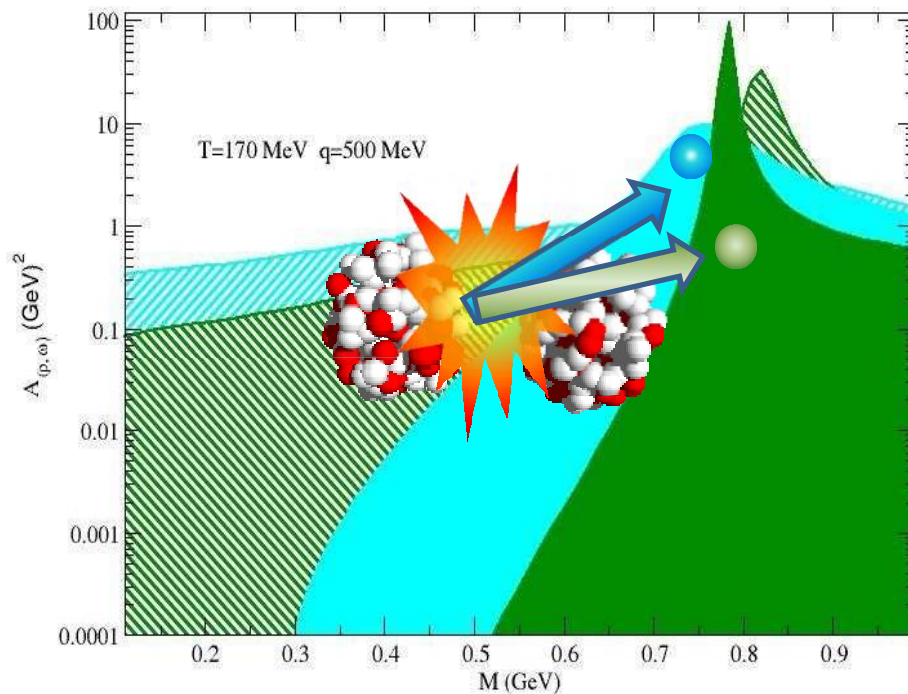


# Vector Meson Spectral Functions in Medium

Sabyasachi Ghosh

## Outline of the talk .....

- ✓ *Motivation + meaning of spectral function*
- ✓ *RTF (essence + calculations)*
- ✓ *Results of in-medium spectral function ( $\rho + \omega$ )*
- ✓ *Application on dilepton*

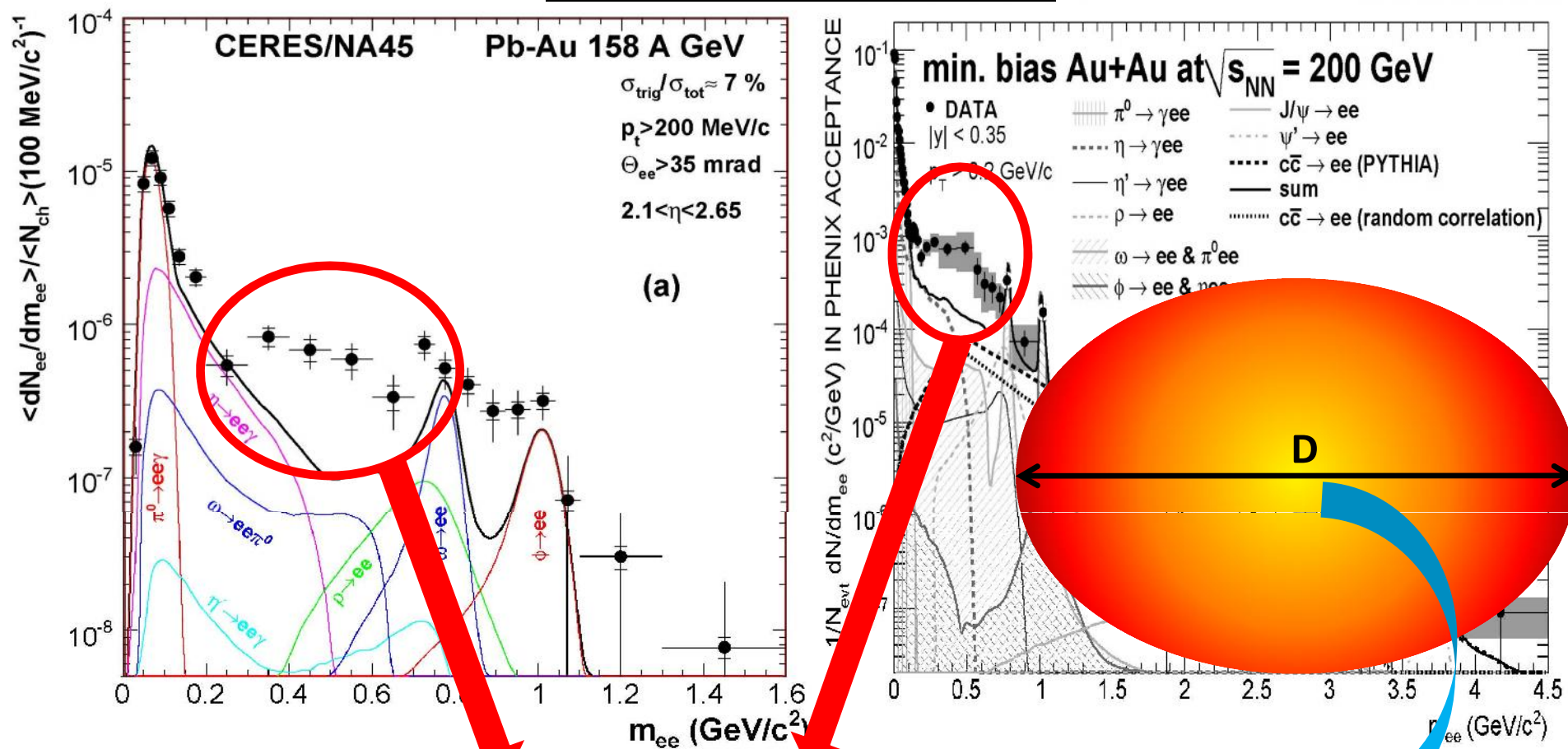


## Thermal Photons and Dileptons in Heavy-Ion Collisions

RIKEN BNL Research Center Workshop  
August 20-22, 2014 at Brookhaven National Laboratory

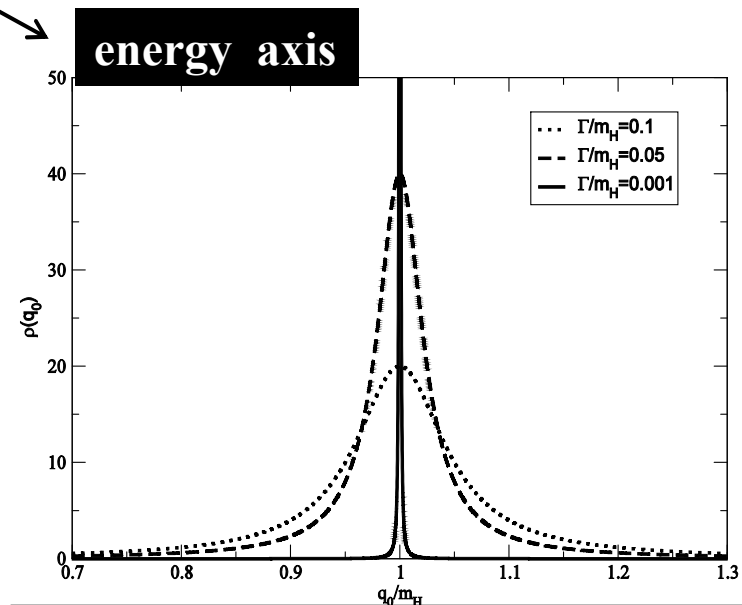
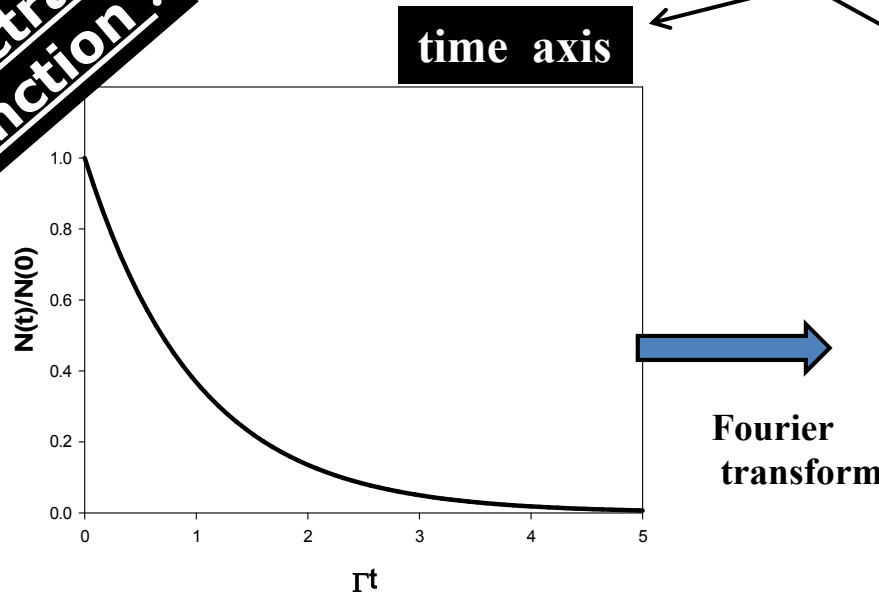


## Experimental motivation :



# Probabilistic amplitude of Unstable particle

**spectral function :**



$$\frac{N(t)}{N(0)} = \left| \frac{\psi(t)}{\psi(0)} \right|^2 = \exp(-\Gamma t)$$

$$\psi(t) \sim \exp(im_H t - \Gamma t / 2)$$

$$\frac{N(q_0)}{N(m_H)} = \left| \frac{\tilde{\psi}(q_0)}{\tilde{\psi}(m_H)} \right|^2 = (\Gamma / 2) \rho(q_0)$$

$$\tilde{\psi}(q_0) \sim \frac{1}{(q_0 - m_H) - i\Gamma / 2}, \rho(q_0) = \text{Im } \tilde{\psi}(q_0)$$

