}/\sigma_{\text{tot}}$  in  $e+A$  ?

- Predictions: ~25-40%?

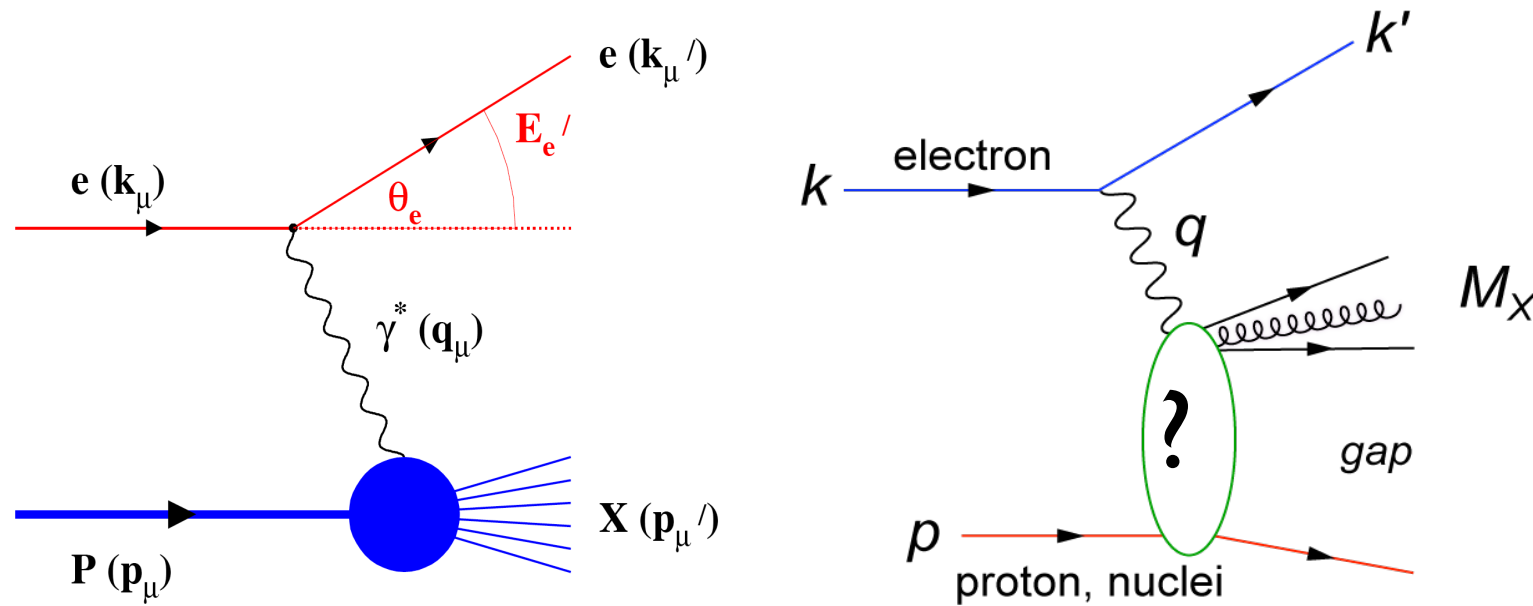
Look inside the “Pomeron”

- Diffractive structure functions

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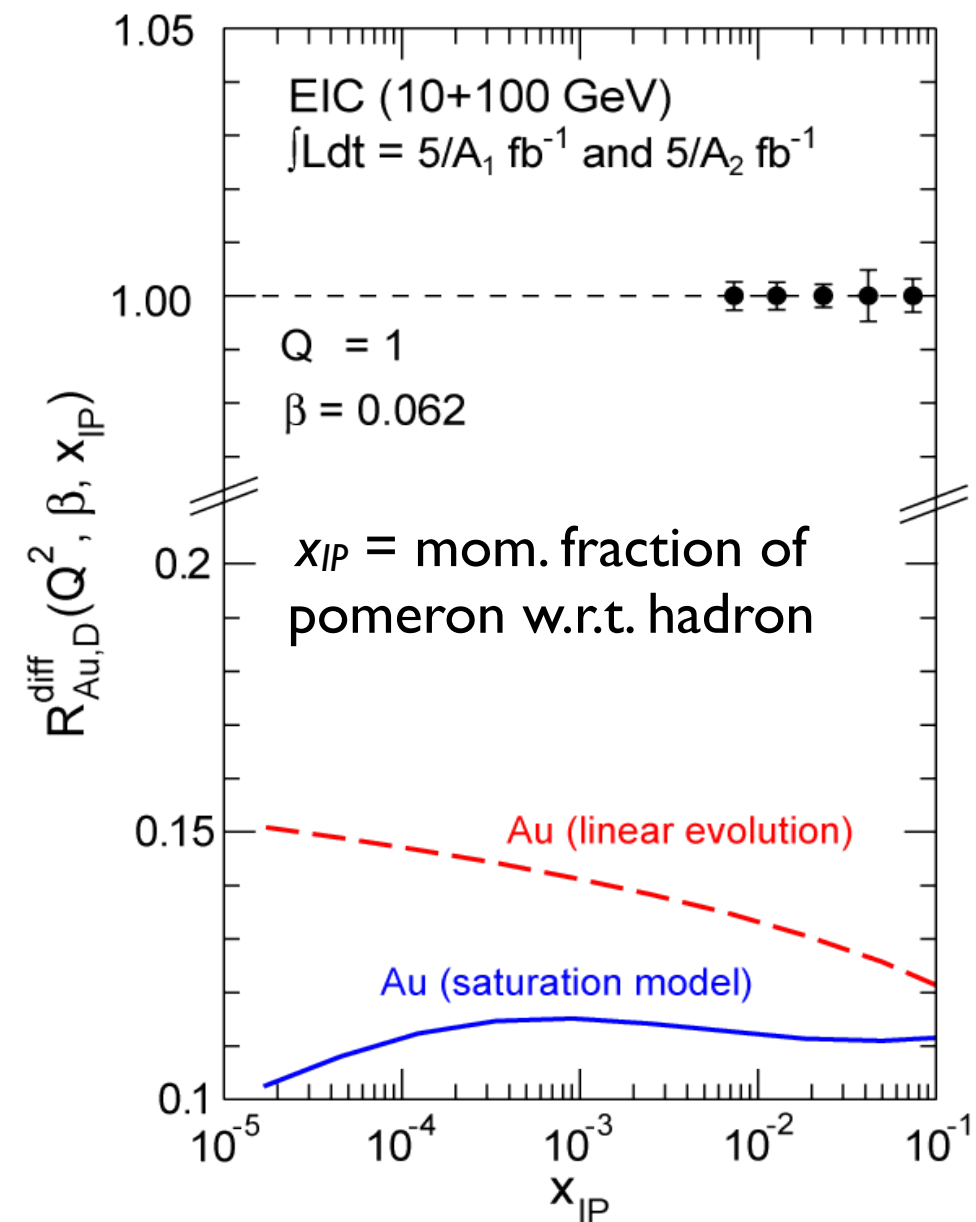
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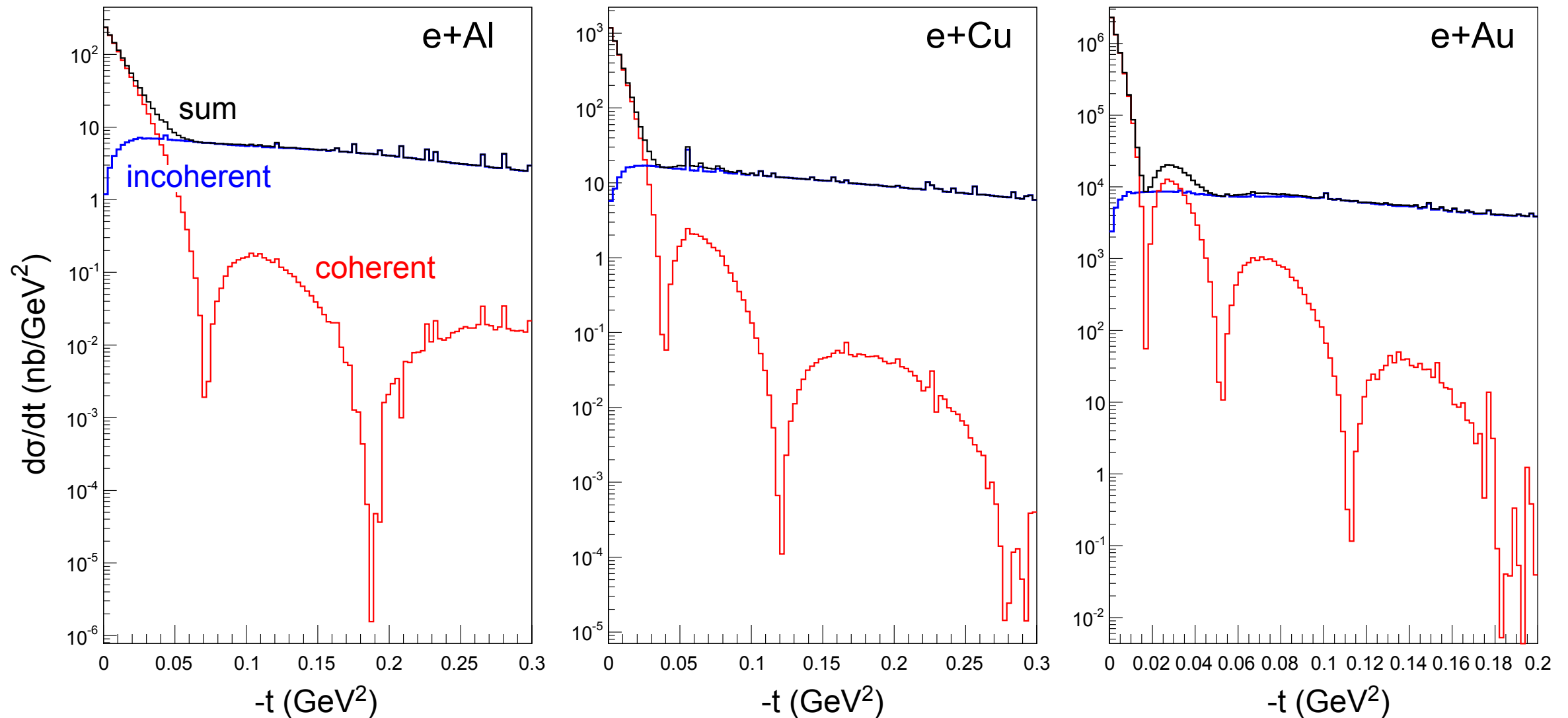
- Diffractive structure functions

Distinguish between **linear evolution** and **saturation** models



# Diffraction in e+A

$$e + A \rightarrow e' + J/\psi + A'$$



- Diffractive cross-section  $\sigma_{\text{diff}}/\sigma_{\text{tot}}$  in e+A predicted to be  $\sim 25\text{-}40\%$
- Process most sensitive to  $xG(x, Q^2)$
- Rich physics program on momentum & spatial gluon distribution
- Coherent vs Incoherent: requires detection of breakup with  $\sim 1\text{-}10^{-4}$  efficiency

# Detecting Nuclear Breakup

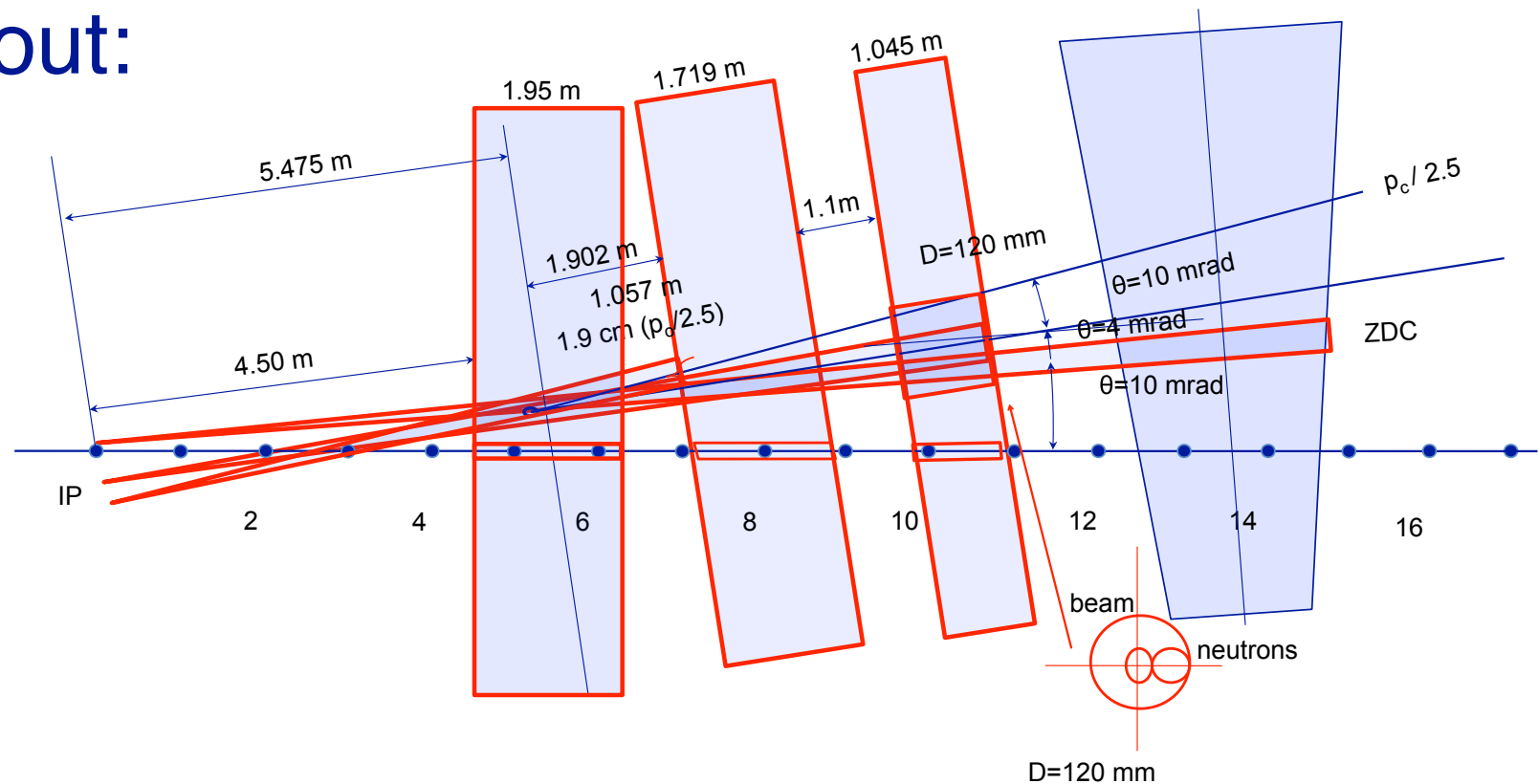
- Detecting **all** fragments  $p_{A'} = \sum p_n + \sum p_p + \sum p_d + \sum p_\alpha \dots$  not possible
- Focus on n emission
  - ➔ Zero-Degree Calorimeter
  - ➔ Requires careful design of IR
- Additional measurements:
  - ▶ Fragments via Roman Pots
  - ▶  $\gamma$  via EMC

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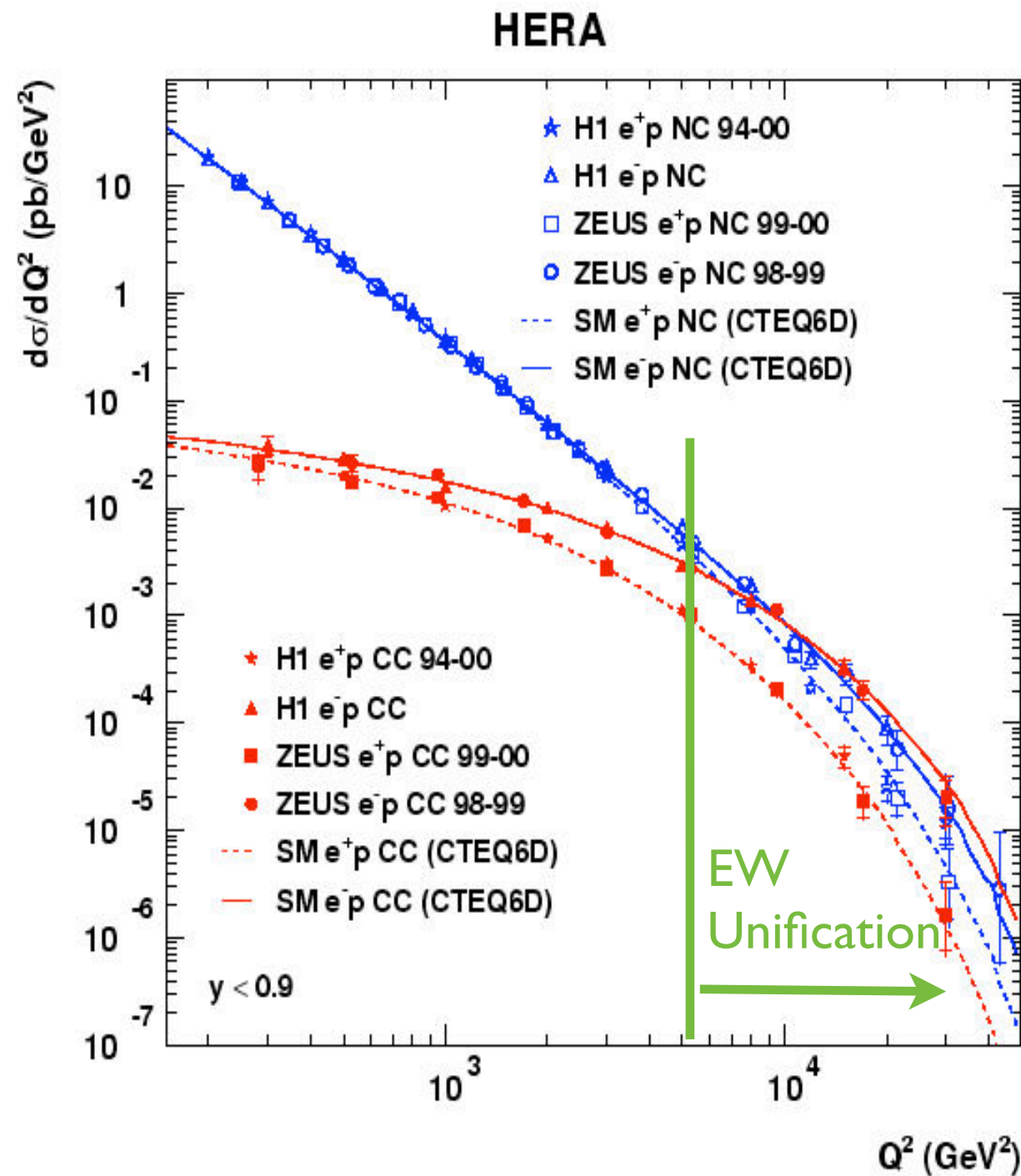
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- Proposed eRHIC IR layout:

- ➡ Need  $\pm X$  mrad opening
- ➡ through triplet for  $n$  and
- ➡ room for ZDC



# EW unification at HERA



- From DIS at HERA:

➡ At small-medium  $Q^2$ ,  $\sigma(\text{NC}) \gg \sigma(\text{CC})$

➡ For  $Q^2 > M_Z^2$  and  $M_W^2$ ,  $\sigma(\text{NC}) \sim \sigma(\text{CC})$

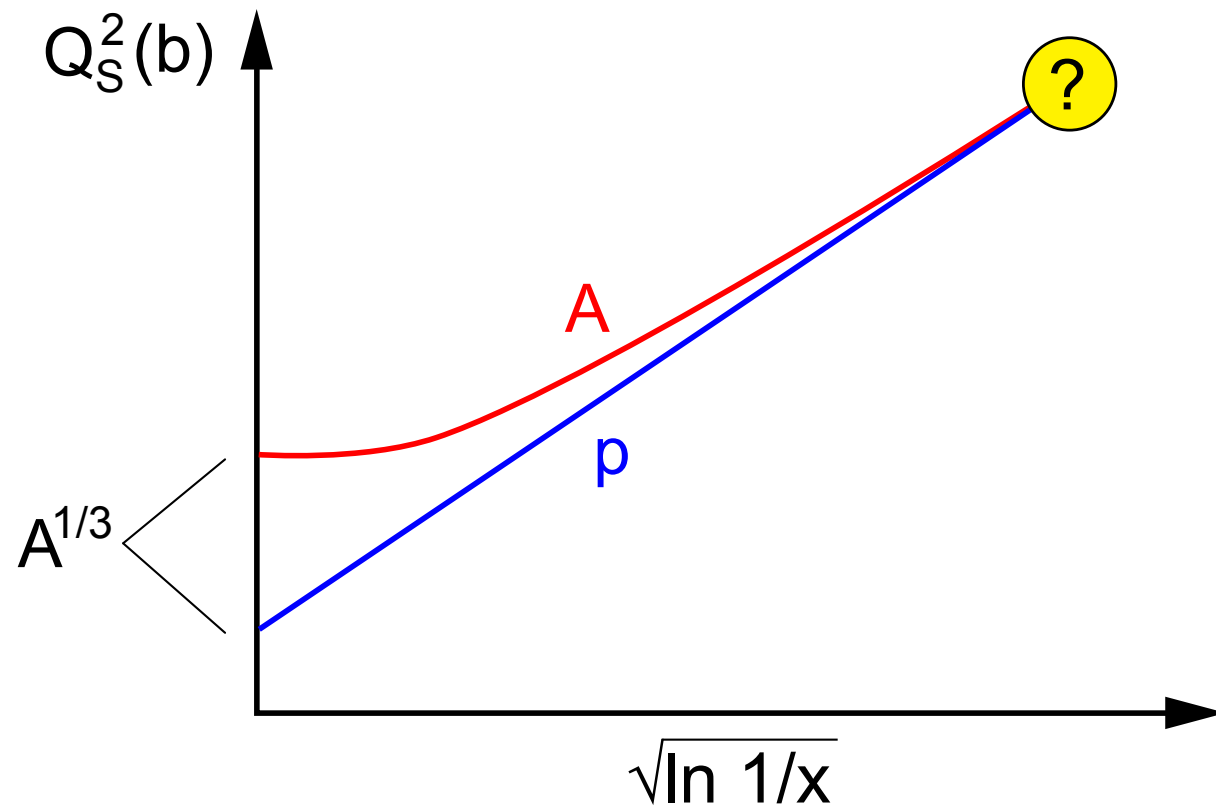
► EW Unification

- Already a textbook figure ...



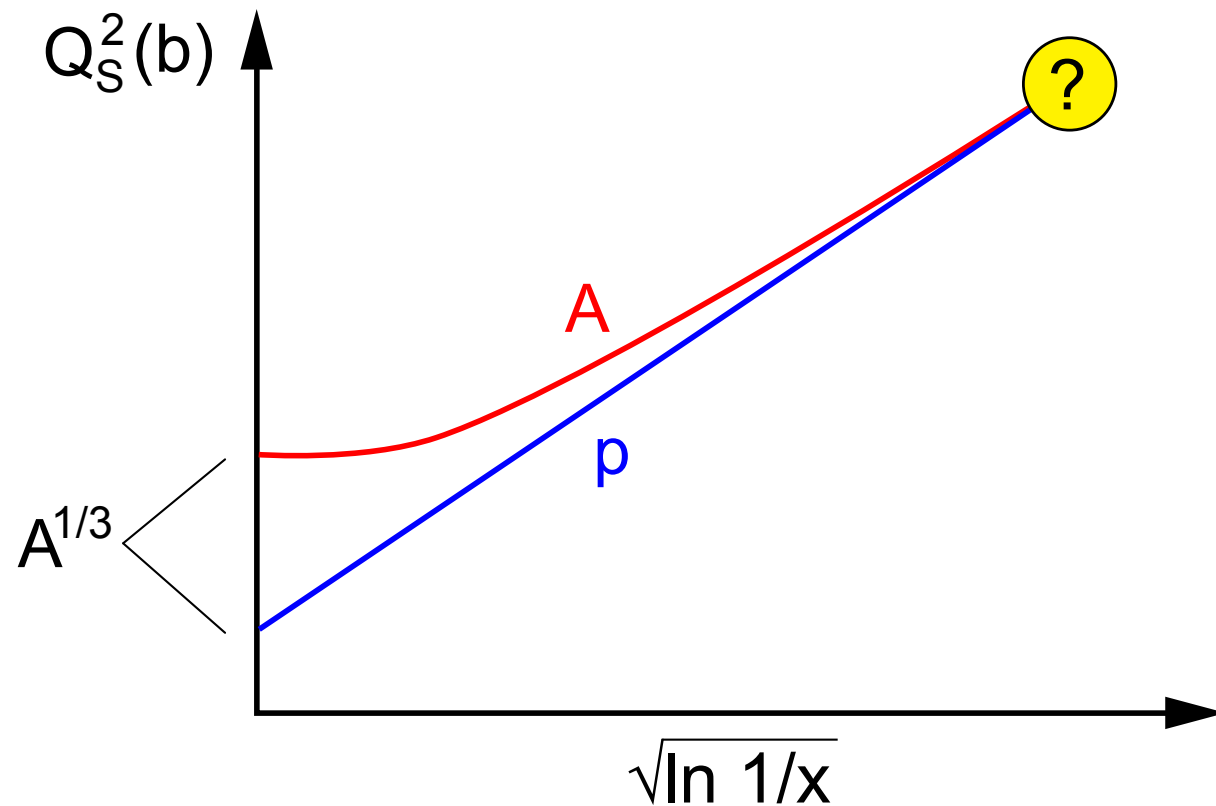
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What about on the parton scale?



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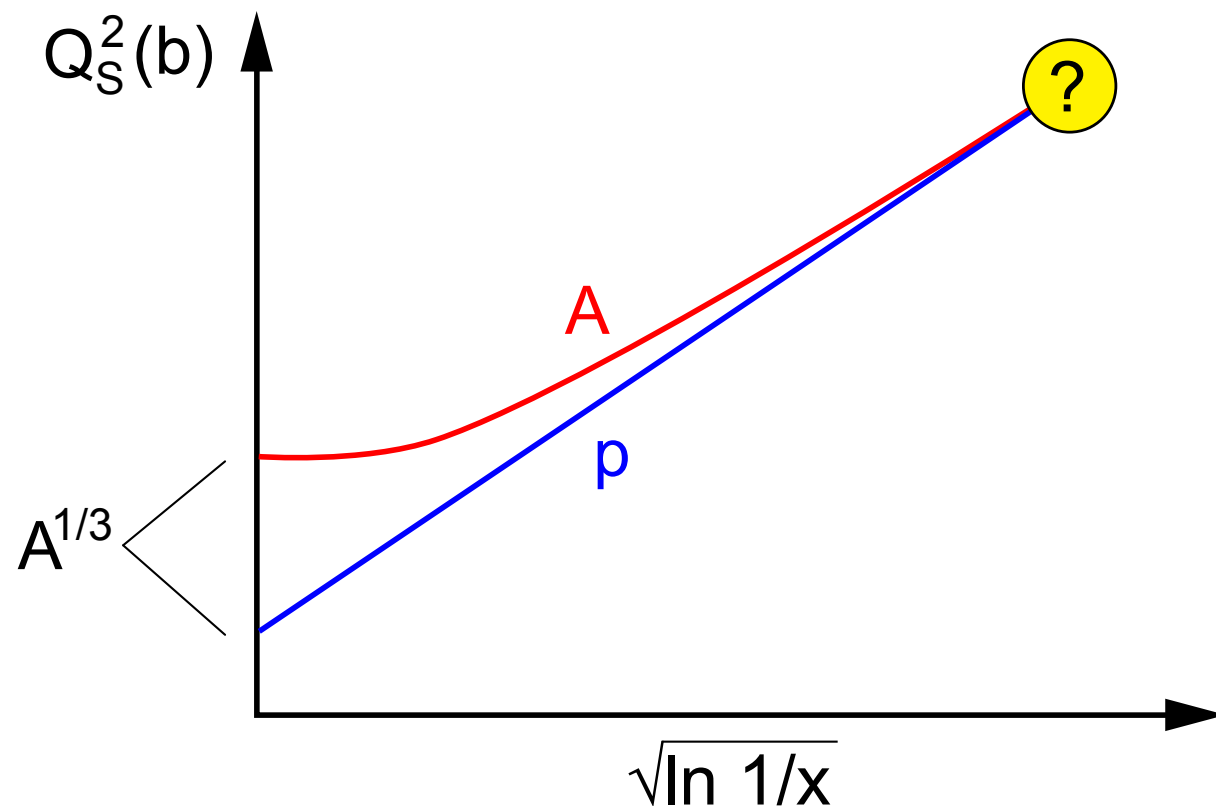


*A.H. Mueller, hep-ph/0301109*

- Small-x running-coupling BFKL QCD evolution predicts:
  - $Q_s$  approaches universal behaviour for **all** hadrons and nuclei
  - No dependence on  $A$ !!
  - Not only functional form  $f(Q_s)$  universal, but even  $Q_s$  itself becomes universal

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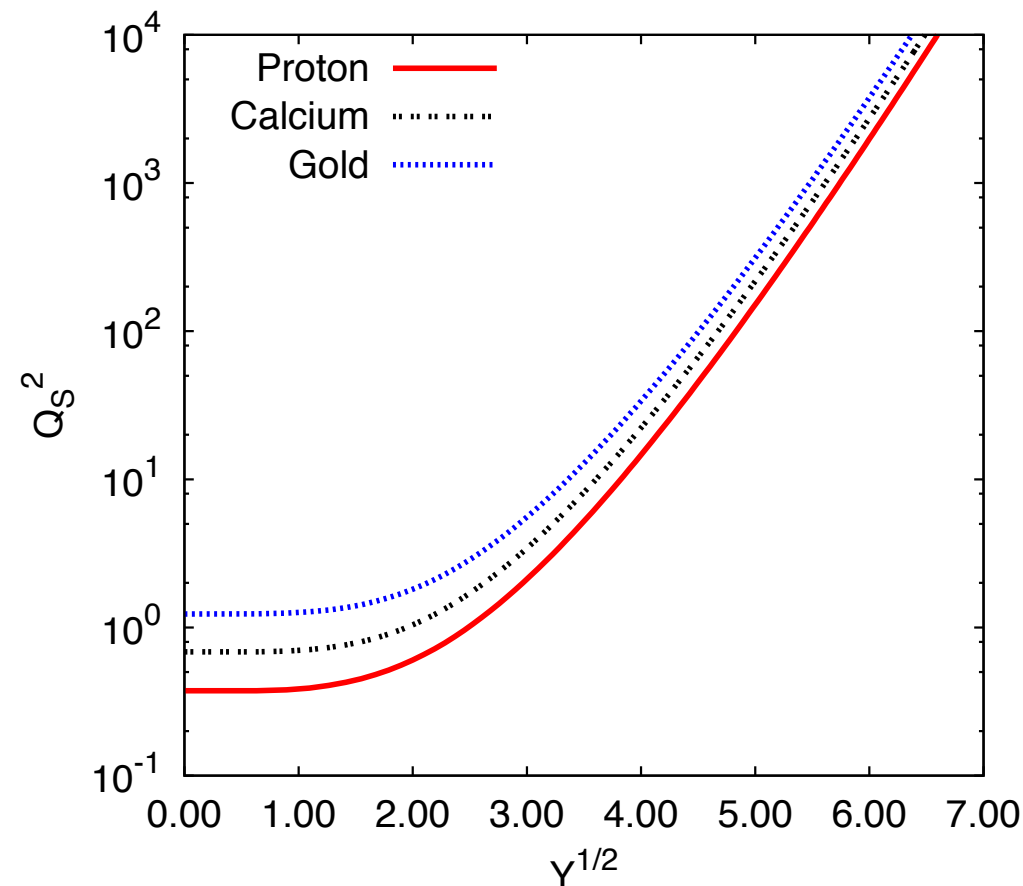
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## Radical View:

- ➡ Nuclei and all hadrons have a component of their wave function with the **same** behaviour
- ➡ This is a conjecture! Needs to be tested

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*Dusling, Gelis, Lappi and Venugopalan arXiv:0911.2720*

