

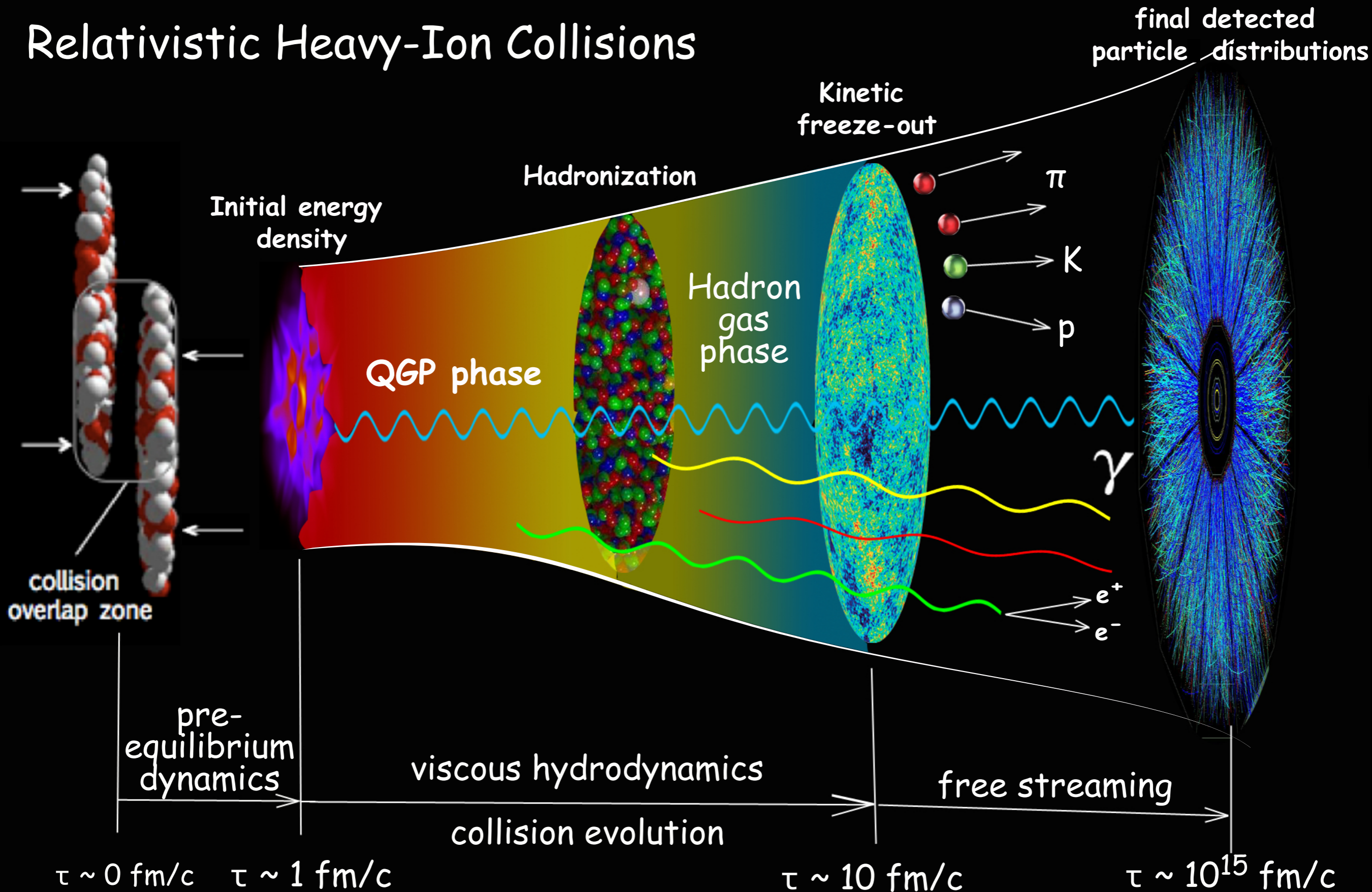
Photon tomography of relativistic heavy-ion collisions

Chun Shen
The Ohio State University

In collaboration with Jean-Francois Paquet, Gabriel Denicol,
Ulrich Heinz, Charles Gale