## QFT definition of spectral function

$$x = (t, \vec{x}) \longrightarrow q = (q_0, \vec{q})$$

$$\lim_{\Gamma \rightarrow 0} \rho(q) \sim \delta(q^2 - m_H^2)$$

$$\sim \text{Im} \frac{1}{q^2 - m_H^2 + i\eta}$$

$$\lim_{\Gamma \rightarrow 0} \rho(q_0) \sim \delta(q_0 - m_H)$$

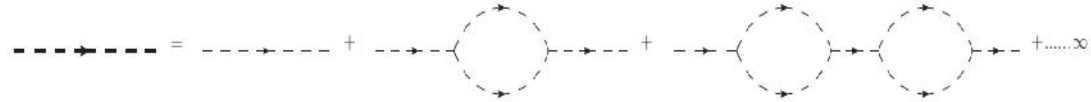
**Propagator**

$$\langle 0 | \psi(x_1) \psi(x_2) | 0 \rangle$$

## Vacuum free propagator

$$\lim_{\Gamma \rightarrow 0} \rho(q) \sim \delta(q^2 - m_H^2)$$

$$\sim \text{Im} \frac{1}{q^2 - m_H^2 + i\eta}$$



## Vacuum interacting propagator

$$D(q^2) = \frac{1}{q^2 - m^2 - \Pi(q^2)}$$

## Essence of RTF

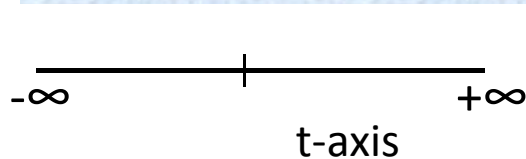
Time evolution  
operator

$$\exp(-i H t) = \exp(-\beta H)$$

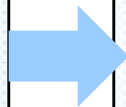
$$t \rightarrow \tau = -i\beta$$

~ density matrix

## Field Theory of vacuum

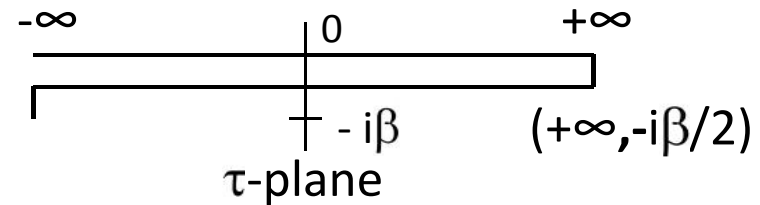


$$\int_{-\infty}^{\infty} dt$$



$$\int_{\text{contour}} d\tau$$

## Field Theory at finite temperature



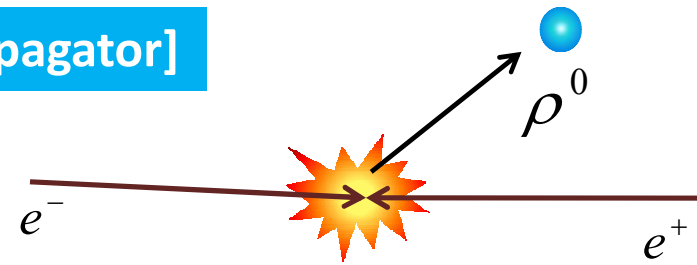
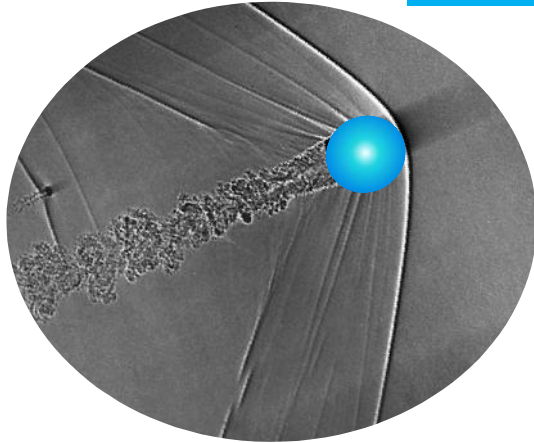
## Diagonalization

$$\bar{D}(q^0, \vec{q}, T) = \frac{1}{q^2 - m^2 - \bar{\Pi}(q^0, \vec{q}, T)}$$

$D^{ab}$  thermal propagator

Thermal self-energy  $\Pi^{ab}$

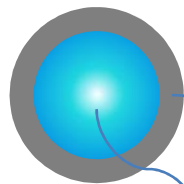
Spectral function = Im [propagator]



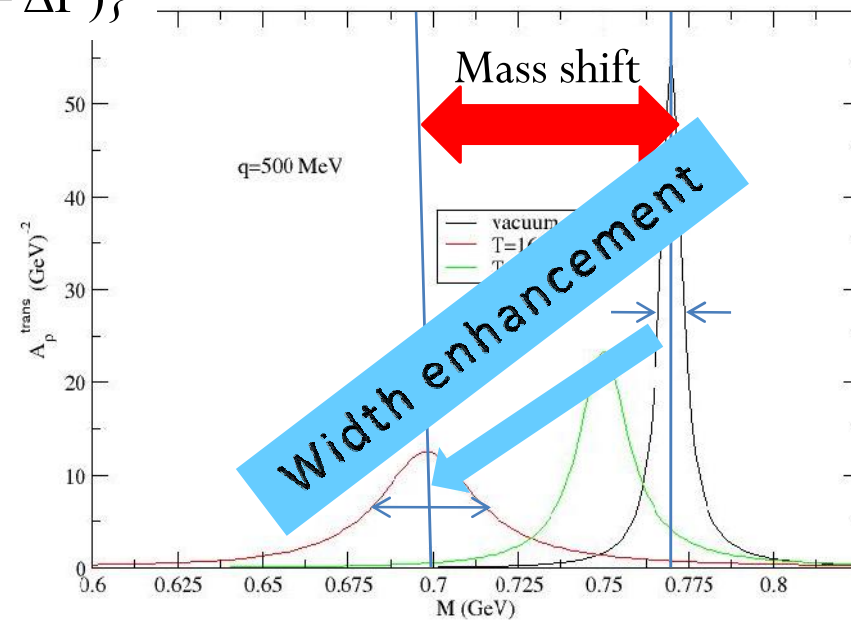
$$A_v(q) = \text{Im} \left[ \frac{1}{\{q^2 - (m_{phy}^2) - iq\Gamma_{decay}\}} \right]$$

$$A_{tot}(q) = \text{Im} \left[ \frac{1}{\{q^2 - (m_{phy} + \Delta m)^2 - iq(\Gamma_{decay} + \Delta\Gamma)\}} \right]$$

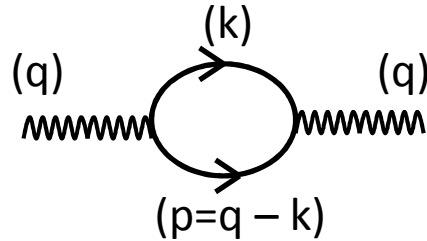
$$\overline{\Pi}_{total} = \overline{\Pi}_{vac} + \overline{\Pi}_{th}$$



$$m_{bare}^2 + \text{Re}\overline{\Pi}_{vac} + \text{Re}\overline{\Pi}_{th} = m_{phy}^2 + \text{Re}\overline{\Pi}_{th} \\ = (m_{phy} + \Delta m)^2$$



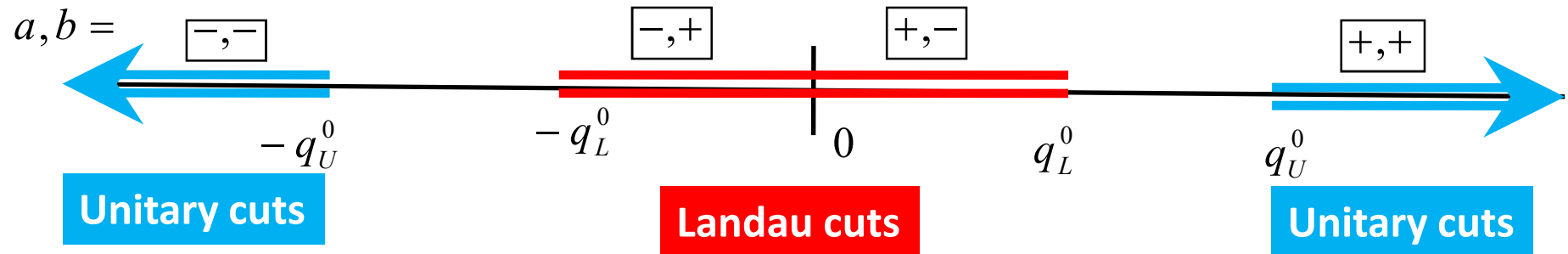
# Self Energy of $\rho$ for mesonic loops :



$$D_{11} = \text{vacuum} + \text{Thermal}$$

$$\Pi_{11}^{\mu\nu}(q, T) = \int \frac{d^4 k}{(2\pi)^4} N^{\mu\nu}(q, k) D_{11}(k, T) D_{11}(p = q - k, T)$$

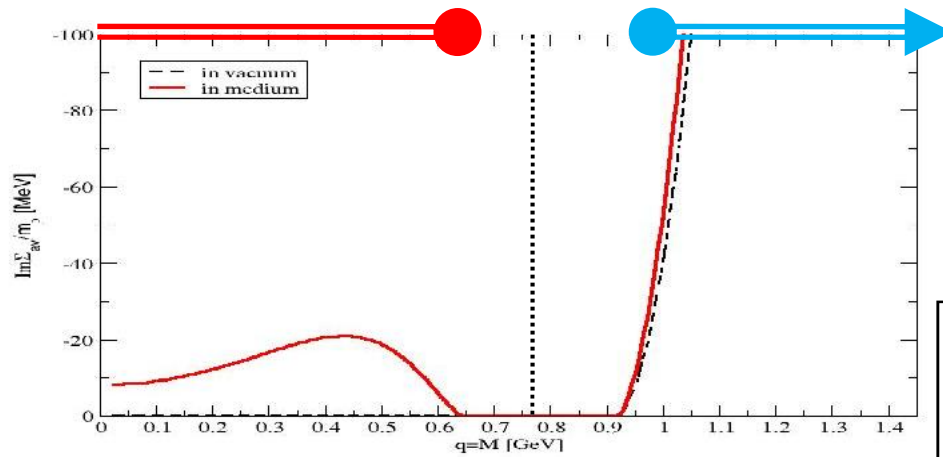
$$\text{Im} \Pi_{11}^{\mu\nu}(q, T) = \int d^3 k \sum_{a,b} (..) \delta(q_0 - a \omega_k - b \omega_p)$$



$$q_L^0 = \sqrt{\vec{q}^2 + (m_p - m_k)^2}$$

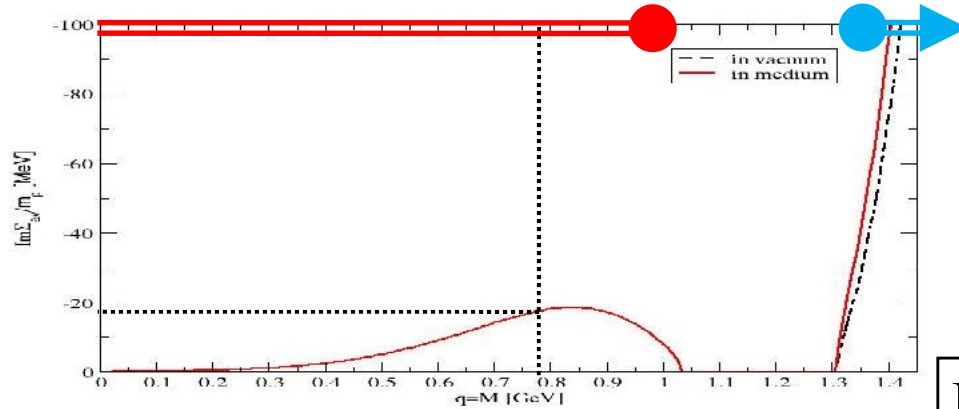
$$q_U^0 = \sqrt{\vec{q}^2 + (m_p + m_k)^2}$$