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# Summary

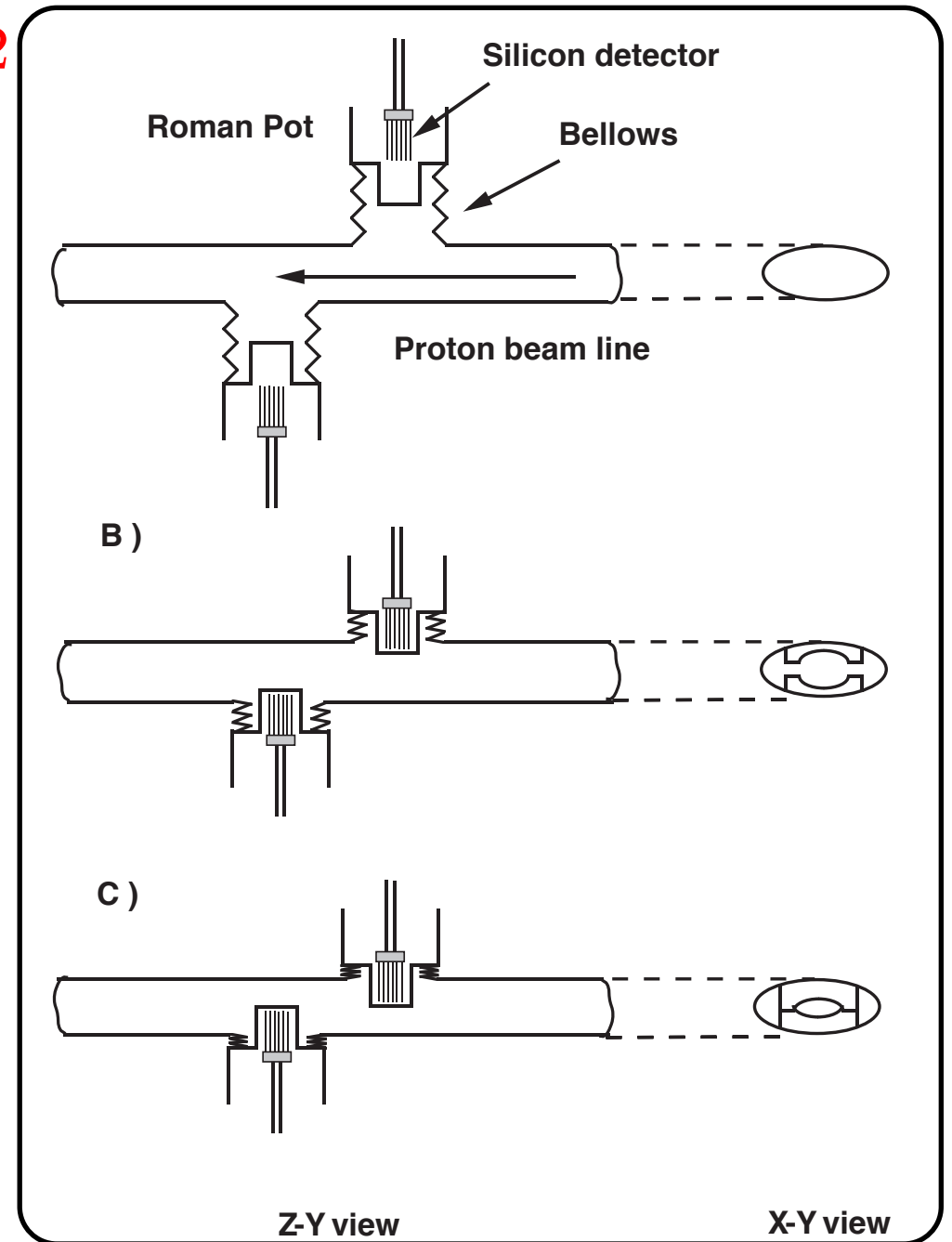
- The e+A programme at an EIC will allow for the unprecedented study on nuclear matter, not well described by linear QCD
- Very little is known about the nuclear structure at low-x
  - ➡ inclusive structure functions and exclusive vector meson production data only available at high-x/light nuclei
  - ➡ diffractive structure functions have never been measured
- Crucial measurements at an EIC:
  - ➡ inclusive and diffractive structure functions
  - ➡ di-hadron correlations: the  $k_T$  dependence of gluon correlations
  - ➡ coherent diffraction: spatial distribution of gluons
  - ➡ incoherent diffraction: spatial correlations between low-x gluons
- However, many experimental challenges need to be resolved
  - ➡ Differentiating coherent and incoherent diffractive events
  - ➡ Radiative corrections for e+A collisions

BACKUP



# How to measure coherent diffraction in e+A ?

$$\frac{d\sigma}{dt} \Big|_{t=0} (\gamma^* \rightarrow M_X A) \propto \alpha^2 [G_A(x, Q^2)]^2$$



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  - ➡  $p_T^{\min} \sim pA\theta_{\min}$ 
    - For beam energies = 100 GeV/n and  $\theta_{\min} = 0.08$  mrad:
- These are large momentum kicks, >> the binding energy ( $\sim 8$  MeV)

species (A)	$p_T^{\min}$ (GeV/c)
d (2)	0.02
Si (28)	0.22
Cu (64)	0.51
In (115)	0.92
Au (197)	1.58
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For large A, nucleus cannot be separated from beam without breaking up

# How else to measure diffraction in $e+A$ ?

Large Rapidity Gap Method:

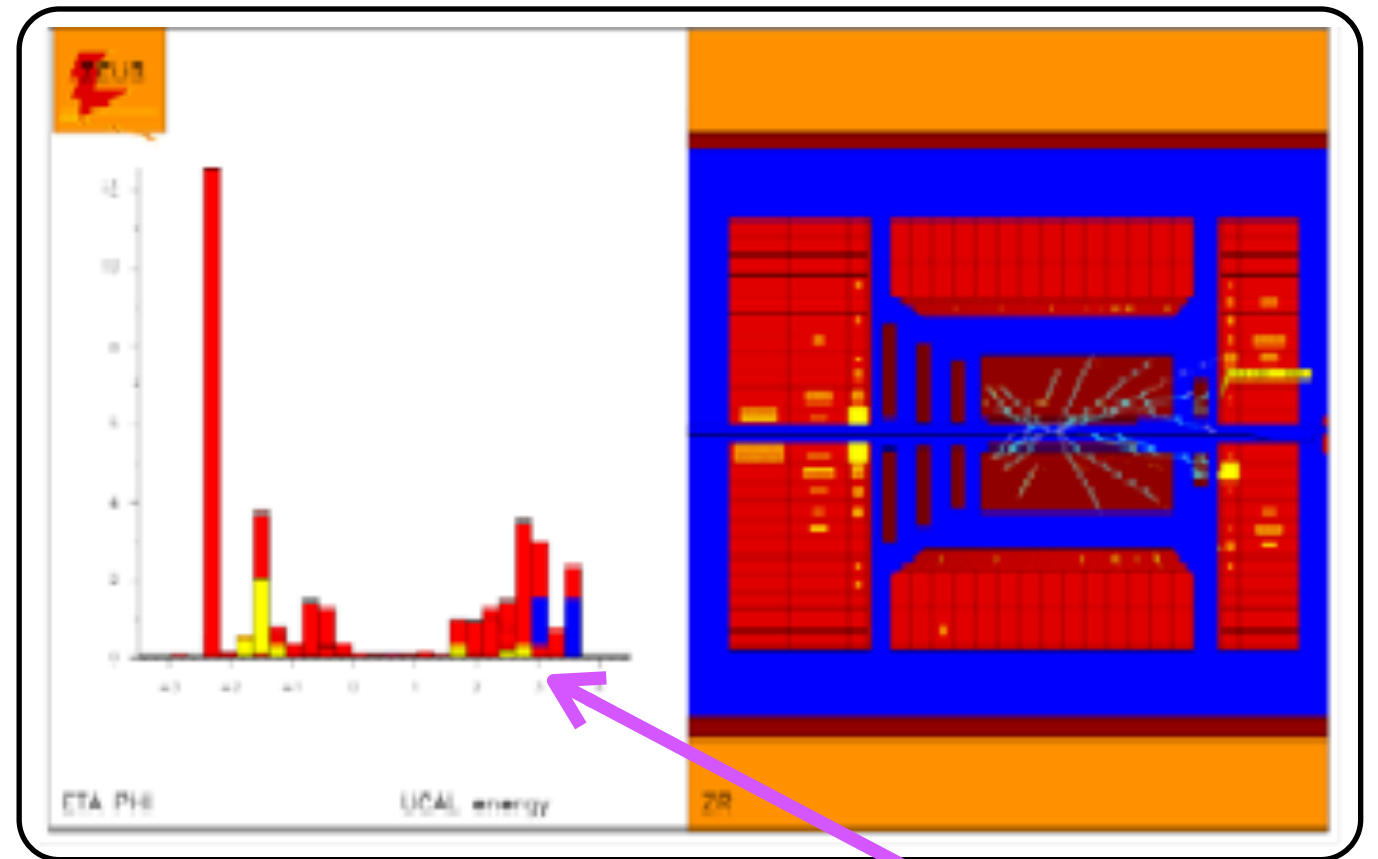
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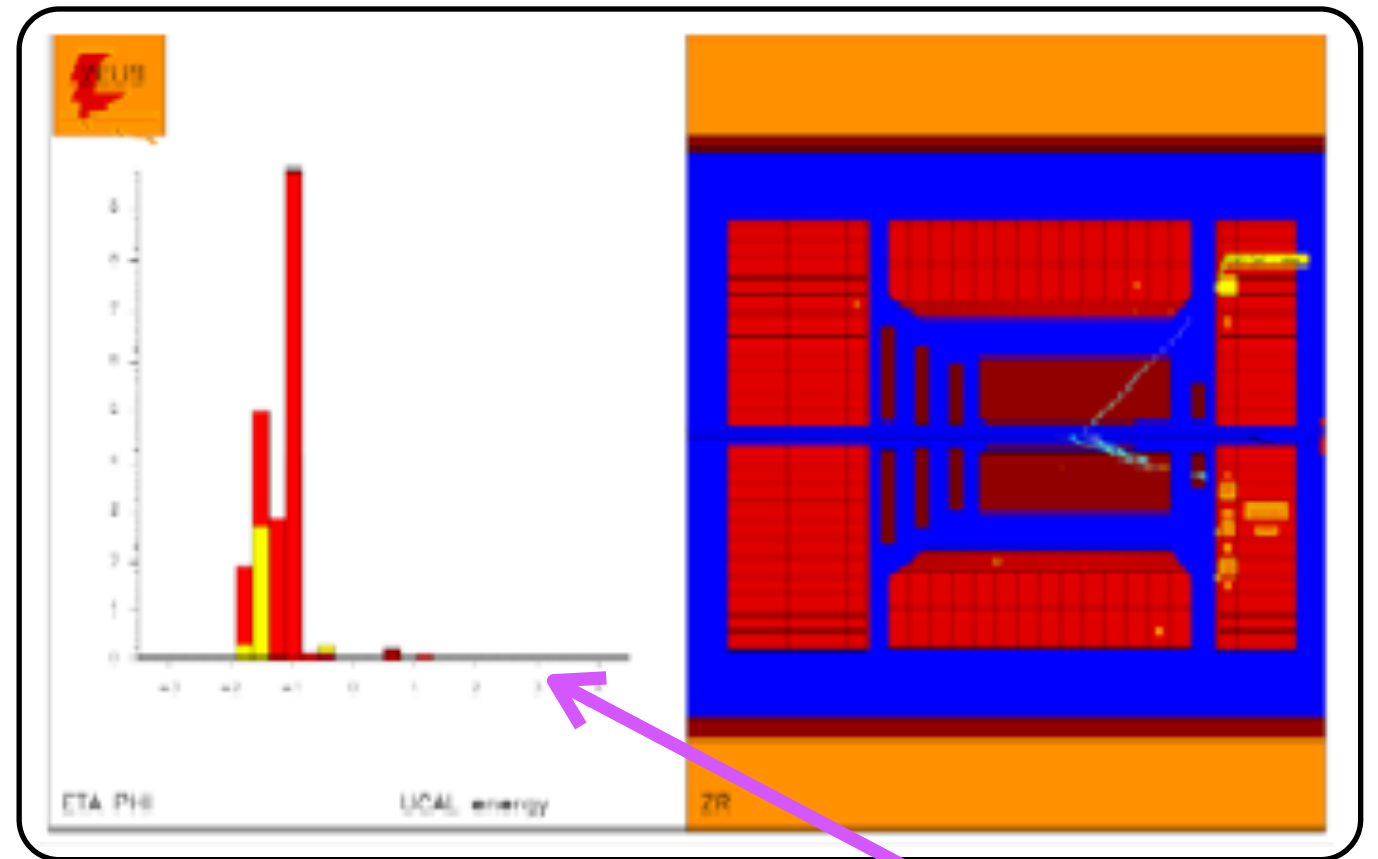


activity in the proton direction

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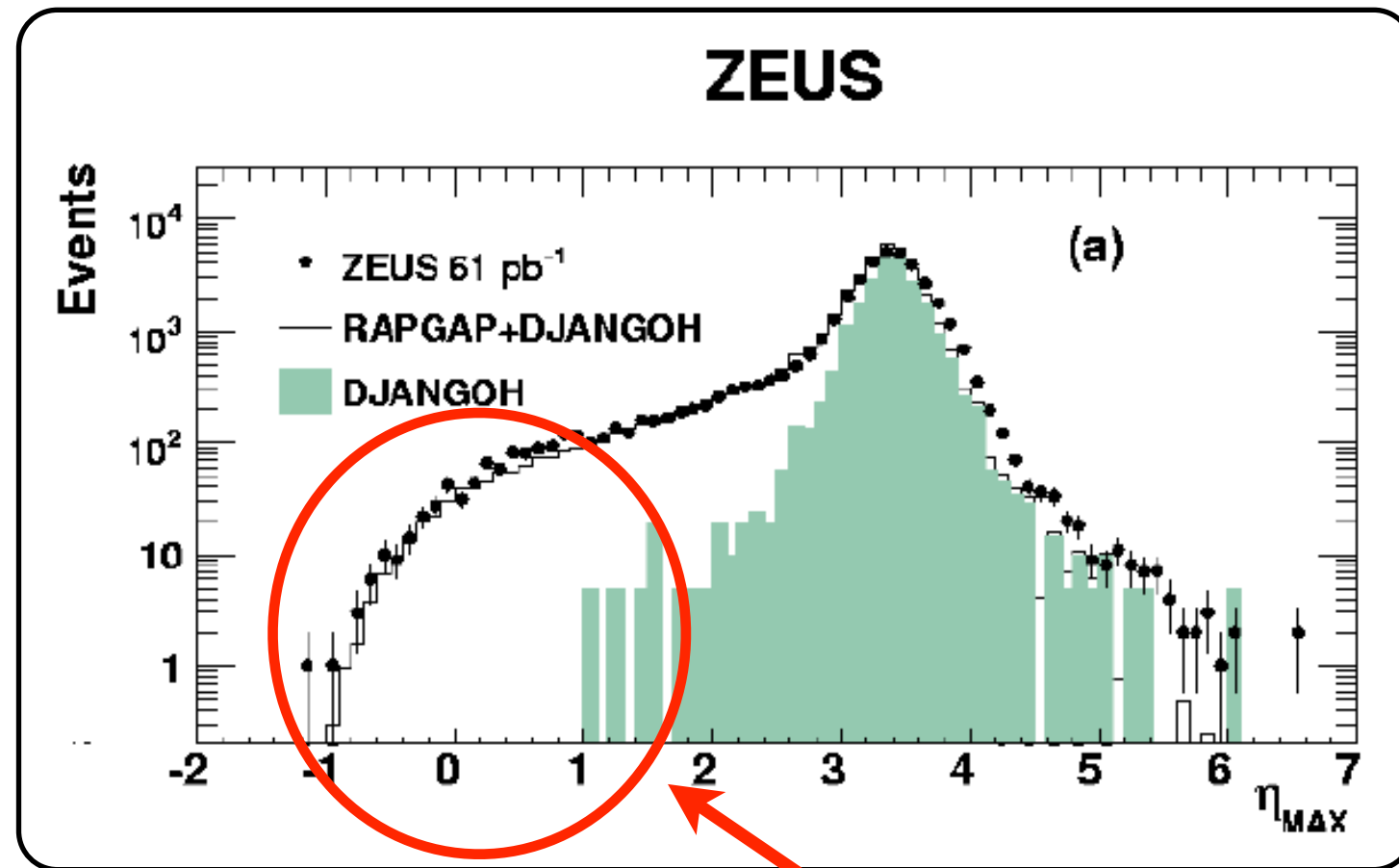


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- At HERA:  $\Delta\eta \sim 7 \Rightarrow$  hadronization reduces this to  $\sim 2.5$

- Pros

➡ Lots of statistics

- Cons

➡ Sensitive to detector acceptance

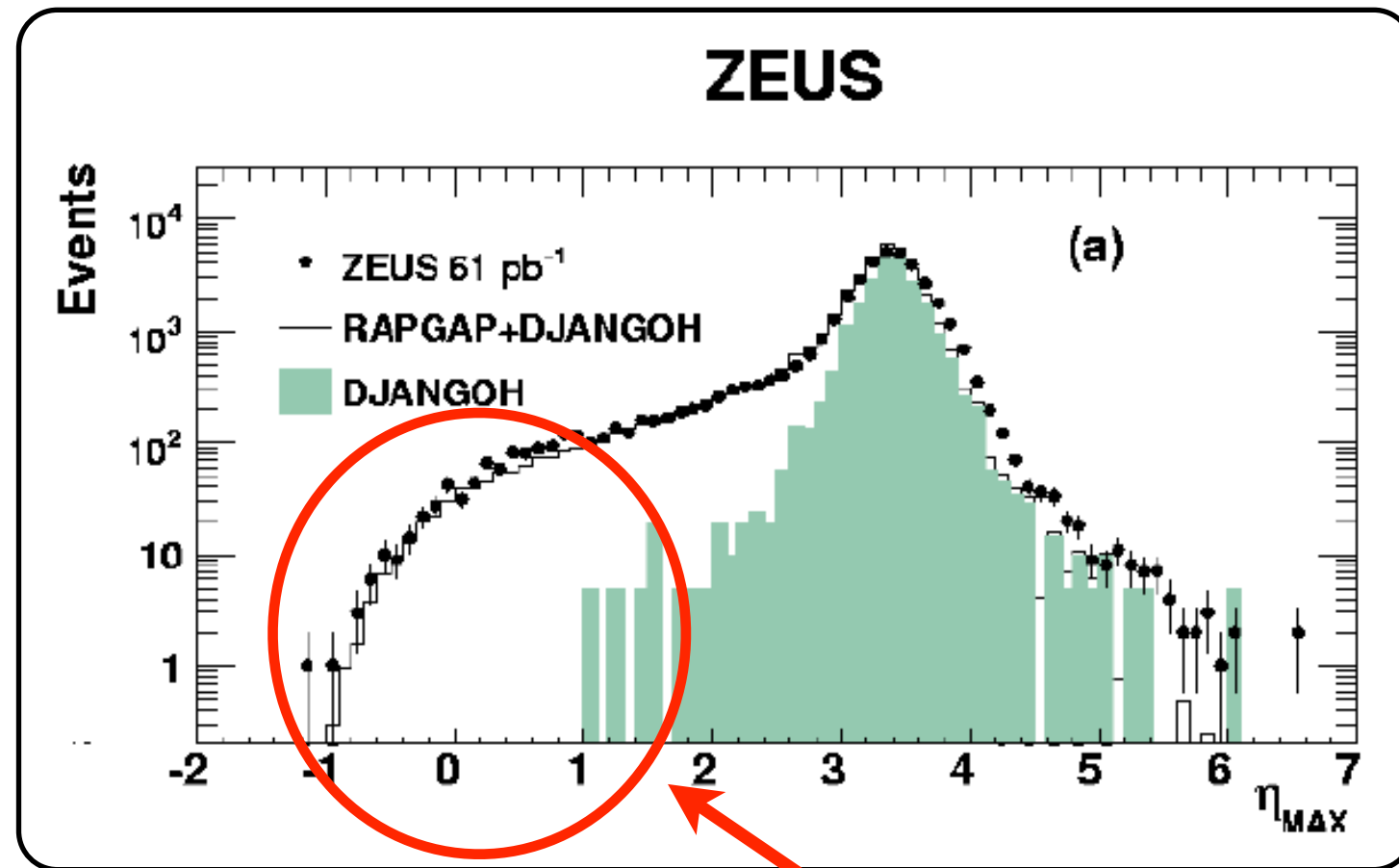
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**Diffractive events**

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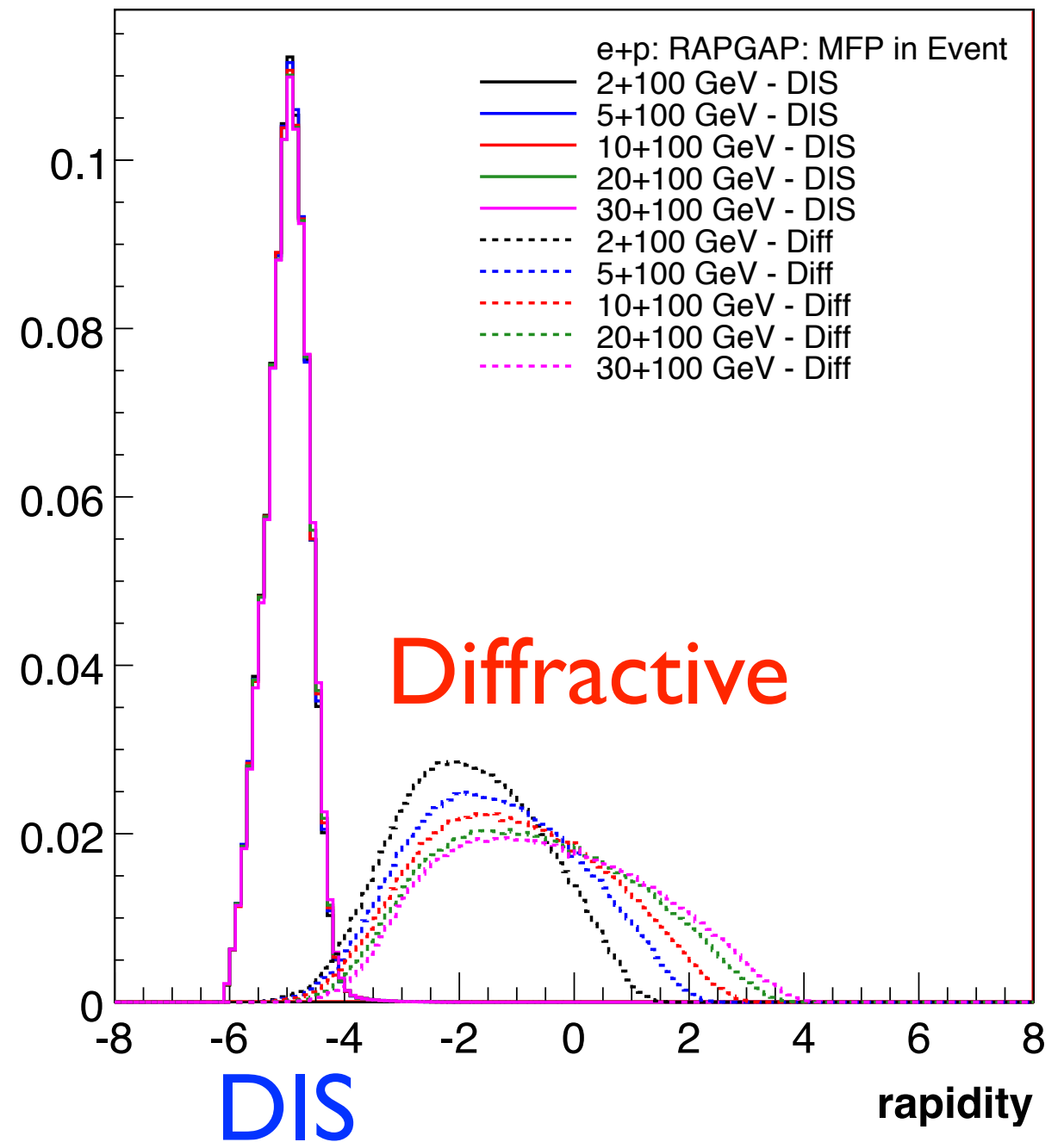
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Can this method  
be used at an EIC?