Little Bang

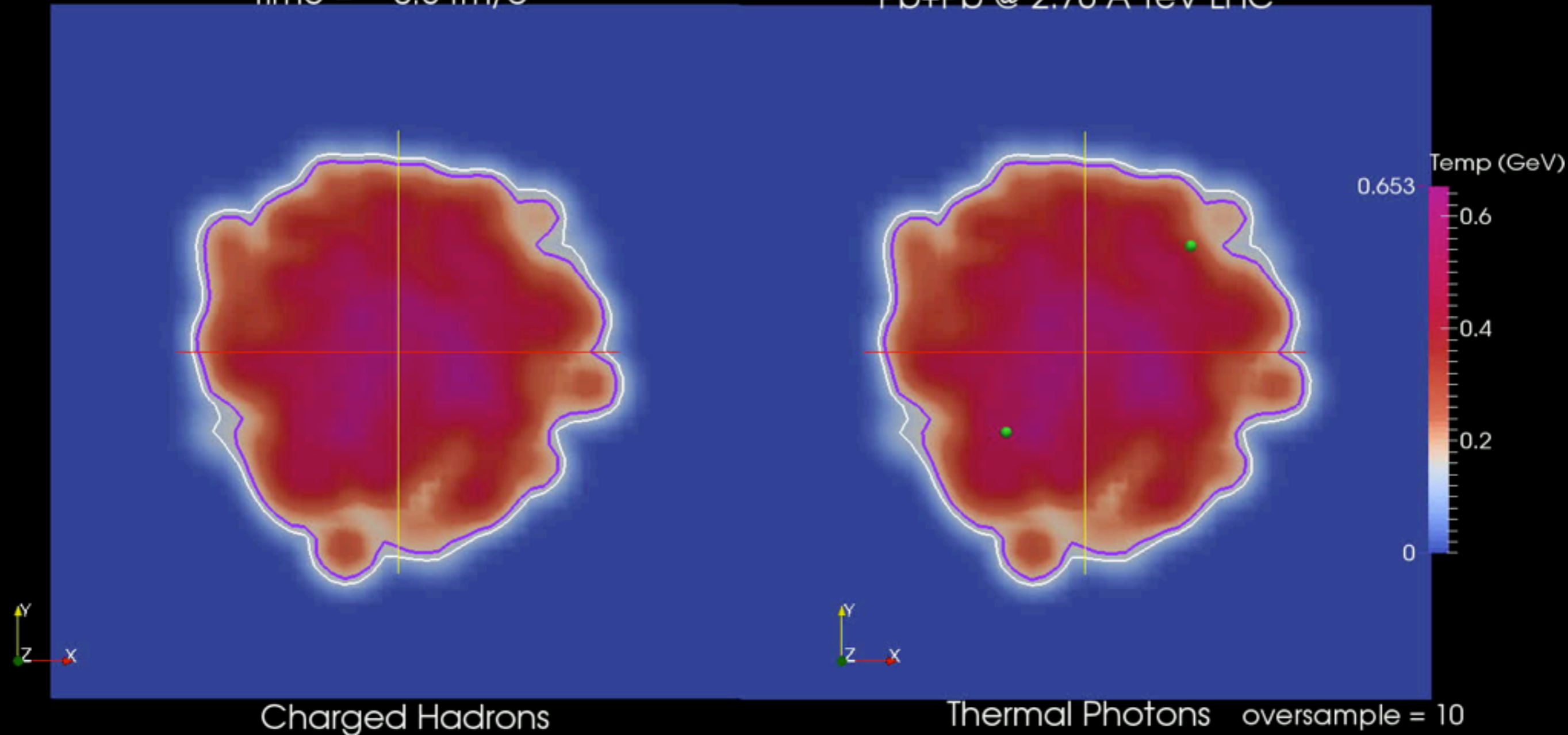
Relativistic Heavy-Ion Collisions



Photons from Heavy-ion Collisions

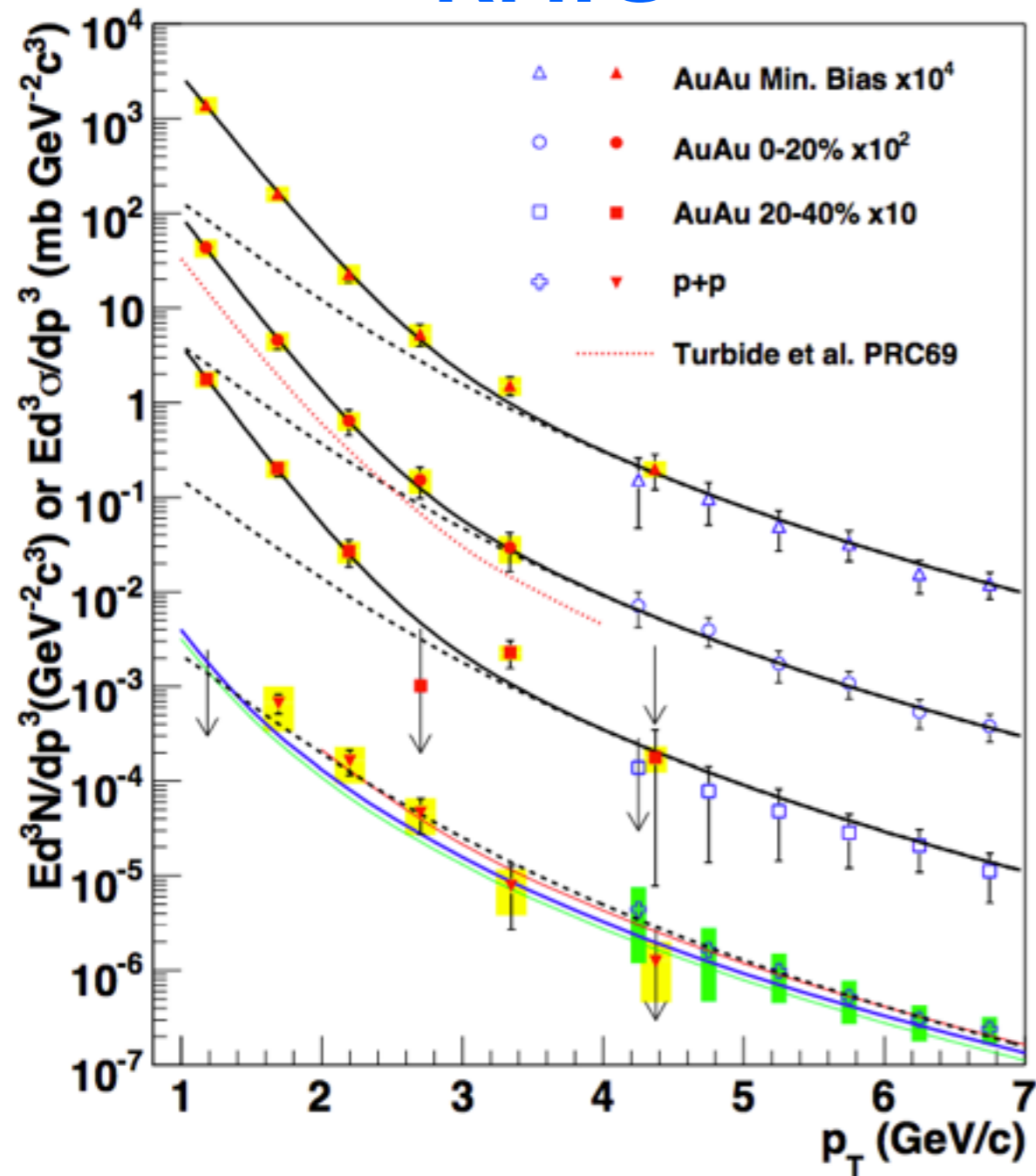
time = 0.6 fm/c

Pb+Pb @ 2.76 A TeV LHC



Fitted T_{eff} from Experiments

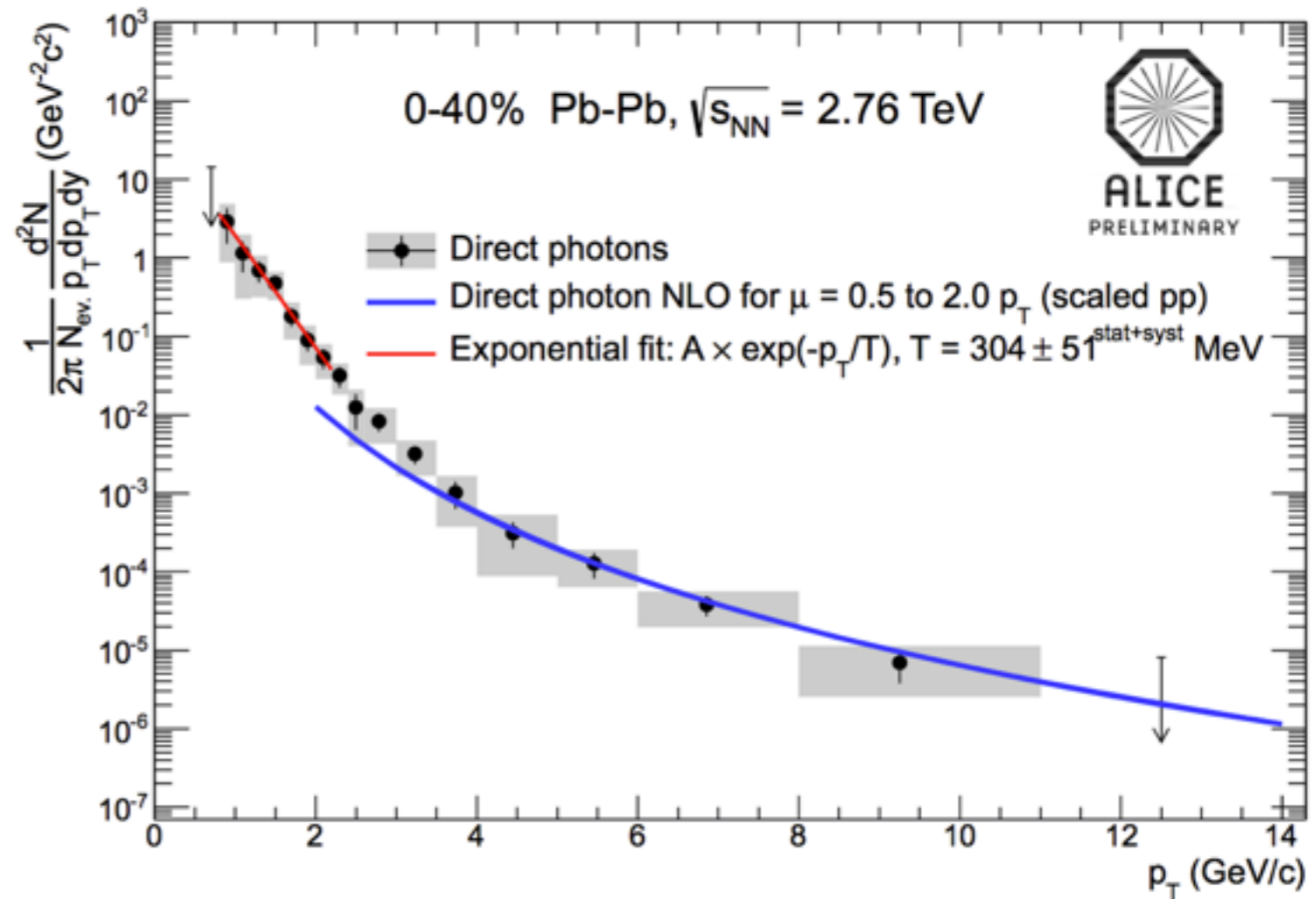
RHIC



0 – 20%

$$T = 221 \pm 19 \pm 19 \text{ MeV}$$

LHC



fit: $A \exp(-p_T/T)$

$$T = 304 \pm 51^{\text{stat+syst}} \text{ MeV}$$

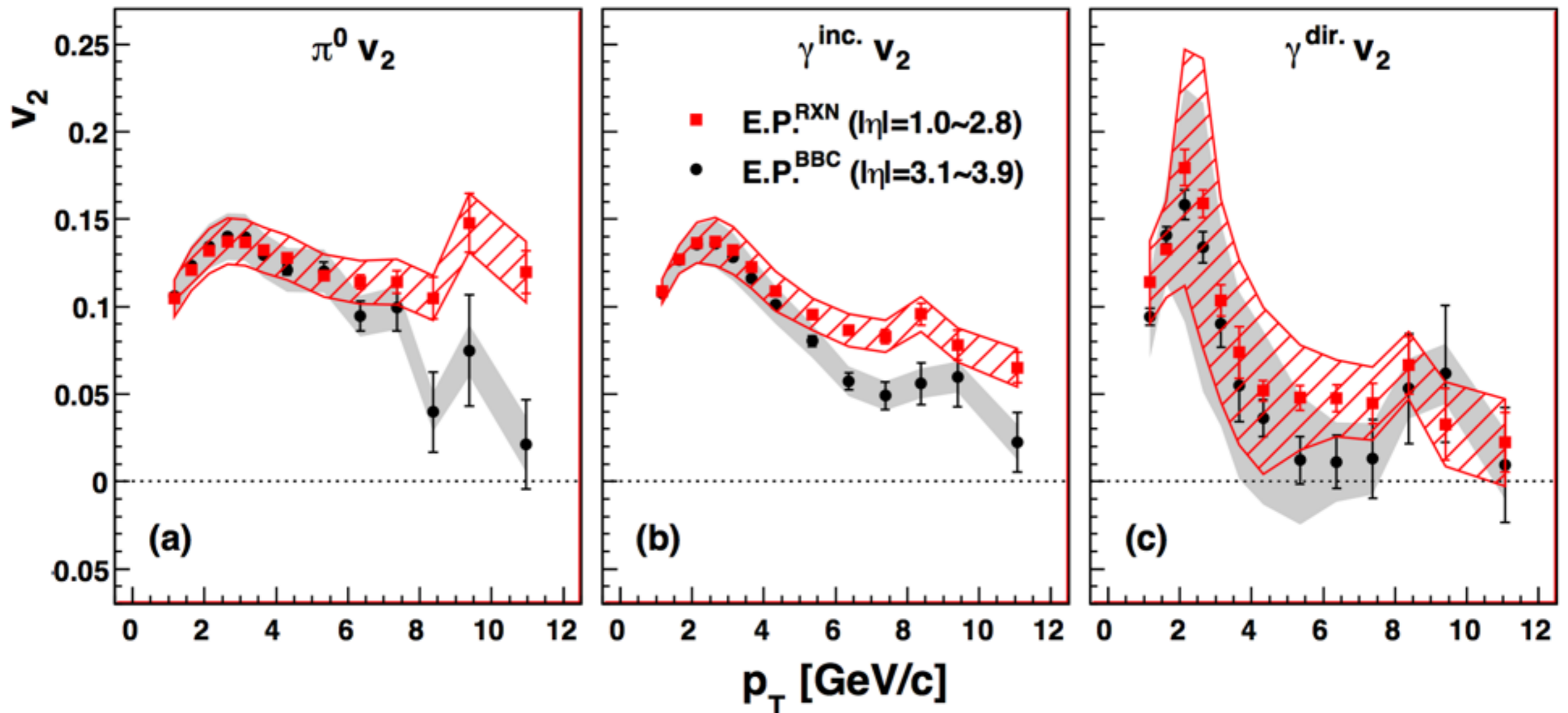


What does this T mean



Photon v_n from Experiment

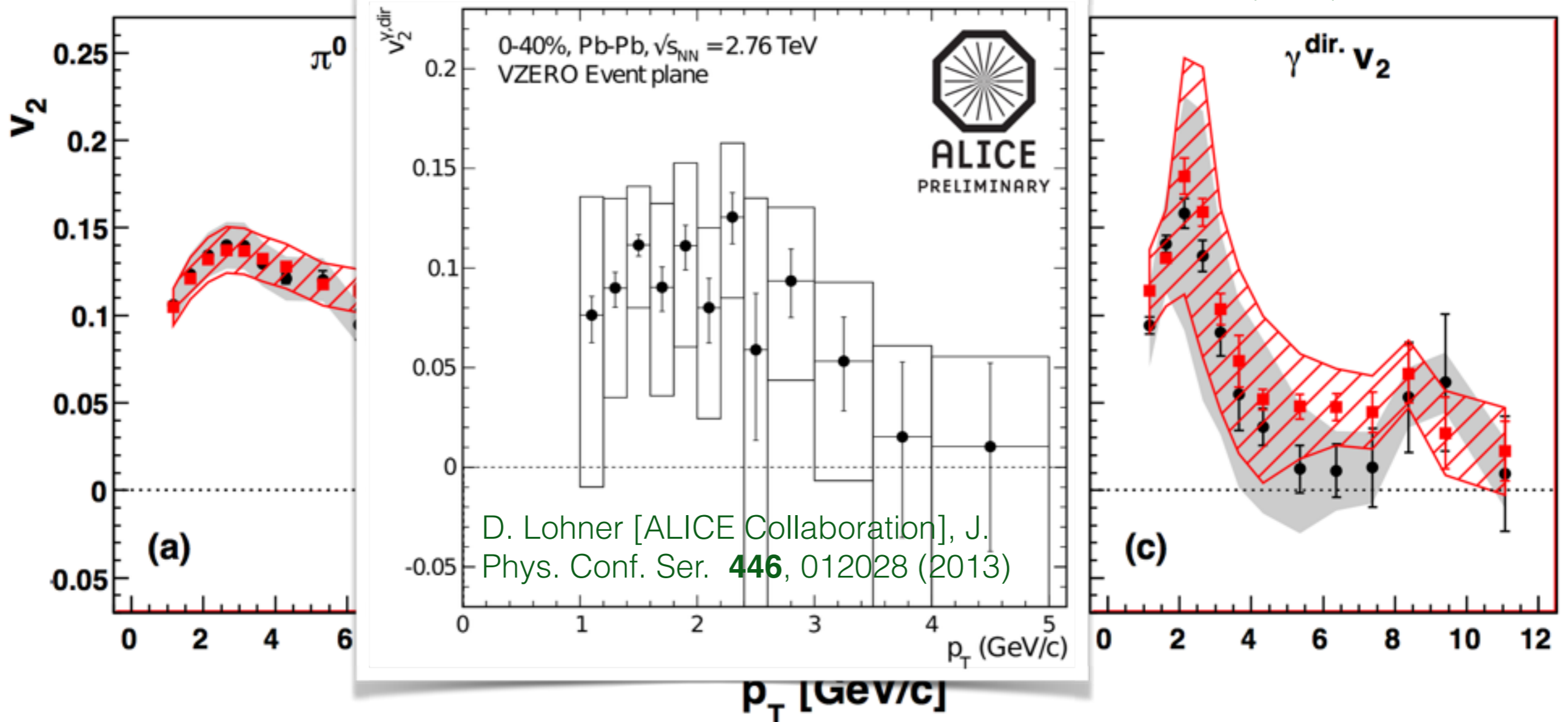
A. Adare *et al.* [PHENIX Collaboration] Phys. Rev. Lett. **109**, 122302 (2012)



- PHENIX measurements show **large** direct photon v_2 at $p_T < 4$ GeV

Photon v_n from Experiment