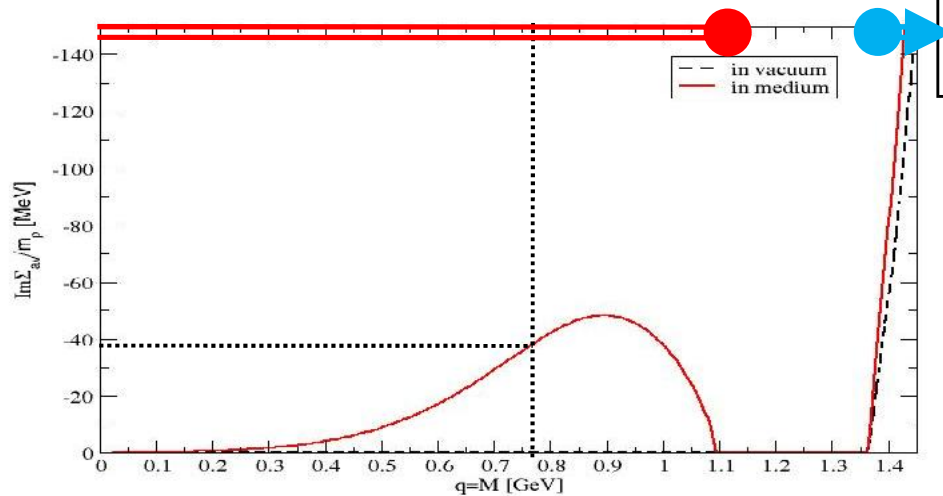
$$\begin{aligned}
 k, p &= \pi \omega(782) \\
 &\pi h_1(1170) \\
 &\pi a_1(1260)
 \end{aligned}$$

$$\frac{\text{Im } \bar{\Pi}_L^M}{m_\rho} \sim \text{vac} \otimes \{n_k(1+n_p) - (1+n_k)n_p\}$$



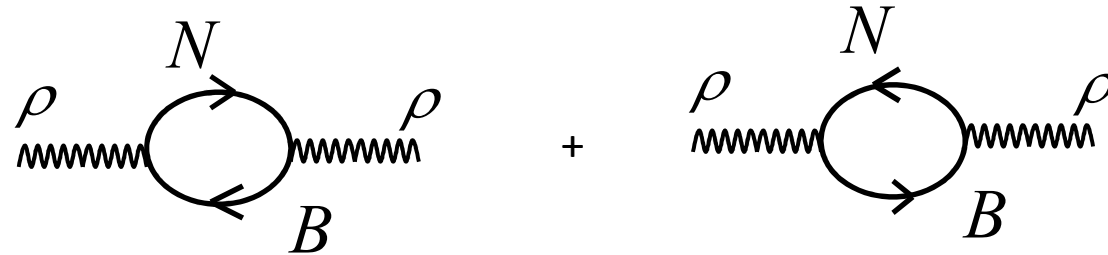
Mesonic collision rate ( $\Gamma_C^M$ )

$$\frac{\text{Im } \bar{\Pi}_U^M}{m_\rho} \sim \text{vac} \otimes \{(1+n_k)(1+n_p) - n_k n_p\}$$

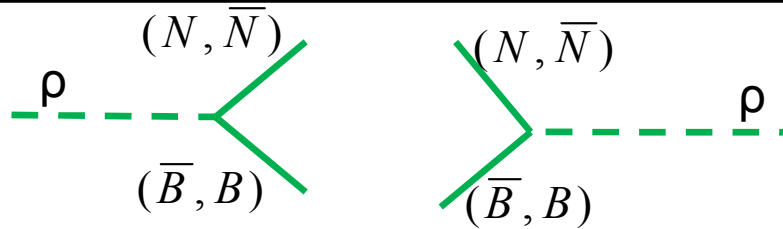


Bose enhancement of decay rate ( $\Gamma_{B.e.}^M$ )

# Self Energy of $\rho$ for baryonic loops :

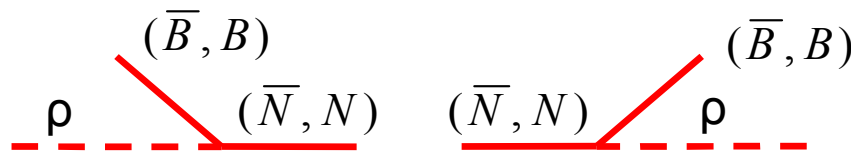


$$\frac{\text{Im } \bar{\Pi}_U^B}{m_\rho} \sim \text{vac} \otimes [ \{ (1 - n_N^+)(1 - n_B^-) - n_N^+ n_B^- \} + \{ (1 - n_N^-)(1 - n_B^+) - n_N^- n_B^+ \} ]$$

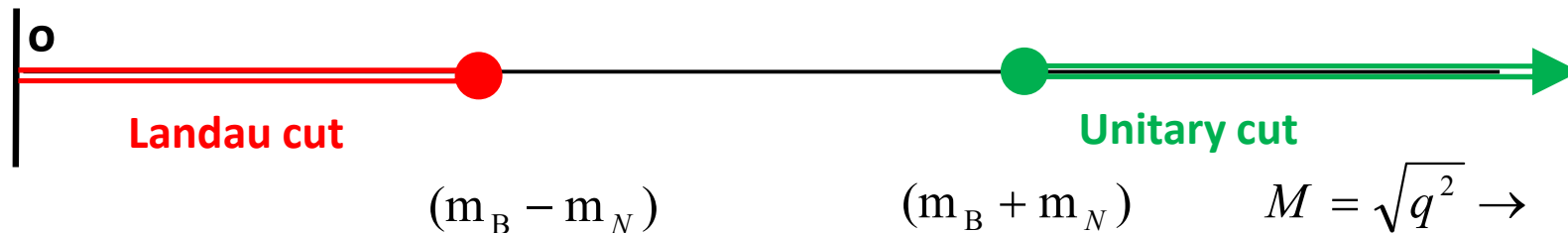


Pauli blocking of decay rate ( $\Gamma_{P.b.}^B$ )

$$\frac{\text{Im } \bar{\Pi}_L^B}{m_\rho} \sim \text{vac} \otimes [ \{ (1 - n_N^-) n_B^- - n_N^- (1 - n_B^-) \} + \{ (1 - n_N^+) n_B^+ - n_N^+ (1 - n_B^+) \} ]$$



Bosonic collision rate ( $\Gamma_C^B$ )





**Landau cut**

$$(m_B - m_N)$$

$$=(0.94-0.94) \text{ GeV} = 0 \text{ GeV [min]}$$

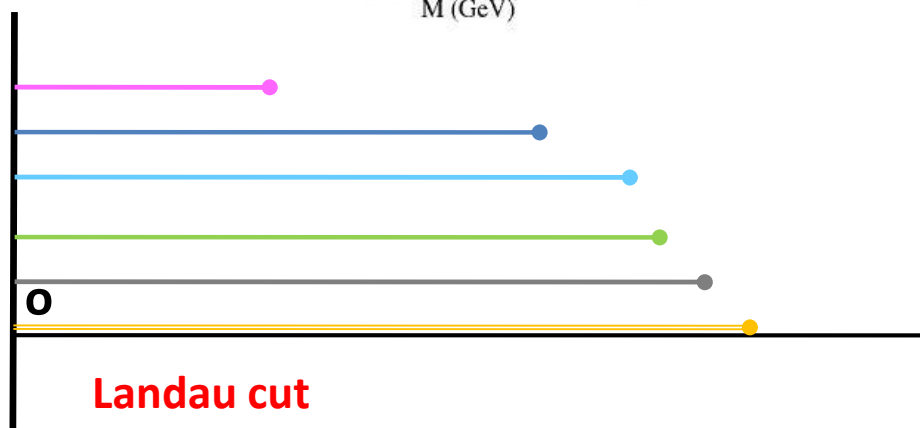
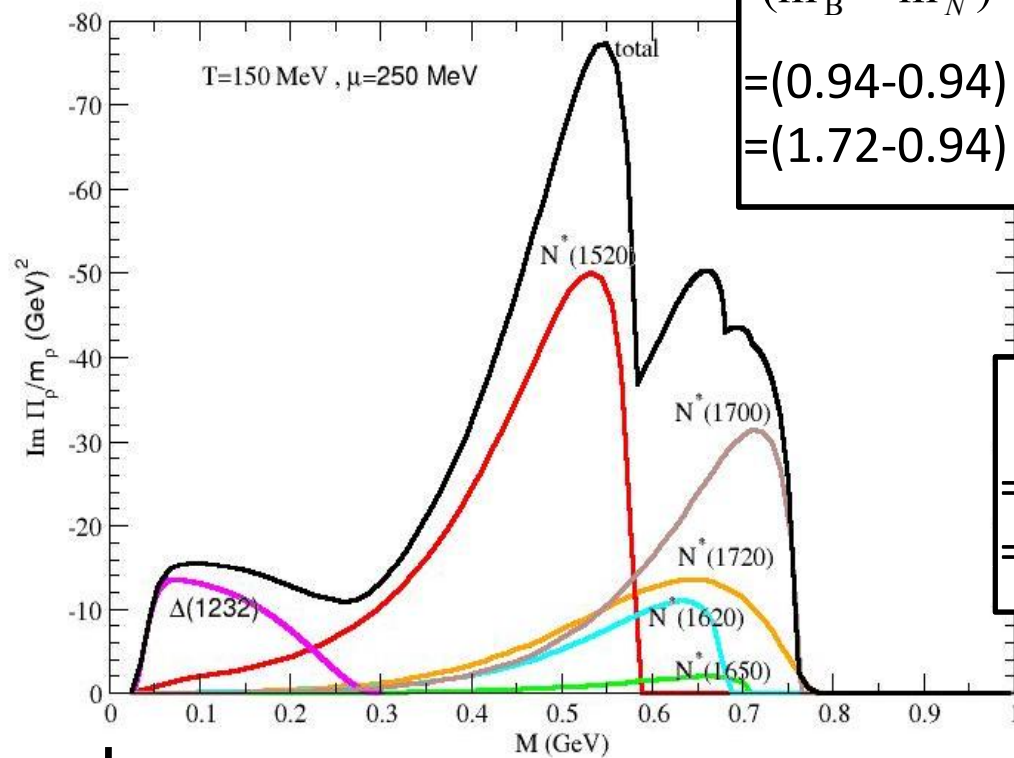
$$=(1.72-0.94) \text{ GeV} = 0.78 \text{ GeV [max]}$$