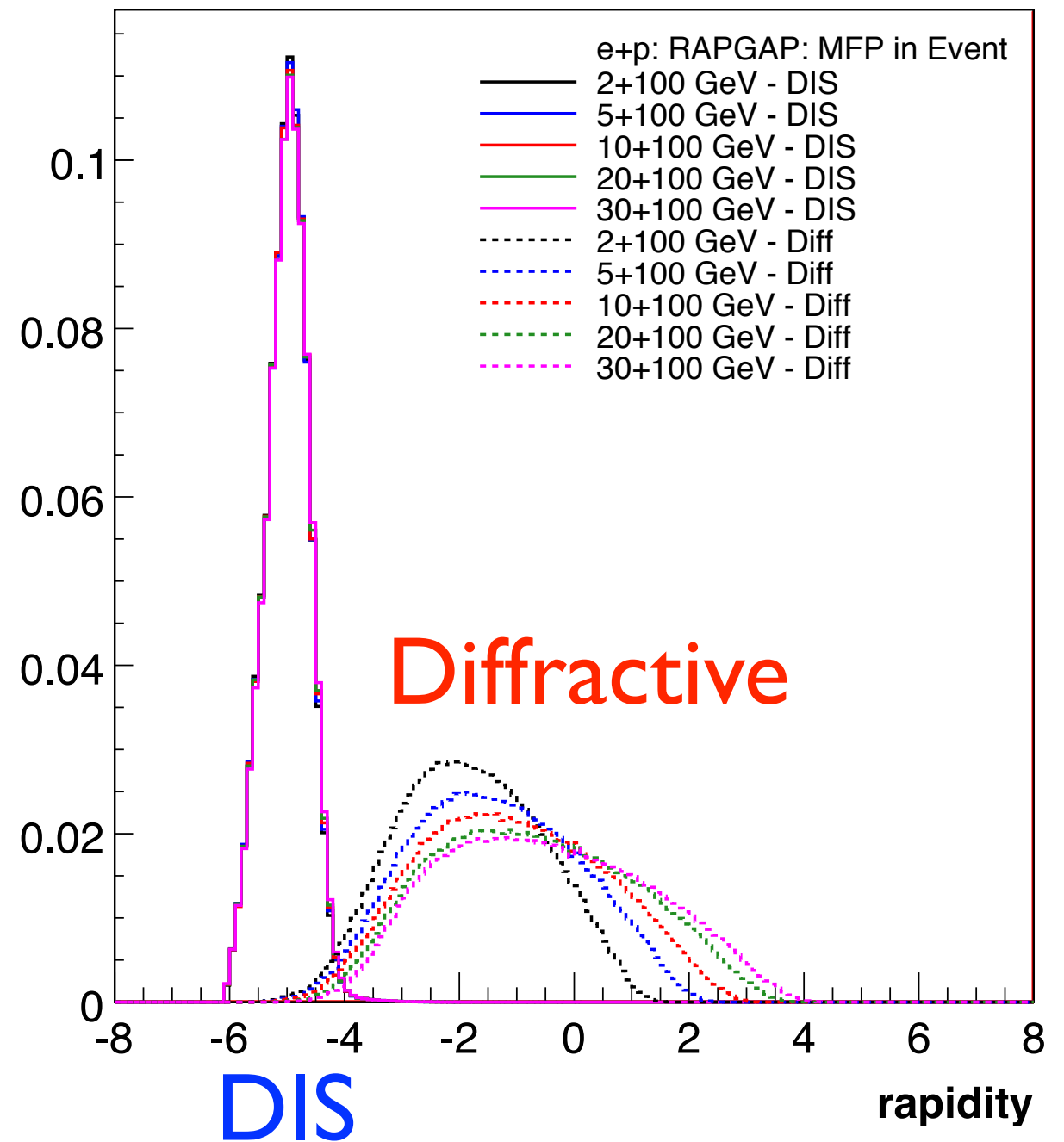


# Large rapidity gaps at eRHIC



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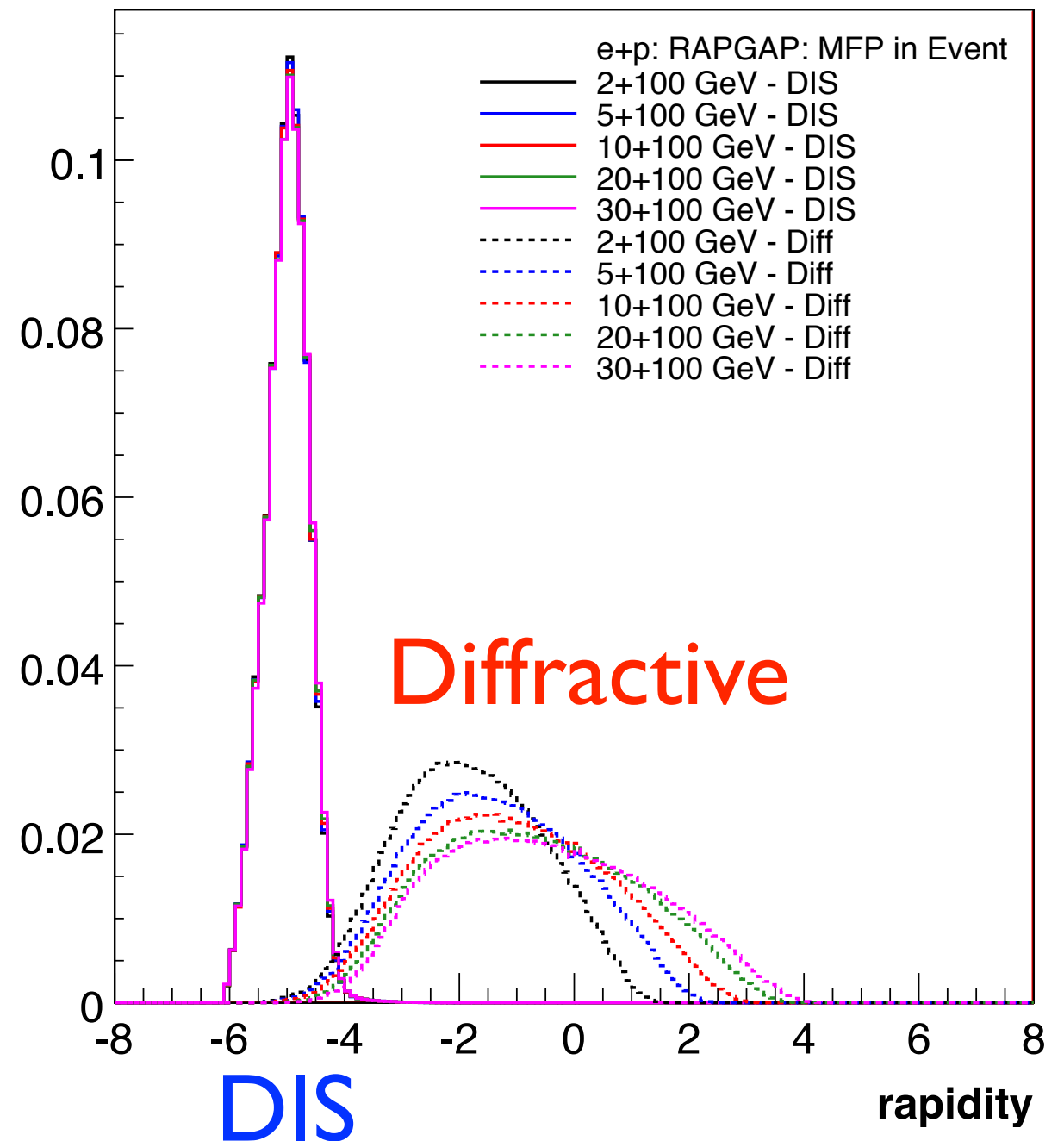
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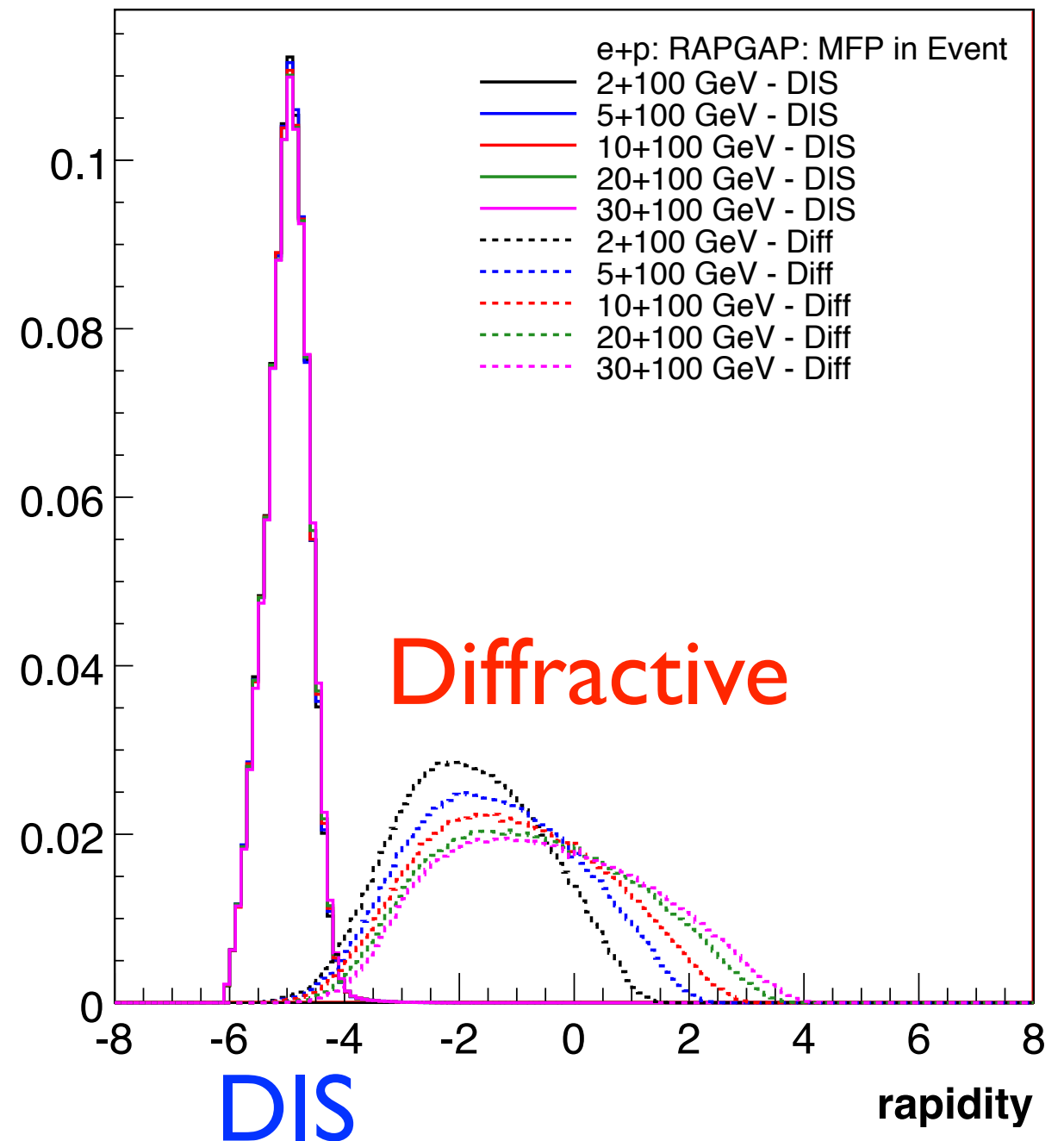
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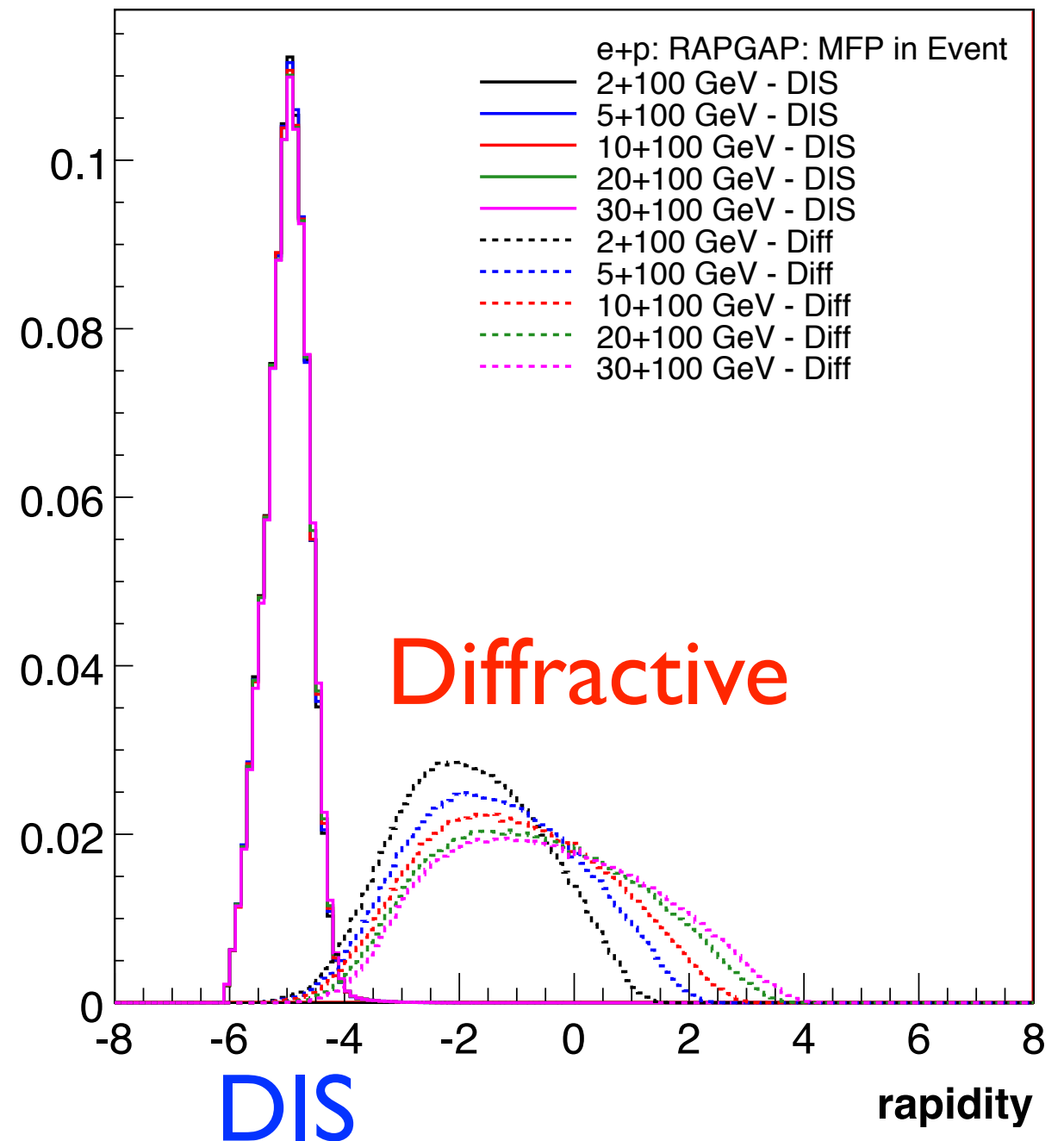
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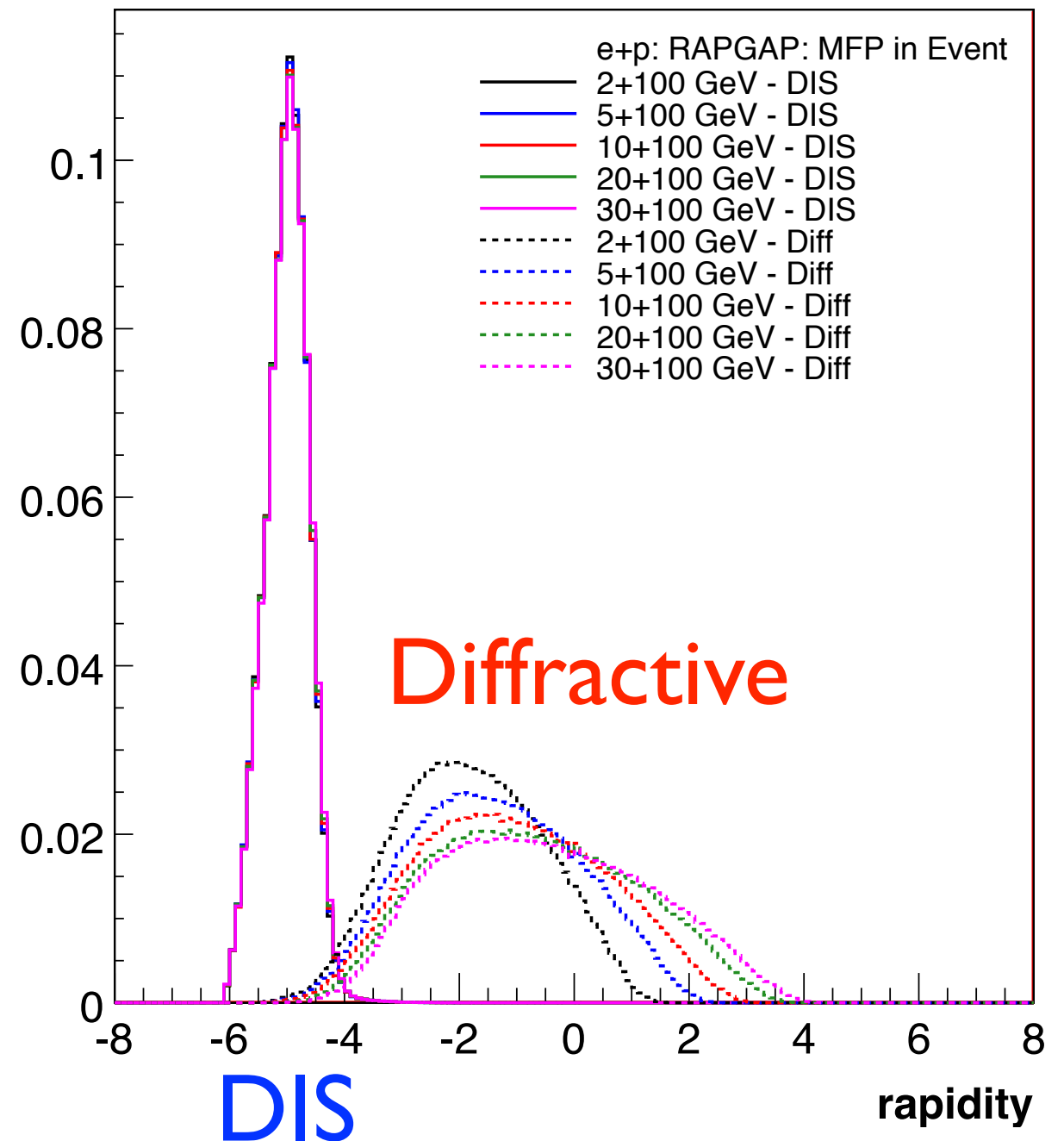
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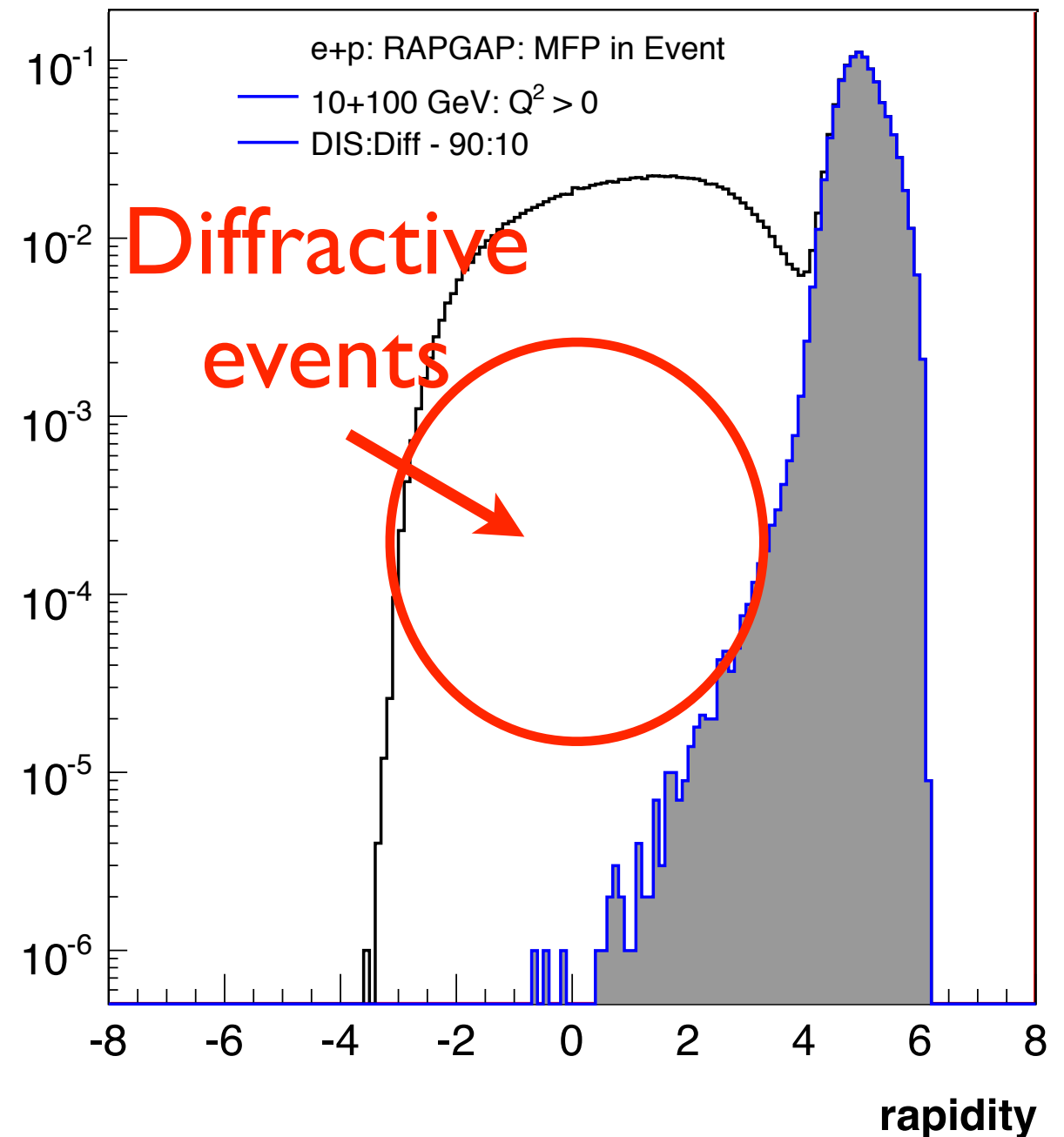




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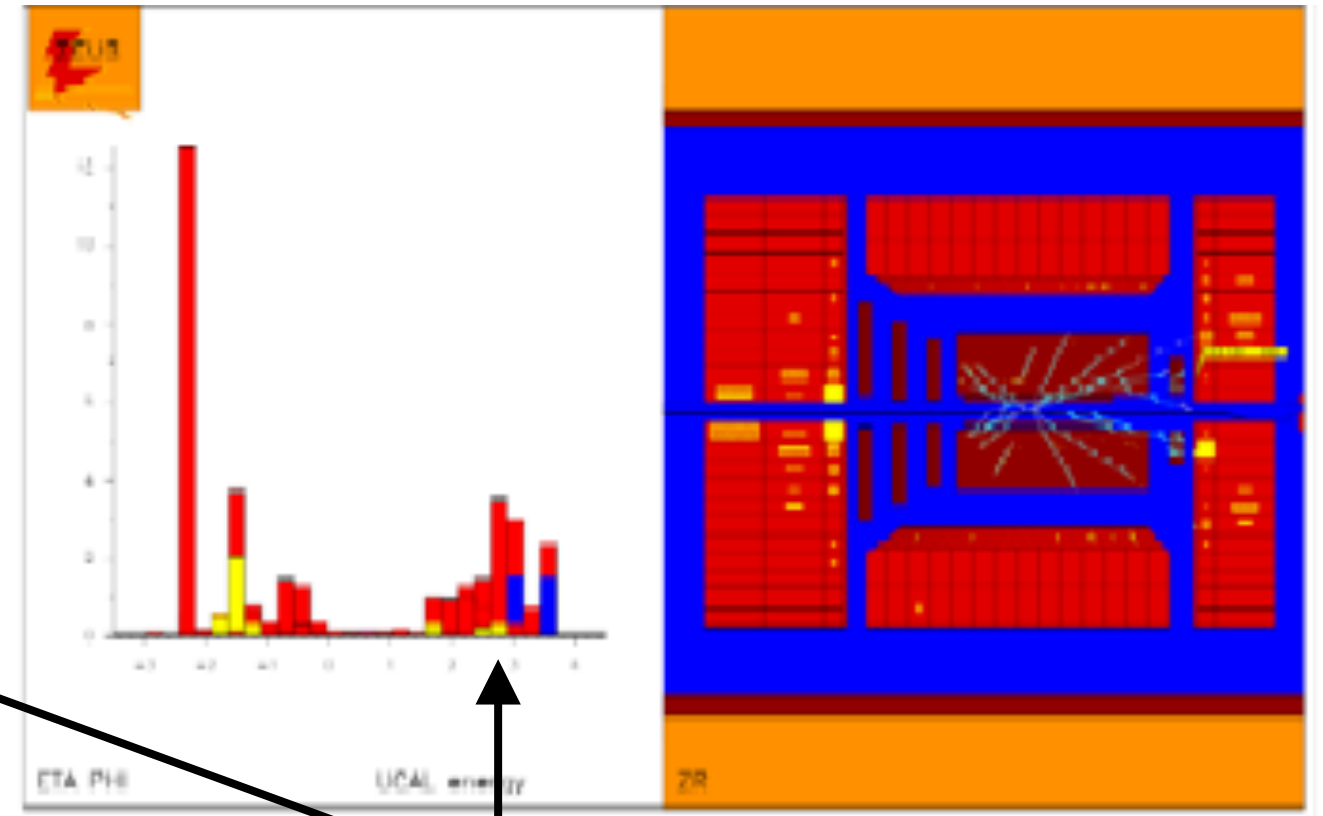
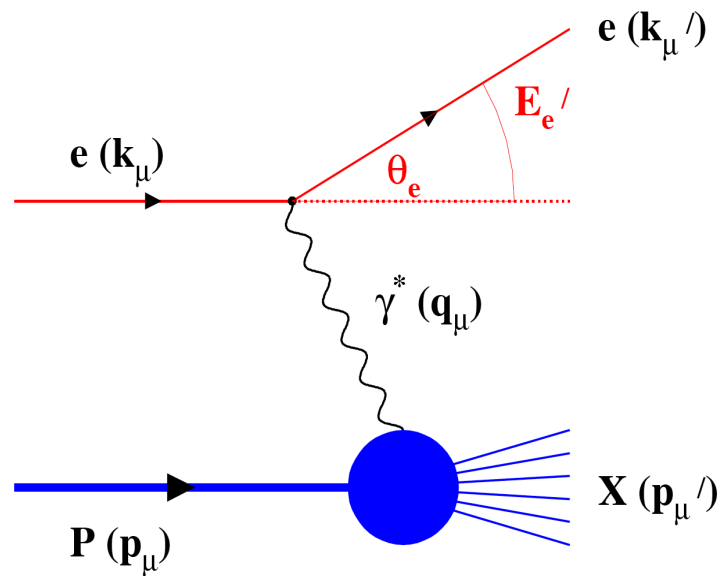
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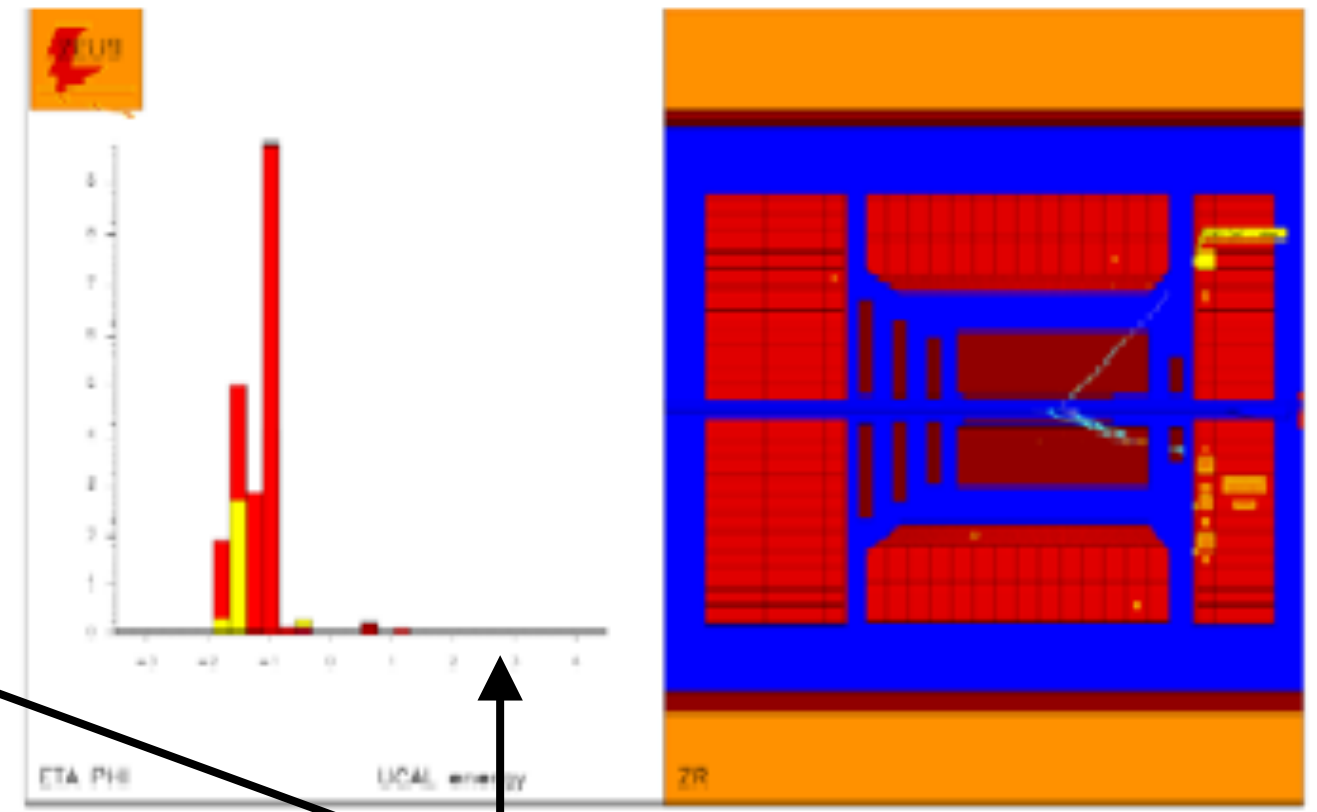
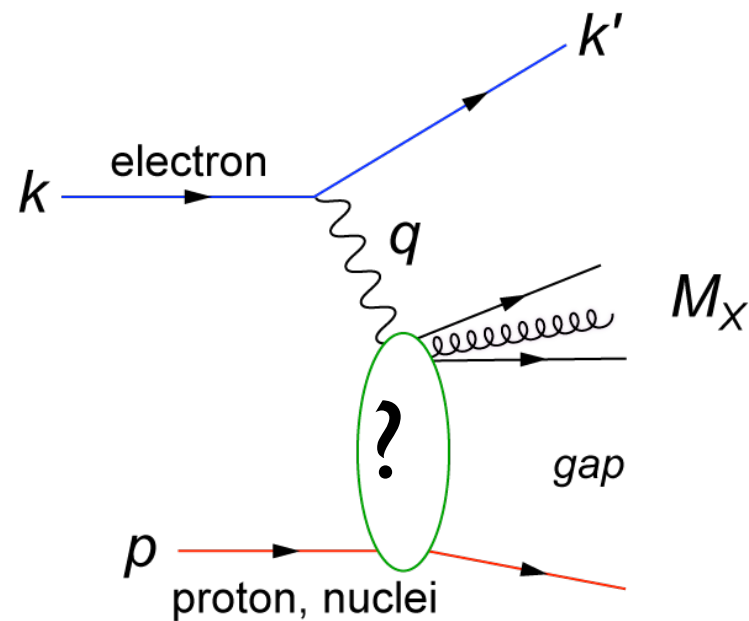
‘Standard DIS event’



Activity in proton direction

# Diffractive Physics in $e+A$

## Diffractive event

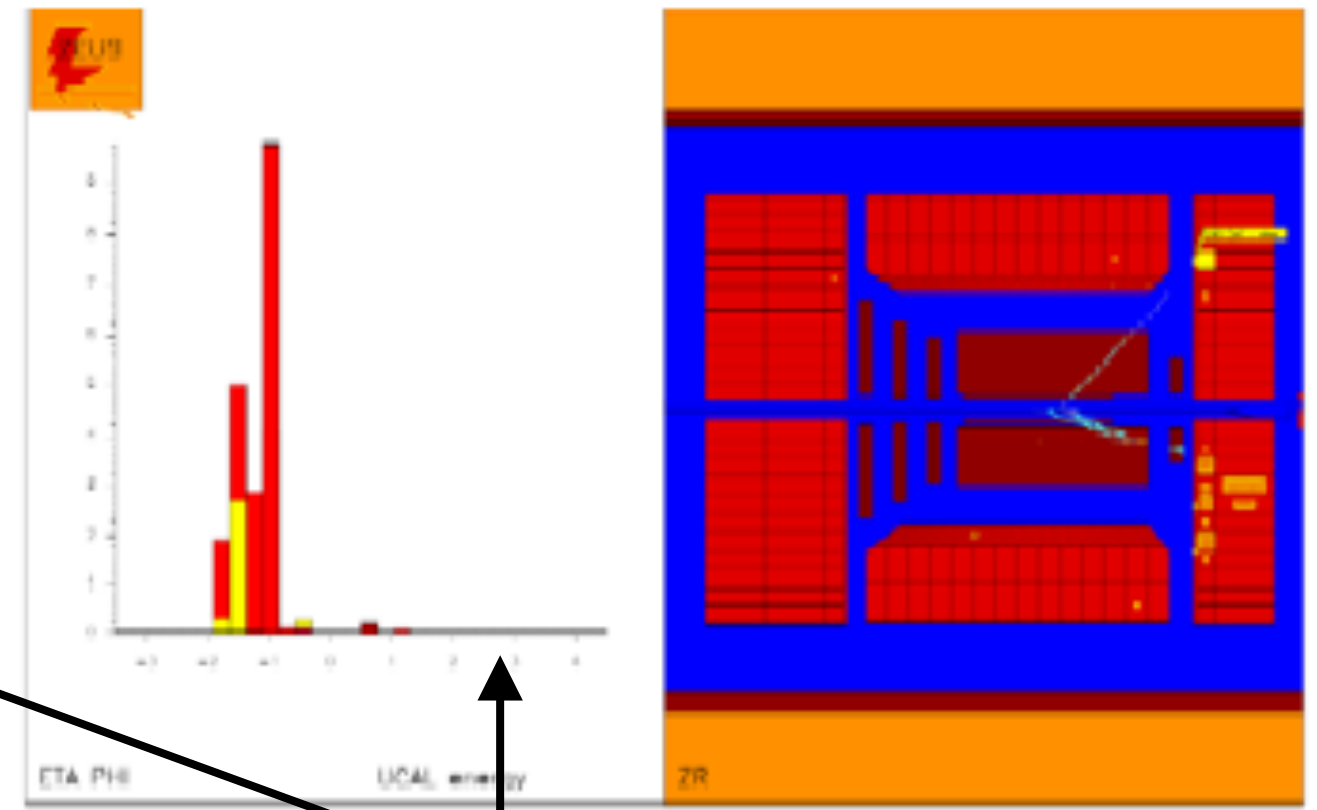
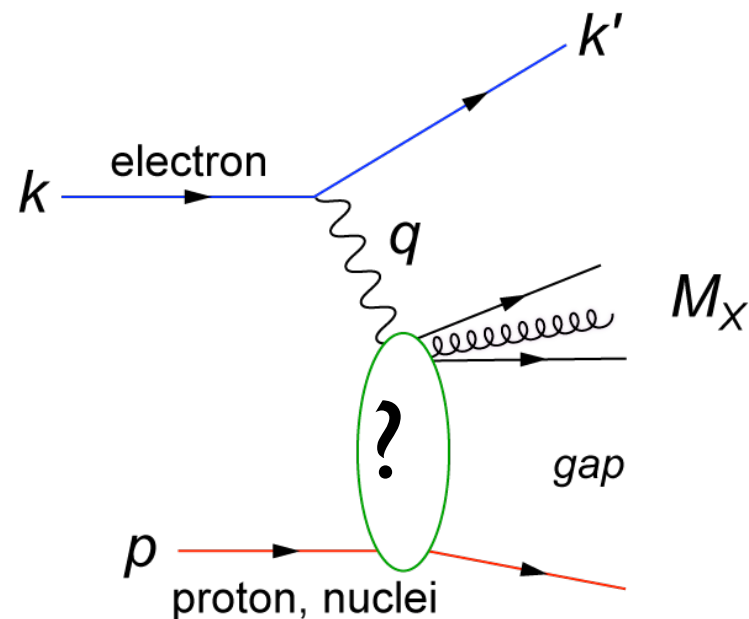


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Activity in proton direction

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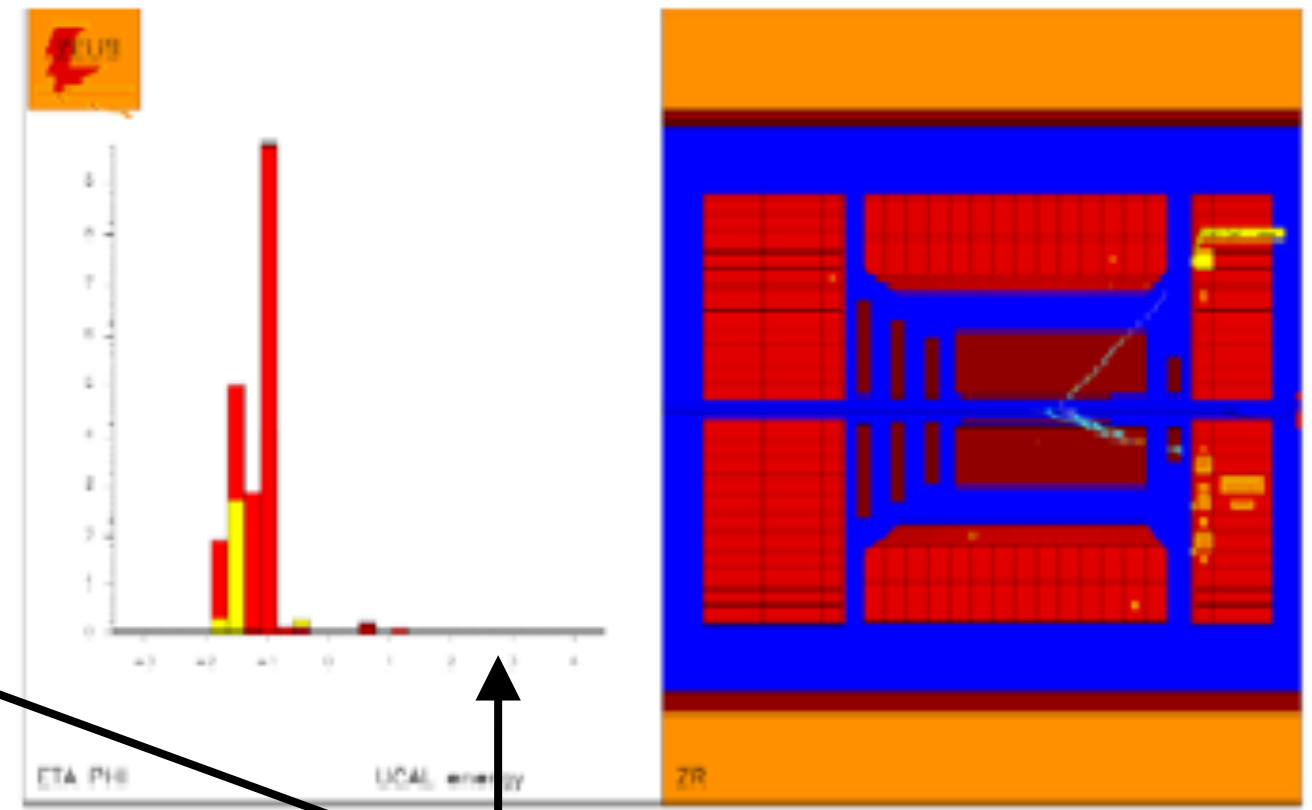
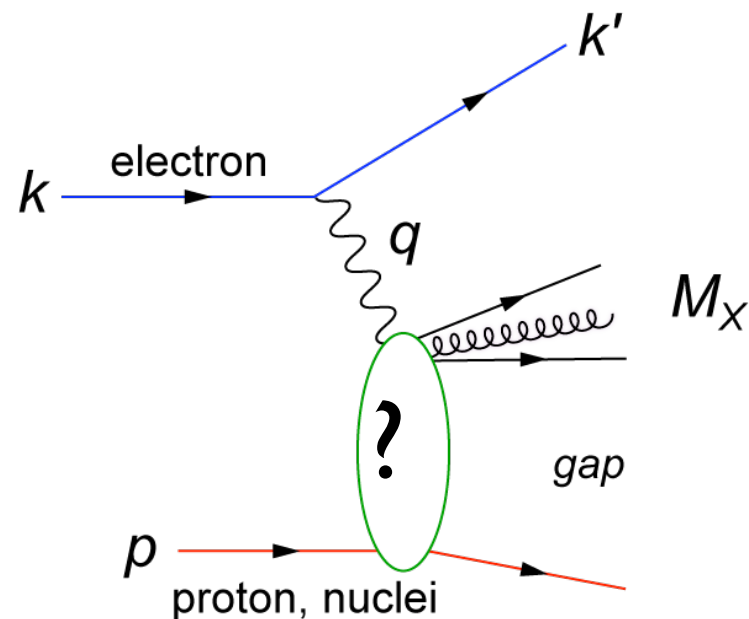