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- PHENIX measurements show **large** direct photon v_2 at $p_T < 4$ GeV
- ALICE also measured similar **large** direct photon elliptic flow at LHC

State-of-the-art hydrodynamic modeling

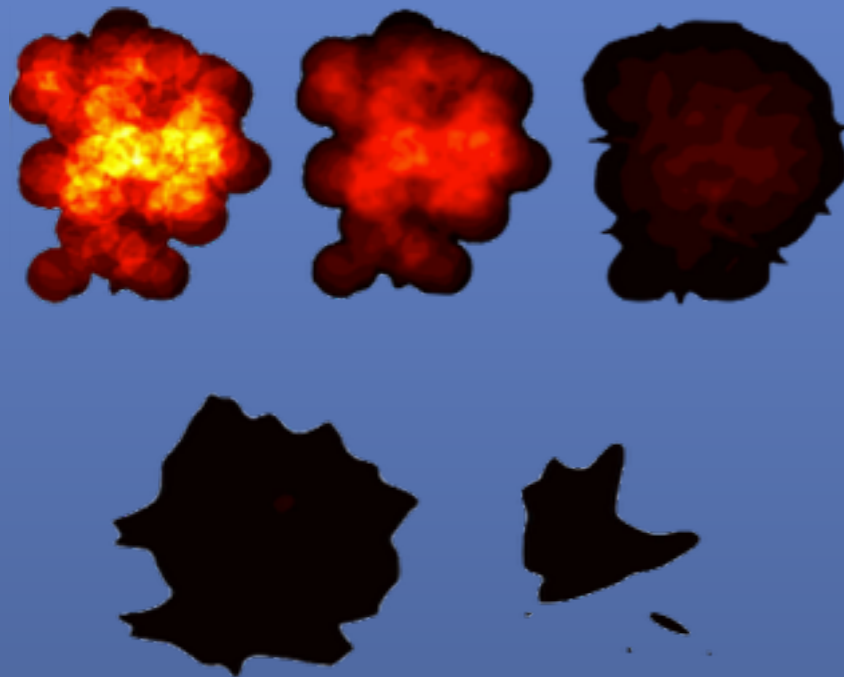
[https://github.com/
chunshen1987/iEBE.git](https://github.com/chunshen1987/iEBE.git)



Initial Condition
Generators
(MC-KLN, MC-Glauber)

Thermal Photon
Emission Rates

Hydrodynamic
Simulations
(VISH2+1)



HydroInfo
Package

$e, s, p, T,$
 $u^\mu, \pi^{\mu\nu}$

Thermal Photon
Interface

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} q_\mu q_\nu}{2(e+p)} a_{\alpha\beta} \Gamma^{\alpha\beta}$$
$$E \frac{dN^\gamma}{d^3p} = \int d^4x q \frac{dR}{d^3q}$$