**Unitary cut**

$$(m_B + m_N)$$

$$=(0.94+0.94) \text{ GeV} = 1.88 \text{ GeV [min]}$$

$$=(1.72+0.94) \text{ GeV} = 2.66 \text{ GeV [max]}$$



$\Delta(1232)$

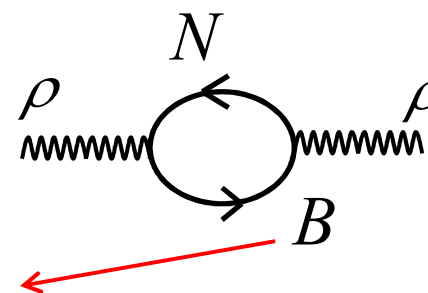
$N^*(1520)$

$\Delta^*(1620)$

$N^*(1650)$

$\Delta^*(1700)$

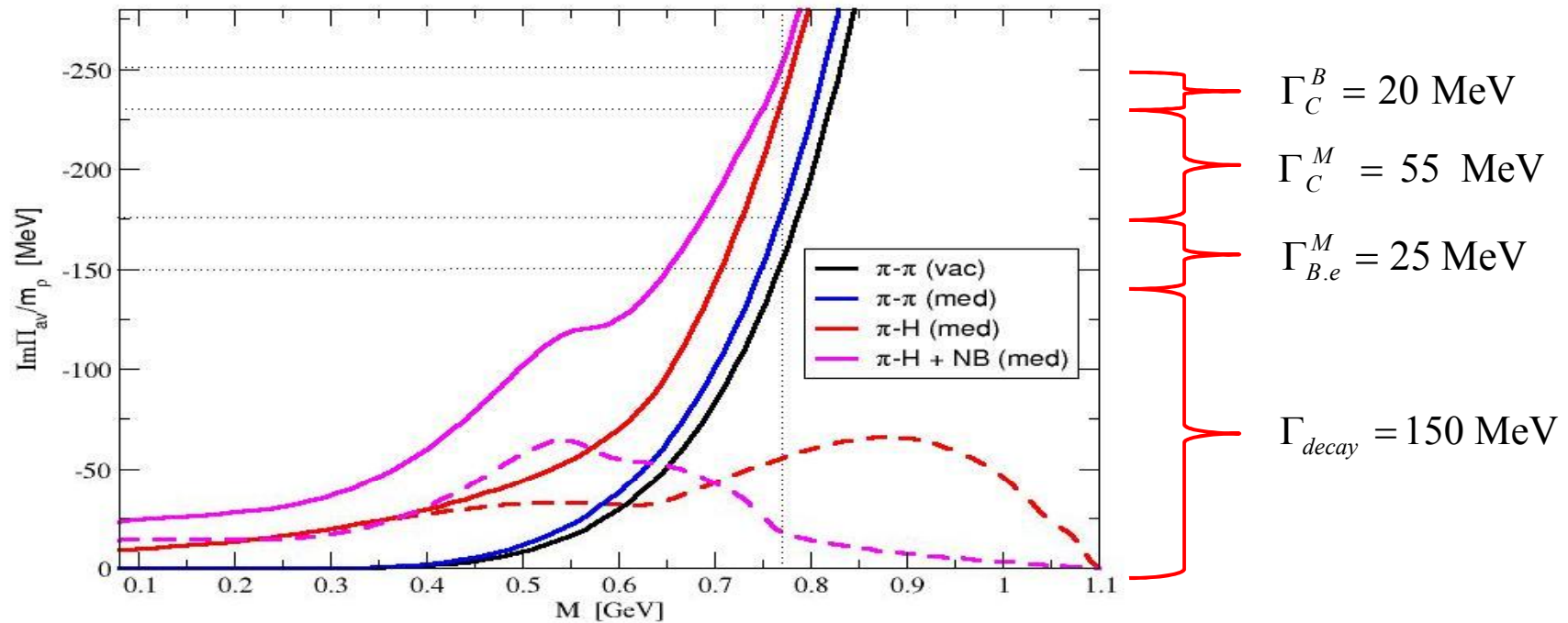
$N^*(1720)$



**Unitary cut**

$$M = \sqrt{q^2} \rightarrow$$

$$T = 150 \text{ MeV} \& \mu_B = 250 \text{ MeV}$$



$$\Gamma_{decay}^{\pi\pi} = 150 \text{ MeV}$$

vacuum part of  $\pi\pi$   
loop (unitary cut)

$$\Gamma_{B.e.}^M = 25 \text{ MeV}$$

thermal part of  $\pi\pi$   
loops (unitary cut)

$$\Gamma_C^M = 55 \text{ MeV}$$

( $\pi H =$   
 $\pi\omega, \pi h_1, \pi a_1$ )  
loops (Landau cut)

$$\Gamma_C^B = 20 \text{ MeV}$$

( $NB =$   
 $NN, N\Delta, NN^*, N\Delta^*$ )  
loops (Landau cut)

# Physical interpretation of imaginary part of in-medium self-energy :

$$\Gamma_{tot} = \Gamma_C^M + \Gamma_{B.e.}^M + \Gamma_C^B = \frac{\text{Im } \Pi(E_q, \vec{q})}{m_\rho}$$



rate at which  $\rho$  try to be thermalized with the thermal bath

$$\frac{1}{\exp(E_q / T) - 1} + c \exp(-\Gamma_{tot} t)$$



$$\frac{1}{\exp(E_q / T) - 1}$$



Thermalized Hadronic matter with mesons(H) and baryons(B)