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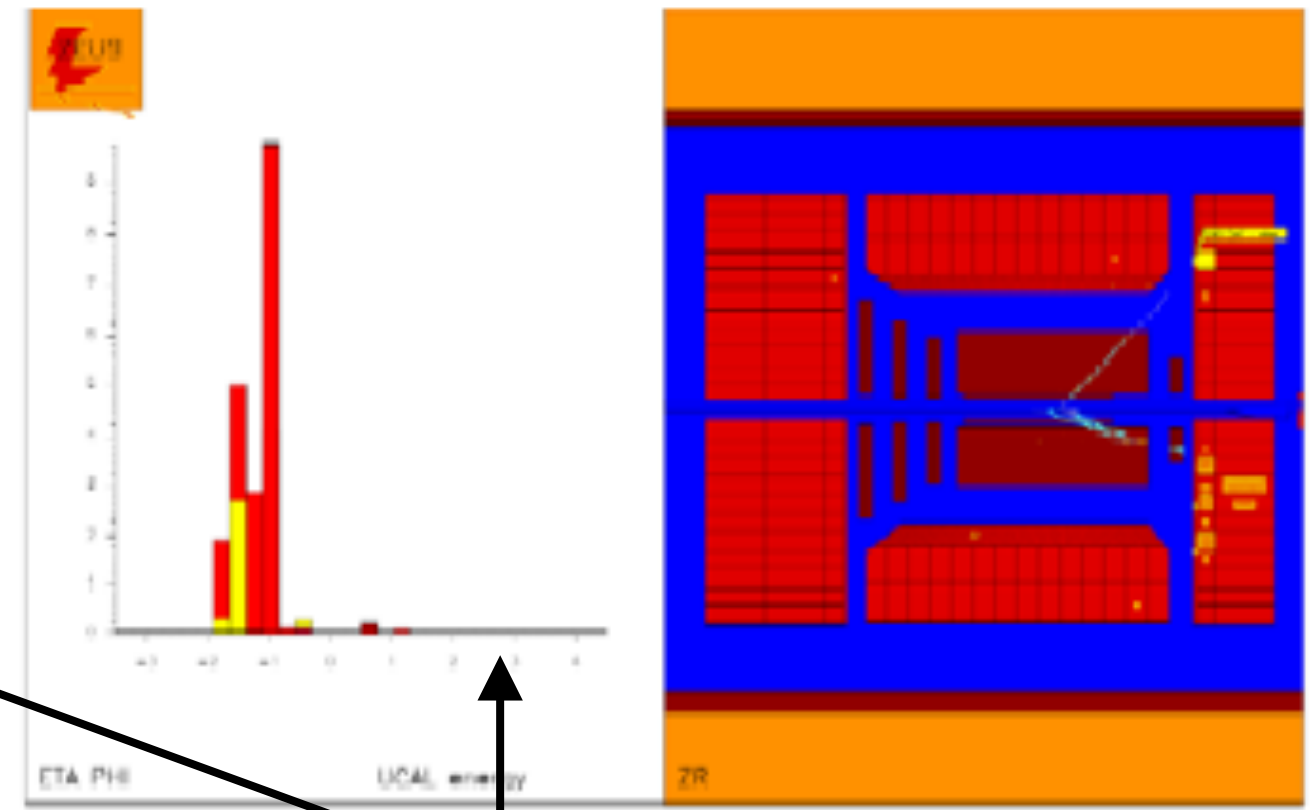
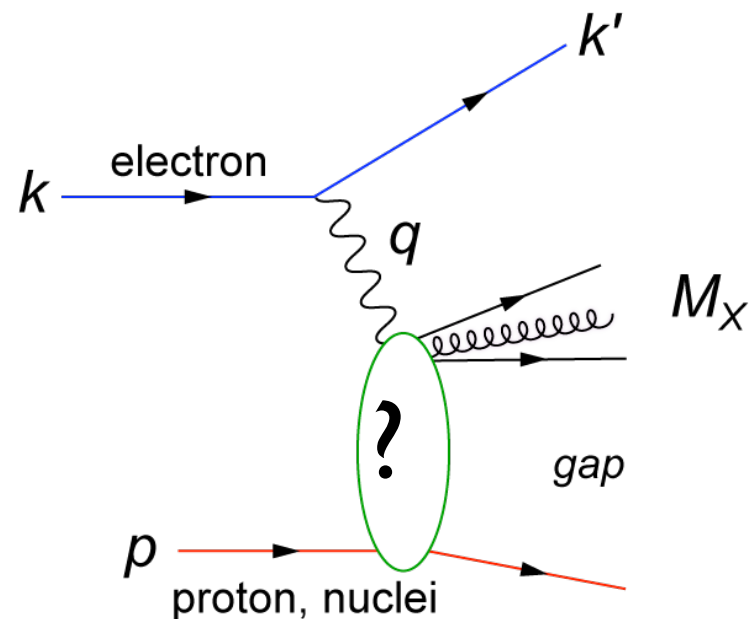


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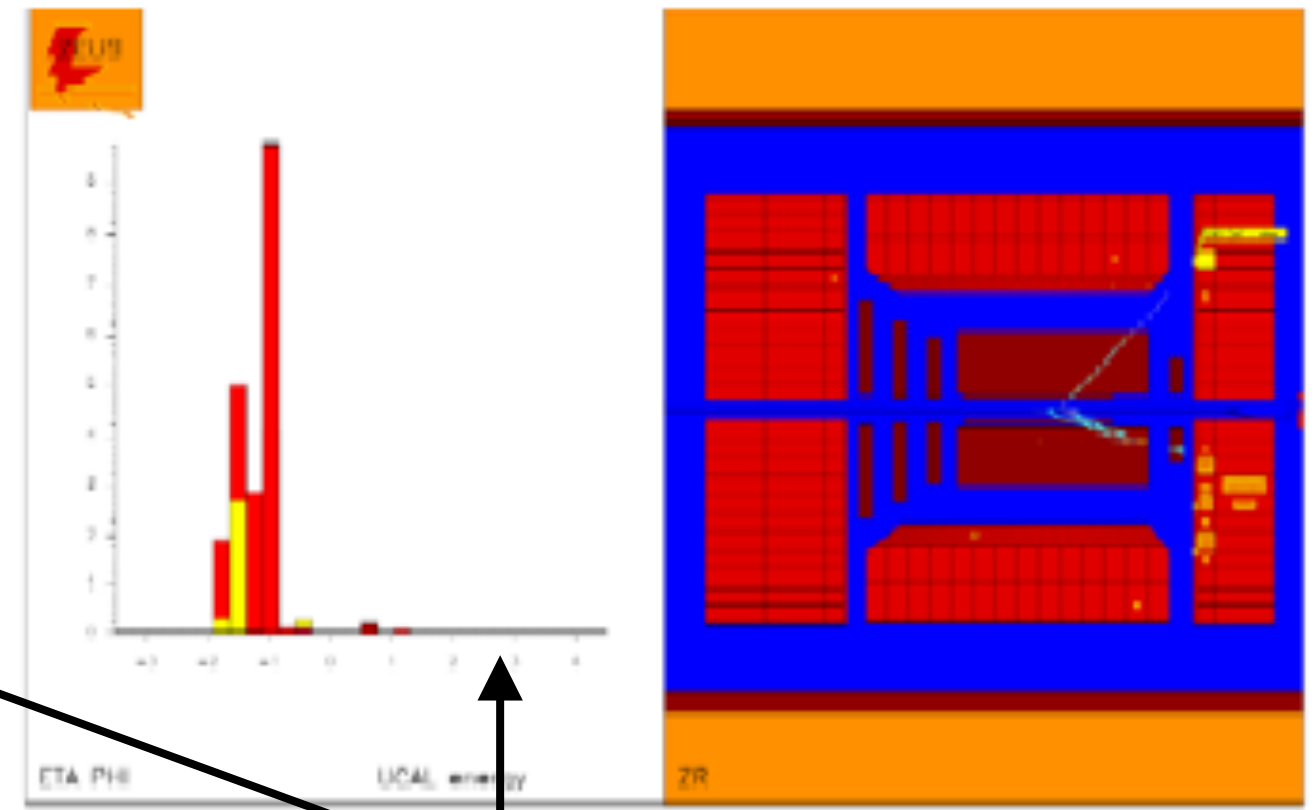
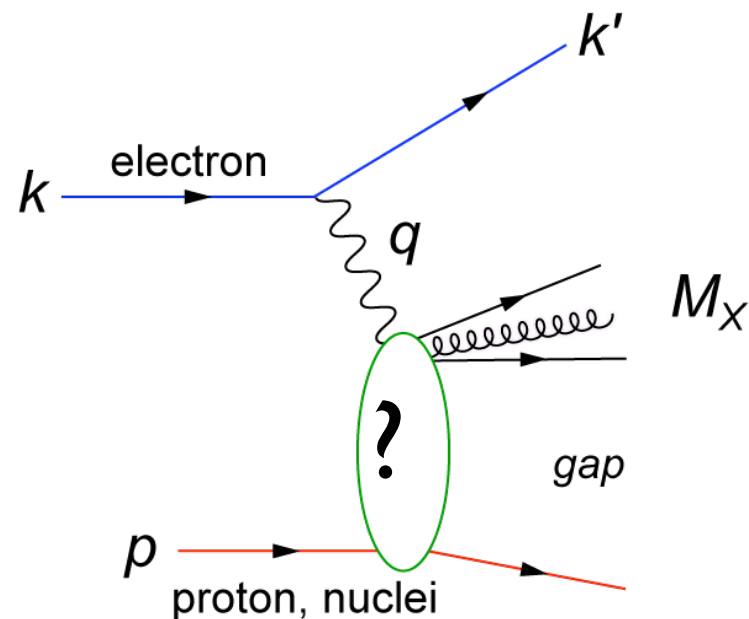
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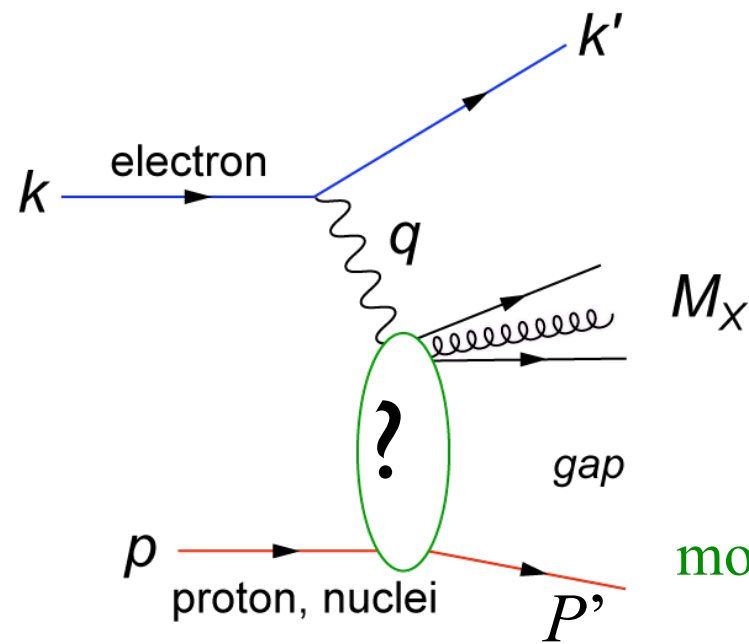


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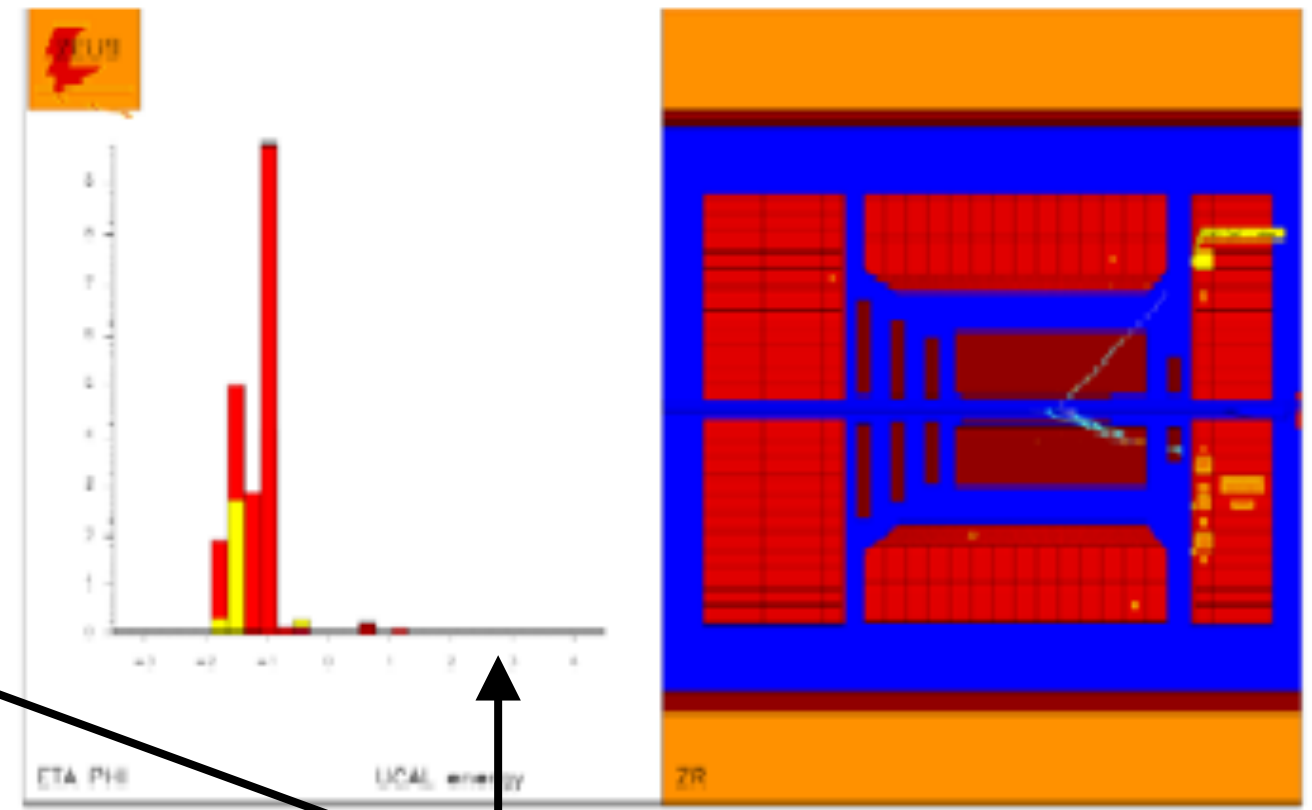
Activity in proton direction

# Diffractive Physics in $e+A$

## Diffractive event



momentum transfer:  
 $t = (P-P')^2$

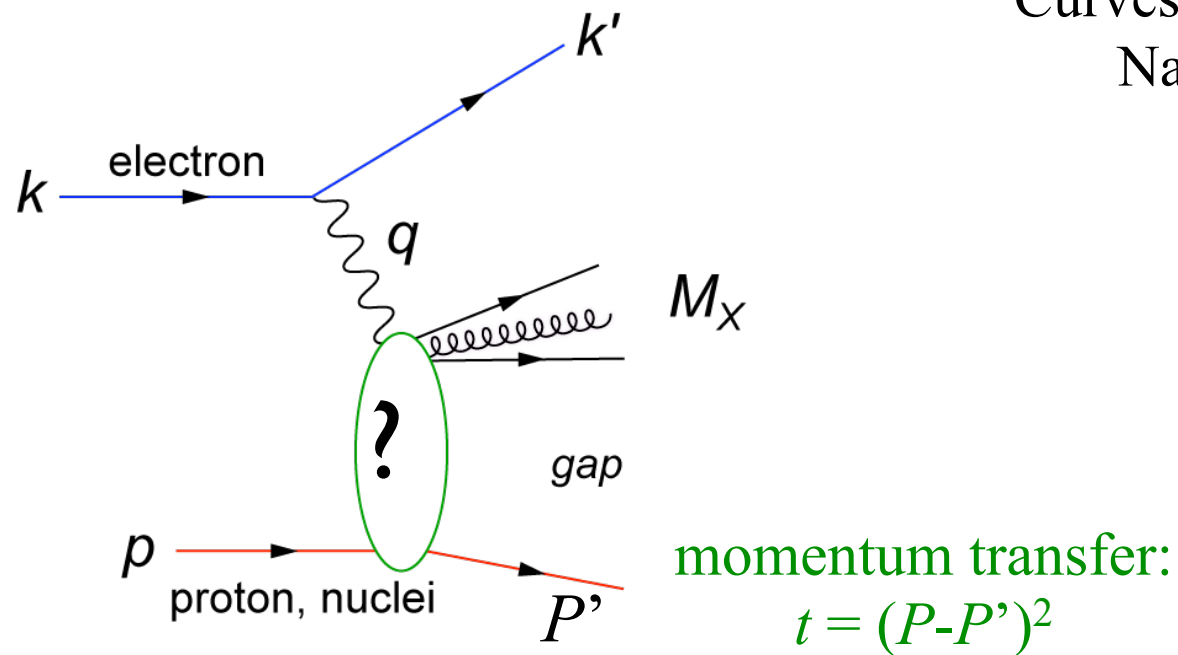


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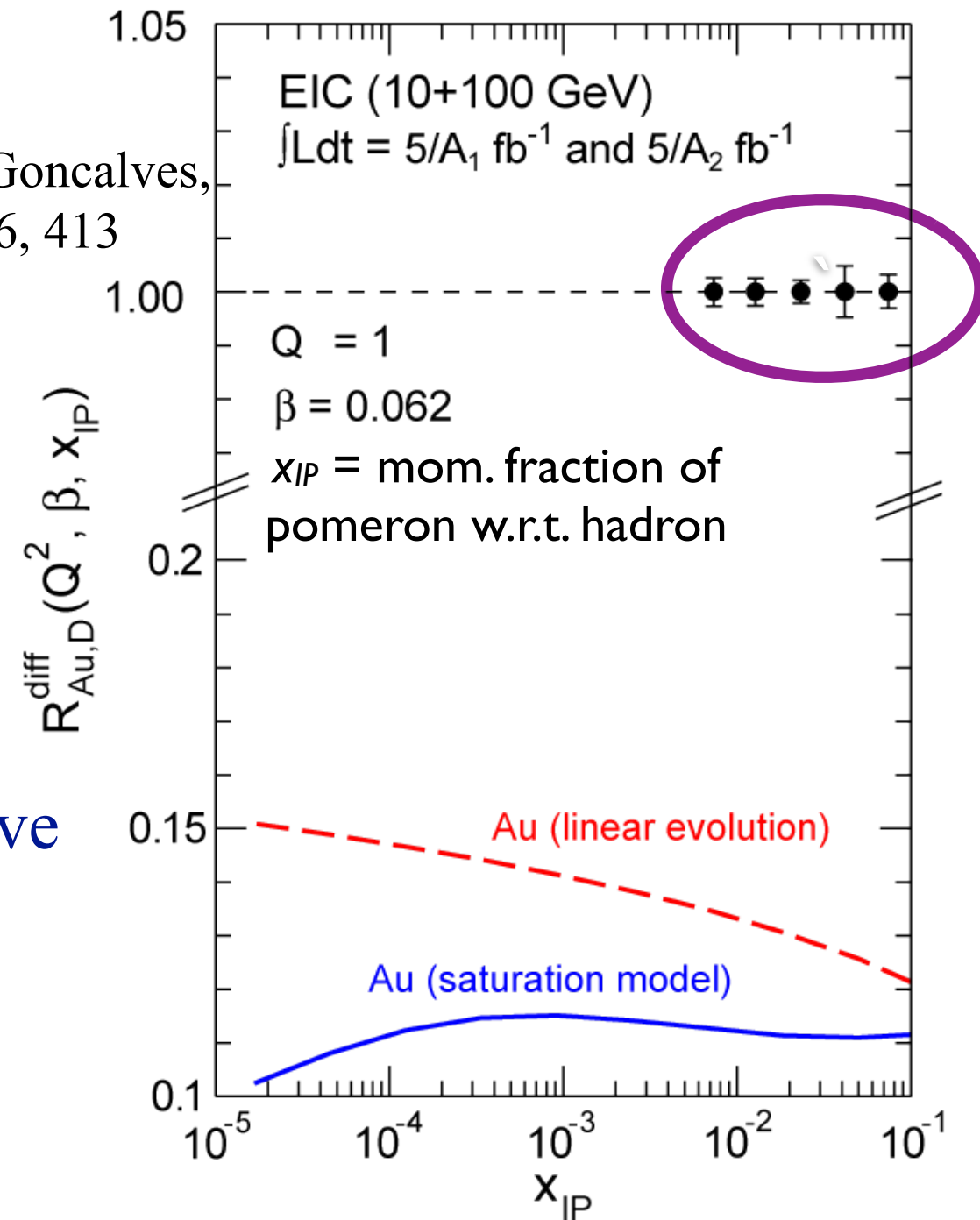
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- Look inside the “Pomeron”
  - ➡ Diffractive structure functions
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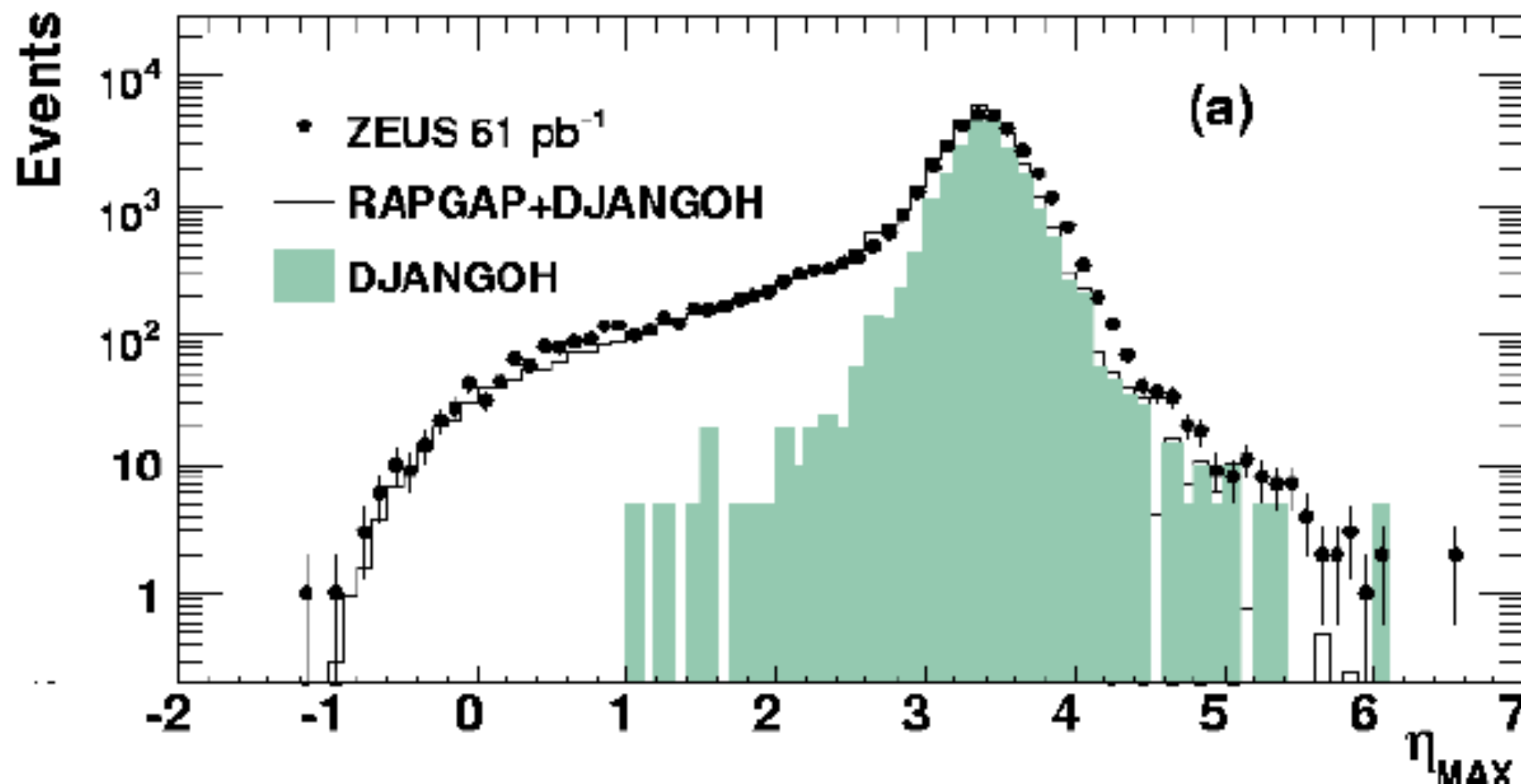
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- Look inside the “Pomeron”  
 ➔ Diffractive structure functions  
 ➔ Diffractive vector meson production:  $d\sigma/dt \sim [xG(x, Q^2)]^2 !!$
- Distinguish between linear evolutions and saturation models



# Diffractive Physics in e+A

- How to measure diffraction in e+A?
  - ➔ Use HERA method of Large Rapidity Gaps
  - ➔ Ideal gap of  $\sim 7.7$  at HERA units reduced to 3-4 due to spread from hadronisation

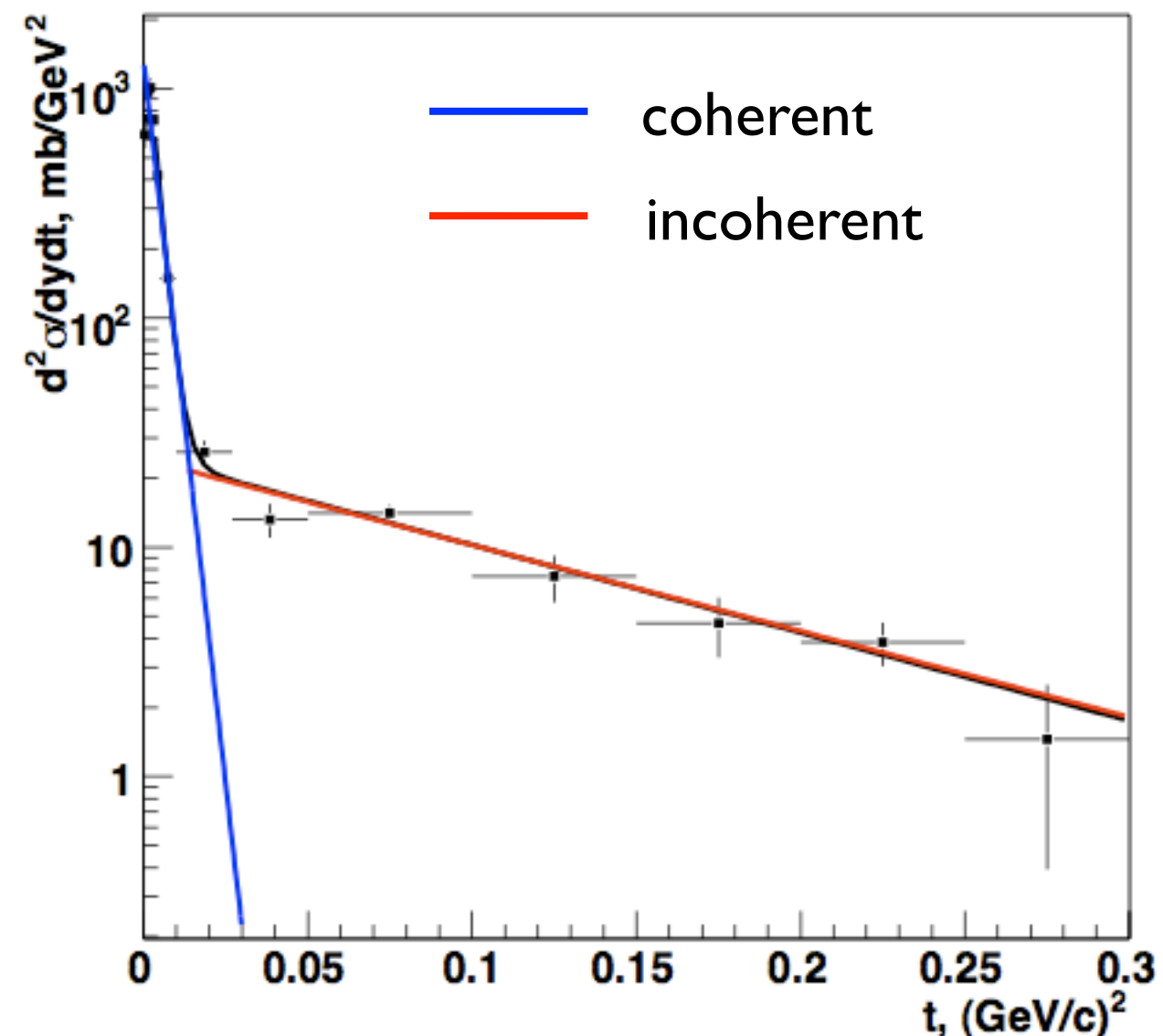
## ZEUS



# Diffraction Physics in e+A

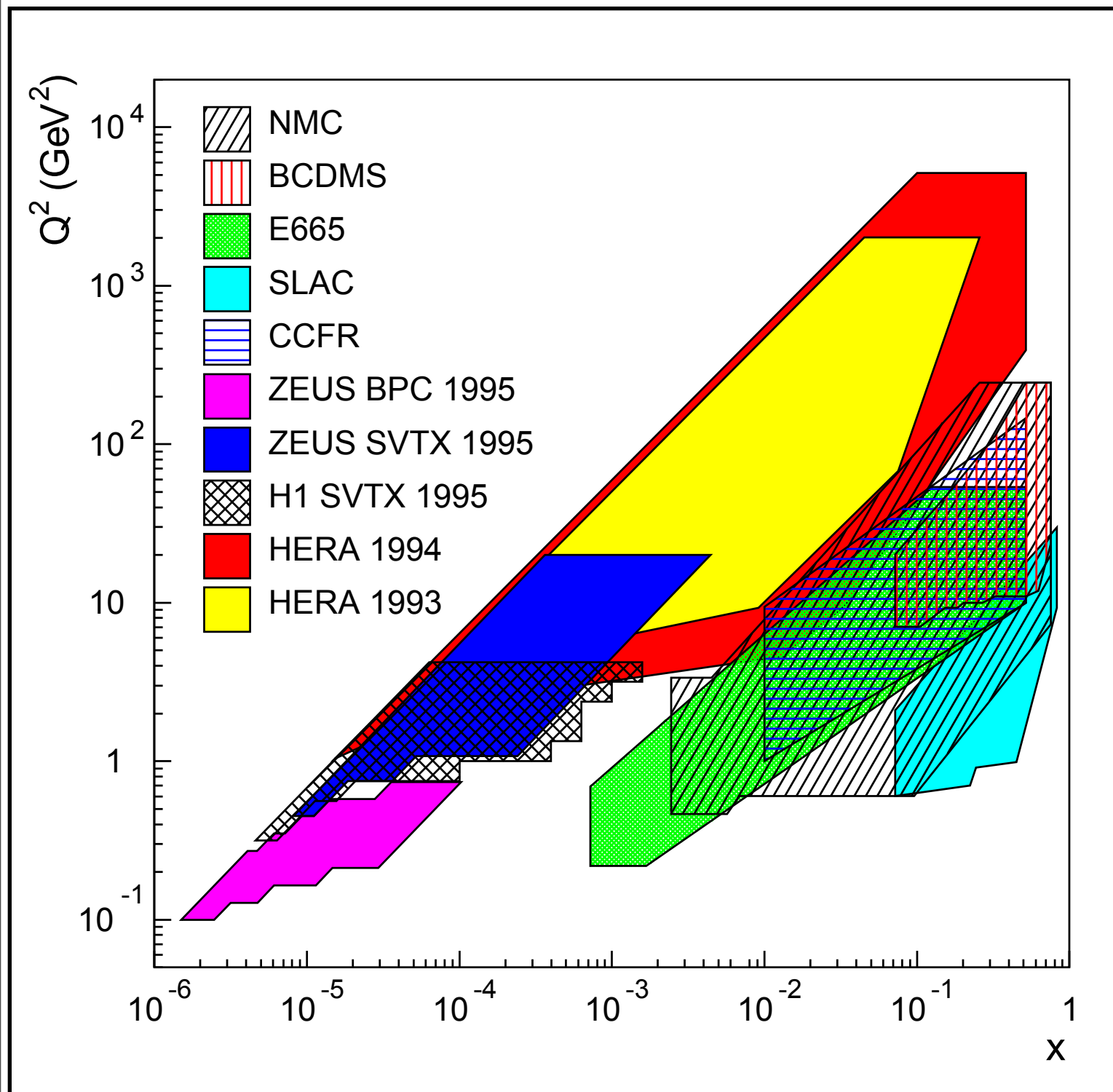
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- Issues with measuring diffractive physics in e+A:
  - ➔  $t$  required for nucleus to break-up is small ( $\sim 30 \text{ MeV}/c^2$ )
  - ➔  $t$  required for nucleus to be measured in detector  $\gg 30 \text{ MeV}/c^2$
  - ➔ To measure  $t$  dependence, must measure exclusive diffraction (e.g. vector mesons -  $t \sim p_T^2$ )