Hadrons spectra &
 V_n

Photon spectrum &
 V_n

State-of-the-art hydrodynamic modeling

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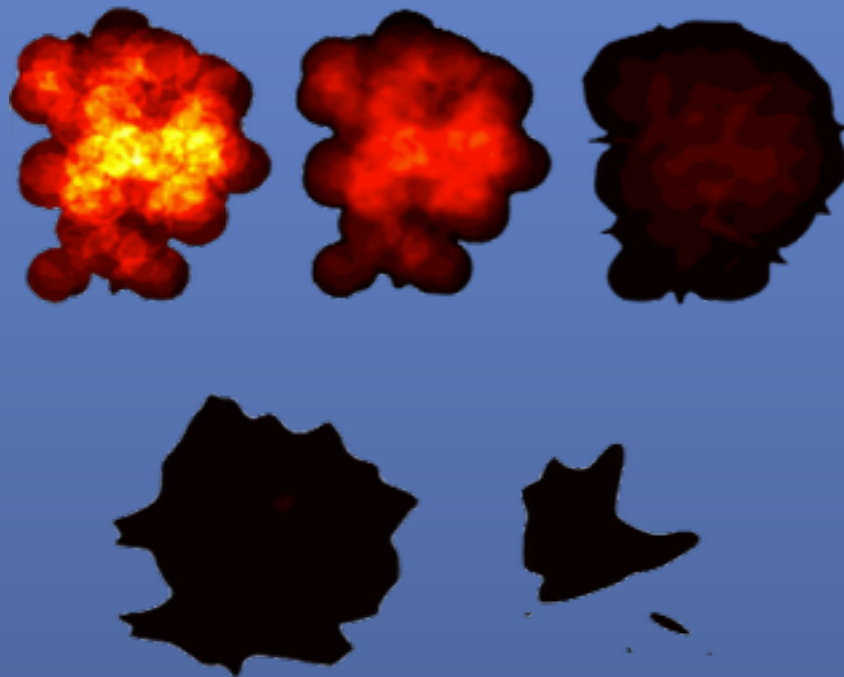
Thermal Photon
Emission Rates

viscous
corrections

Hydrodynamic
Simulations
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HydroInfo
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$e, s, p, T,$
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$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} q_\mu q_\nu}{2(e+p)} a_{\alpha\beta} \Gamma^{\alpha\beta}$$

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viscous
corrections

Hadrons spectra &
 V_n

Photon spectrum &
 V_n

Viscous Photon Emission Rates: General Formalism

Thermal photon emission rates can be calculated by

$$E_q \frac{dR}{d^3q} = \int \frac{d^3p_1}{2E_1(2\pi)^3} \frac{d^3p_2}{2E_2(2\pi)^3} \frac{d^3p_3}{2E_3(2\pi)^3} \frac{1}{2(2\pi)^3} |\mathcal{M}|^2 \\ \times f_1(p_1^\mu) f_2(p_2^\mu) (1 \pm f_3(p_3^\mu)) (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - q)$$

With

$$f(p^\mu) = f_0(E) + f_0(E)(1 \pm f_0(E)) \frac{\pi^{\mu\nu} \hat{p}_\mu \hat{p}_\nu}{2(e+p)} \chi\left(\frac{p}{T}\right)$$

We can expand photon emission rates around the thermal equilibrium:

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e+p)} a_{\alpha\beta} \Gamma^{\alpha\beta},$$

$$a_{\mu\nu} = \frac{3}{2(u \cdot \hat{q})^4} \hat{q}_\mu \hat{q}_\nu + \frac{1}{(u \cdot \hat{q})^2} u_\mu u_\nu + \frac{1}{2(u \cdot \hat{q})^2} g_{\mu\nu} - \frac{3}{2(u \cdot \hat{q})^3} (\hat{q}_\mu u_\nu + \hat{q}_\nu u_\mu).$$

Viscous Photon Emission Rates: General Formalism

Thermal photon emission rates can be calculated by

$$E_q \frac{dR}{d^3q} = \int \frac{d^3p_1}{2E_1(2\pi)^3} \frac{d^3p_2}{2E_2(2\pi)^3} \frac{d^3p_3}{2E_3(2\pi)^3} \frac{1}{2(2\pi)^3} |\mathcal{M}|^2 \\ \times f_1(p_1^\mu) f_2(p_2^\mu) (1 \pm f_3(p_3^\mu)) (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - q)$$

With

$$f(p^\mu) = \Gamma_0(q, T) + a_{\alpha\beta} \Gamma^{\alpha\beta}(q, T) \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{(e+p)} \chi\left(\frac{p}{T}\right)$$

We can expand $f(p^\mu)$ calculated in fluid local rest frame and the thermal equilibrium:

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e+p)} a_{\alpha\beta} \Gamma^{\alpha\beta},$$

$$a_{\mu\nu} = \frac{3}{2(u \cdot \hat{q})^4} \hat{q}_\mu \hat{q}_\nu + \frac{1}{(u \cdot \hat{q})^2} u_\mu u_\nu + \frac{1}{2(u \cdot \hat{q})^2} \dots$$

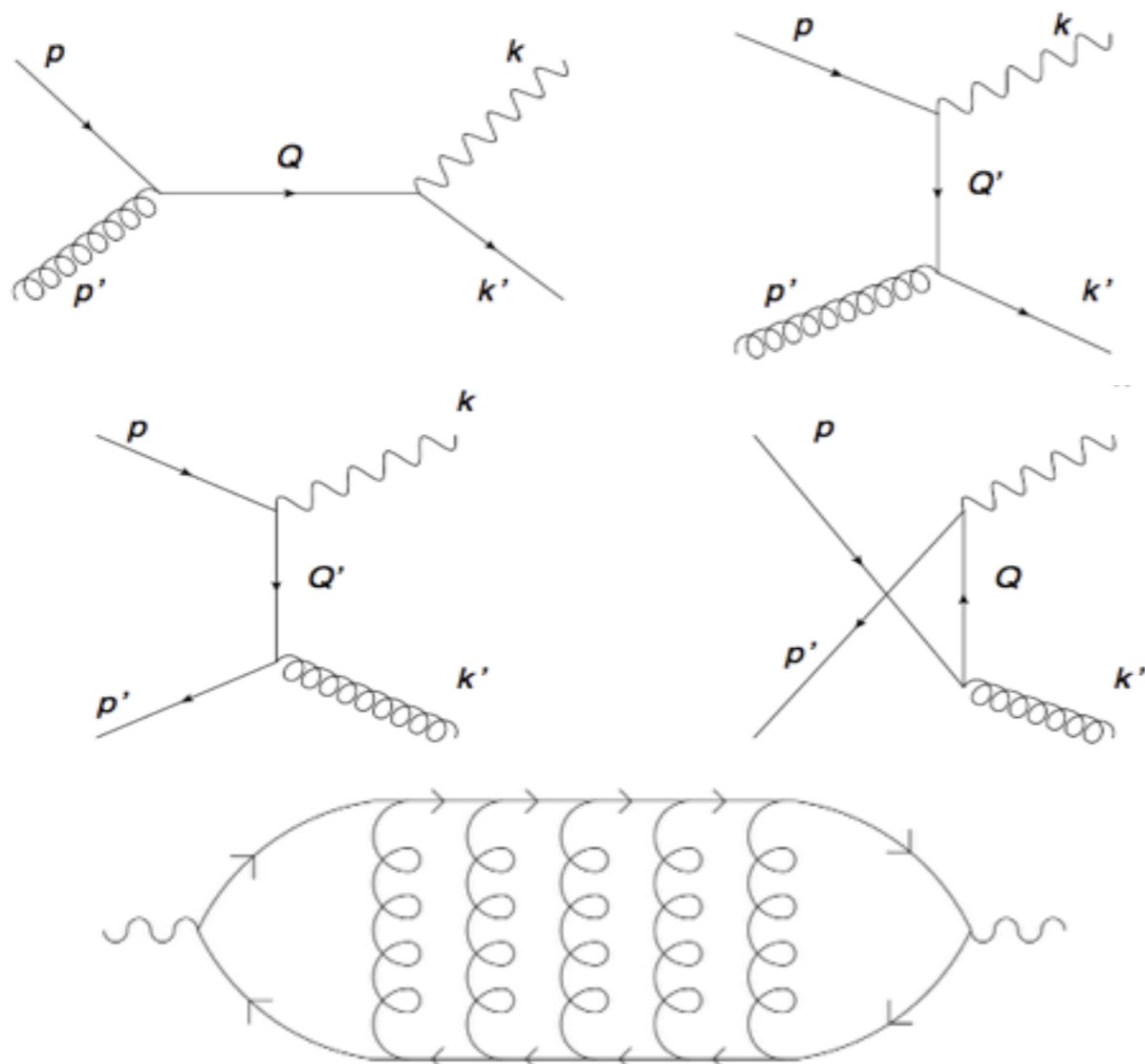
calculated in lab frame

Viscous Photon Emission Rates

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e+p)} a_{\alpha\beta} \Gamma^{\alpha\beta}$$

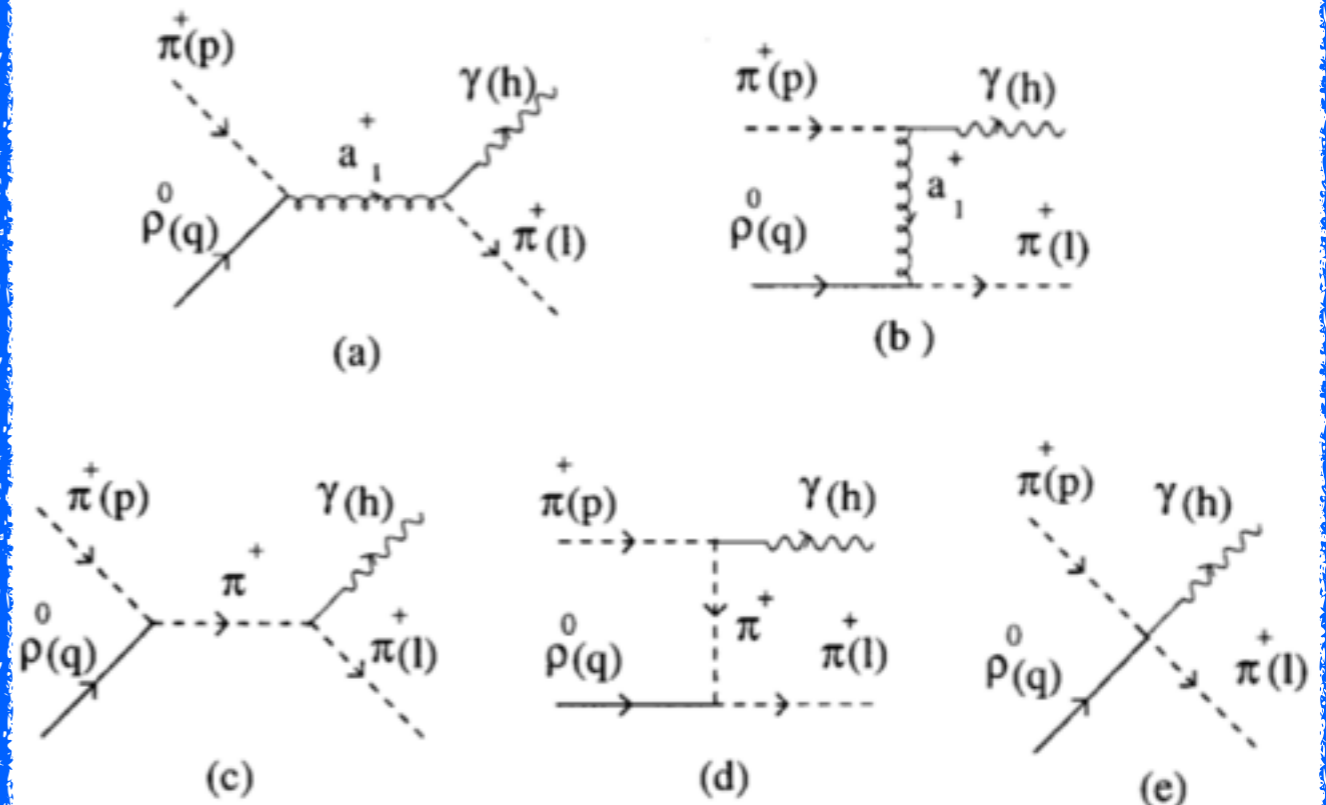
Equilibrium rates

QGP (AMY 2001)



Hadron Gas (TRG 2004)

(TRG 2004)