$\pi(140)$

$\Delta(1232)$

$\omega(782)$

$N^*(1520)$

$h_1(1170)$

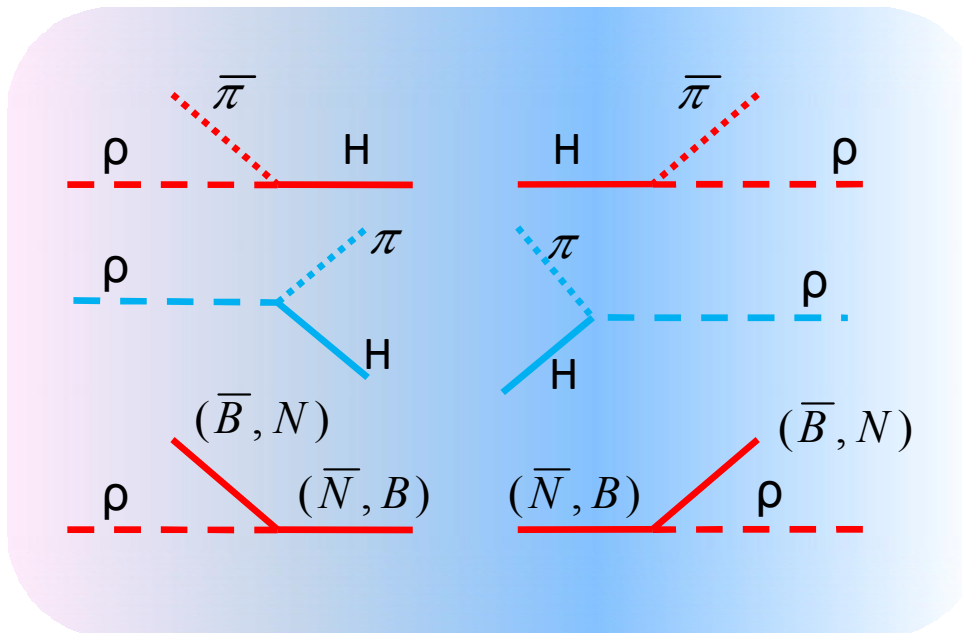
$\Delta^*(1620)$

$a_1(1260)$

$N^*(1650)$

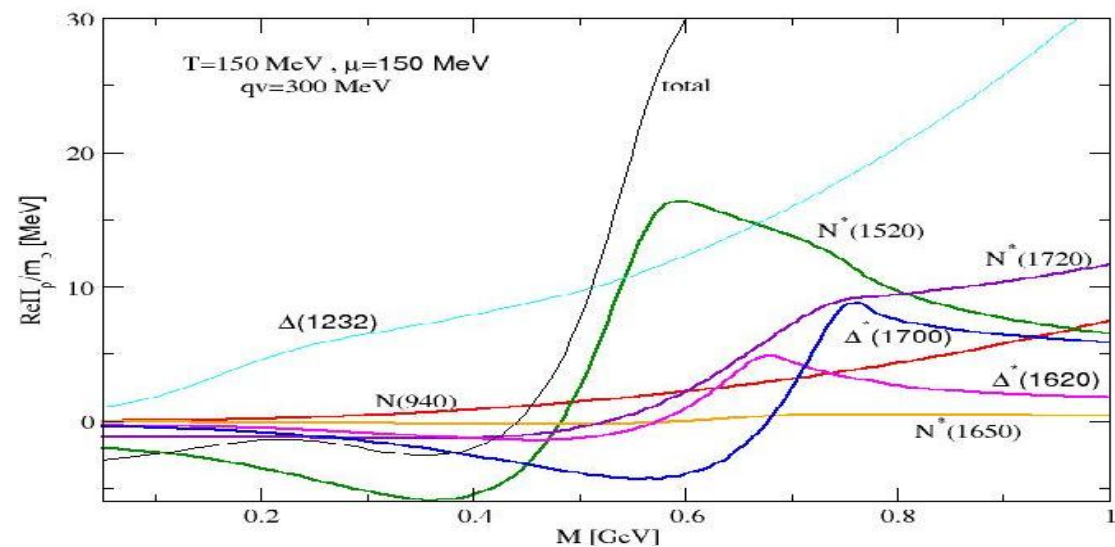
$\Delta^*(1700)$

$N^*(1720)$

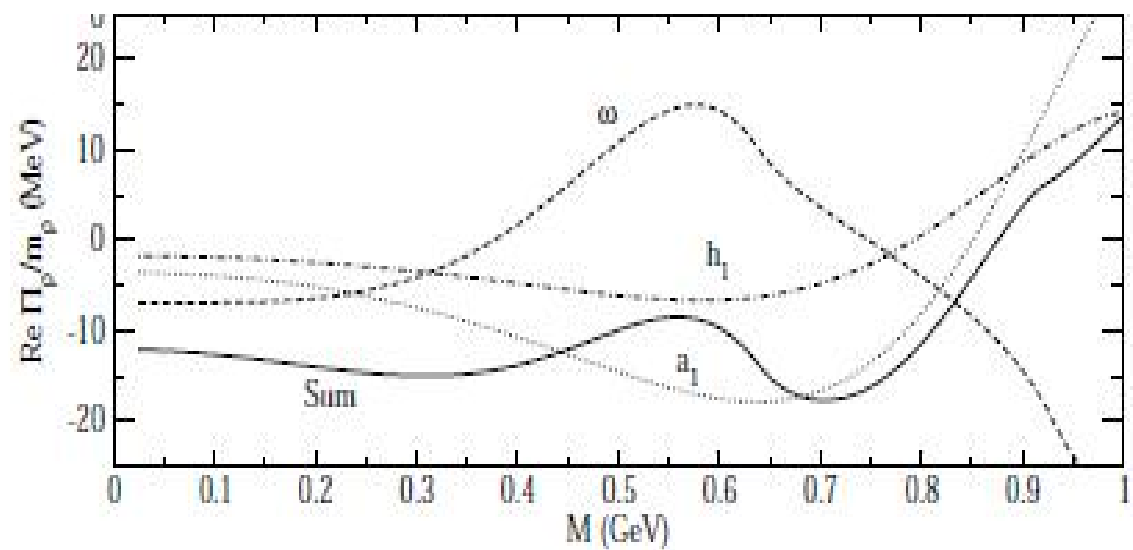


Real part  
of self-energy

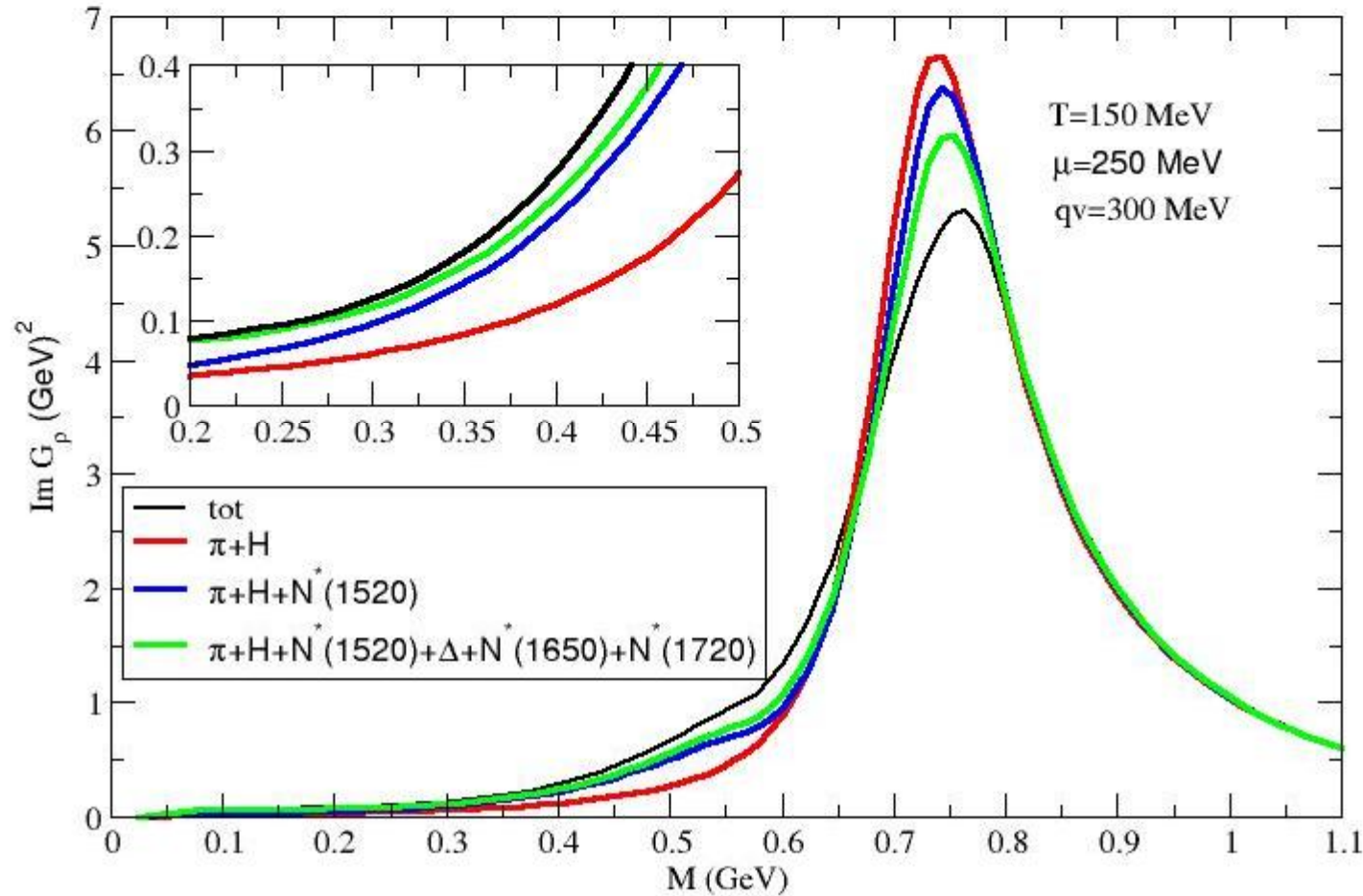
Baryonic loops



Mesonic loops

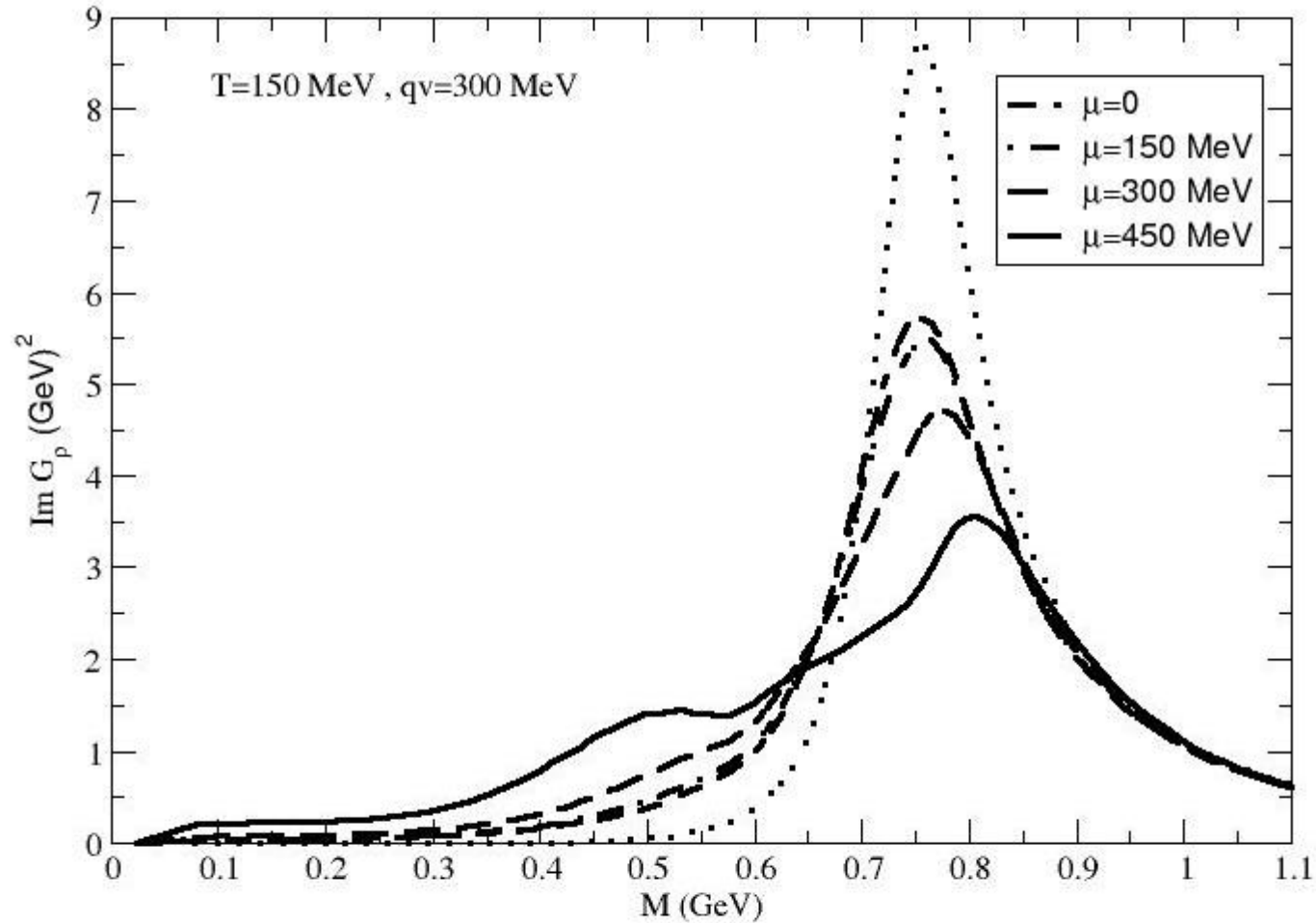


Effect of various loops on low mass invariant mass space  
in  $\rho$  spectral function :



$\pi\omega$  ,  $NN^*(1520)$  &  $N\Delta(1232)$

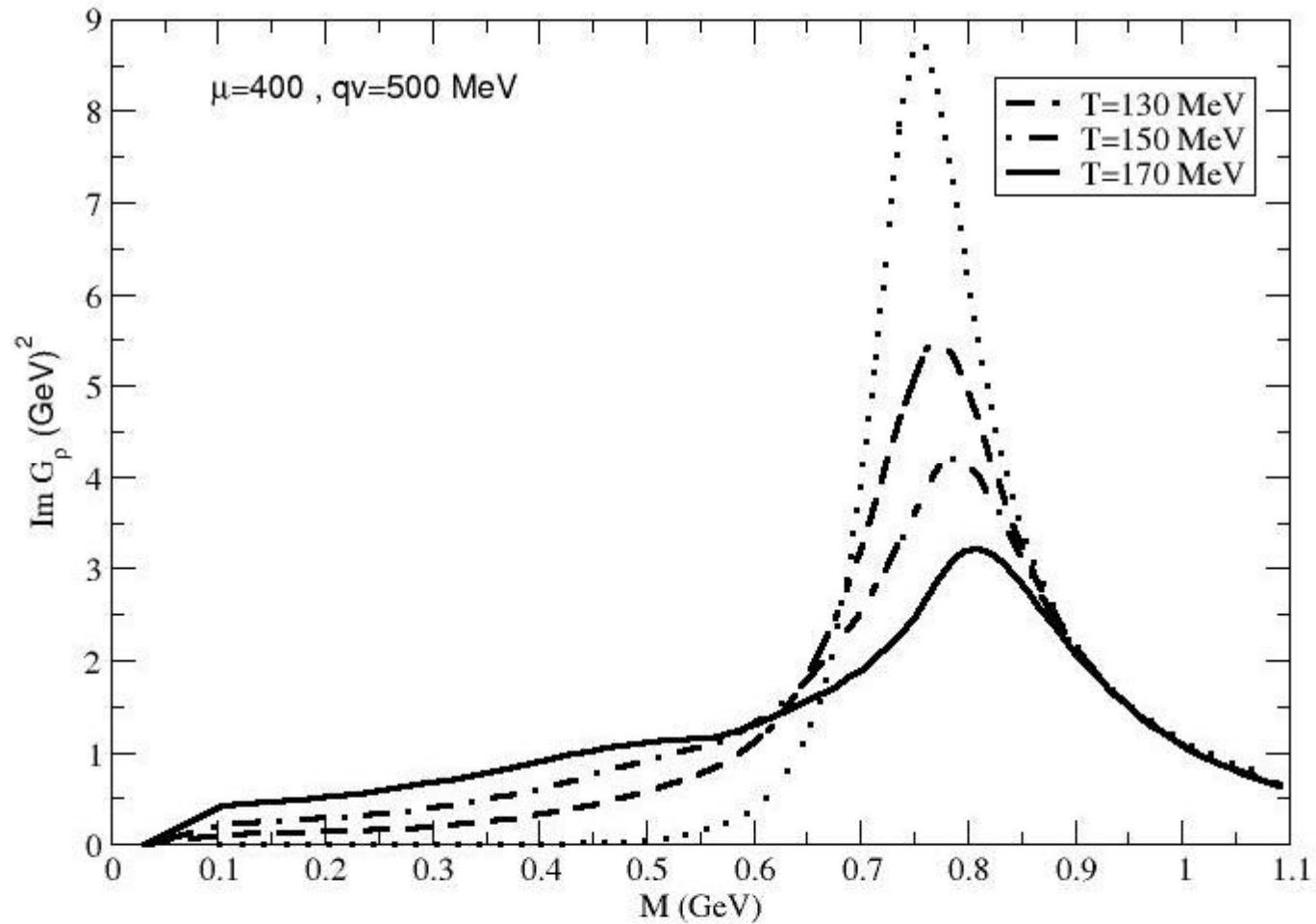
Effect of baryonic chemical potential on  $\rho$  spectral function in low mass region:



$$A_\rho(q^0, \vec{q}, T, \mu_B)$$

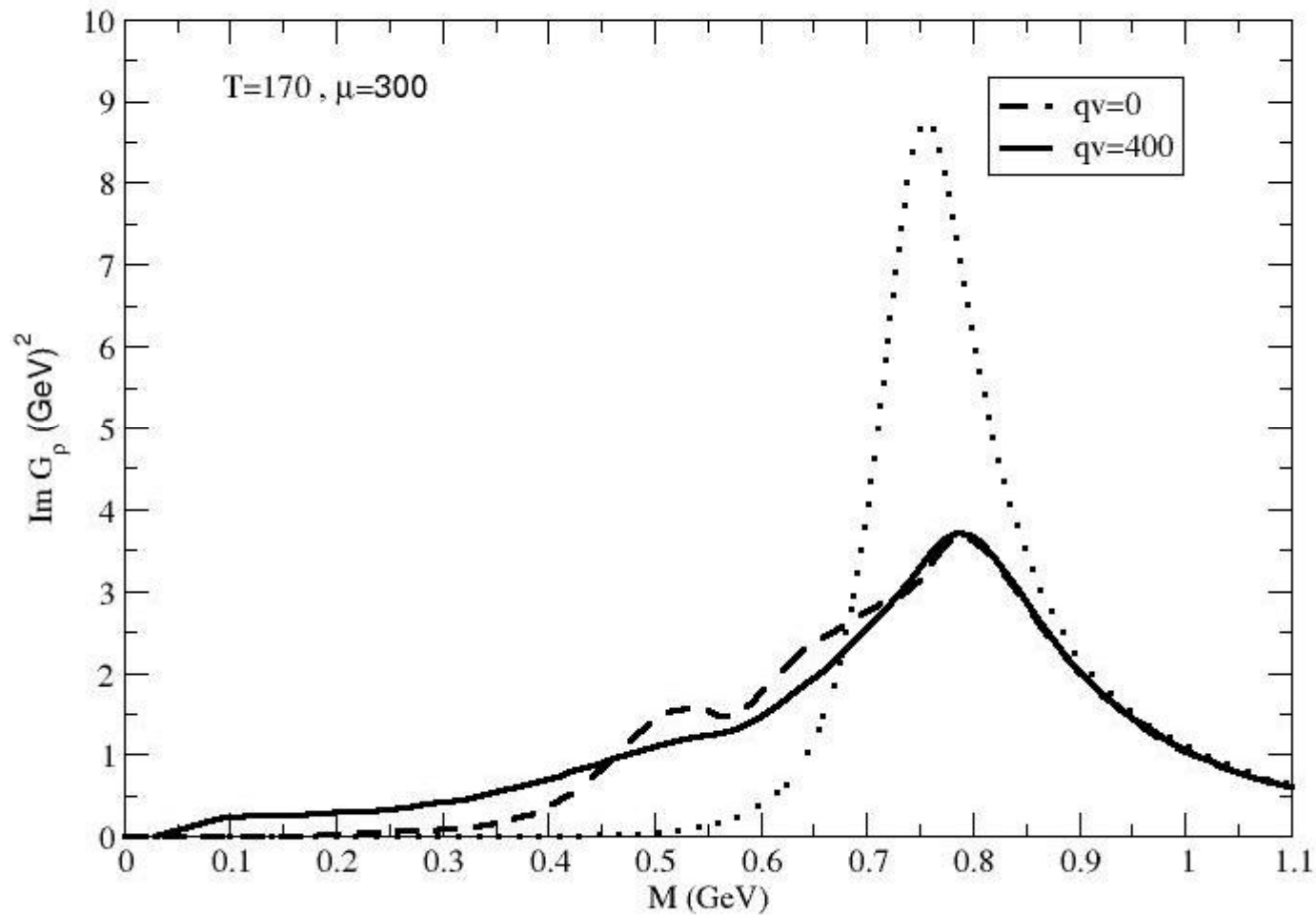


Effect of temperature on  $\rho$   
spectral function in low mass region:



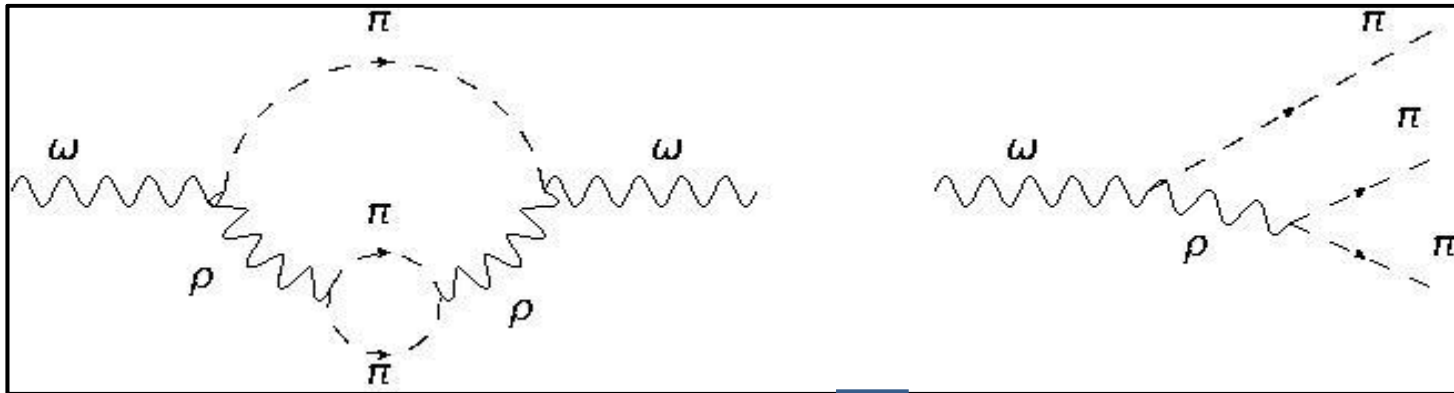
$$A_\rho(q^0, \vec{q}, T, \mu_B)$$

Effect of momentum of  $\rho$  in off mass shell on its spectral function in low mass region:

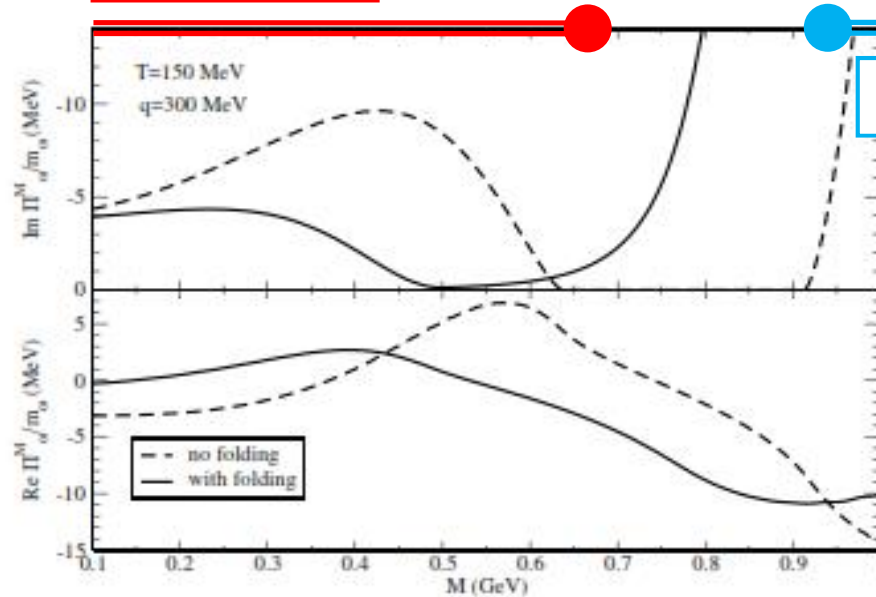


$$A_\rho(q^0, \vec{q}, T, \mu_B)$$

# Self Energy of $\omega$ for mesonic loops :



Landau cut

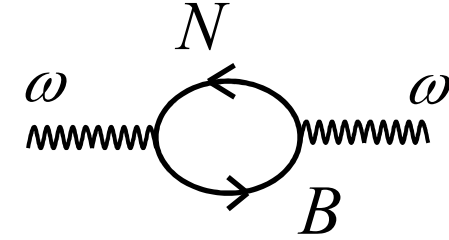
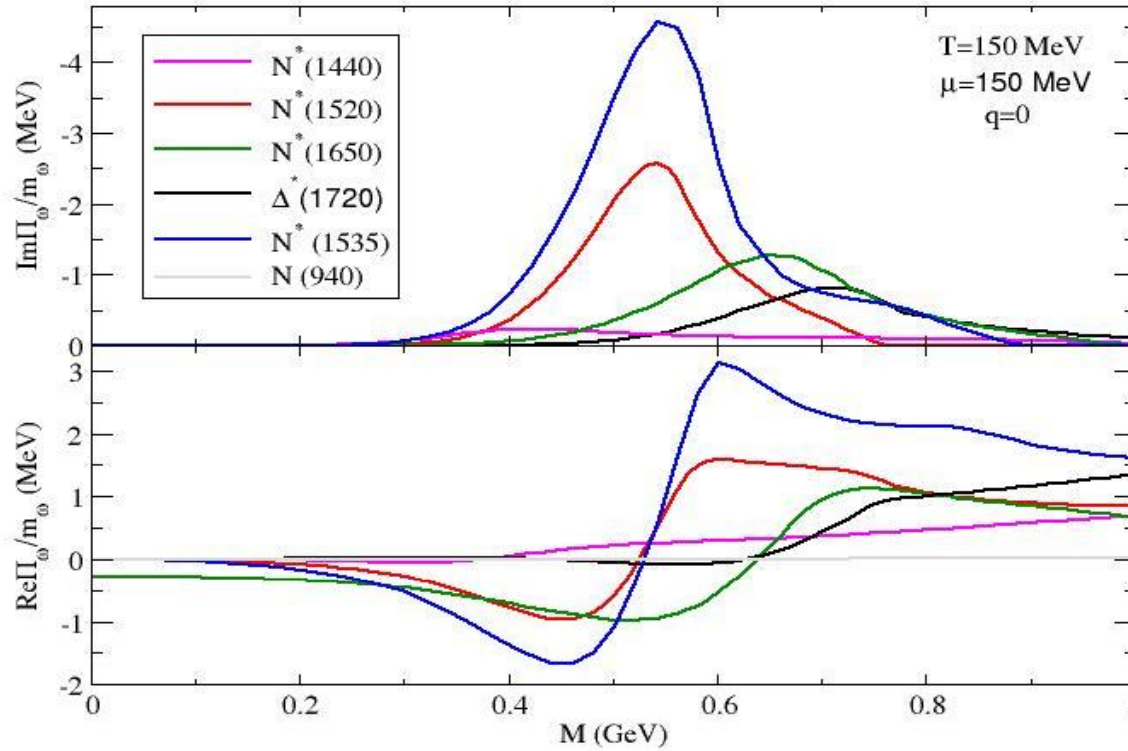


Unitary cut

$$\bar{\Pi}_M^{\mu\nu}(q) = \frac{1}{N_\rho} \int_{4m_\pi^2}^{(q-m_\pi)^2} dM^2 [\bar{\Pi}_{(\rho\pi)}^{\mu\nu}(q, M)] A_\rho(M)$$

$$\mathcal{L}_{int} = \frac{g_m}{F_\pi} \epsilon_{\mu\nu\lambda\sigma} (\partial^\nu \omega^\mu \bar{\rho}^\lambda - \omega^\mu \partial^\nu \bar{\rho}^\lambda) \cdot \partial^\sigma \vec{\pi}$$

# Self Energy of $\omega$ for baryonic loops :



S. Ghosh & S. Sarkar  
Eur. Phys. J. A 49 (2013) 97

P. Muehlich, V. Shklyar, S. Leupold, U. Mosel, M. Post, Nucl. Phys. A **780**, 187 (2006).

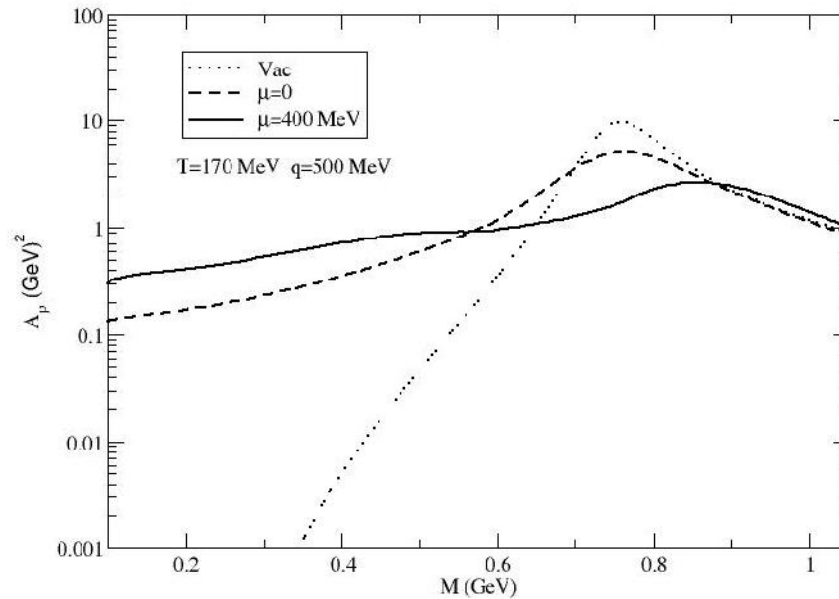
$$\mathcal{L} = -[\bar{\psi}_R(g_1\gamma_\mu - \frac{g_2}{2m_N}\sigma_{\mu\nu}\partial^\nu)\psi_N\omega^\mu + h.c.] \quad J_R^P = \frac{1}{2}^+$$

$$\mathcal{L} = i[\bar{\psi}_R\gamma^5(g_1\gamma_\mu - \frac{g_2}{2m_N}\sigma_{\mu\nu}\partial^\nu)\psi_N\omega^\mu + h.c.] \quad J_R^P = \frac{1}{2}^-$$