## STAR - UPC Collisions

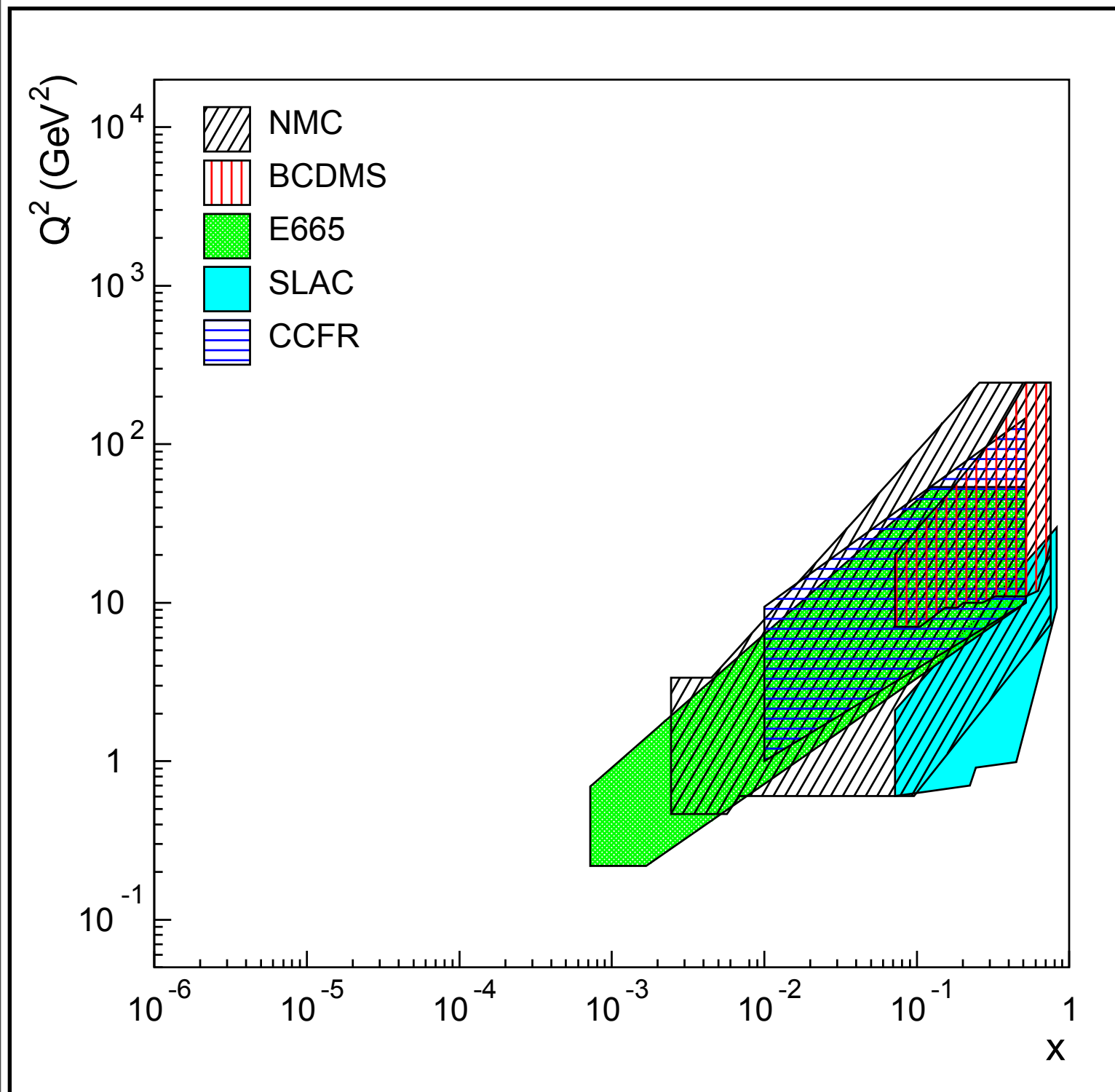


# Requirements for an Electron Ion Collider

Well mapped in  $e+p$



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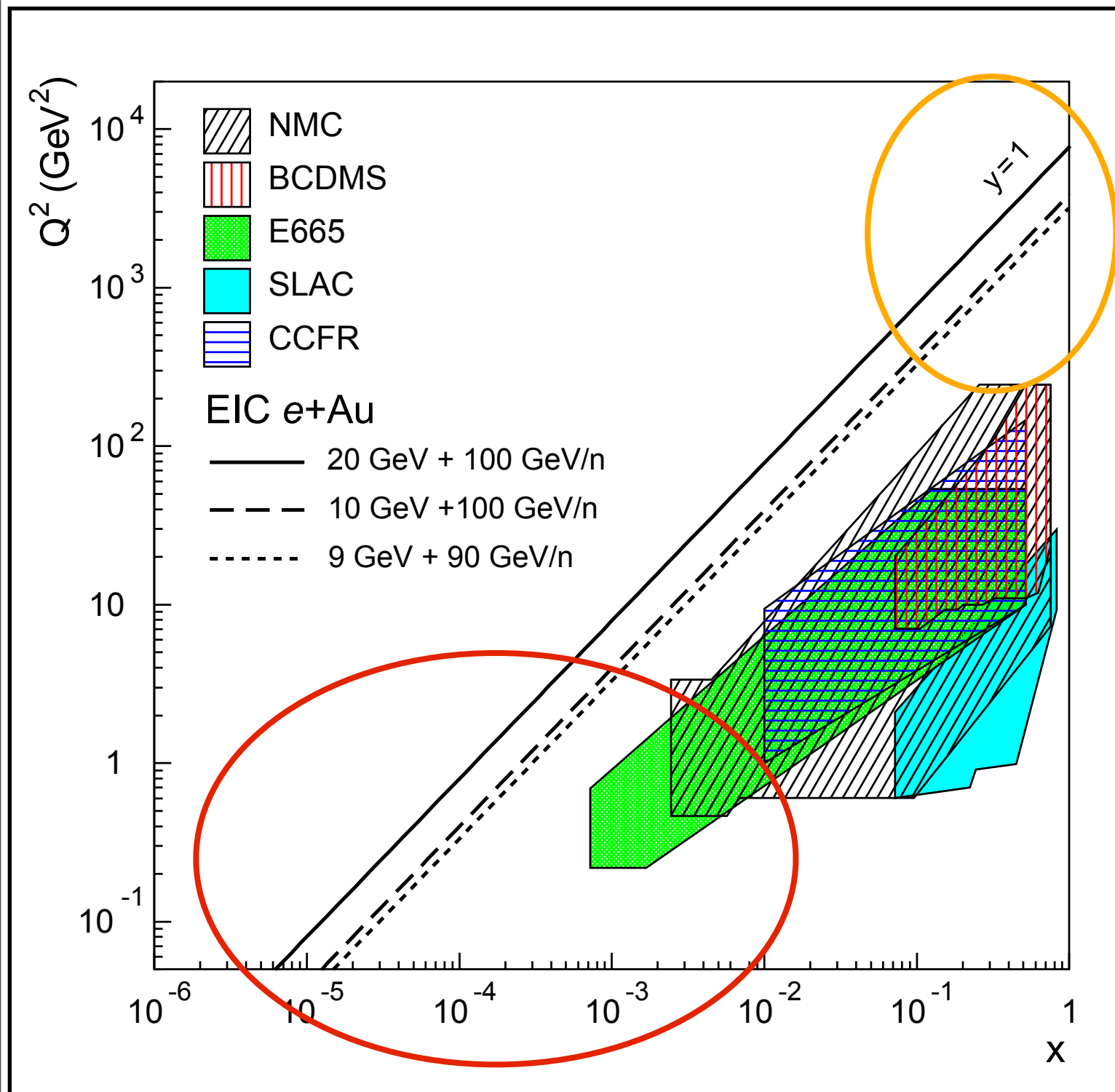
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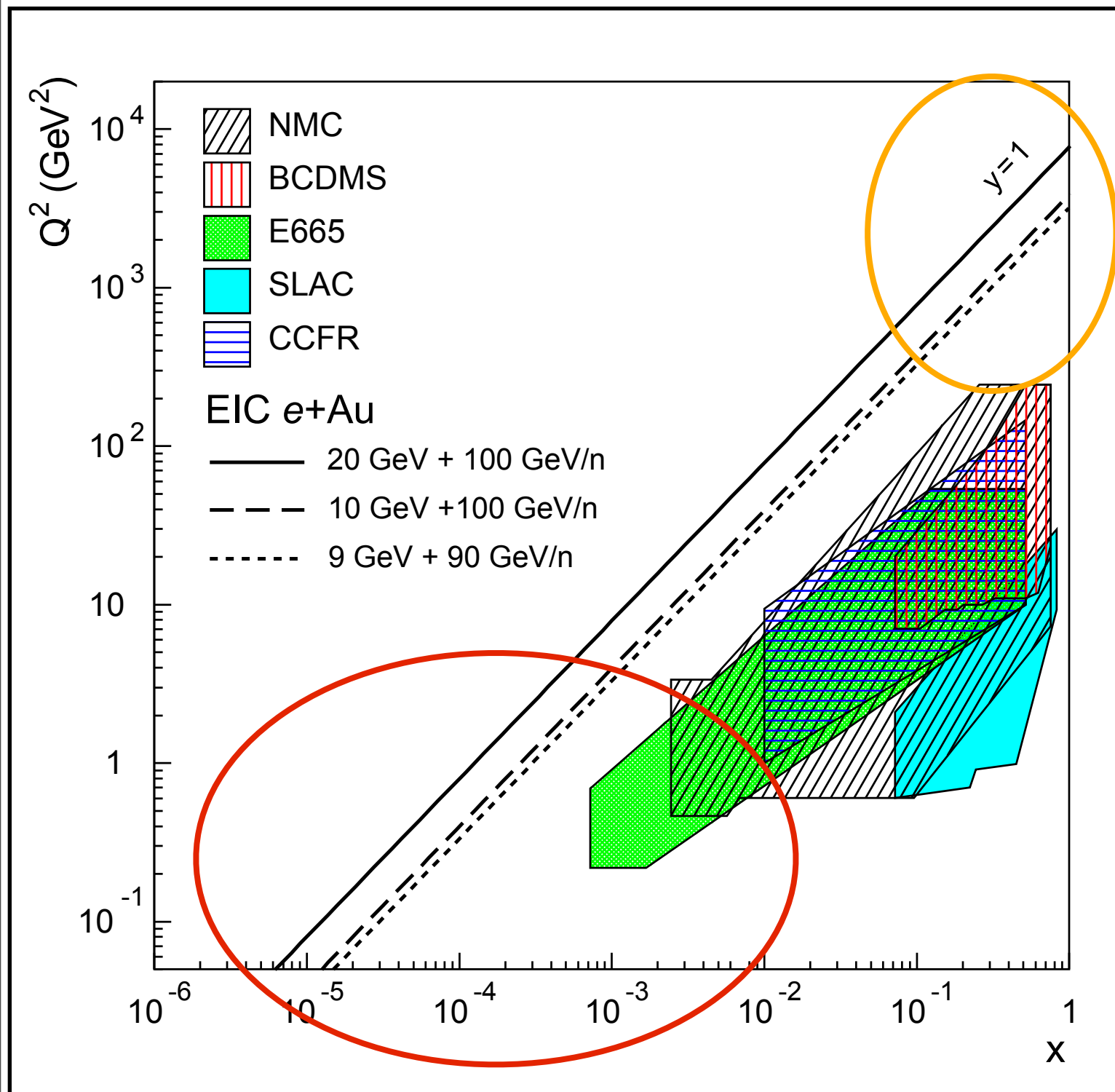
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## Electron Ion Collider:

- $\mathcal{L}(\text{EIC}) > 100 \times \mathcal{L}(\text{HERA})$
- Electrons
  - $E_e = 3 - 20$  GeV
  - polarized
- Hadron Beams
  - $E_A = 100$  GeV
  - $A = p \rightarrow U$
  - polarized  $p$  & light ions

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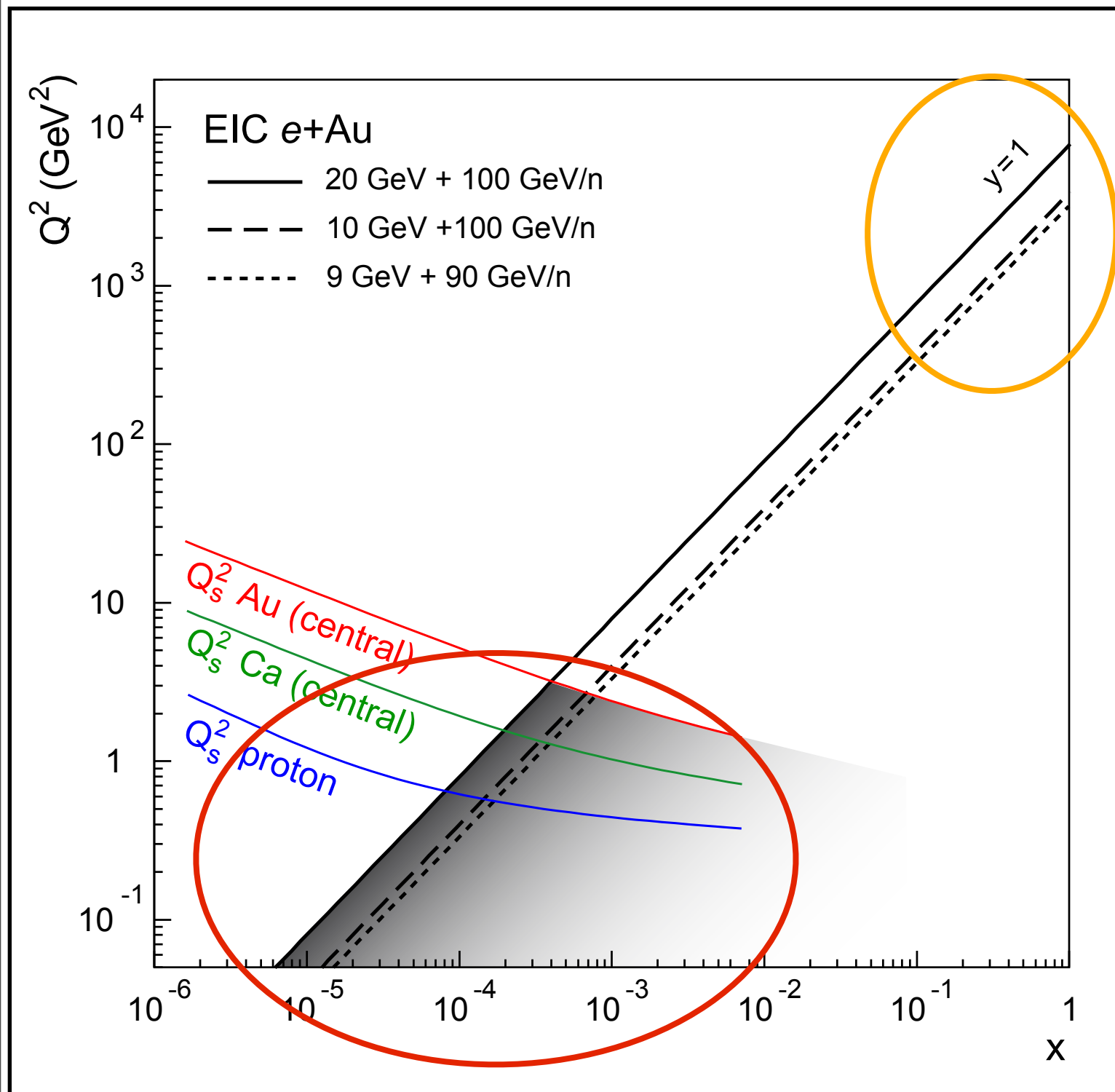
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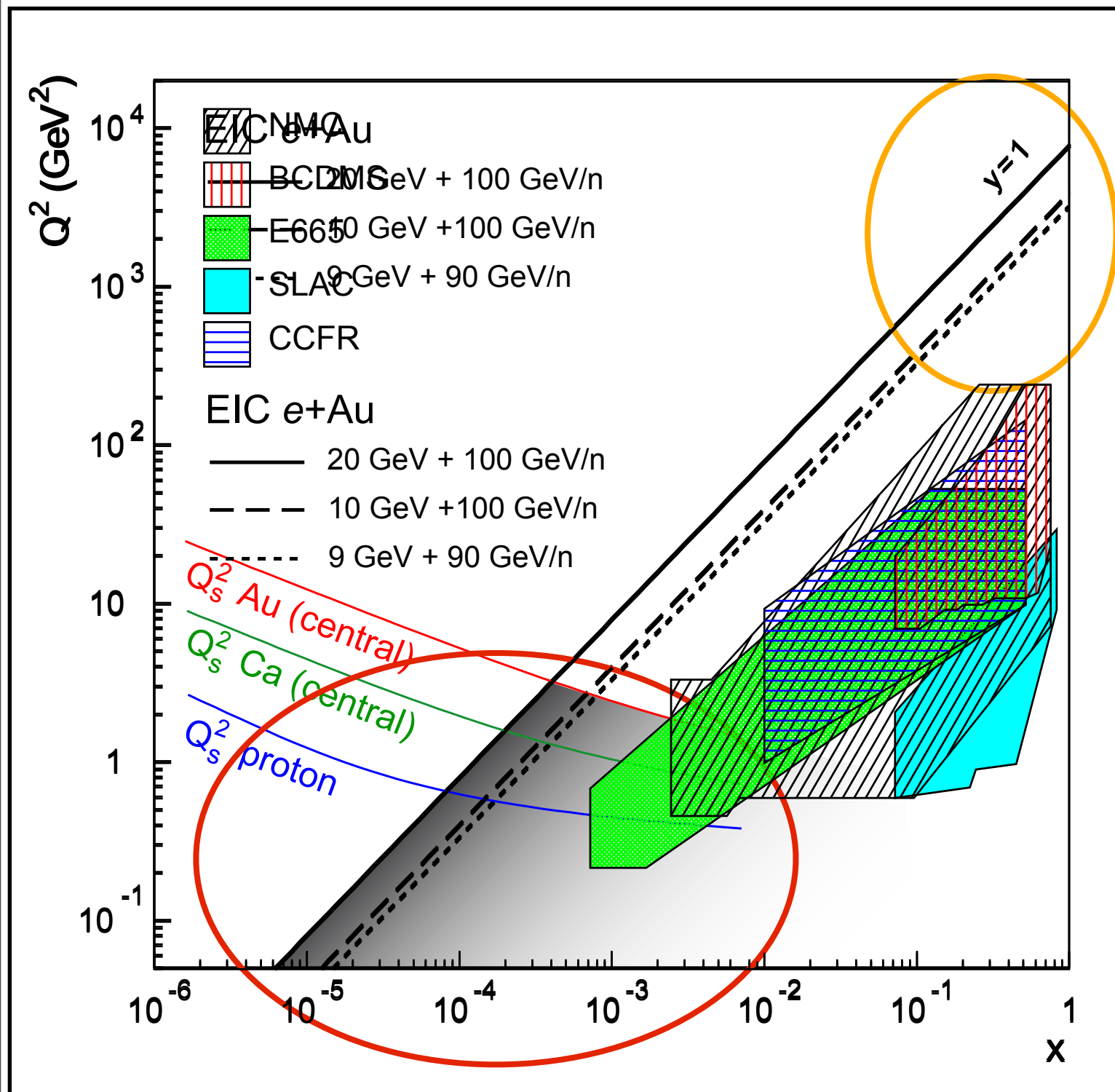
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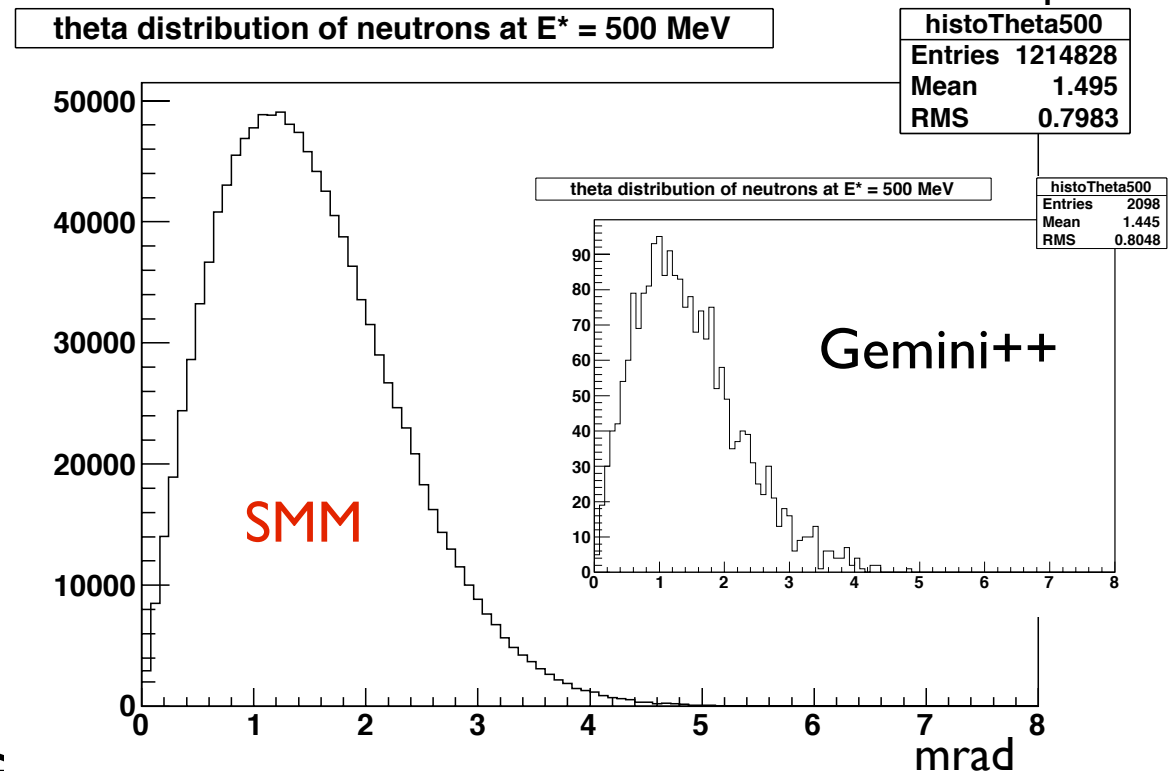
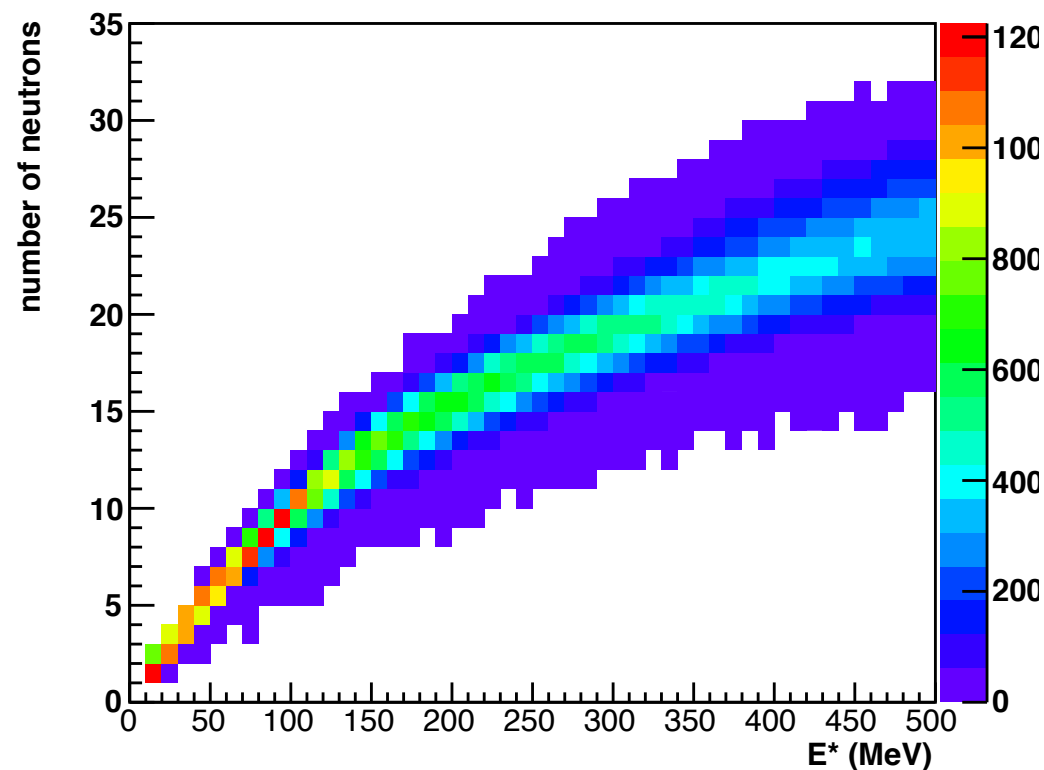
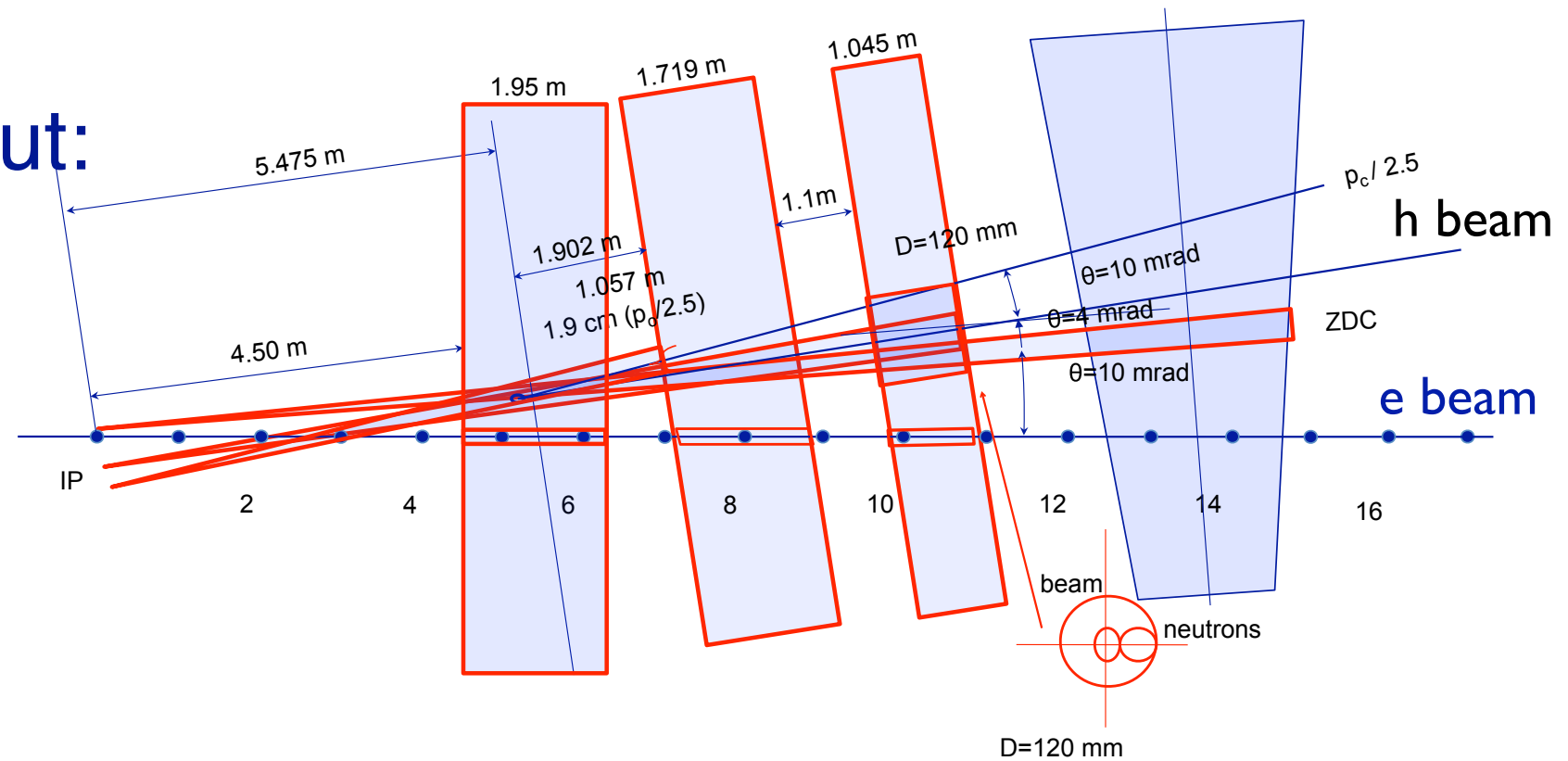
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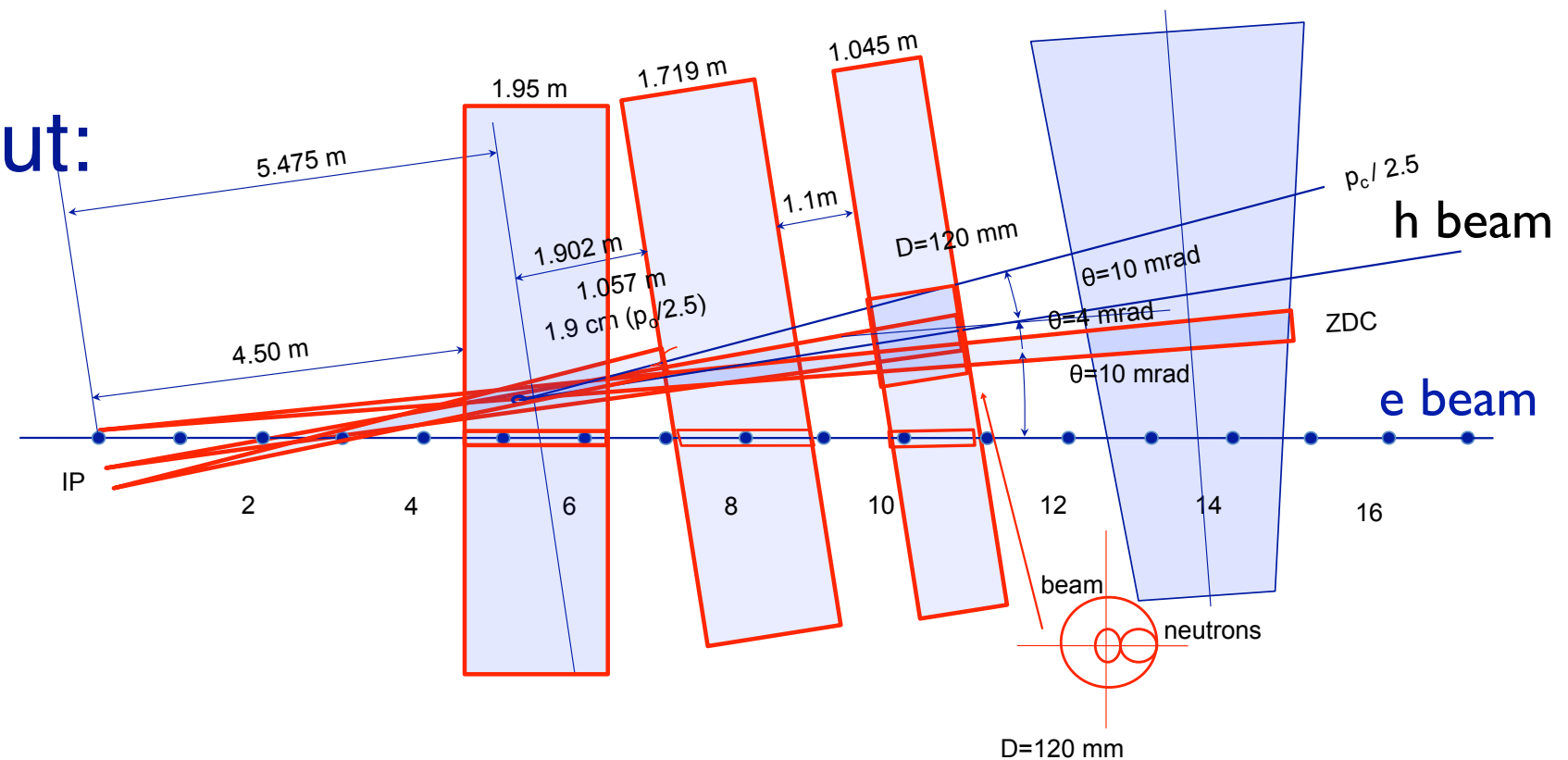
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- Here eRHIC IR layout:
- Need  $\pm X$  mrad opening
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All breakup simulators show **it works**:

- For  $E_{\text{tot}}^* \geq 10$  MeV and 2.5 mrad  $n$  acceptance we have rejection power of at least  $10^5$ .
- Separating incoherent from coherent diffractive events is possible at a collider with  $n$ -detection via ZDCs alone



# Detecting Nuclear Breakup

➔ Detecting **all** fragments  $p_{A'} = \sum p_n + \sum p_p + \sum p_d + \sum p_\alpha \dots$  not possible

➔ Focus on n emission

▶ Zero-Degree Calorimeter

▶ Requires careful design of IR

• Additional measurements:

▶ Fragments via Roman Pots

▶  $\gamma$  via EMC

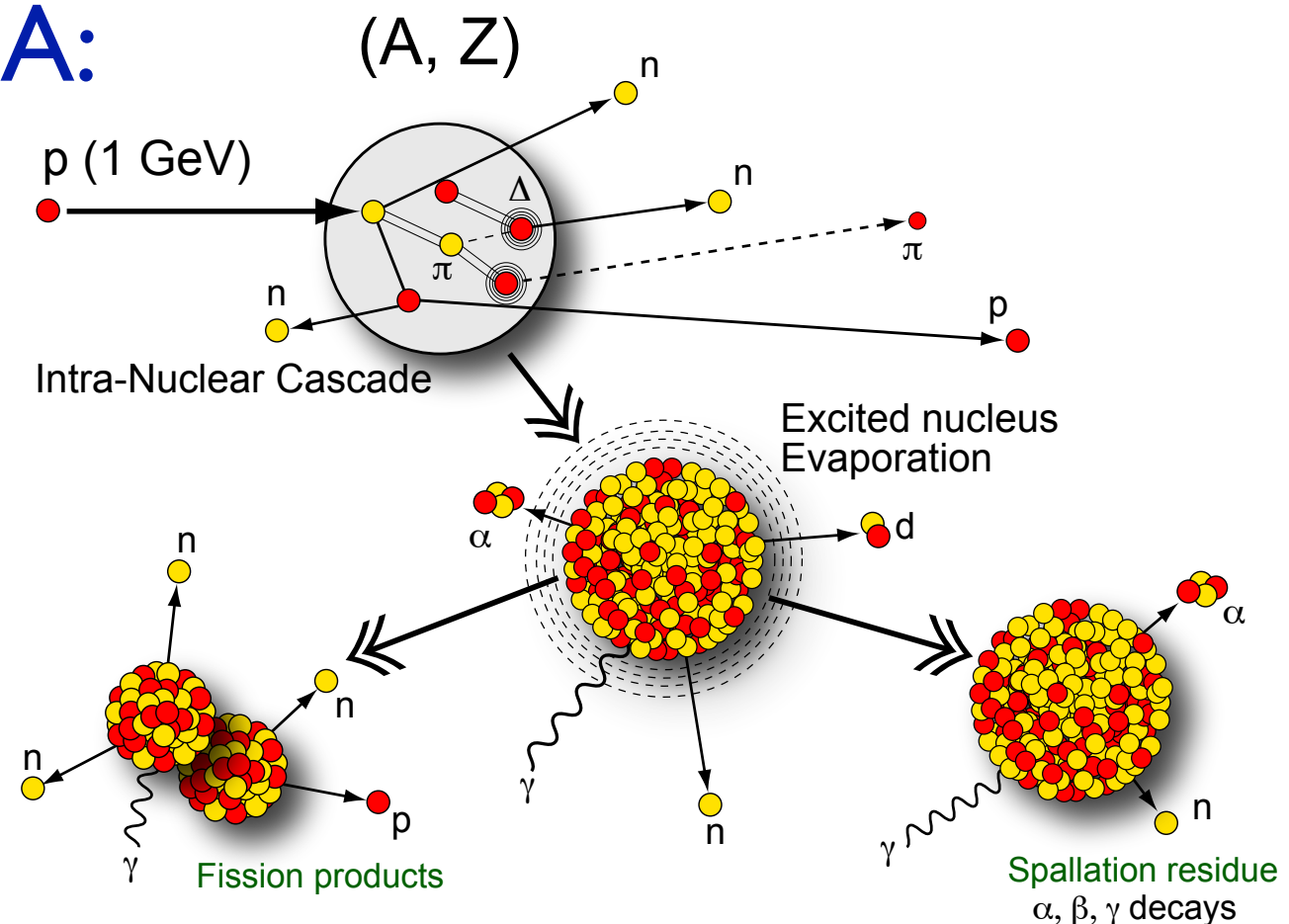
## Traditional modeling done in pA:

### Intra-Nuclear Cascade

- Particle production
- Remnant Nucleus ( $A, Z, E^*, \dots$ )
- ISABEL, INCL4

### De-Excitation

- Evaporation
- Fission
- Residual Nuclei
- Gemini++, SMM, ABLA (all no  $\gamma$ )





# Key Measurements in $e+A$

- **Momentum distribution of gluons  $xG(x, Q^2)$** 
  - ➔ Extract via scaling violation in  $F_2$ :  $\delta F_2 / \delta \ln Q^2$
  - ➔ Direct measurement:  $F_L \sim xG(x, Q^2)$  (requires  $\sqrt{s}$  scan)
  - ➔ 2+1 jet rates
  - ➔ Inelastic vector meson production (e.g.  $J/\psi$ )
  - ➔ Diffractive vector meson production  $\sim [xG(x, Q^2)]^2$

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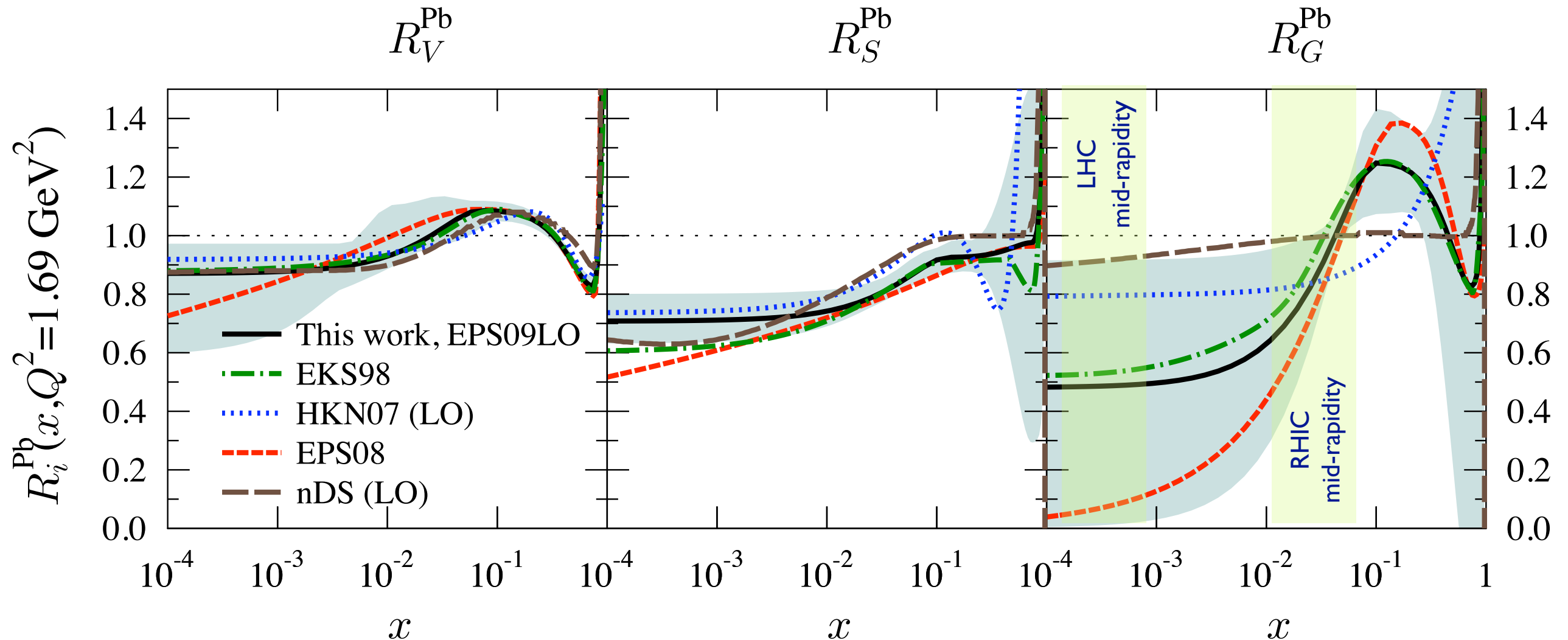
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  - ➔ Exclusive final states (e.g. vector meson production  $\rho$ ,  $J/\psi$ )
  - ➔ Deep Virtual Compton Scattering (DVCS) -  $\sigma \sim A^{4/3}$
  - ➔  $F_2, F_L$  for various  $A$  and impact parameter dependence

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- **Interaction of fast probes with *gluonic* medium?**
  - ➔ Hadronization, Fragmentation
  - ➔ Energy loss (charm, bottom!)

See next talk by A. Accardi

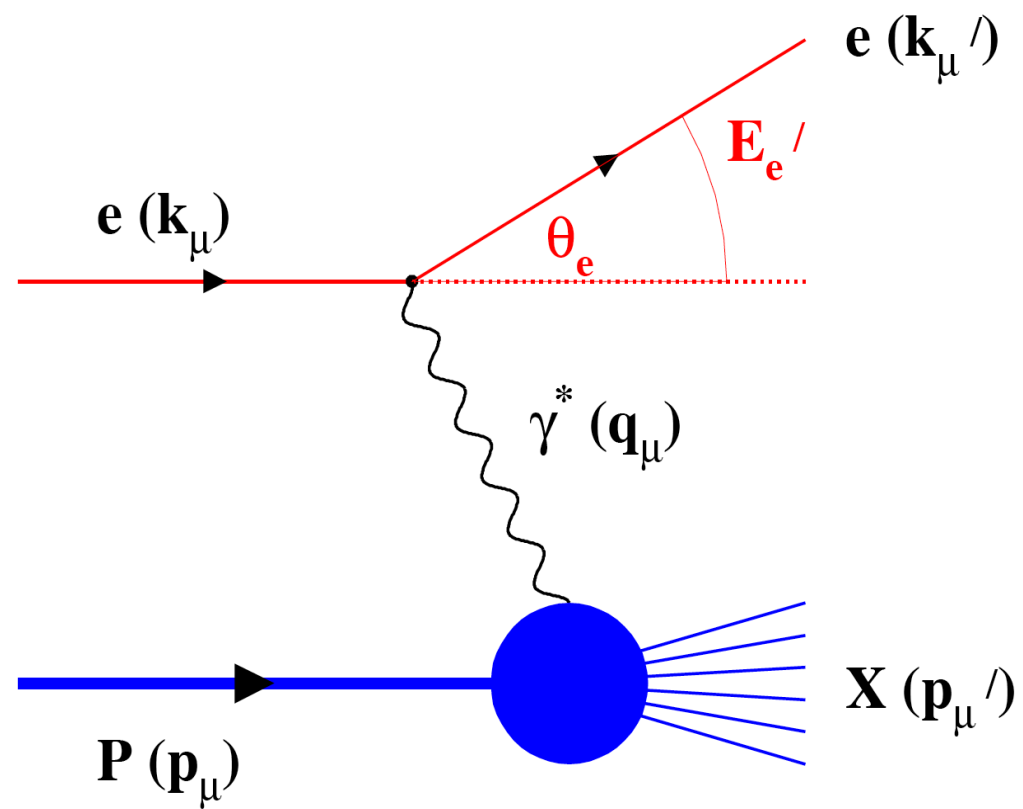
# What about gluons in nuclei?



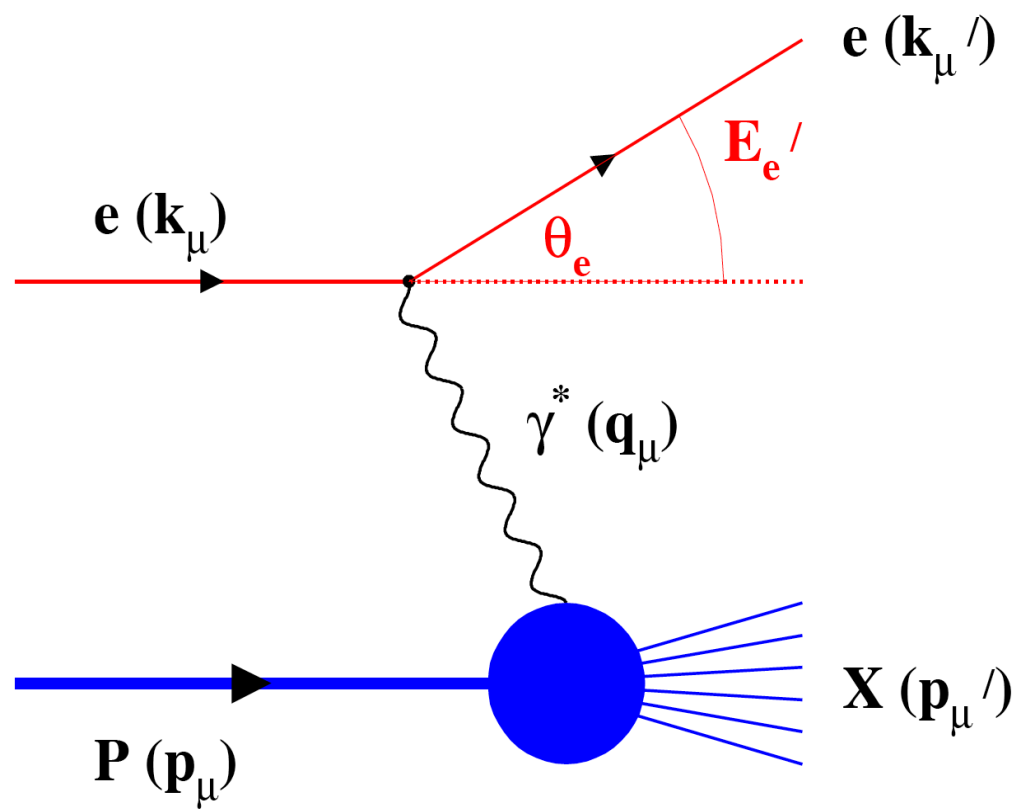
The distribution of valence and sea quarks are relatively well known in nuclei - theories agree well

Large discrepancies exist in the gluon distributions from models for mid-rapidity LHC and forward RHIC rapidities !!

# DIS Kinematics



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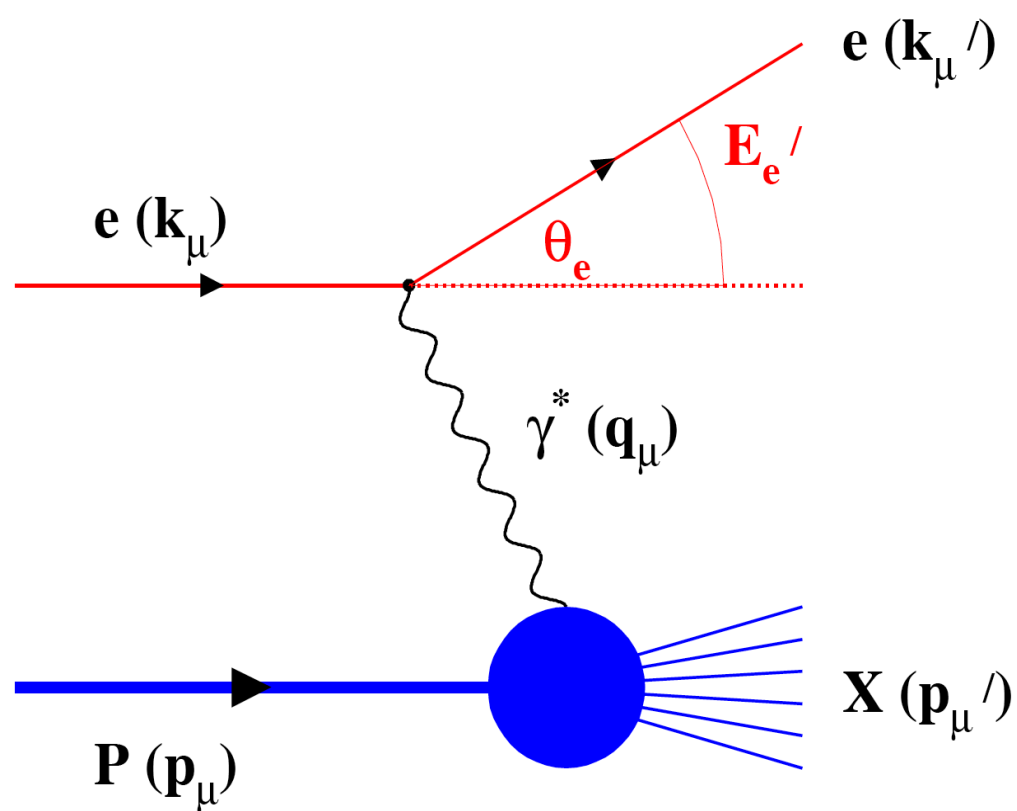


$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

$$Q^2 = 4E_e E'_e \sin^2\left(\frac{\theta'_e}{2}\right)$$

Measure of  
resolution  
power or  
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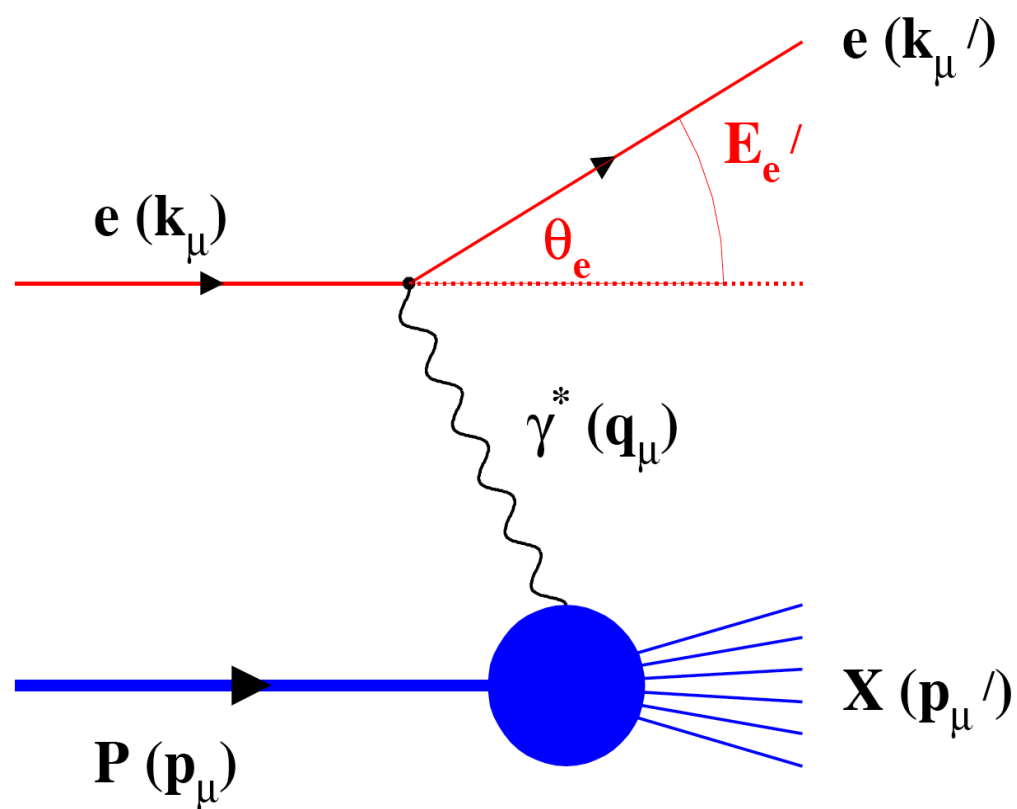
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Measure of inelasticity



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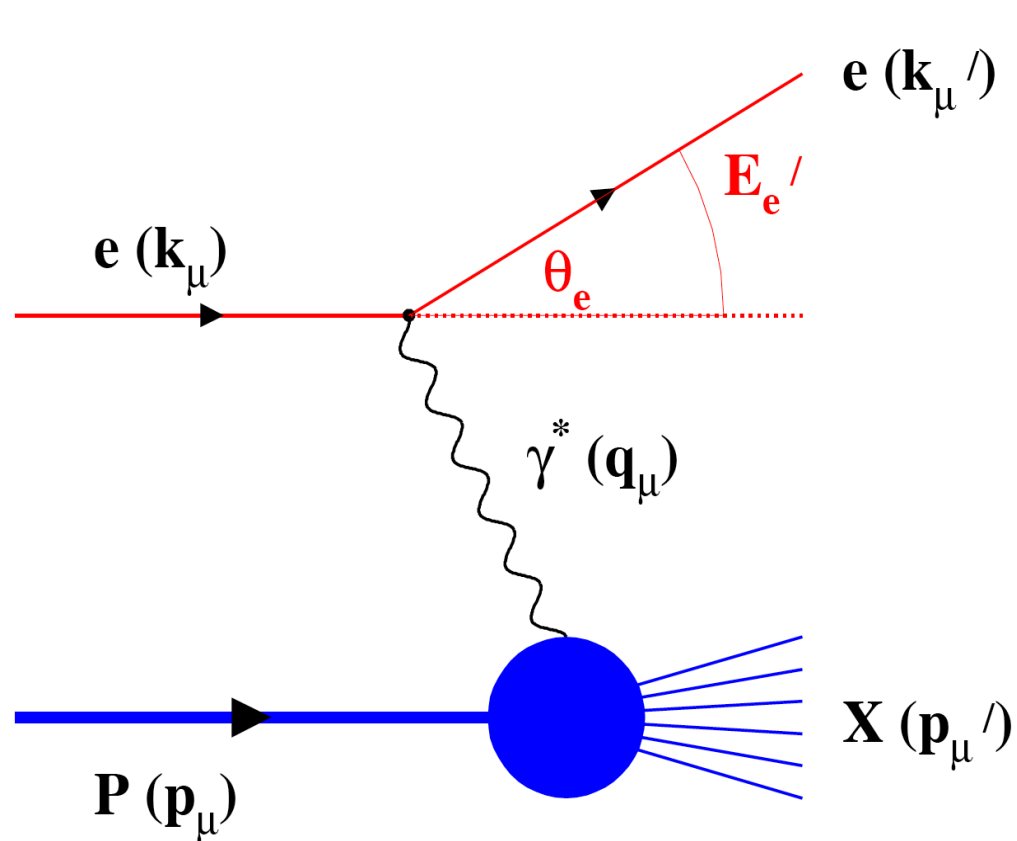
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Measure of  
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Measure of  
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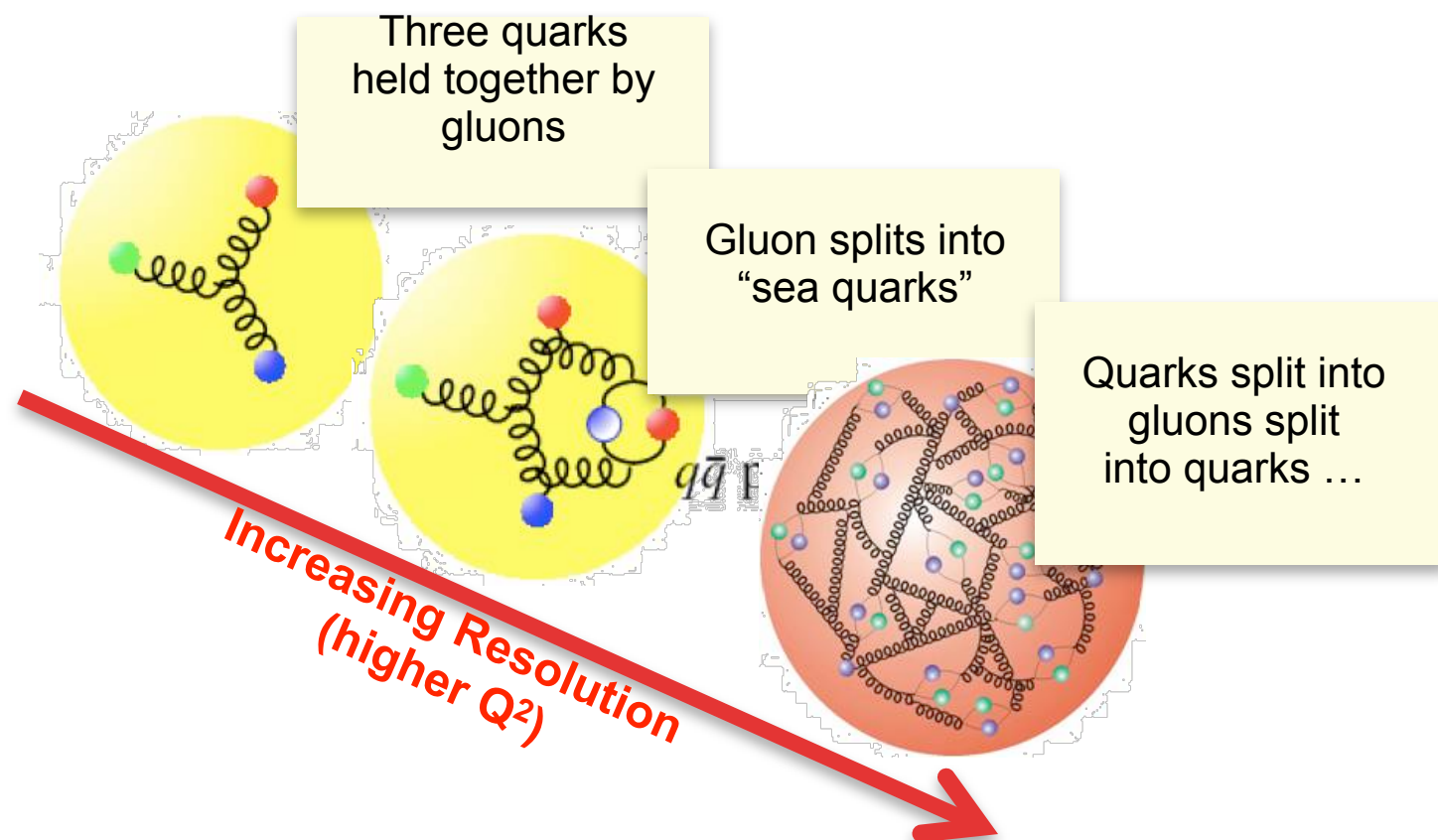
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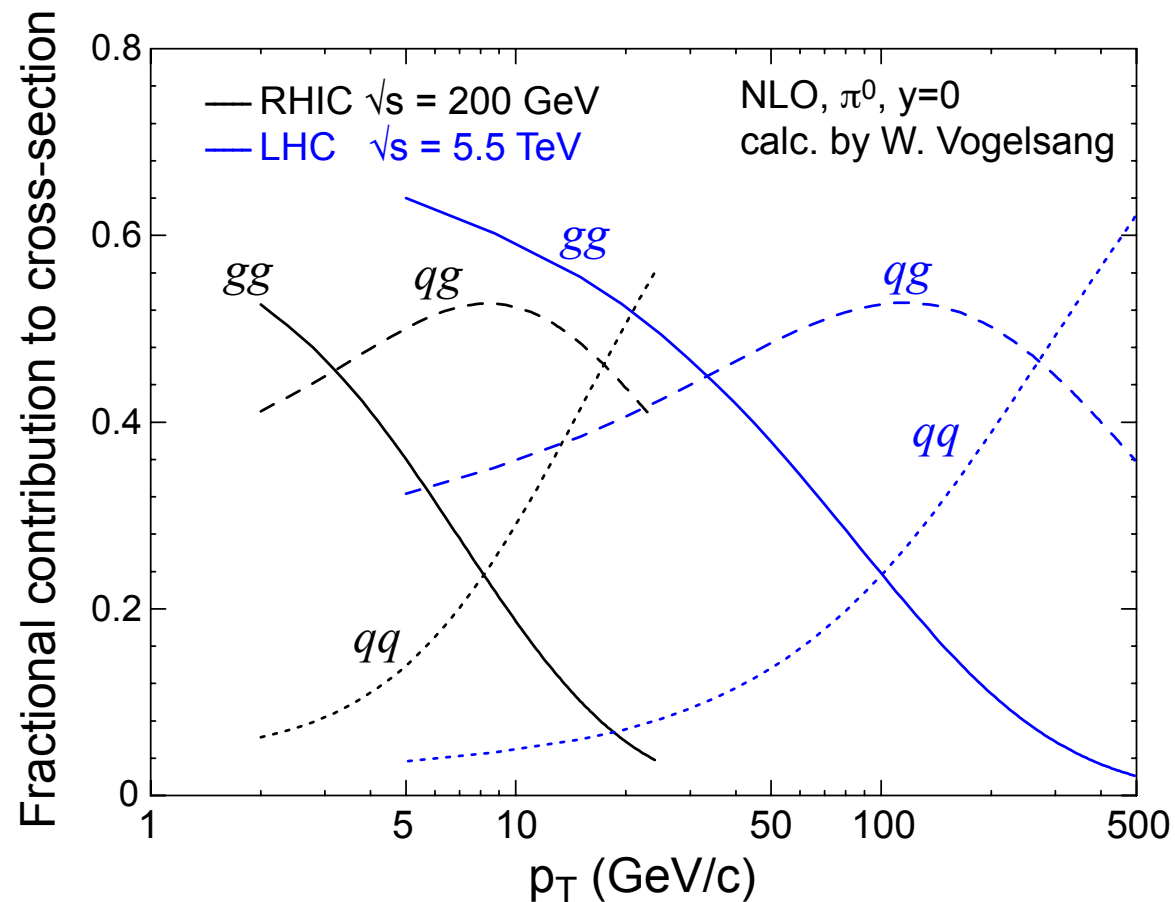
$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