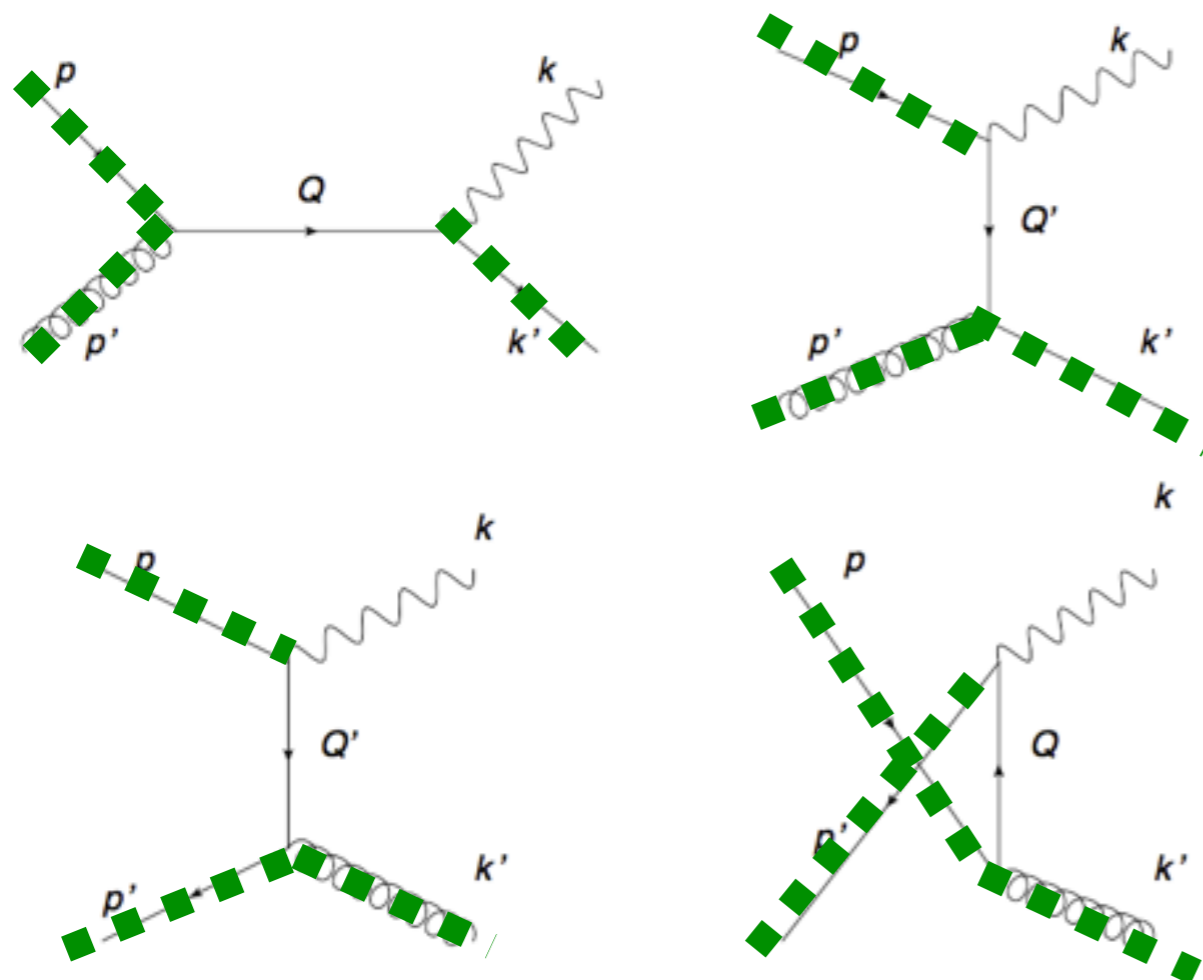
Viscous Photon Emission Rates

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Equilibrium rates

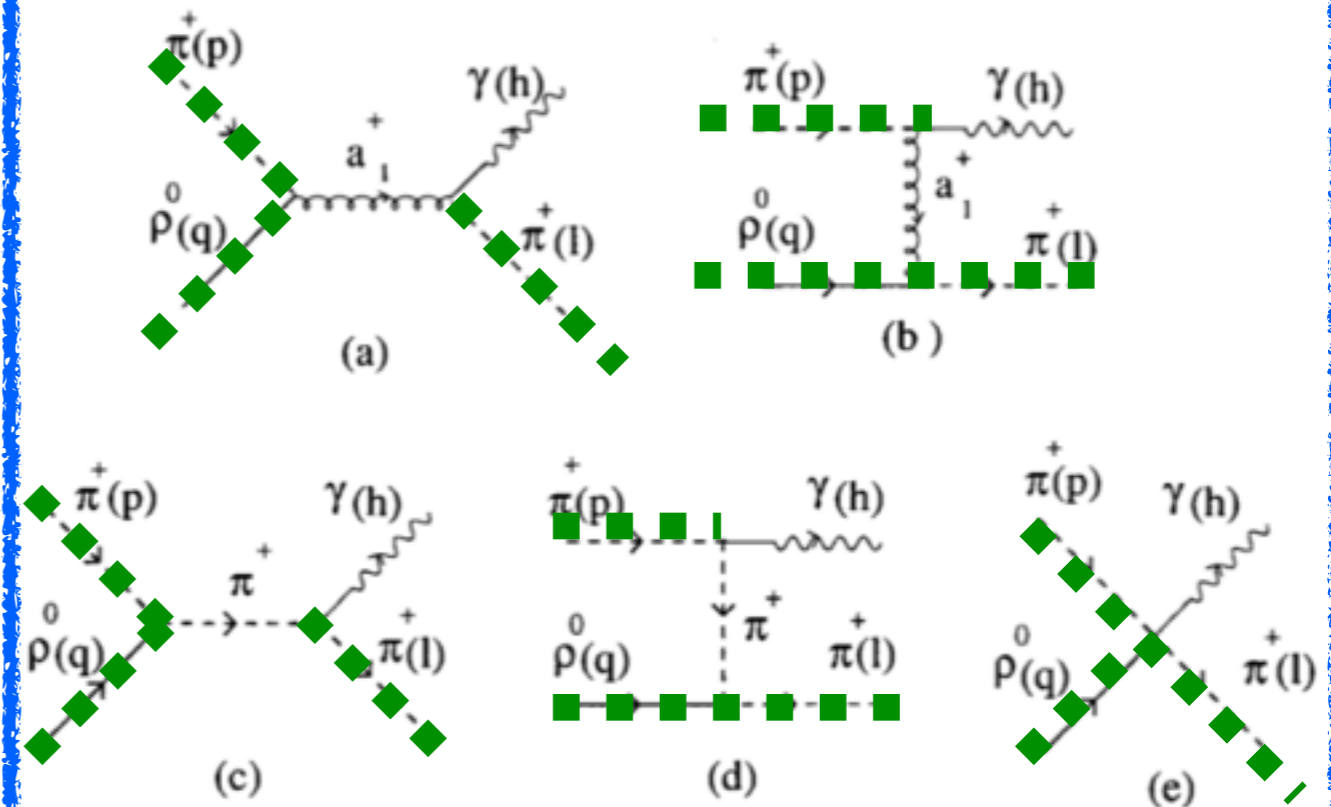
off-equilibrium δf corrections

QGP



Dusling NPA839 (2010) 70

Hadron Gas



Dion et al. PRC84 (2011) 064901

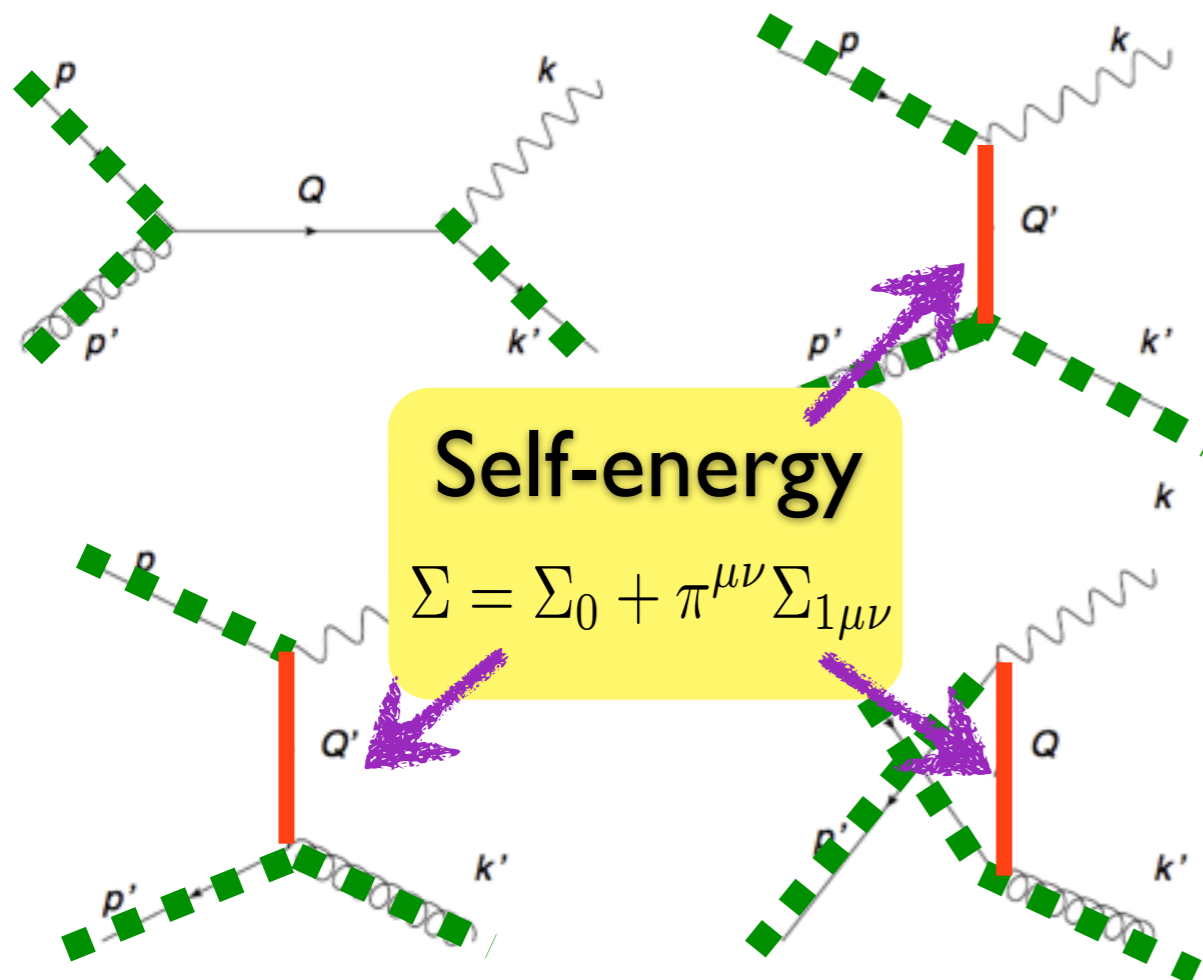
Viscous Photon Emission Rates

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Equilibrium rates

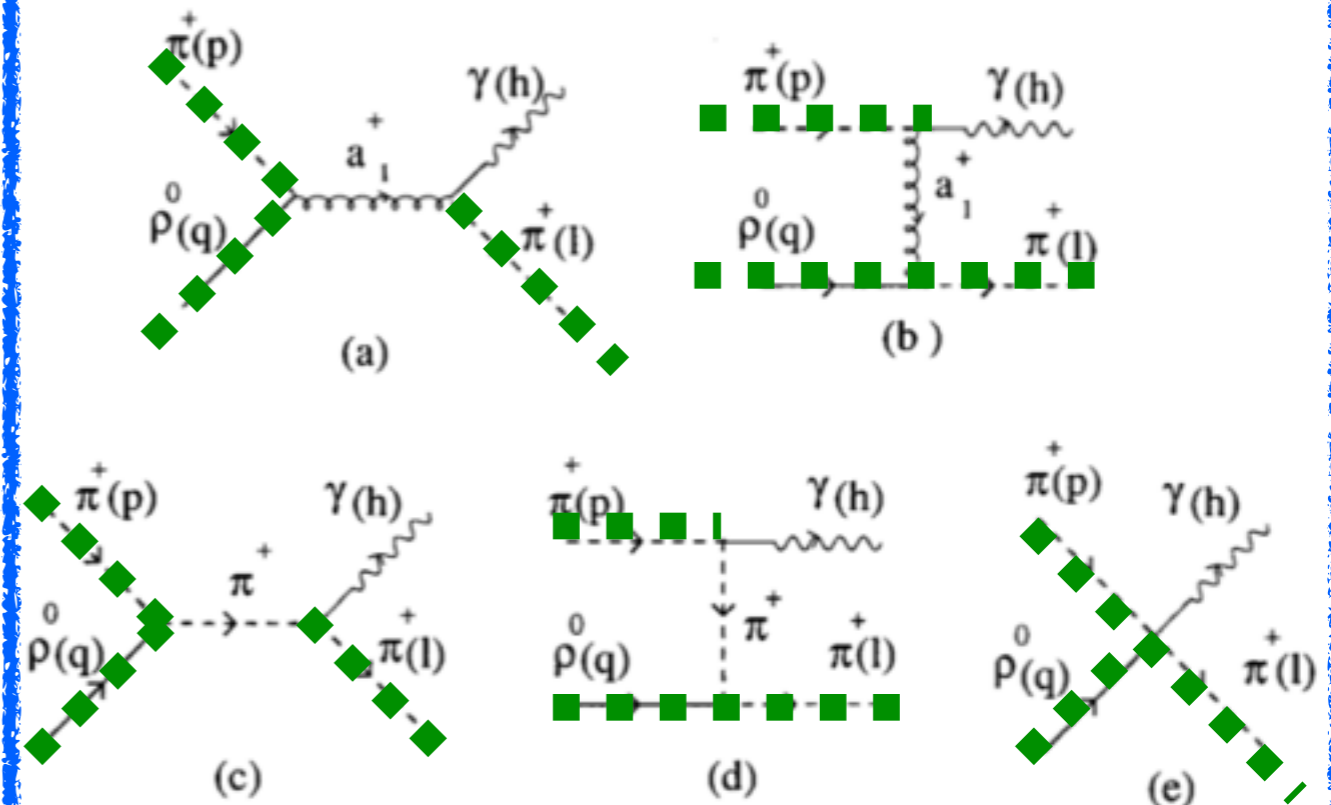
off-equilibrium δf corrections

QGP

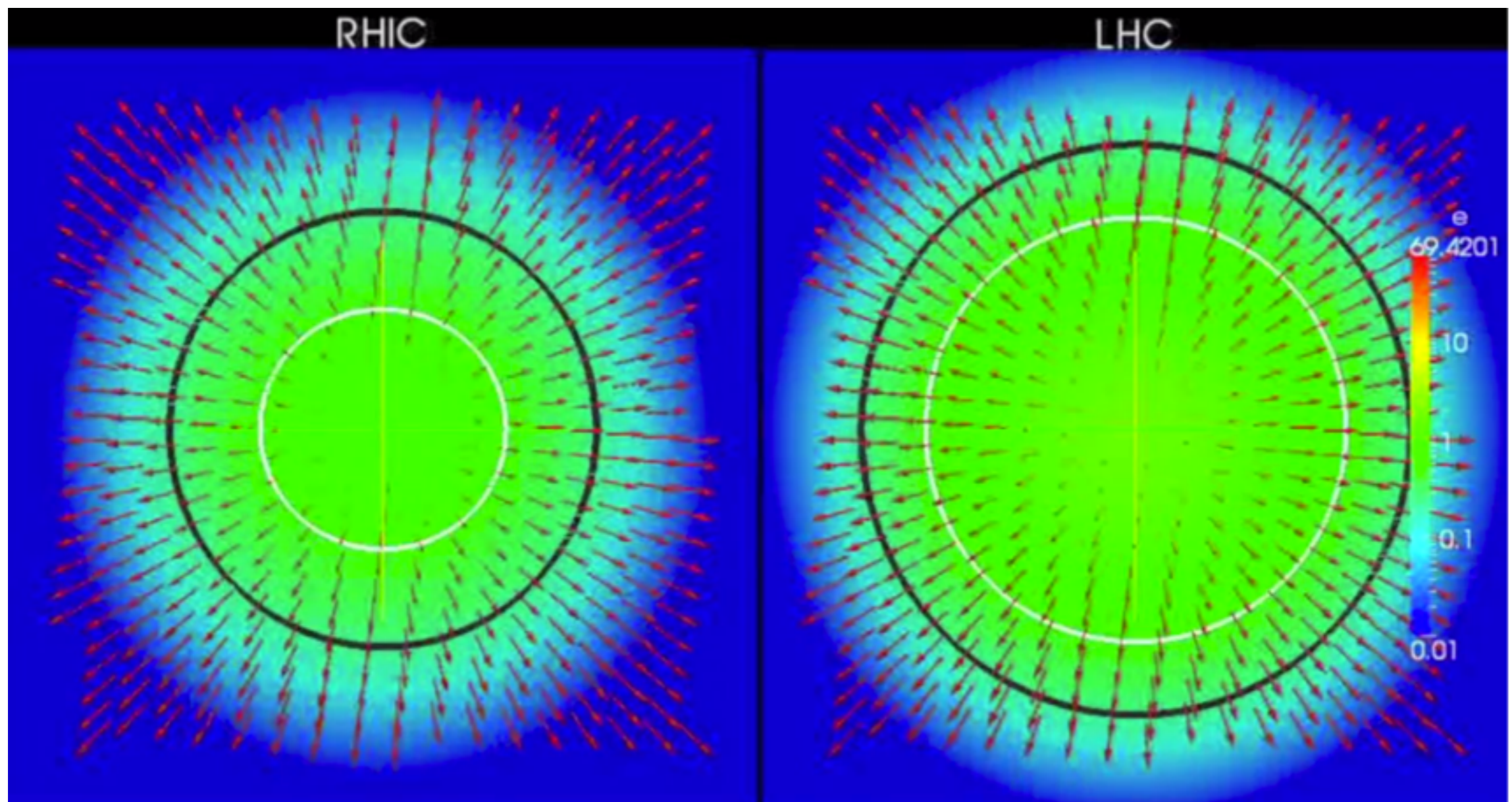


Shen, Paquet et al. (2014)