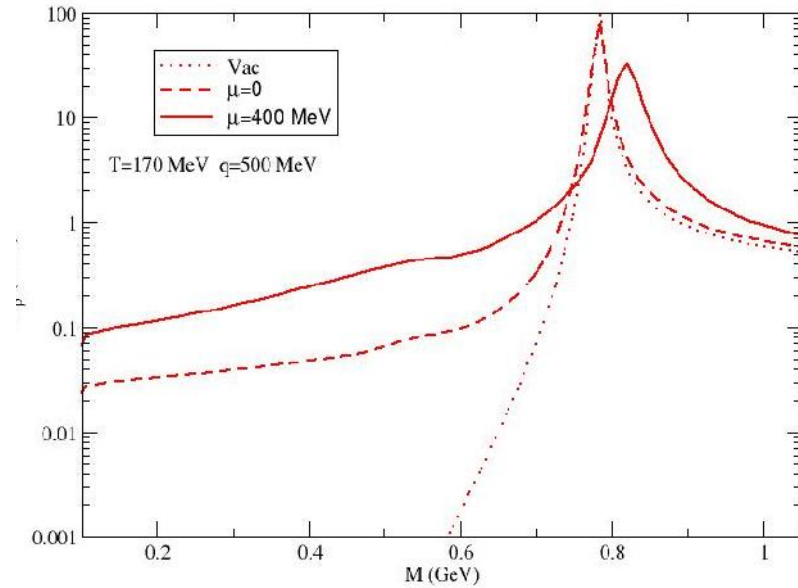
$$\mathcal{L} = -i[\bar{\psi}_R^\mu\gamma^5(\frac{g_1}{2m_N}\gamma^\alpha i\frac{g_2}{4m_N^2}\partial_N^\alpha + i\frac{g_3}{4m_N^2}\partial_\omega^\alpha)(\partial_\alpha^\omega\mathcal{O}_{\mu\nu} - \partial_\mu^\omega\mathcal{O}_{\alpha\nu})\psi_N\omega^\nu + h.c.] \quad J_R^P = \frac{3}{2}^+$$

$$\mathcal{L} = -[\bar{\psi}_R^\mu(\frac{g_1}{2m_N}\gamma^\alpha i\frac{g_2}{4m_N^2}\partial_N^\alpha + i\frac{g_3}{4m_N^2}\partial_\omega^\alpha)(\partial_\alpha^\omega\mathcal{O}_{\mu\nu} - \partial_\mu^\omega\mathcal{O}_{\alpha\nu})\psi_N\omega^\nu + h.c.] \quad J_R^P = \frac{3}{2}^-$$

## $\rho$ meson spectral function

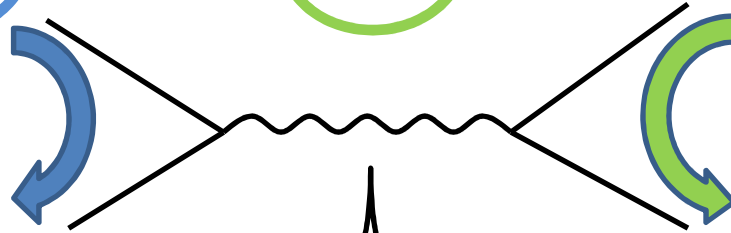


## $\omega$ meson spectral function



## Formalism of dilepton :

$$L_{\text{int}} = e \underbrace{J_{\mu}^h(x)}_{\text{blue circle}} A^{\mu}(x) + e \underbrace{J_{\mu}^l(x)}_{\text{green circle}} A^{\mu}(x)$$



$$S_{\text{fi}} \sim \left\langle F \left| \int d^4x d^4y \underbrace{J_{\mu}^h(x) A^{\mu}(x)}_{\text{blue box}} \underbrace{A^{\nu}(y) J_{\nu}^l(y)}_{\text{green box}} \right| I \right\rangle$$

$$\frac{dN}{d^4x d^4q} = (\dots) \underbrace{L_{\mu\nu}}_{\text{green box}} \underbrace{W^{\mu\nu}}_{\text{blue box}}$$

L. D. McLerran and T. Toimela,  
Phys. Rev. D **31**, 545 (1985).

H. A. Weldon,  
Phys. Rev. D **42**, 2384 (1990).

current correlator  $W_{\mu\nu}$  is defined by

$$W_{\mu\nu}(q_0, \vec{q}) = \int d^4x e^{iq \cdot x} \langle [J_\mu^{em}(x), J_\nu^{em}(0)] \rangle$$

$$\begin{aligned} J_\mu^h &= \frac{1}{2}(\bar{u}\gamma_\mu u - \bar{d}\gamma_\mu d) + \frac{1}{6}(\bar{u}\gamma_\mu u + \bar{d}\gamma_\mu d) + \dots \\ &= J_\mu^V + J_\mu^S + \dots \\ &= J_\mu^\rho + J_\mu^\omega/3 + \dots \end{aligned}$$

$$W_{\mu\nu} = 2\epsilon(q_0)F_\rho^2 m_\rho^2 \text{Im} \bar{D}_{\mu\nu}^\rho + 2\epsilon(q_0)F_\omega^2 m_\omega^2 \text{Im} \bar{D}_{\mu\nu}^\omega + \dots$$

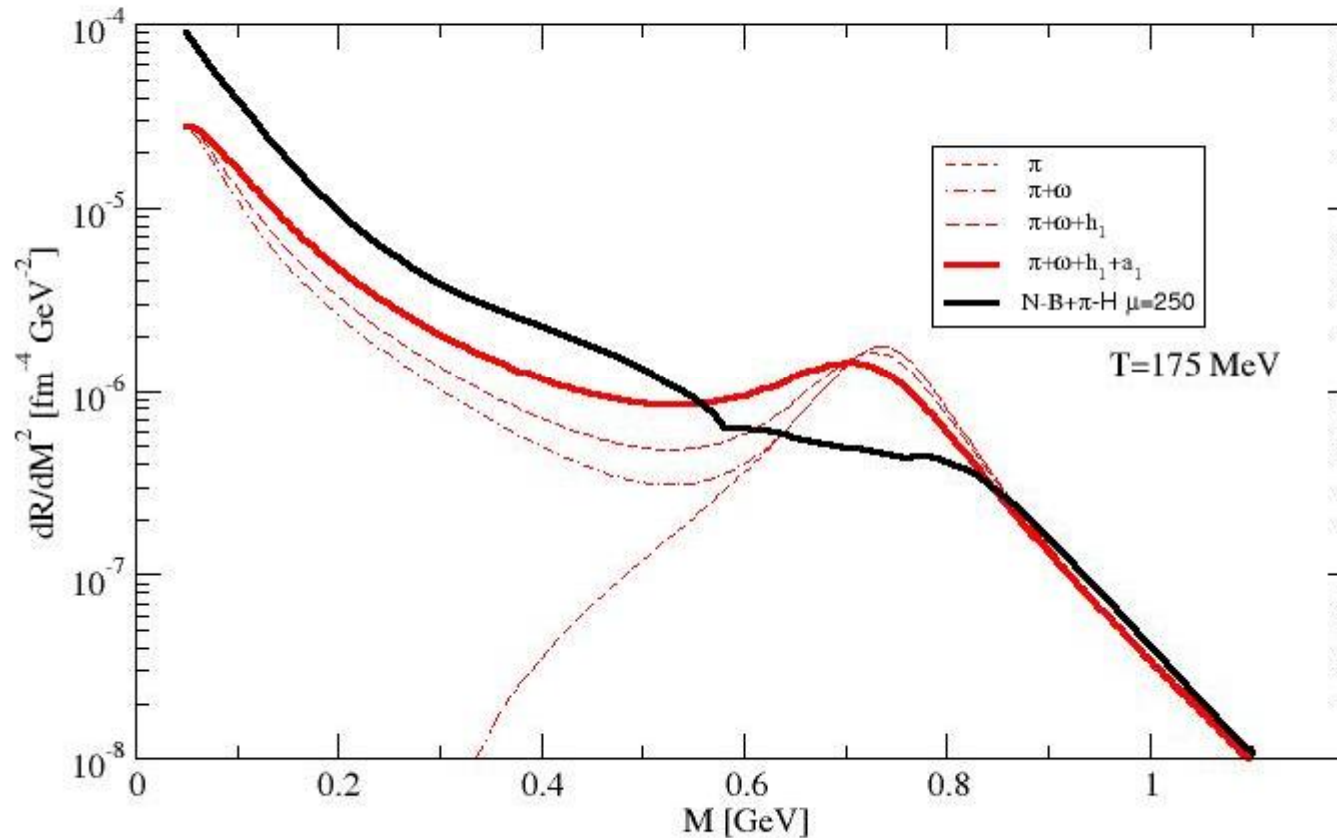
$$\bar{D}_{\mu\nu}(q) = -\frac{P_{\mu\nu}}{q^2 - m_\rho^2 - \bar{\Pi}_t(q)} - \frac{Q_{\mu\nu}/q^2}{q^2 - m_\rho^2 - q^2 \bar{\Pi}_l(q)} - \frac{q_\mu q_\nu}{q^2 m_\rho^2}$$

$$F_R^2 = \frac{3m_R \Gamma_{R \rightarrow e^+ e^-}}{4\pi\alpha^2} \quad F_R = 0.156 \text{ GeV}, 0.046 \text{ GeV} \text{ for } \rho, \omega$$

**Contribution of  $\omega$  is down by a factor  $\sim 10$**



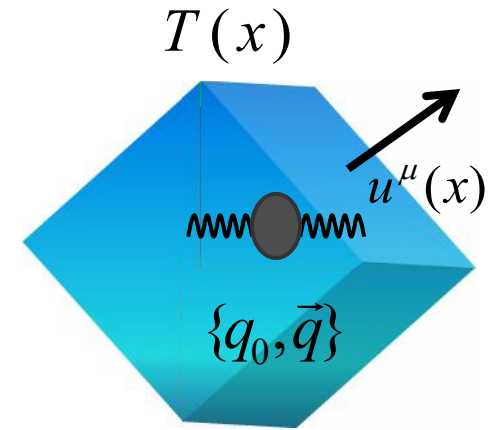
Effect of mesonic as well as baryonic medium modification of  $\rho$  on dilepton rate in low mass region :



$$\frac{dR}{dM^2 d^2 q_T dy} = \frac{dR}{d^4 q} = (\dots) L_{\mu\nu} A_{\rho}^{\mu\nu}(q^0, \vec{q}, T, \mu_B)$$