Measure of momentum fraction of struck quark



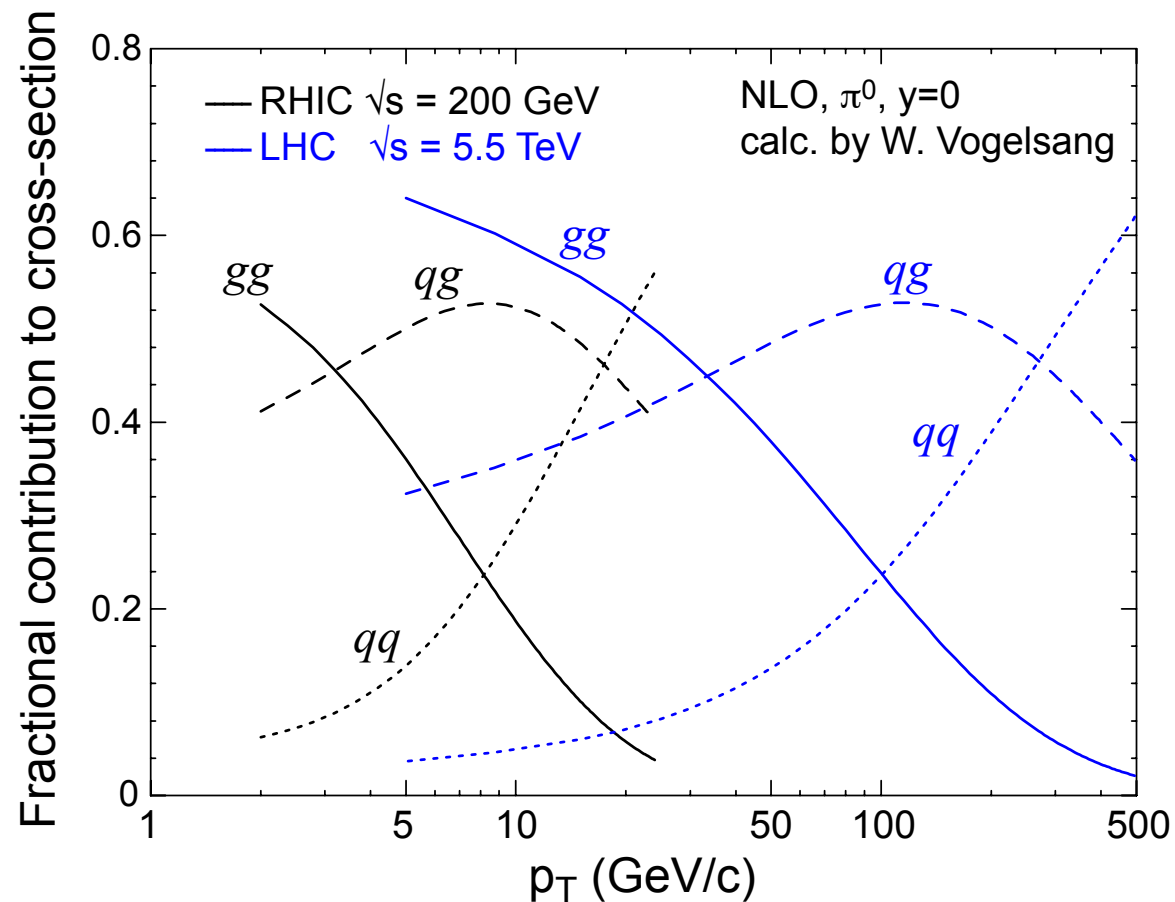
# The role of Glue in Heavy-Ion collisions

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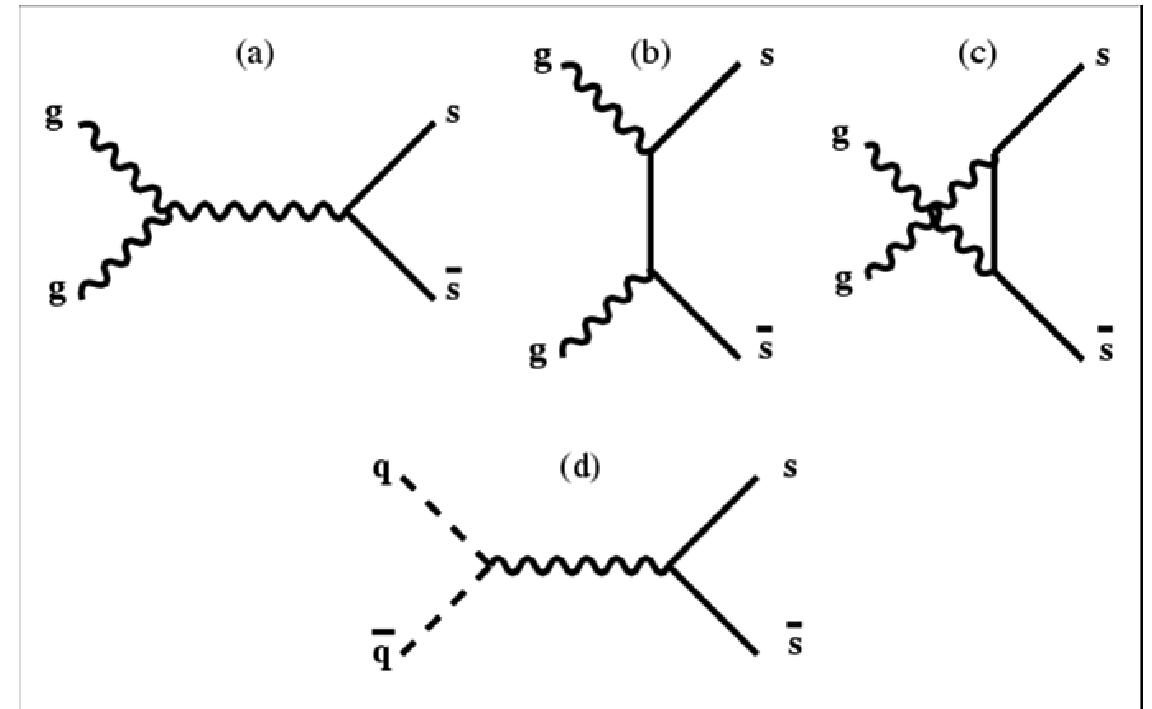


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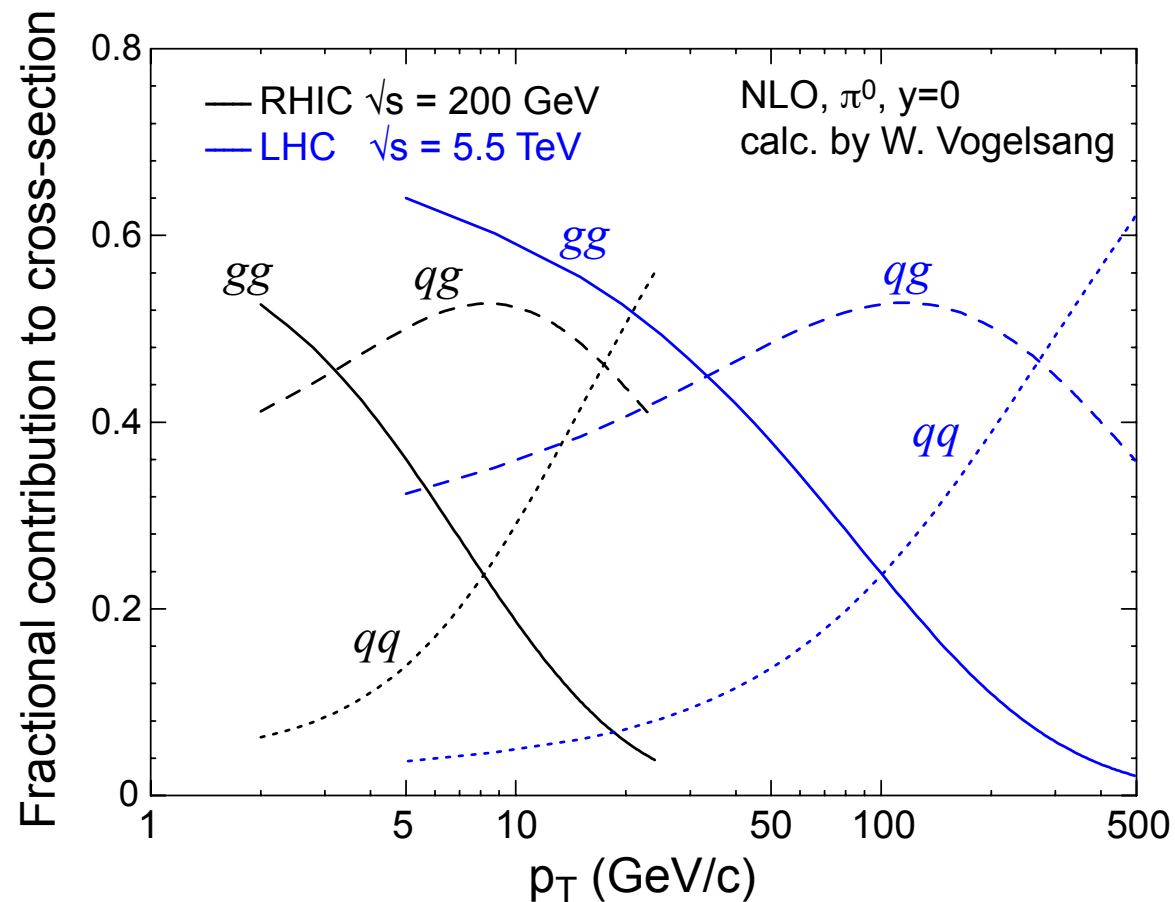


## Strangeness Production

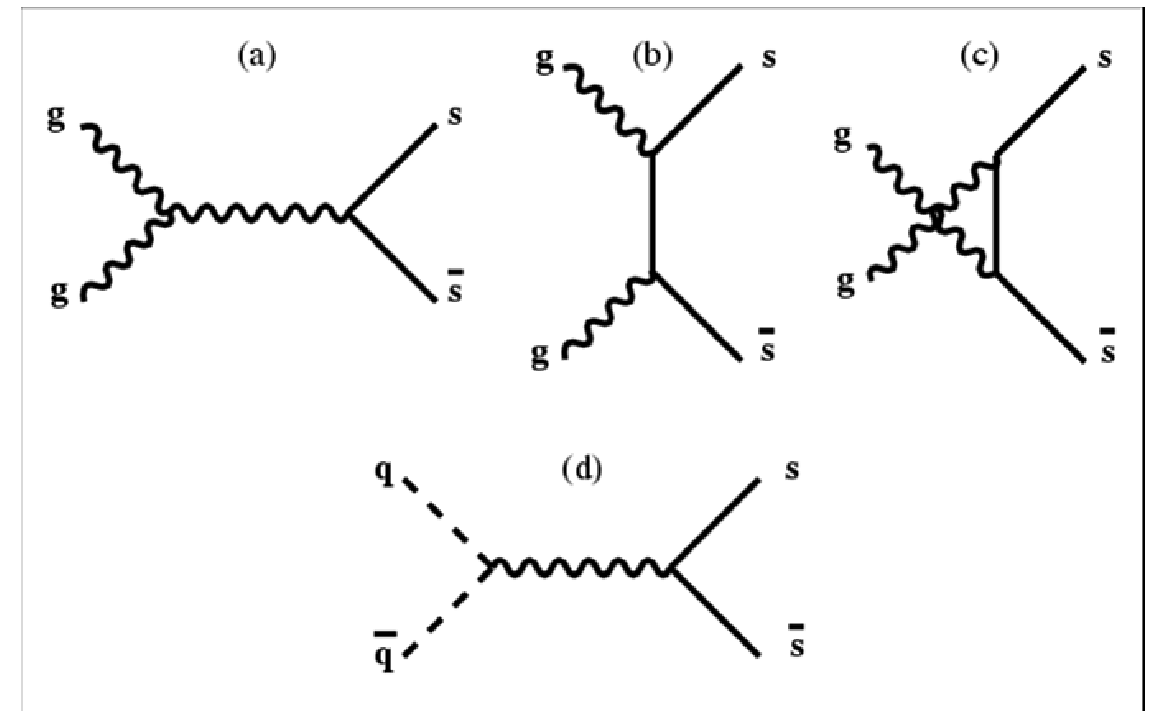


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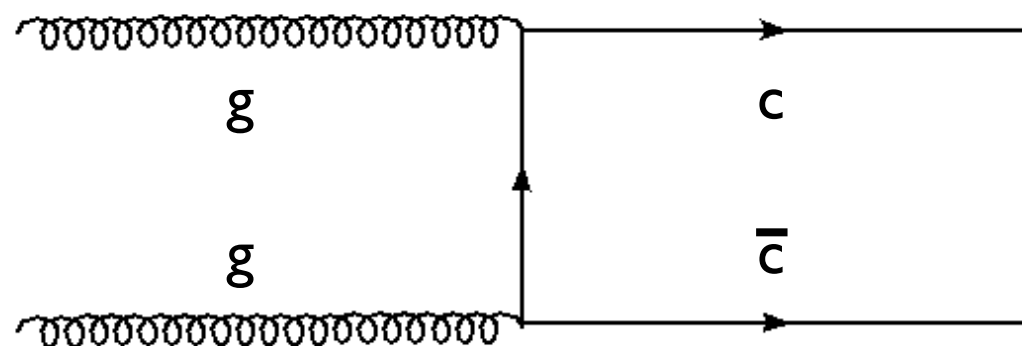
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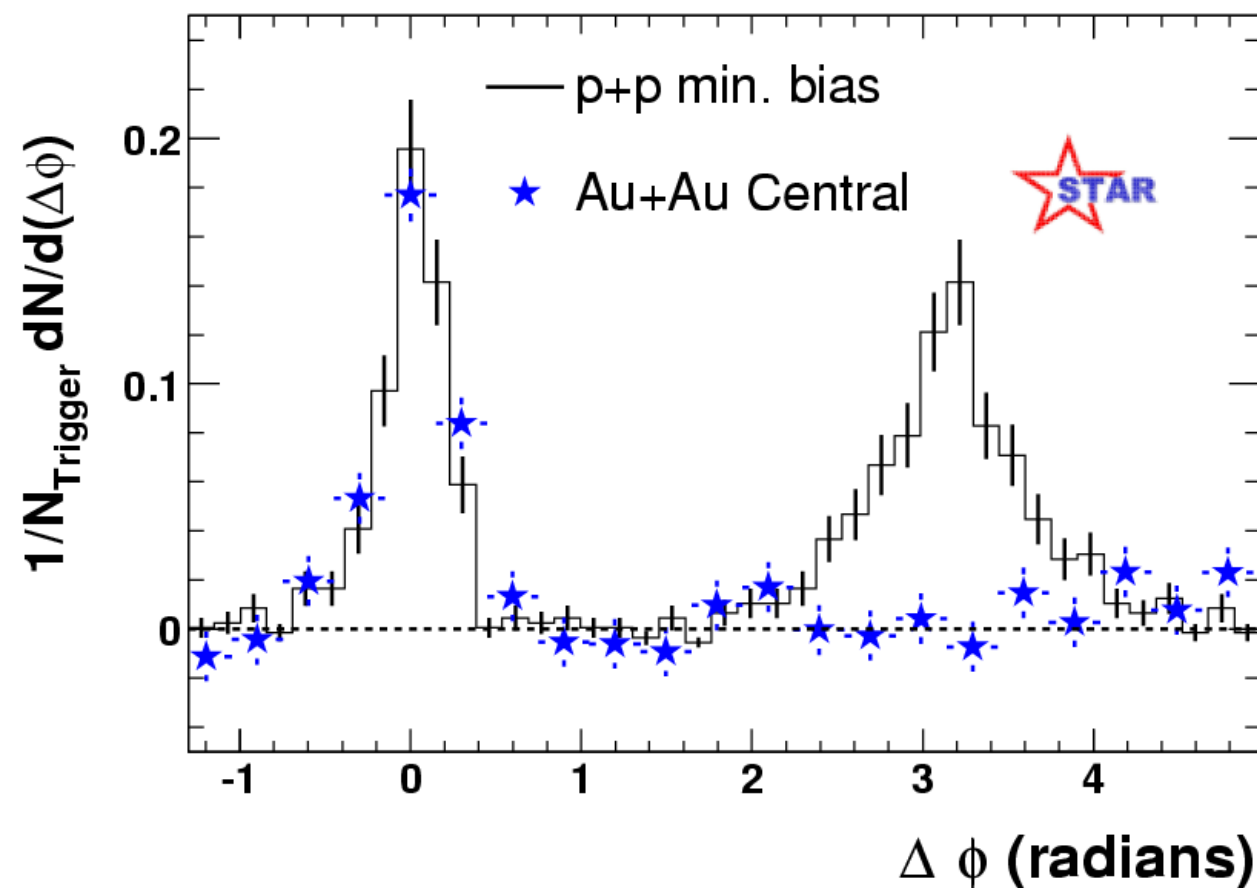
## Strangeness Production



## Heavy Flavour Production

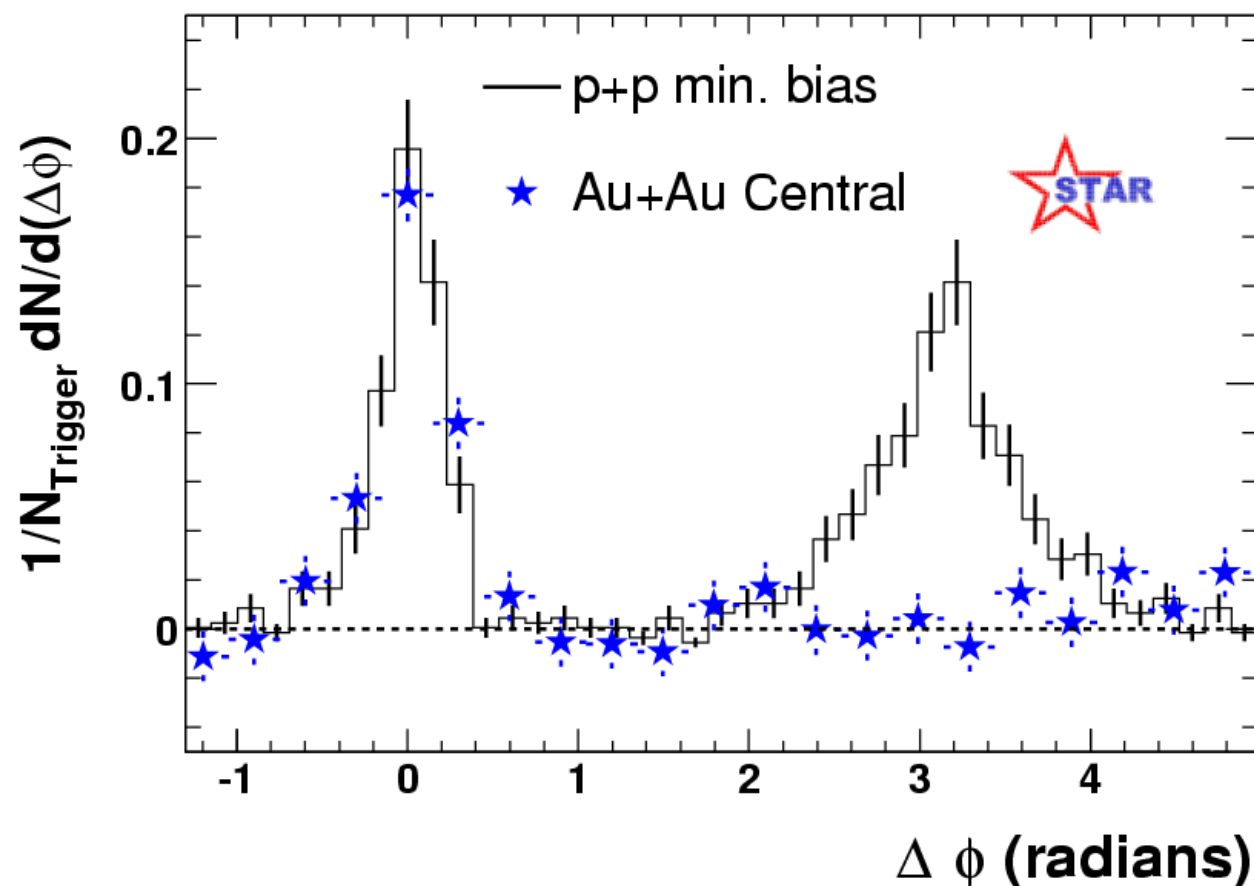


# Jet suppression: final or initial state effect?



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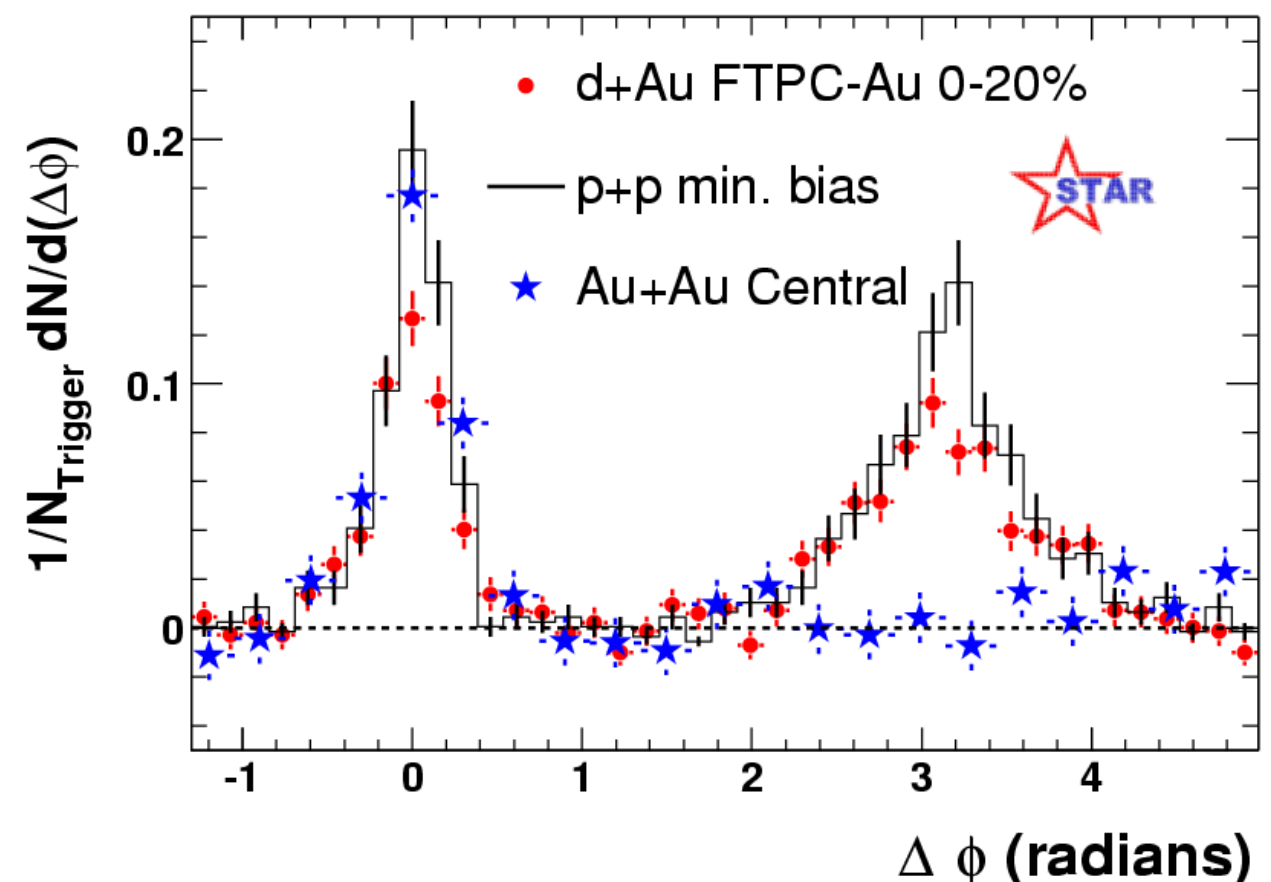
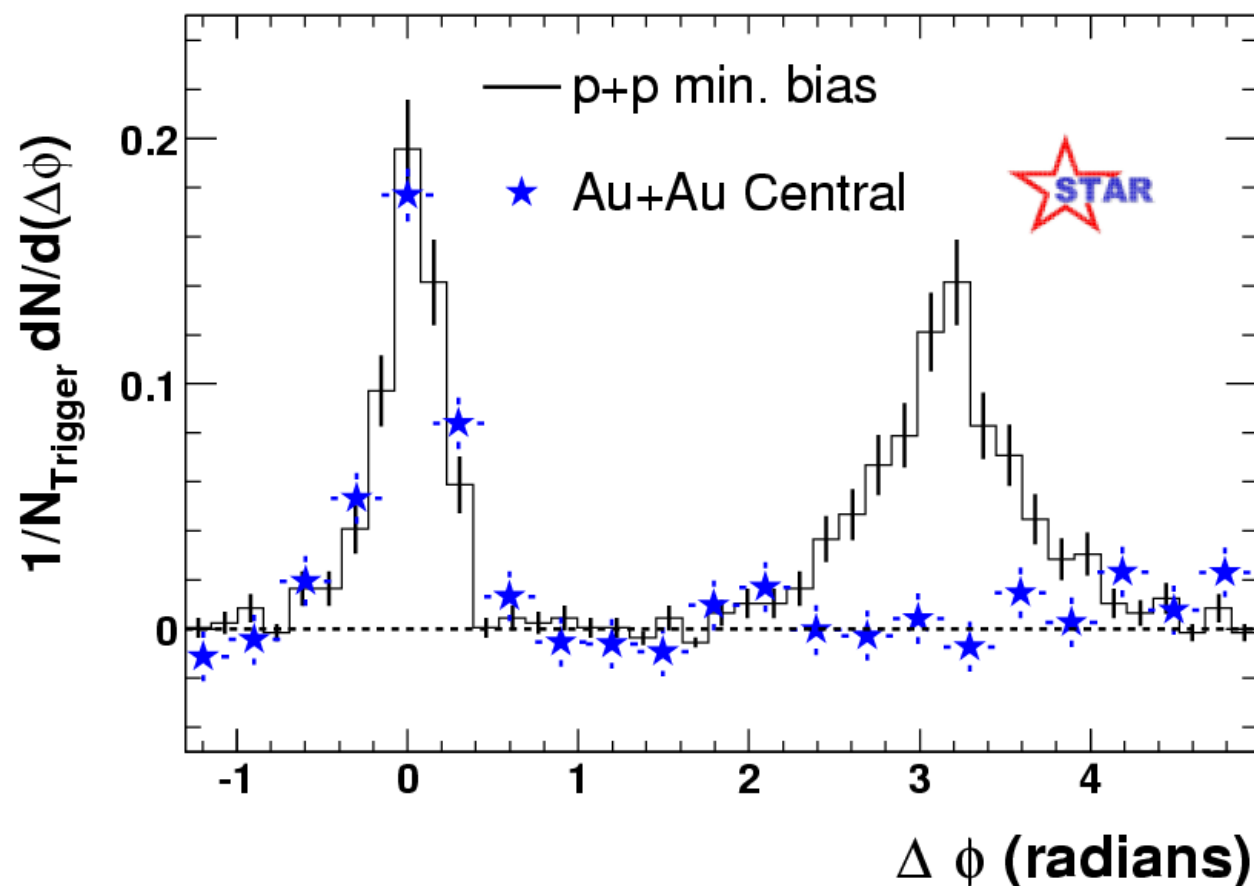
- In d+Au collisions, deconfinement is not expected
  - ▶ Measure correlations in d+Au collisions to determine if this is an initial or a final state effect





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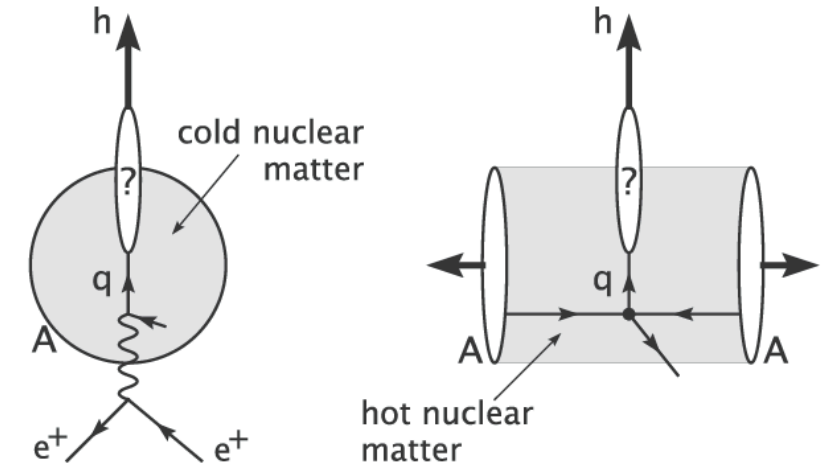
- In d+Au collisions, deconfinement is not expected
  - ▶ Measure correlations in d+Au collisions to determine if this is an initial or a final state effect
- No suppression is observed in d+Au collisions at **mid-rapidity** at RHIC
  - ▶ Jet suppression a final state effect?