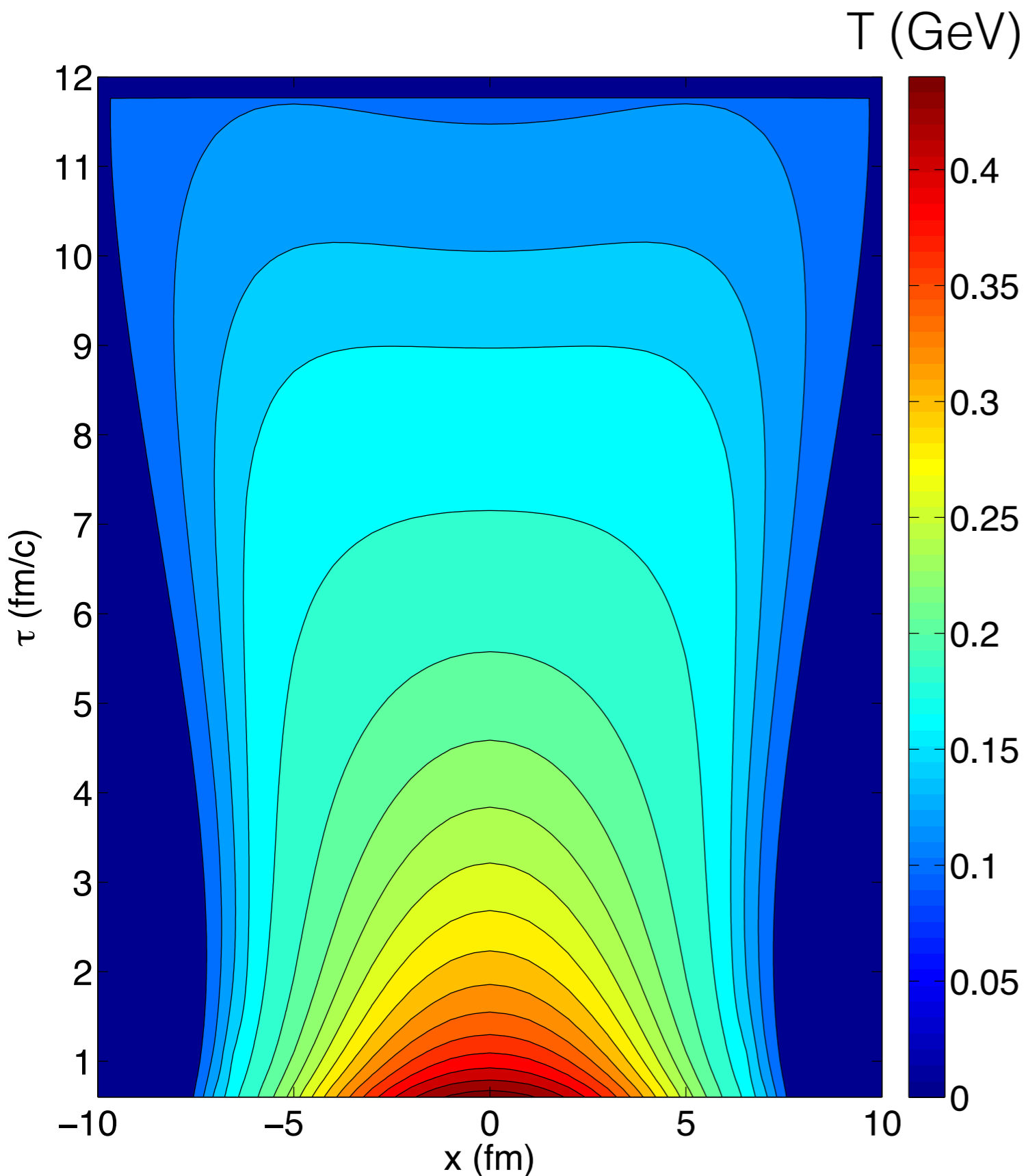
Hadron Gas



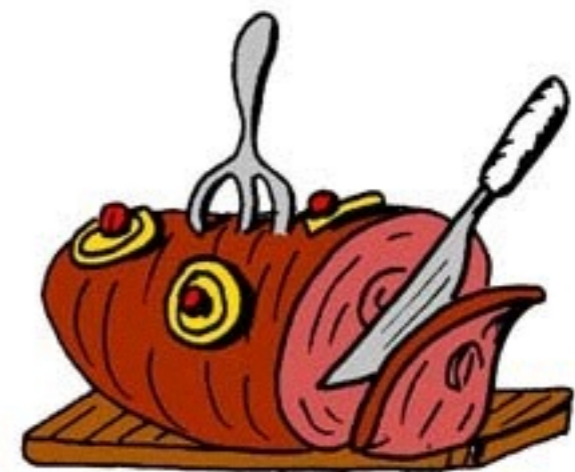
Photon spectra and radial flow



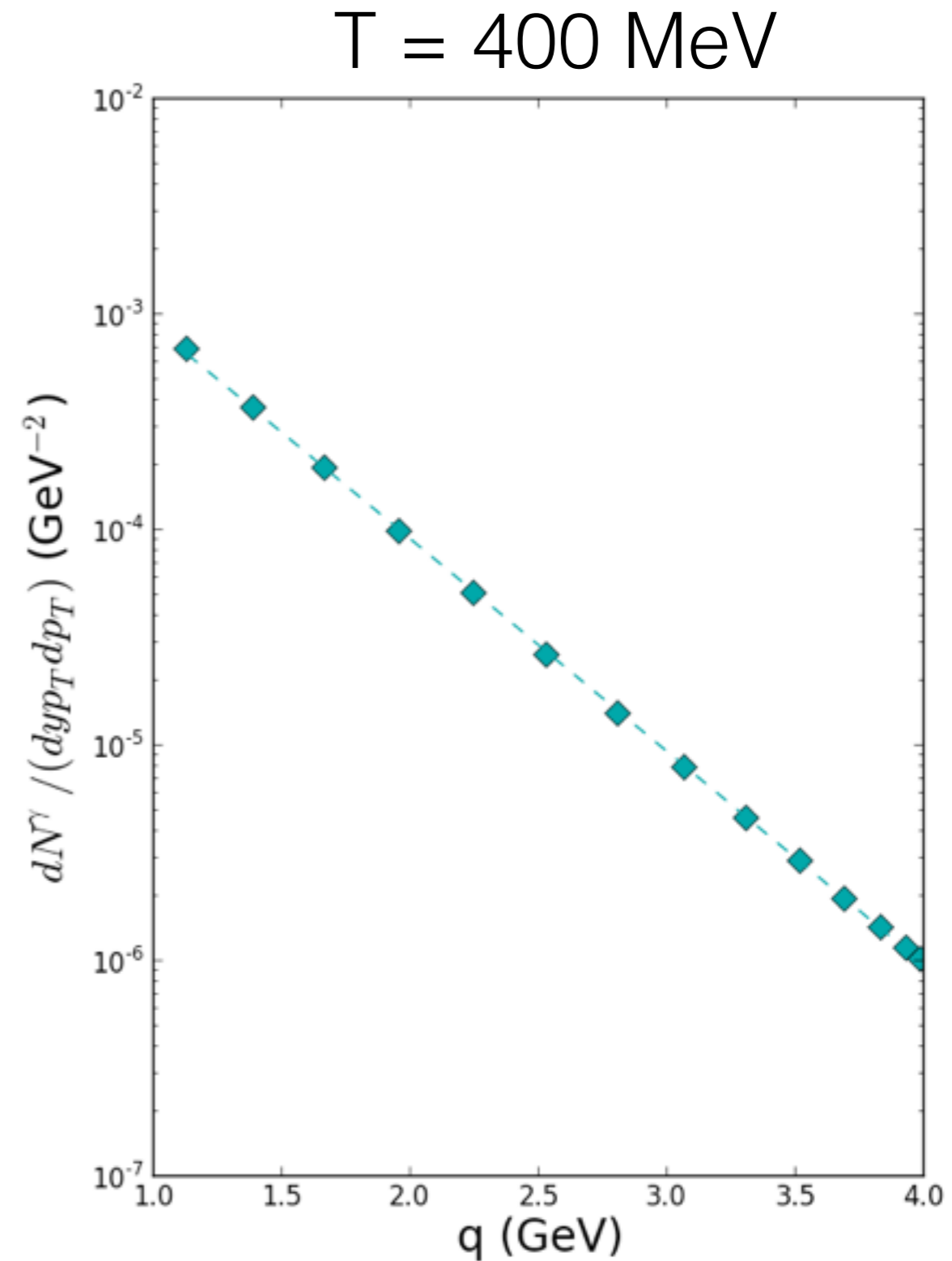
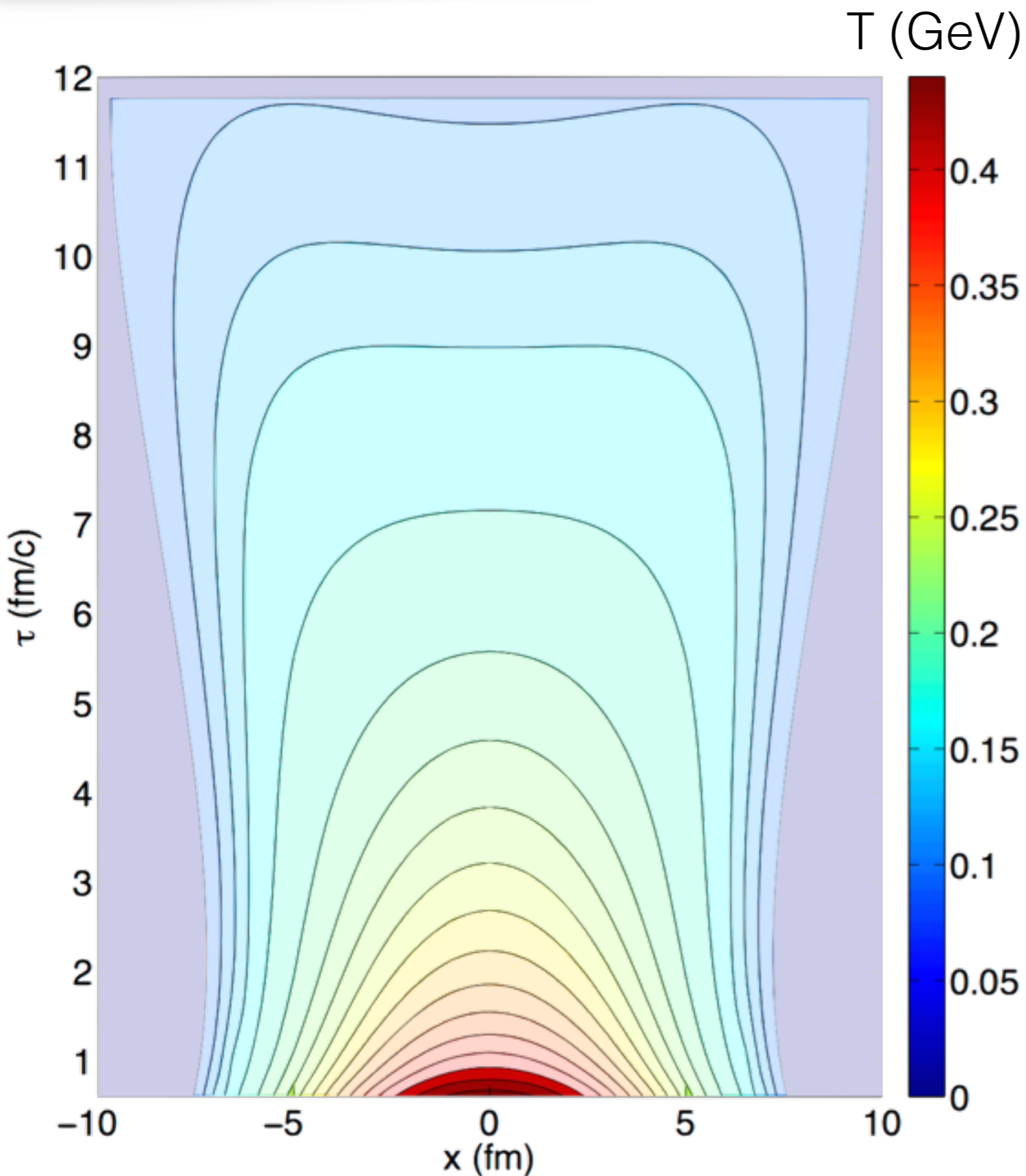
Slope of Photon Spectrum



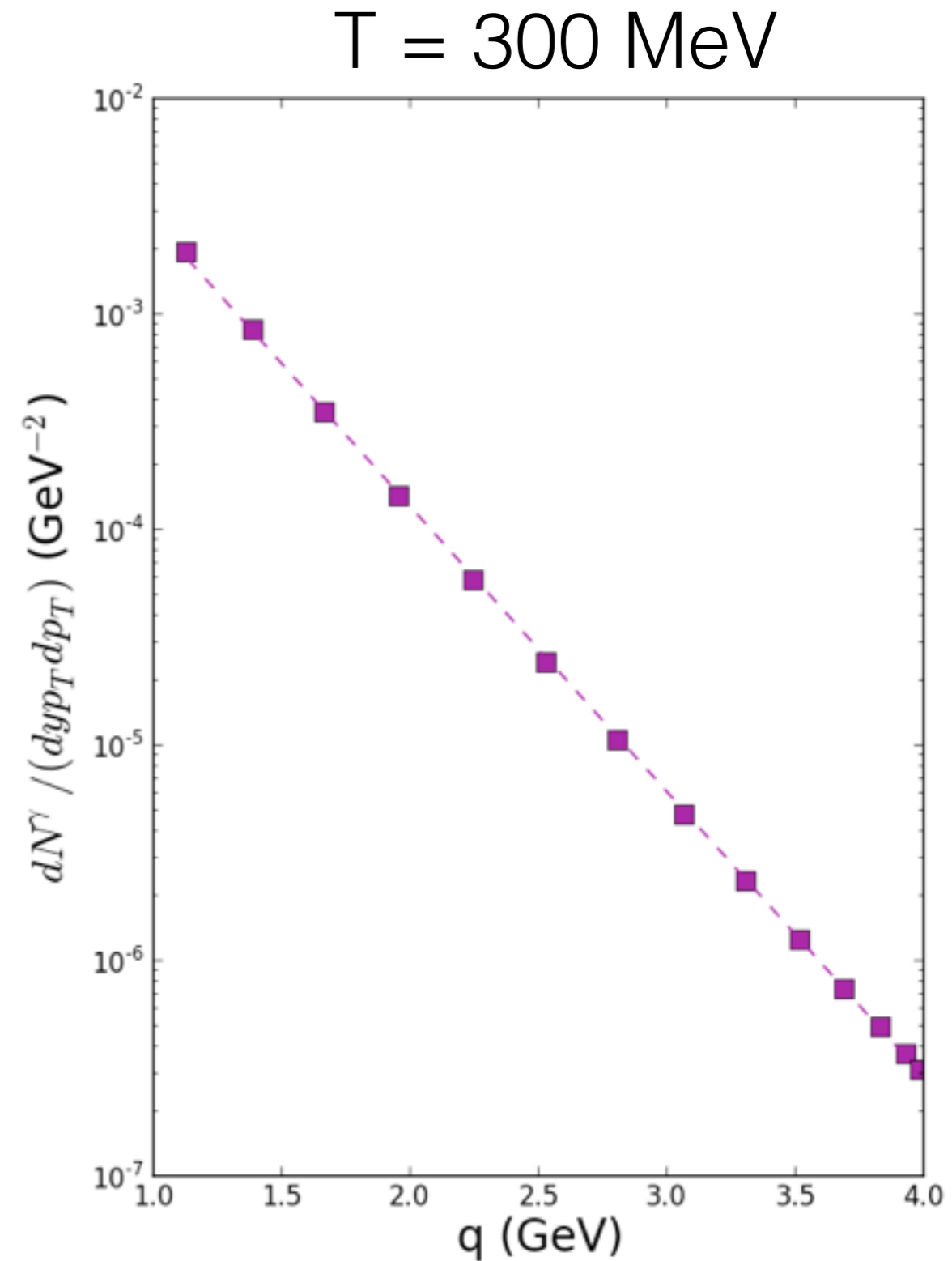
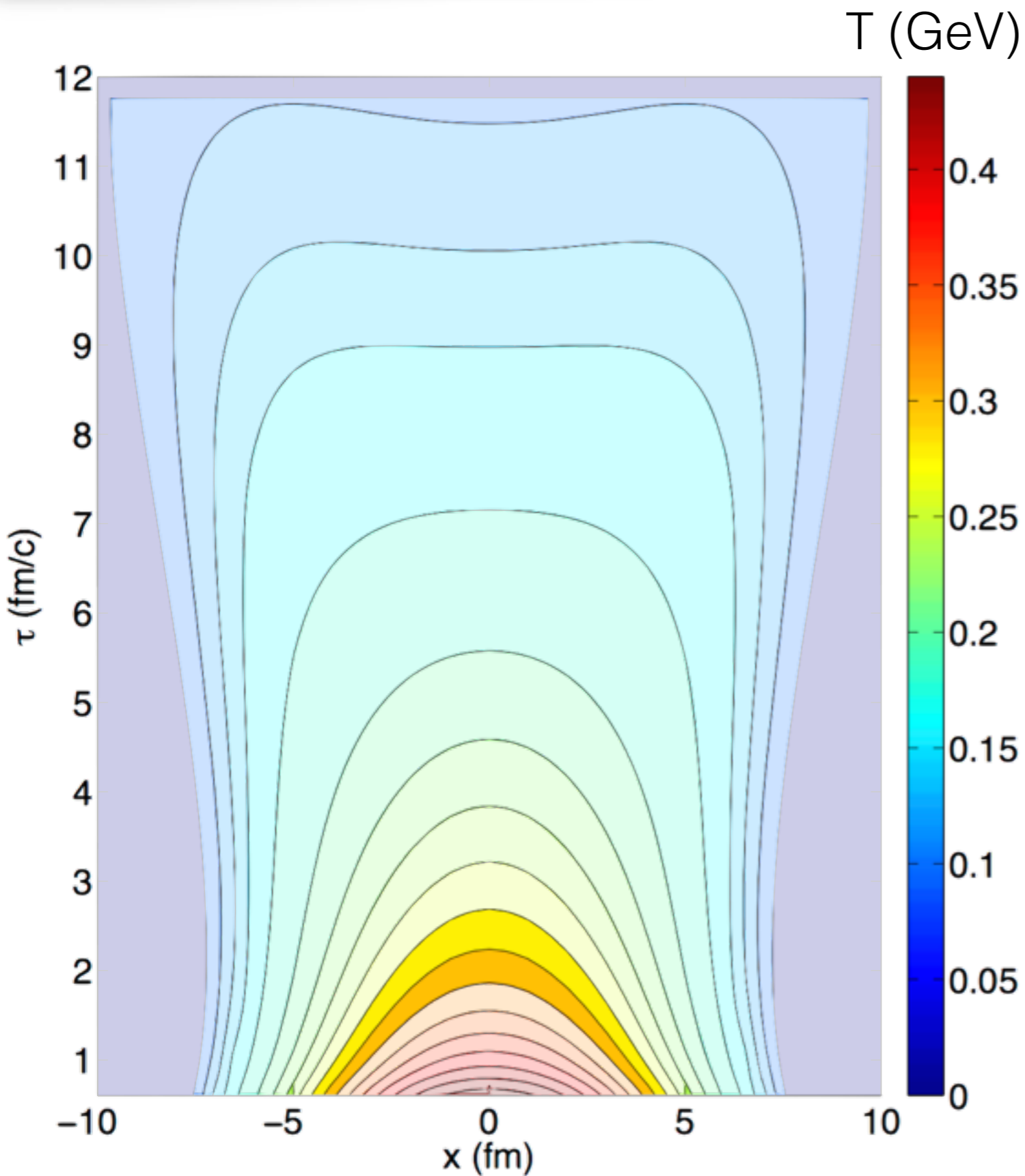
Slicing the hydrodynamic medium