*Dilepton production in  
transverse momentum and  
invariant mass space :*

$$\frac{dN\{q_0, \vec{q}, T\}}{d^4x d^4q}$$



Fluid element

$$\frac{dN\{q_\mu u^\mu(x), \vec{q}, T(x)\}}{d^4x d^4q} \otimes d^4x$$

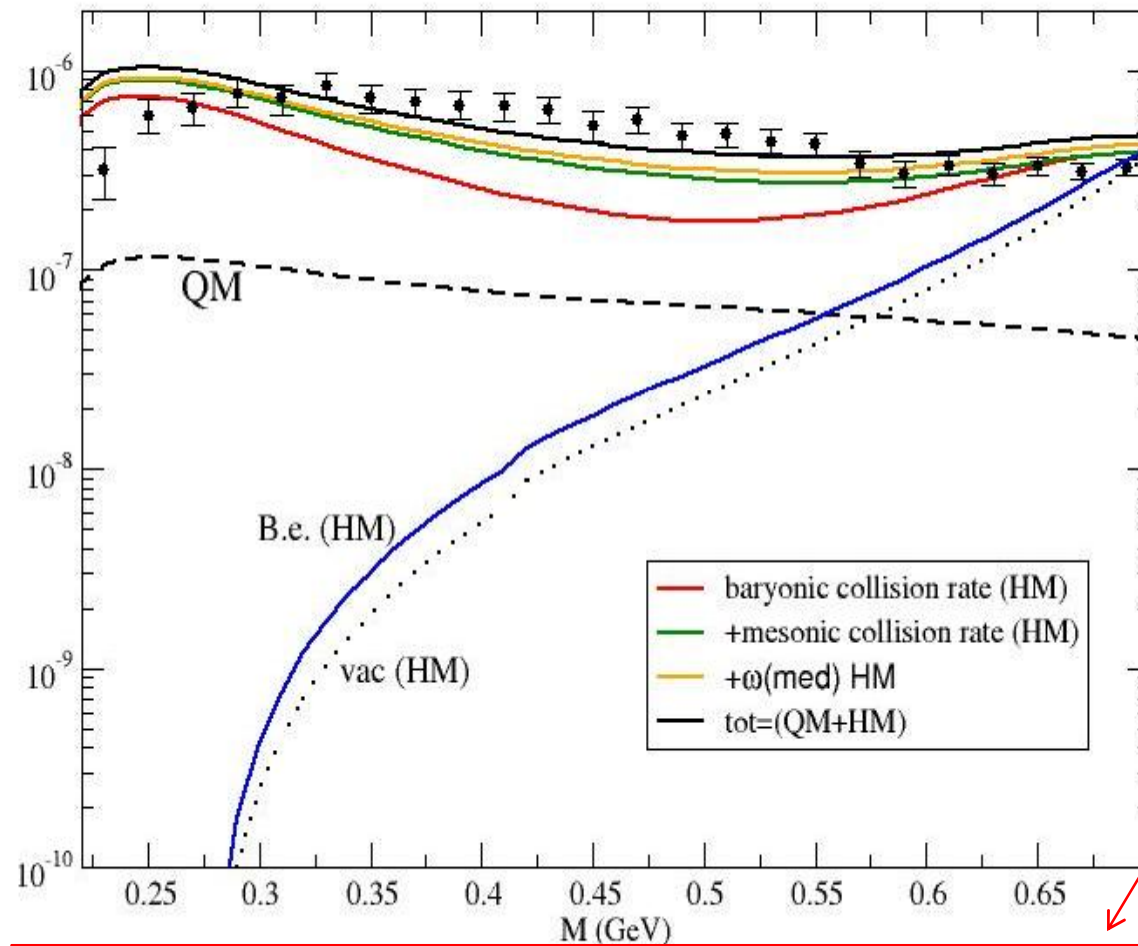
Invariant mass spectra

$$\frac{dN(M)}{dM} = \int d^4x (M dy d^2\vec{q}_T) \left[ \frac{dN\{q_\mu u^\mu(x), \vec{q}, T(x)\}}{d^4x d^4q} \right]$$

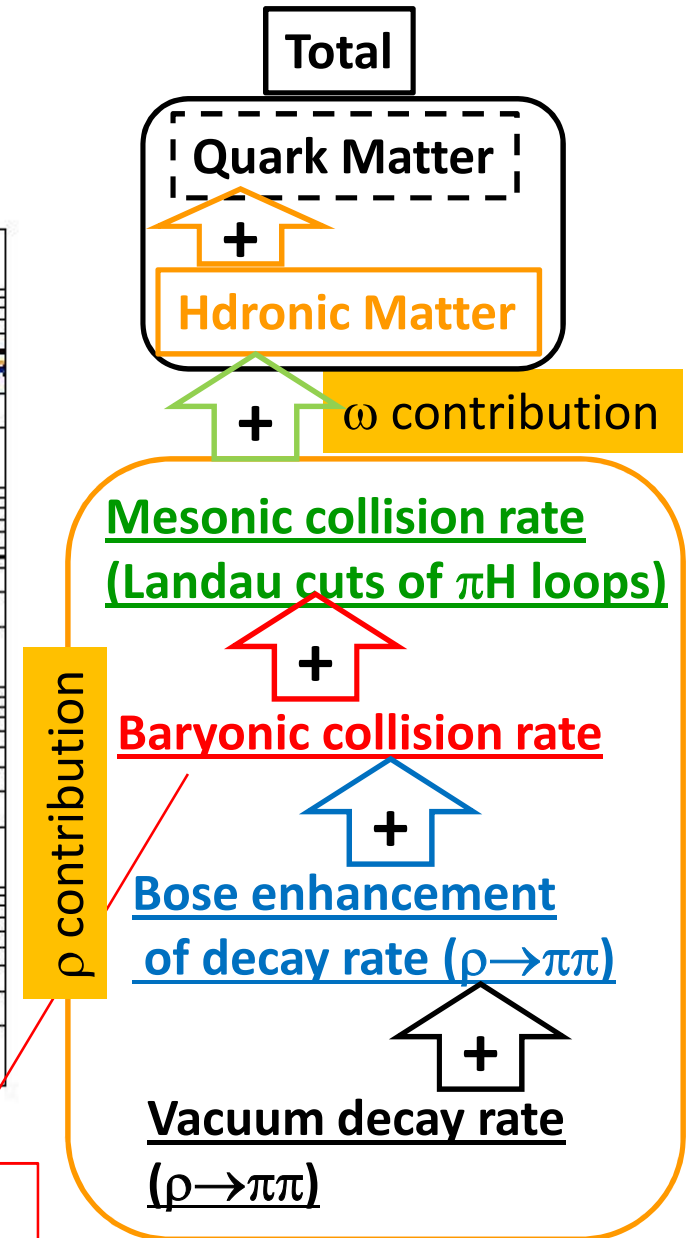
Transverse momentum spectra

$$\frac{dN(\vec{q}_T)}{\vec{q}_T d\vec{q}_T} = \int d^4x (2\pi M dM dy) \left[ \frac{dN\{q_\mu u^\mu(x), \vec{q}, T(x)\}}{d^4x d^4q} \right]$$

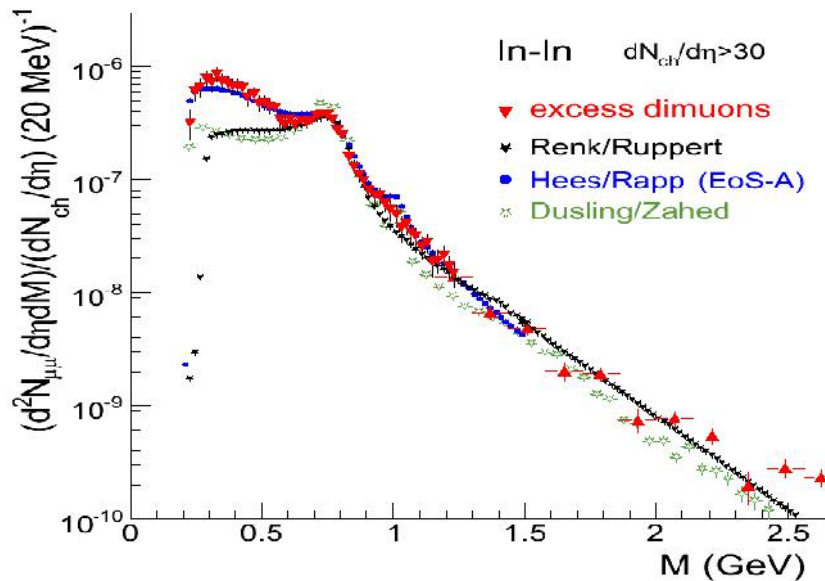
# Understanding low mass enhancement in the language of Thermal Field Theory :



Baryon part from  
Eltesky et. al. [Phys. Rev. C 64, (2001) 035202 ]

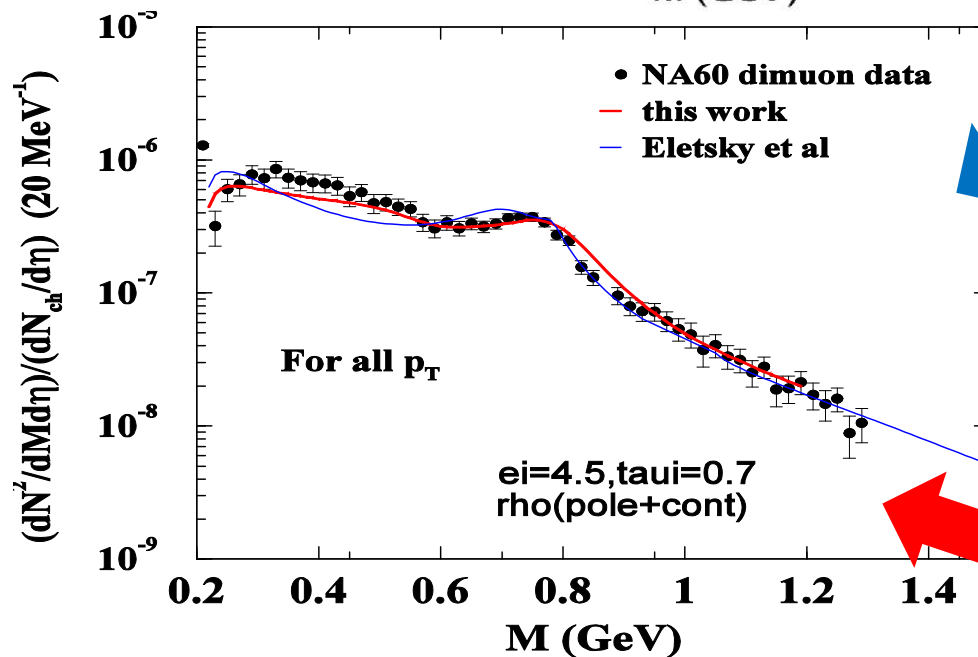


# Low mass enhancement at SPS :



J K Nayak, J Alam, T Hirano, S Sarkar and B Sinha  
**Phys.Rev. C85 (2012) 064906**

**Meson loop self-energies**  
 (S.Ghosh, S.Mallik. S.Sarkar  
 Eur. Phys. C 70, (2010) 251)  
 + **Baryon part from Eltesky et. al.**  
 (Phys. Rev. C 64, (2001) 035202)



S. Sarkar & S. Ghosh  
**J.Phys.Conf.Ser. 374 (2012) 012010**

**Meson (S.Ghosh, S.Mallik. S.Sarkar**  
 Eur. Phys. C 70, (2010) 251)  
 + **Baryon (S.Ghosh, S.Sarkar**  
 Nucl. Phys. A 870, (2011) 94)  
**-loop self-energies**

## Collaborators at VECC (India)

Analytic structure of  $\rho$  meson propagator at finite temperature

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# THANK U...