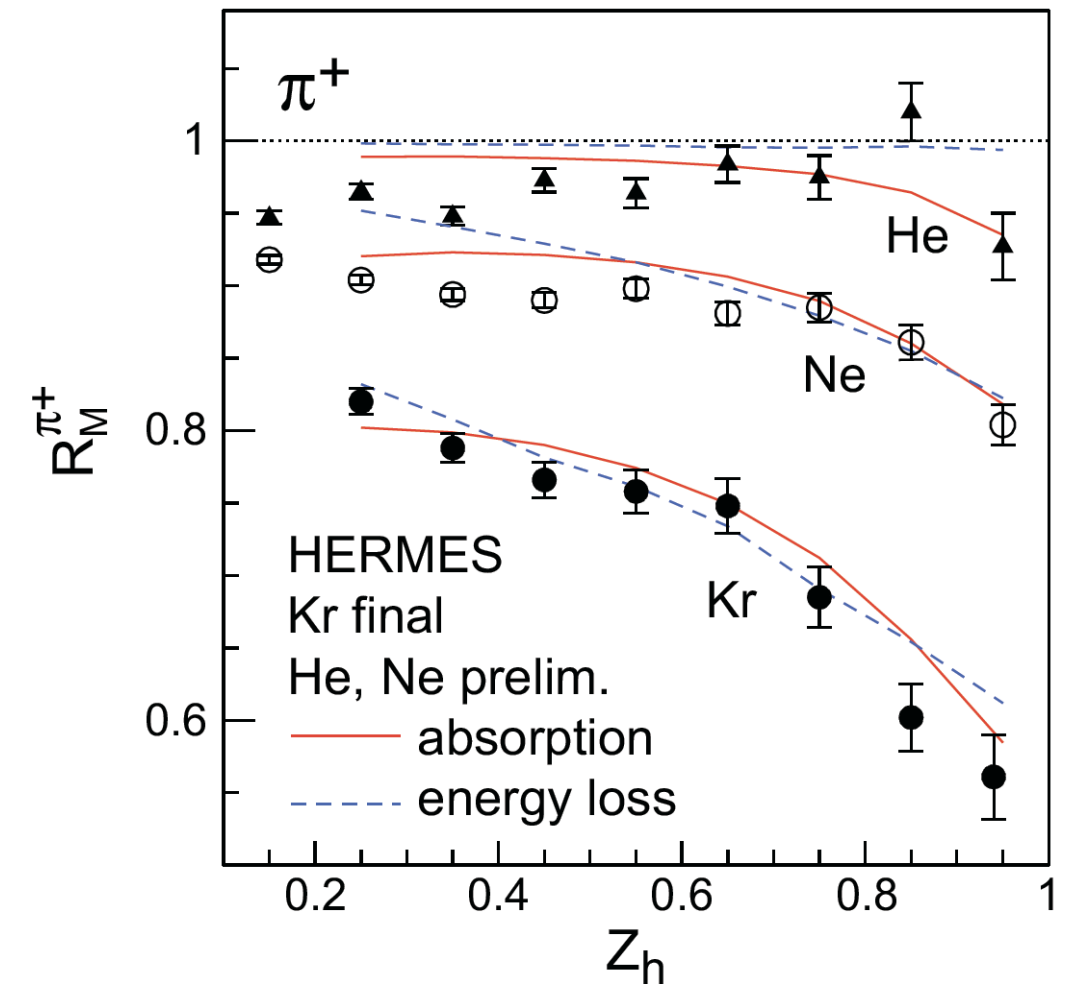
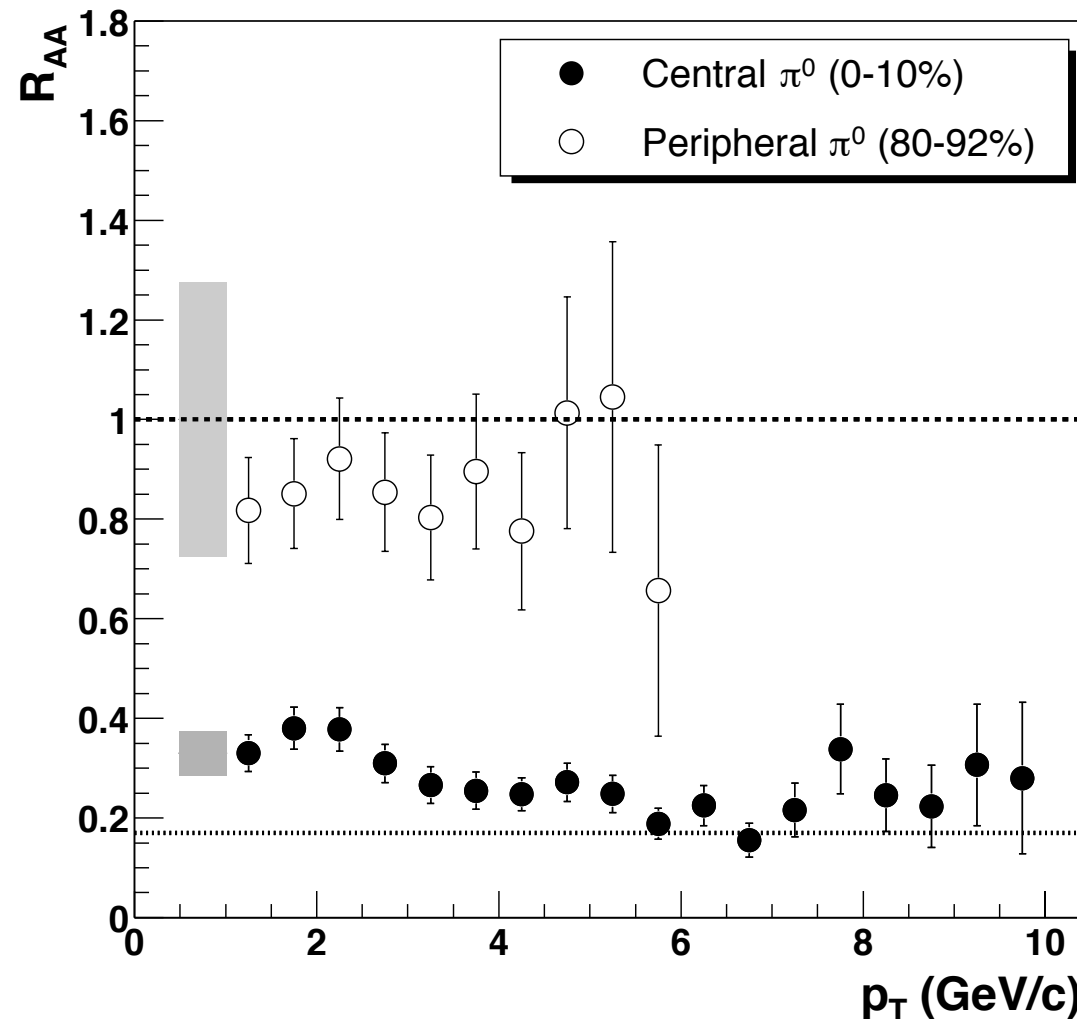
# Interaction of fast probes with gluonic medium

- nDIS:

- ➔ Clean measurement in 'cold' nuclear matter
- ➔ Suppression of high- $p_T$  hadrons analogous to, but weaker than at RHIC



PHENIX expt: Phys.Rev.Lett.91:072301 (2003)



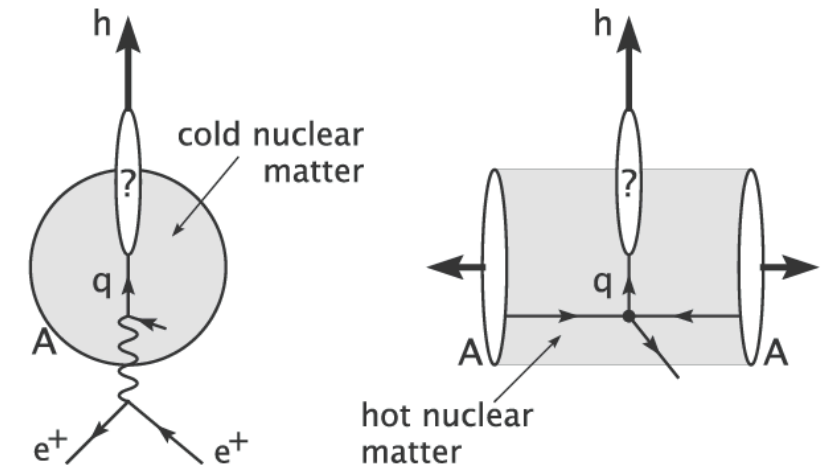
$\nu$  = virtual photon energy

$Z_h = E_h/\nu$

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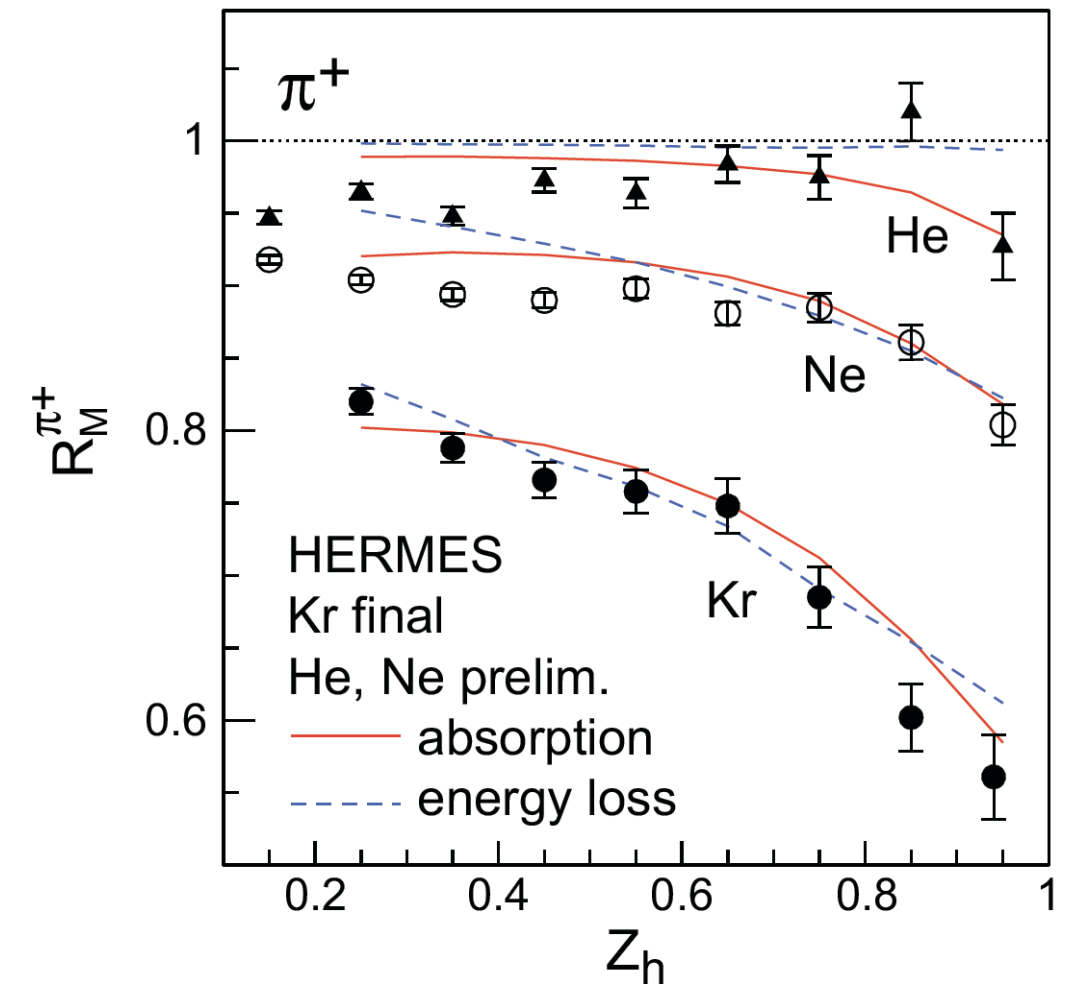
- ➔ Clean measurement in 'cold' nuclear matter
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- Fundamental question:

- When do partons get colour neutralized?

Parton energy loss vs. (pre)hadron absorption



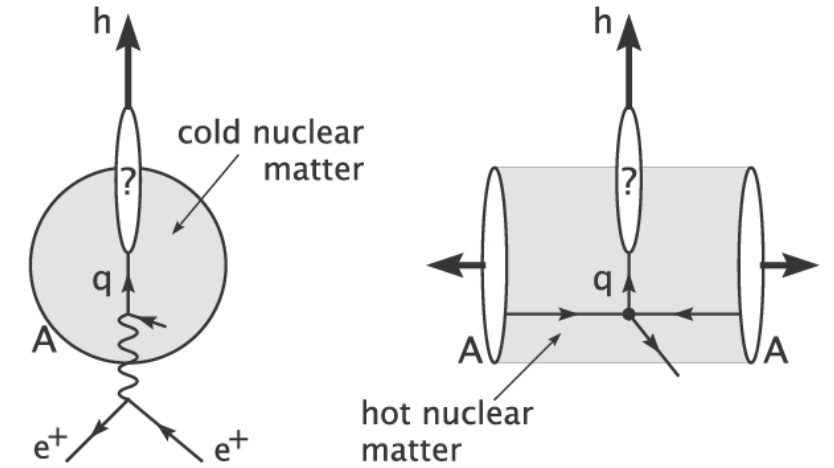
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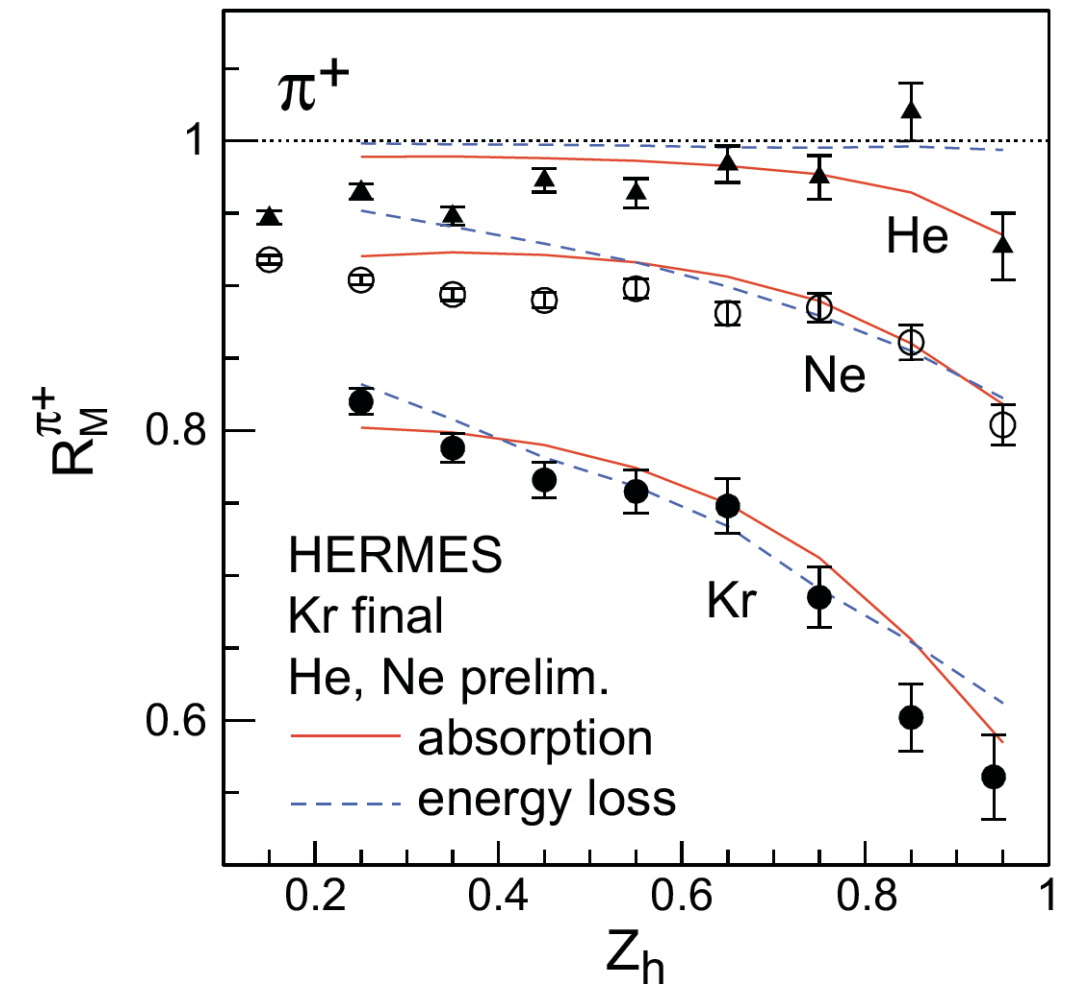
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Energy transfer in lab rest frame:

EIC:  $10 < \nu < 1600$  GeV

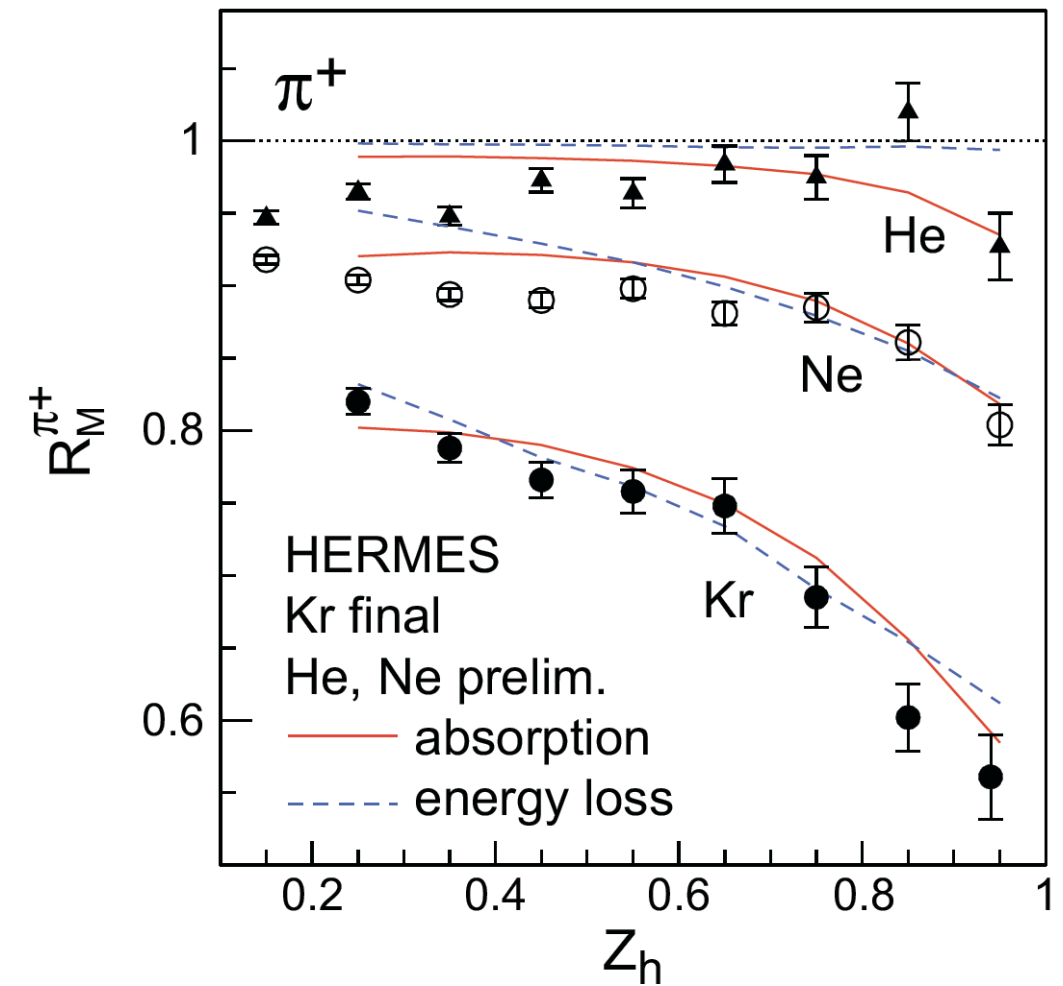
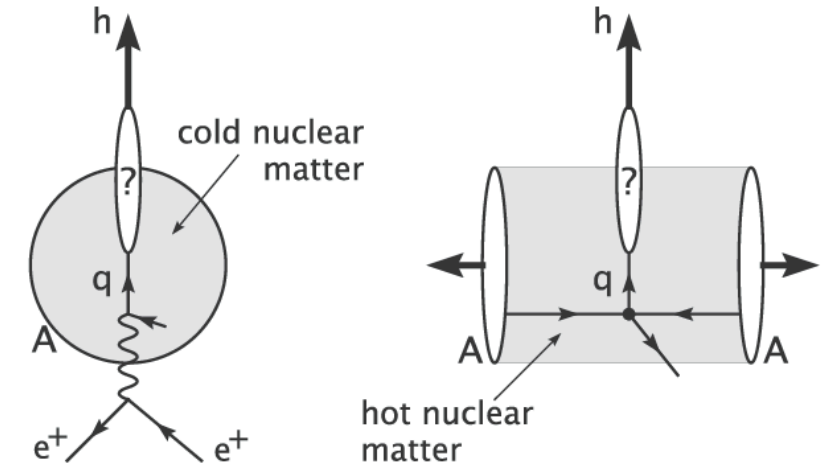
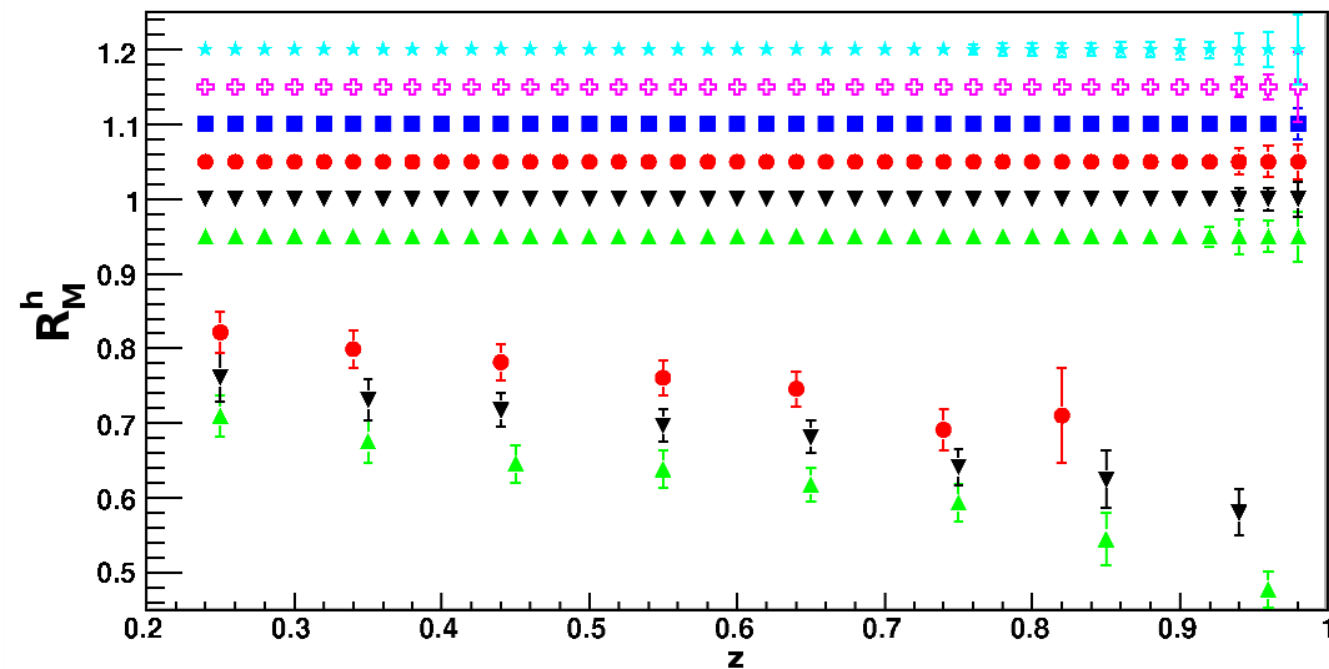
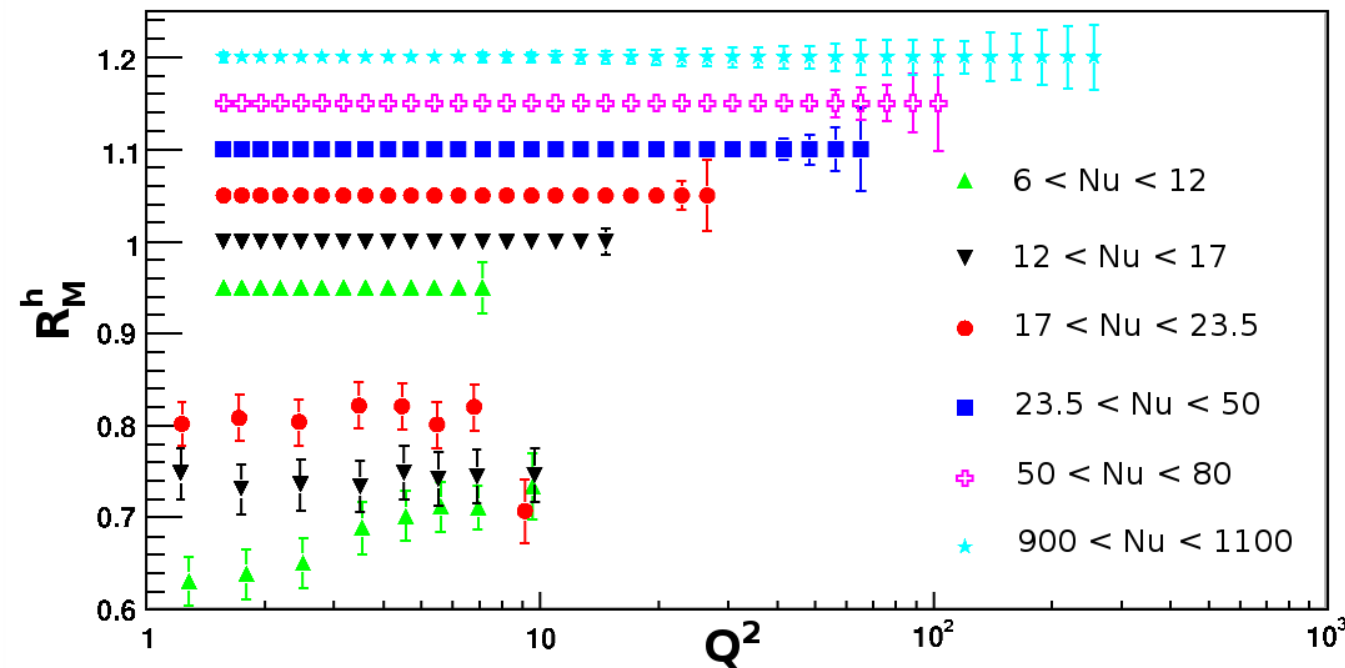
HERMES: 2-25 GeV



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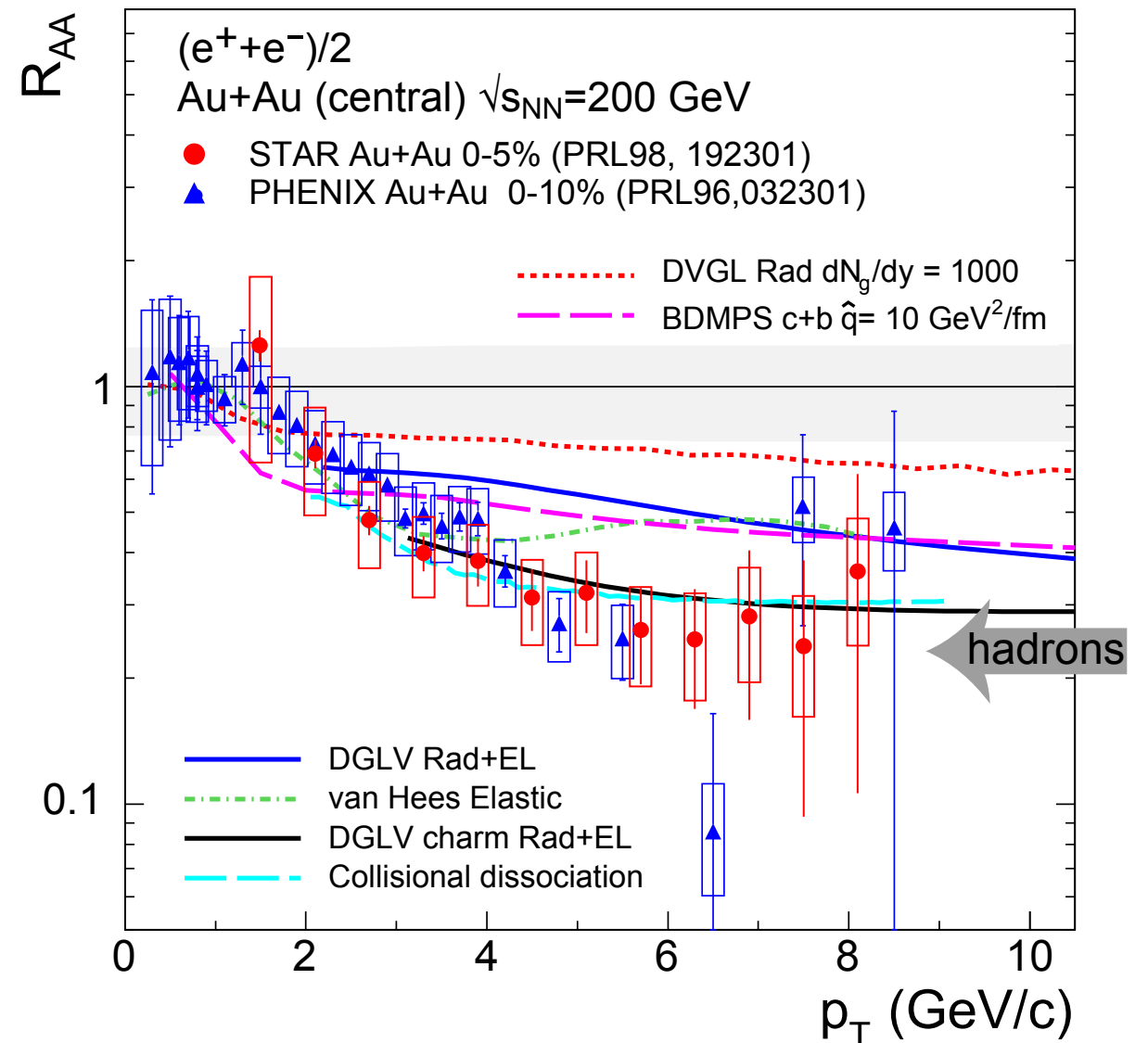
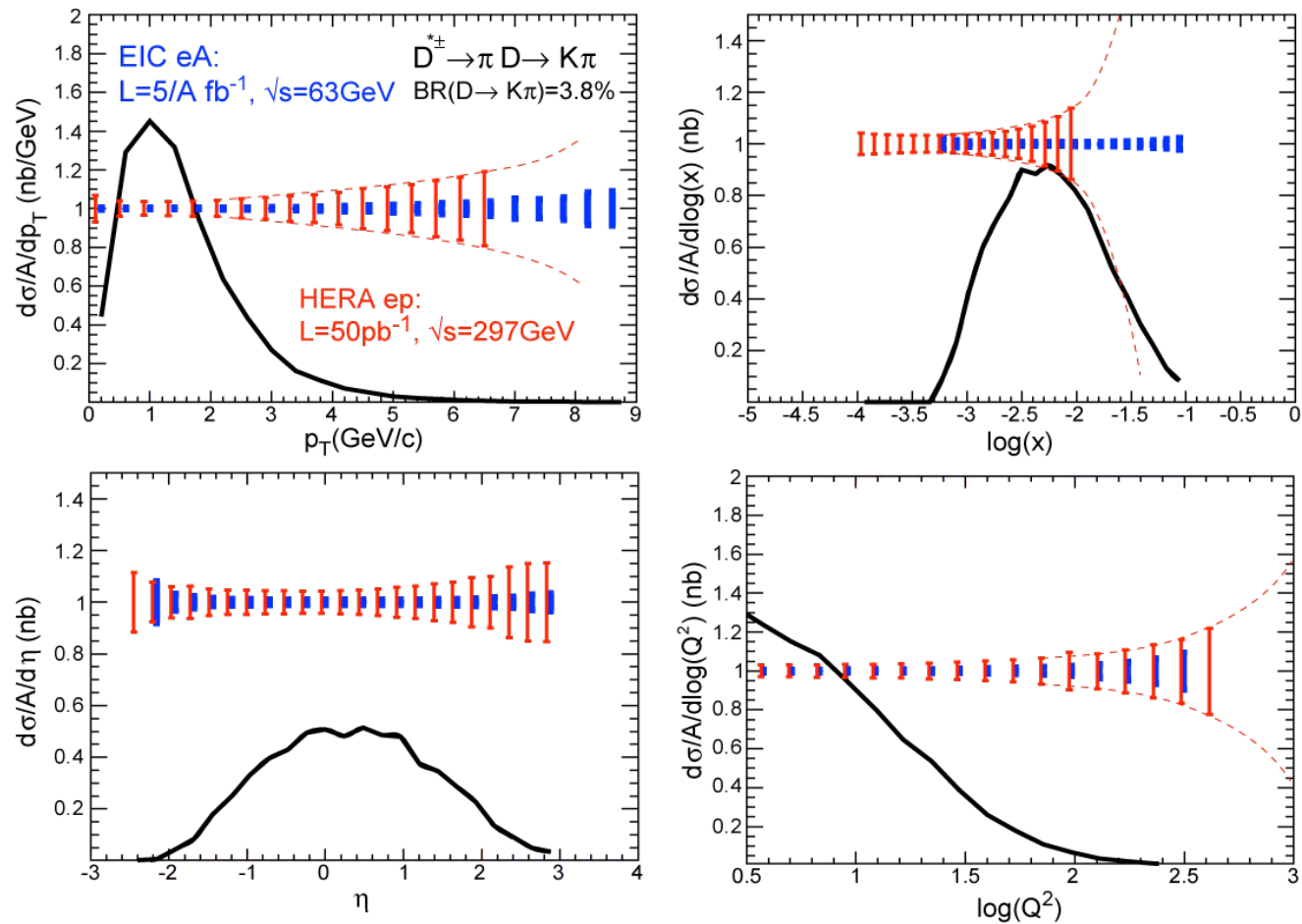
# Interaction of fast probes with gluonic medium



$\nu$  = virtual photon energy

$z_h = E_h/\nu$

# Charm at an EIC



- ➔ EIC allows **multi-differential measurements** of heavy flavour
- ➔ covers and **extends energy range of previous experiments**, allowing for **the study of a wide range of formation lengths**