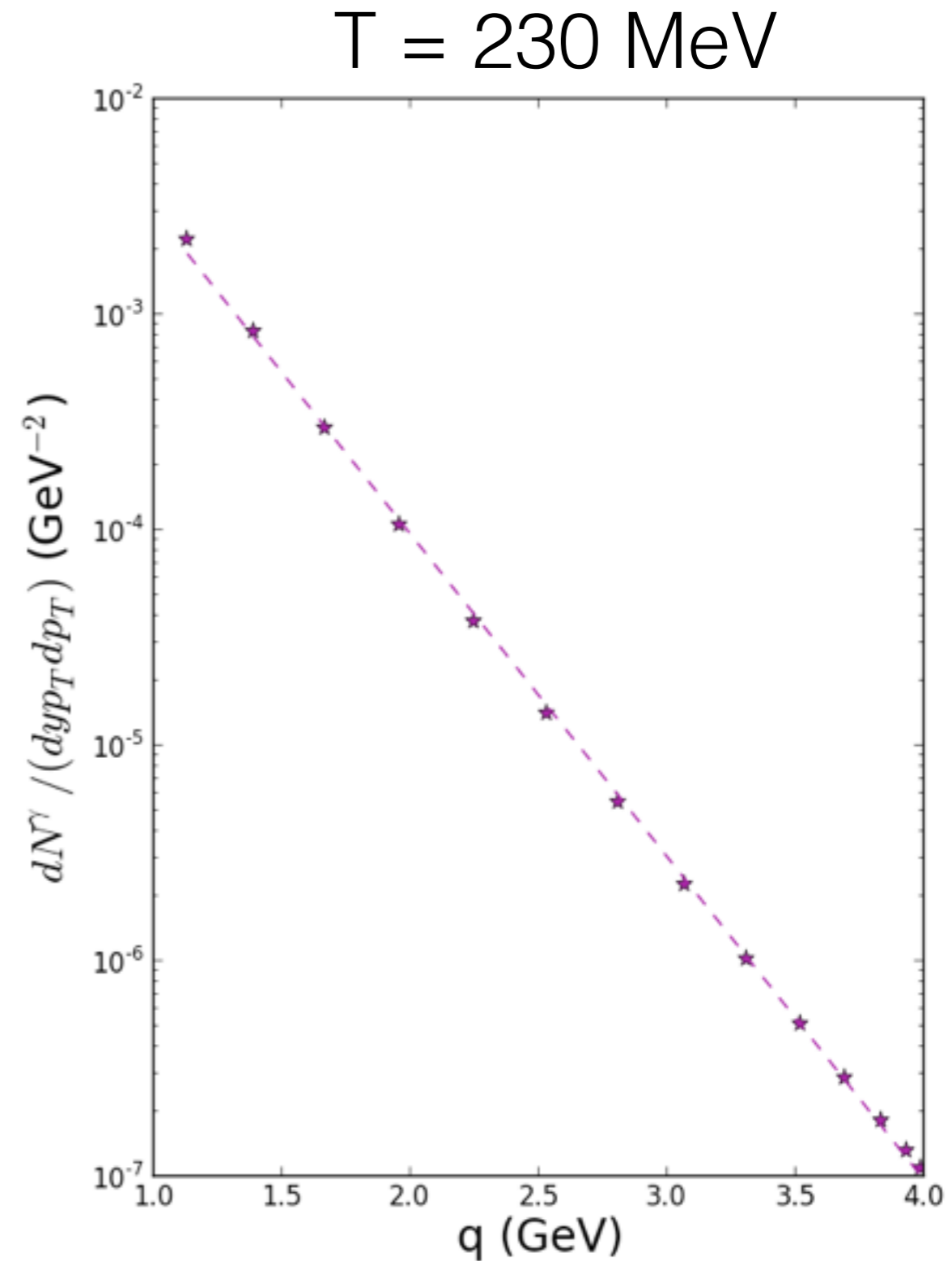
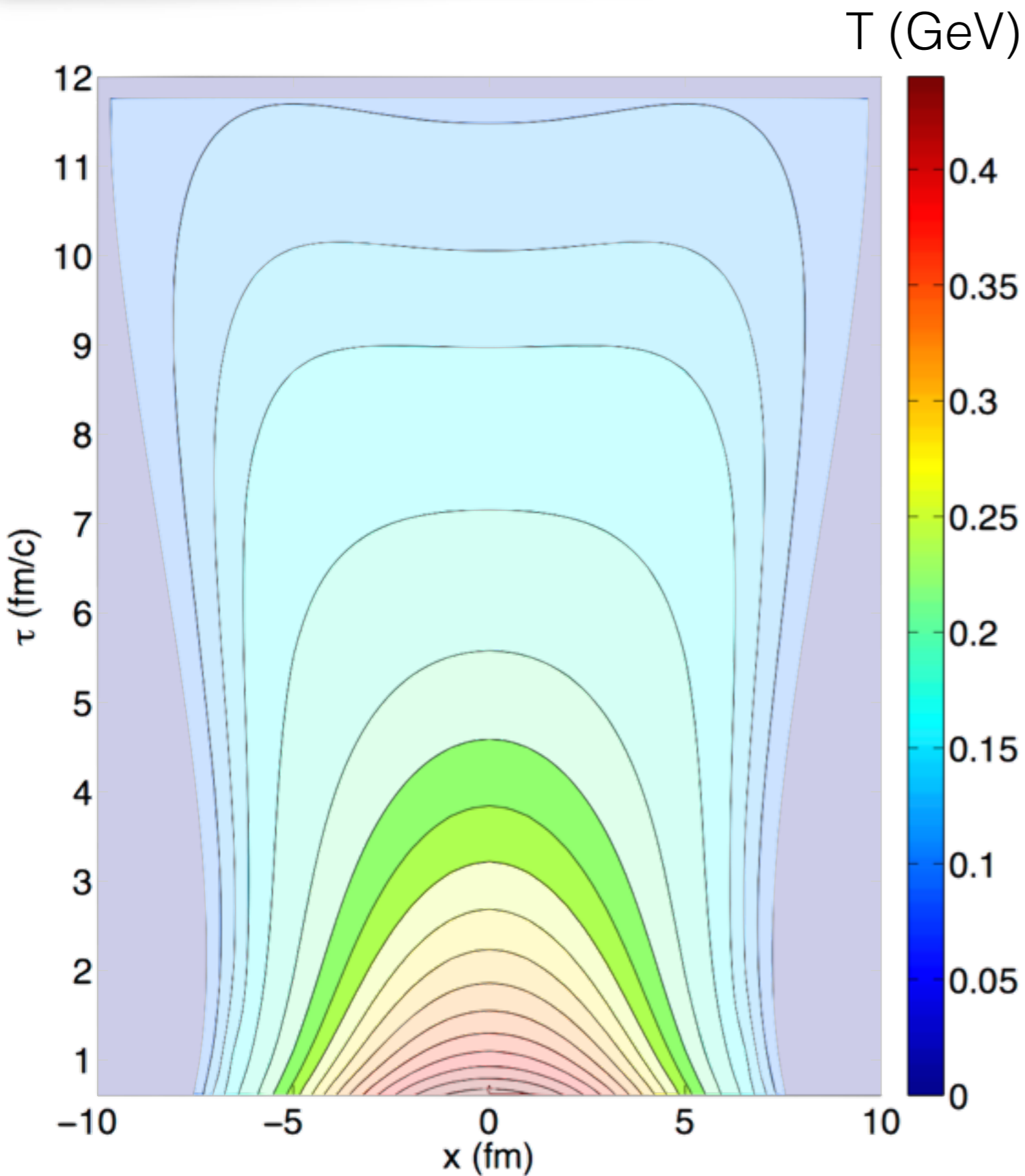
Slope of Photon Spectrum



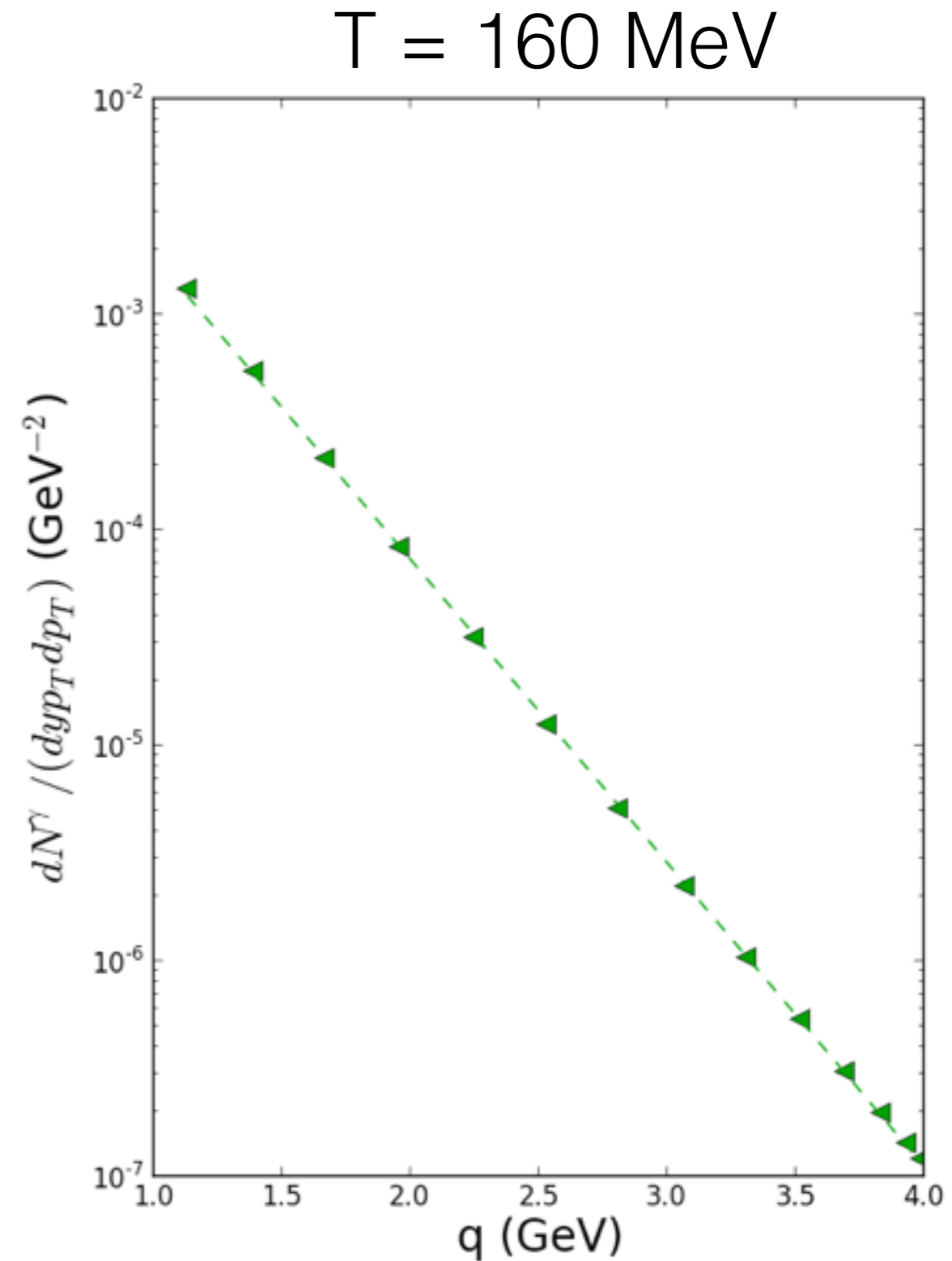
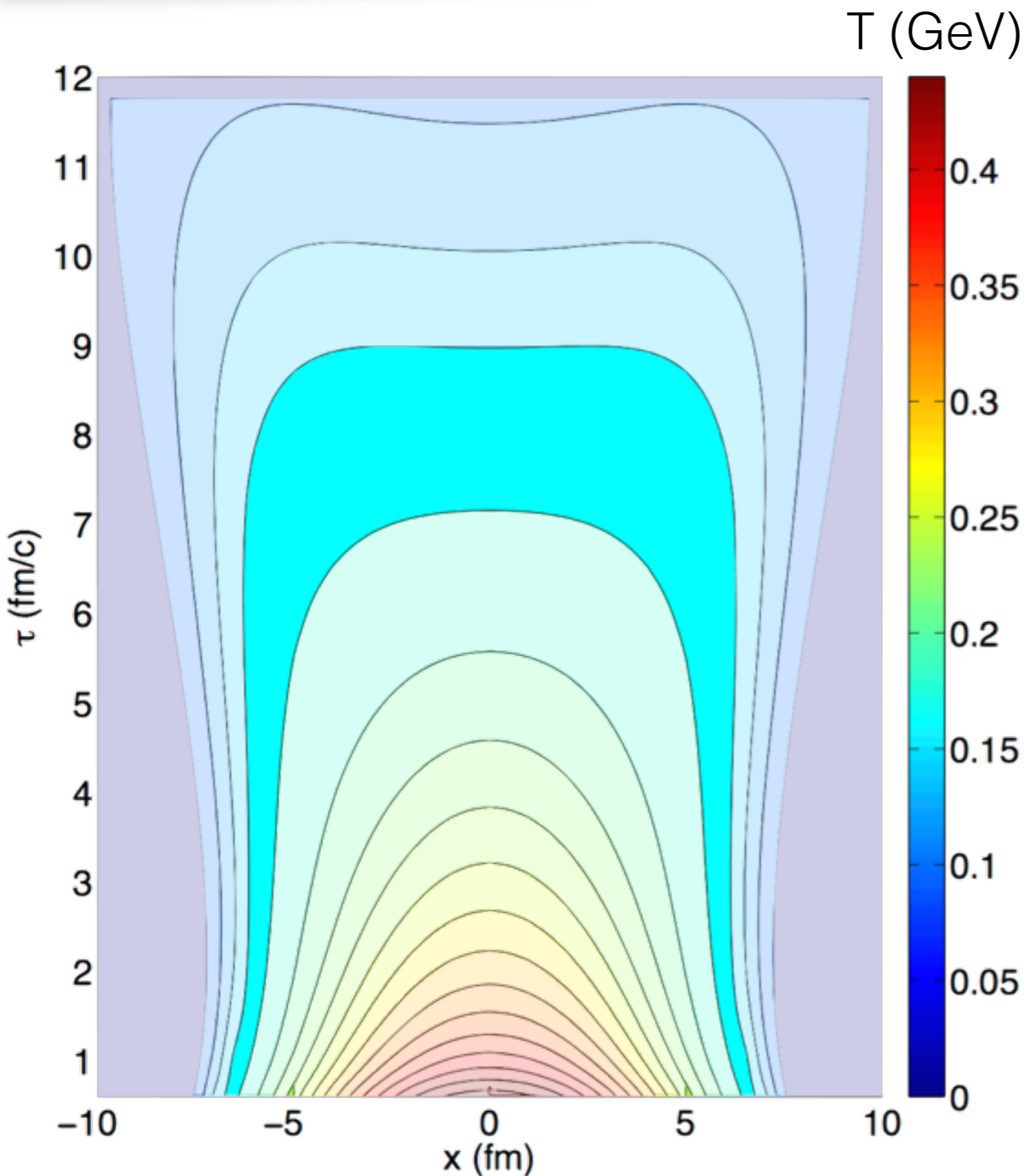
Slope of Photon Spectrum



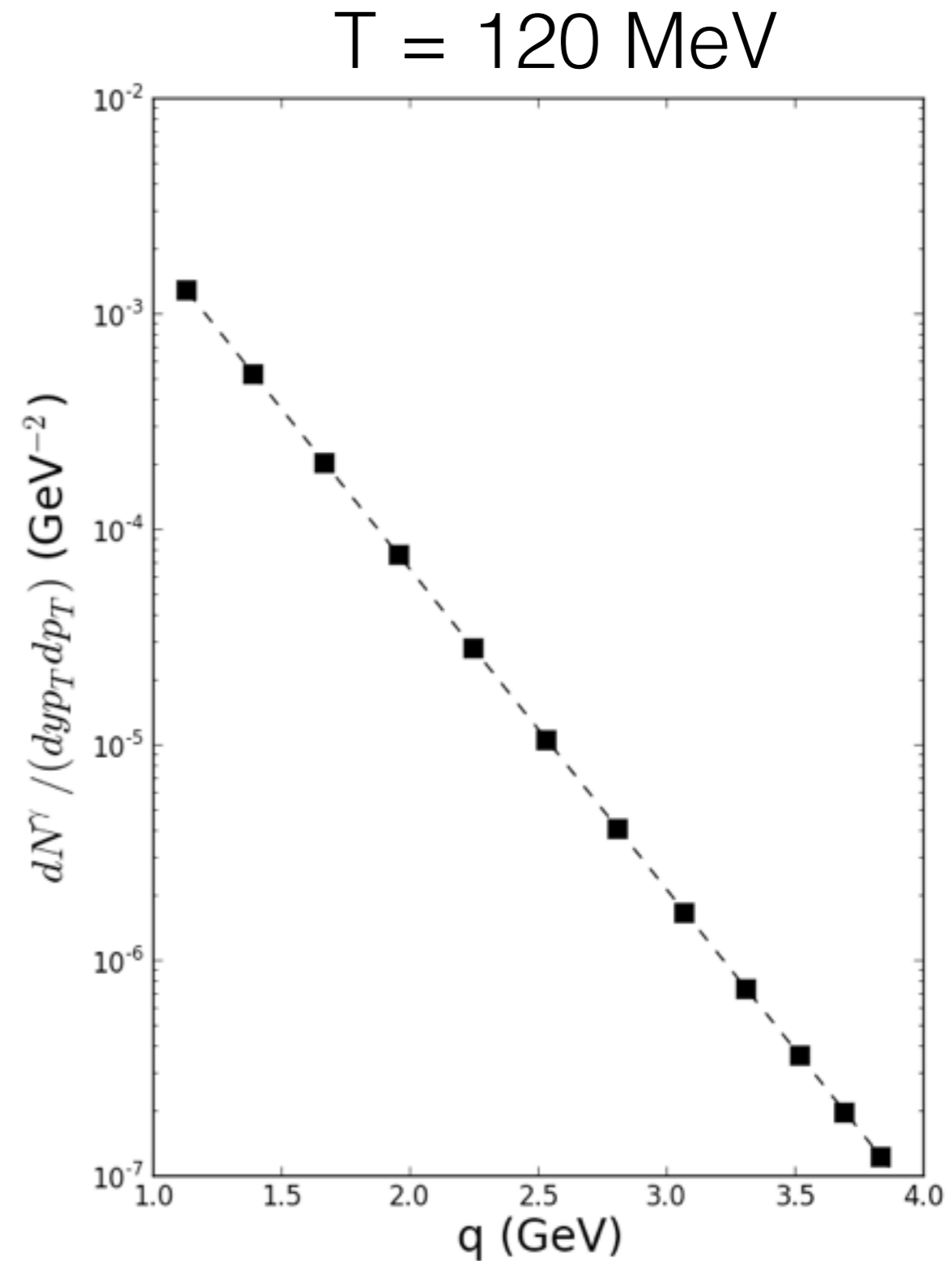
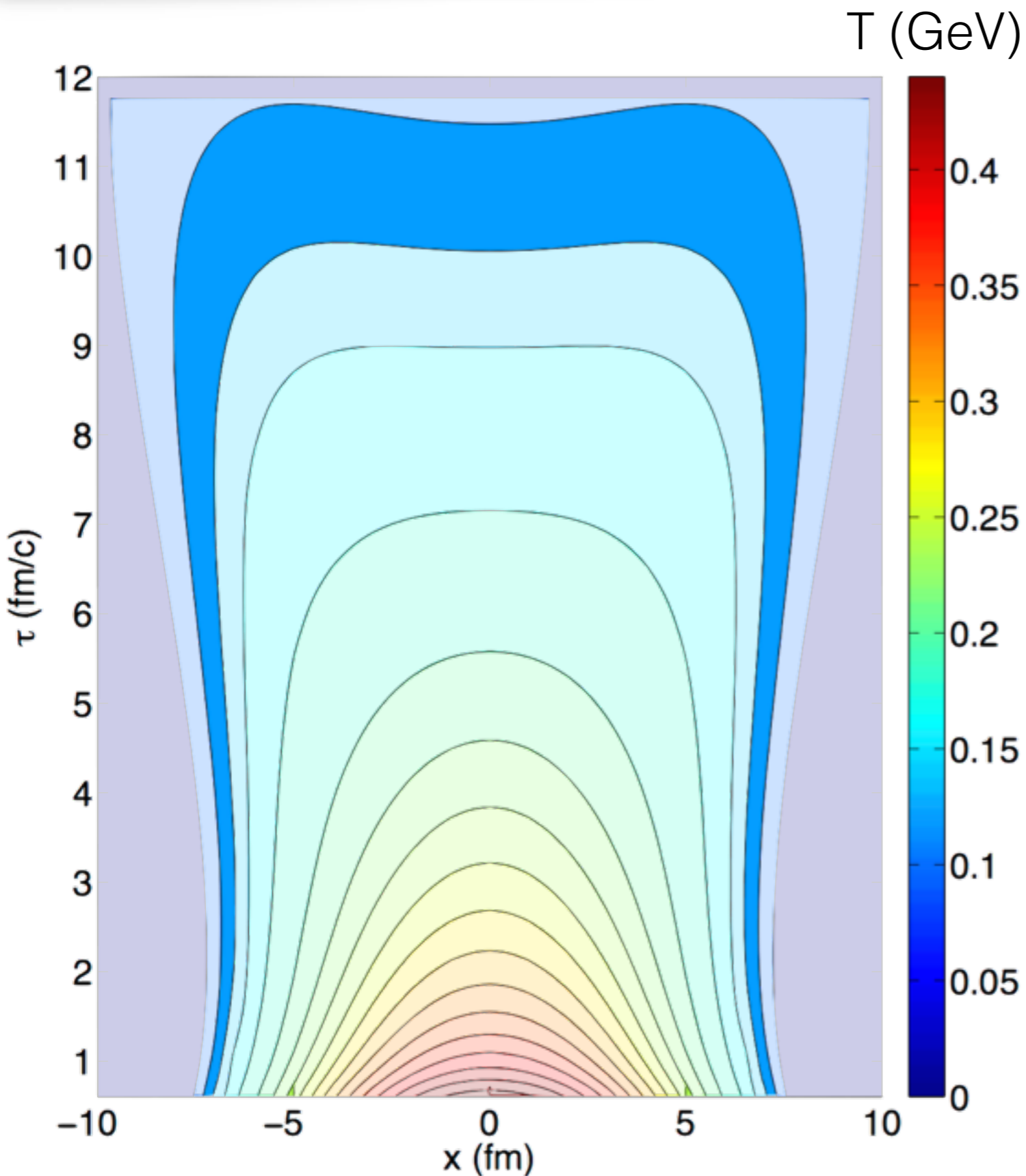
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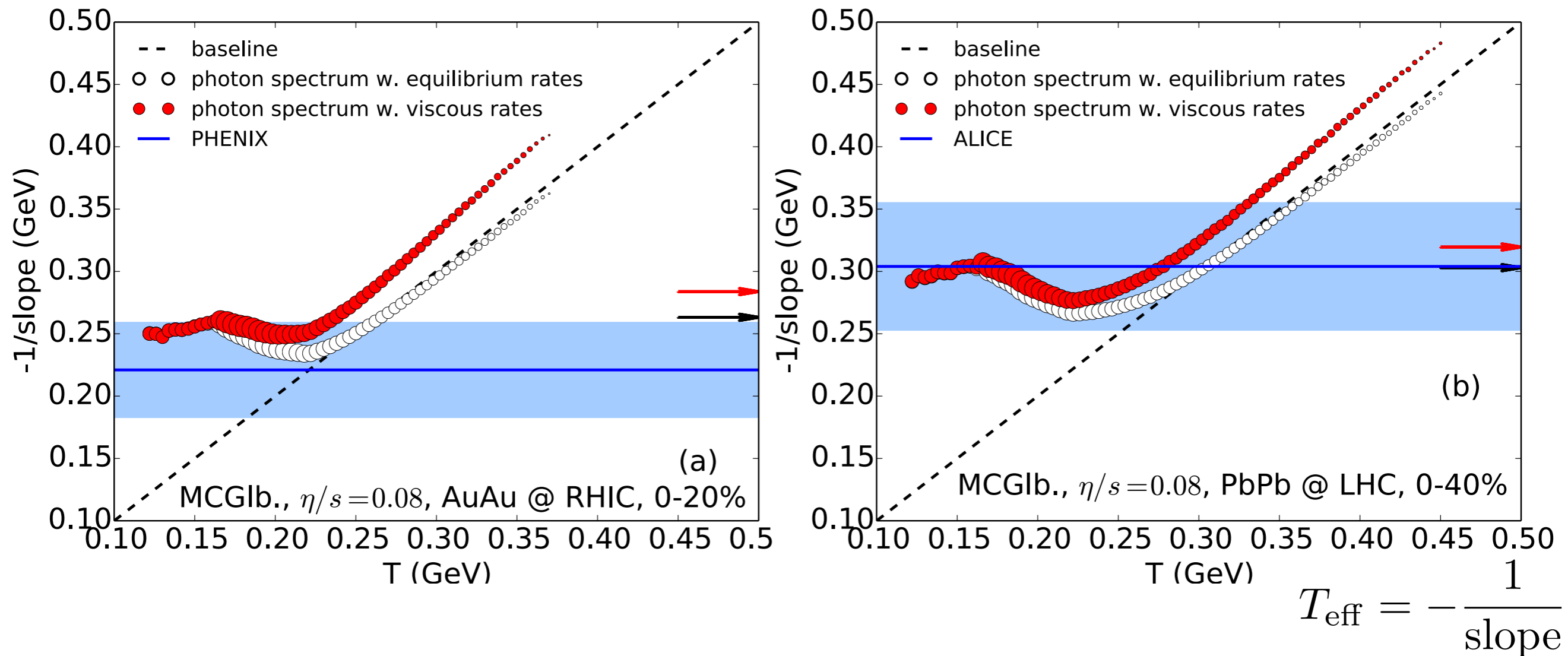


Slope of Photon Spectrum



Fitted T_{eff} vs. True Temperature

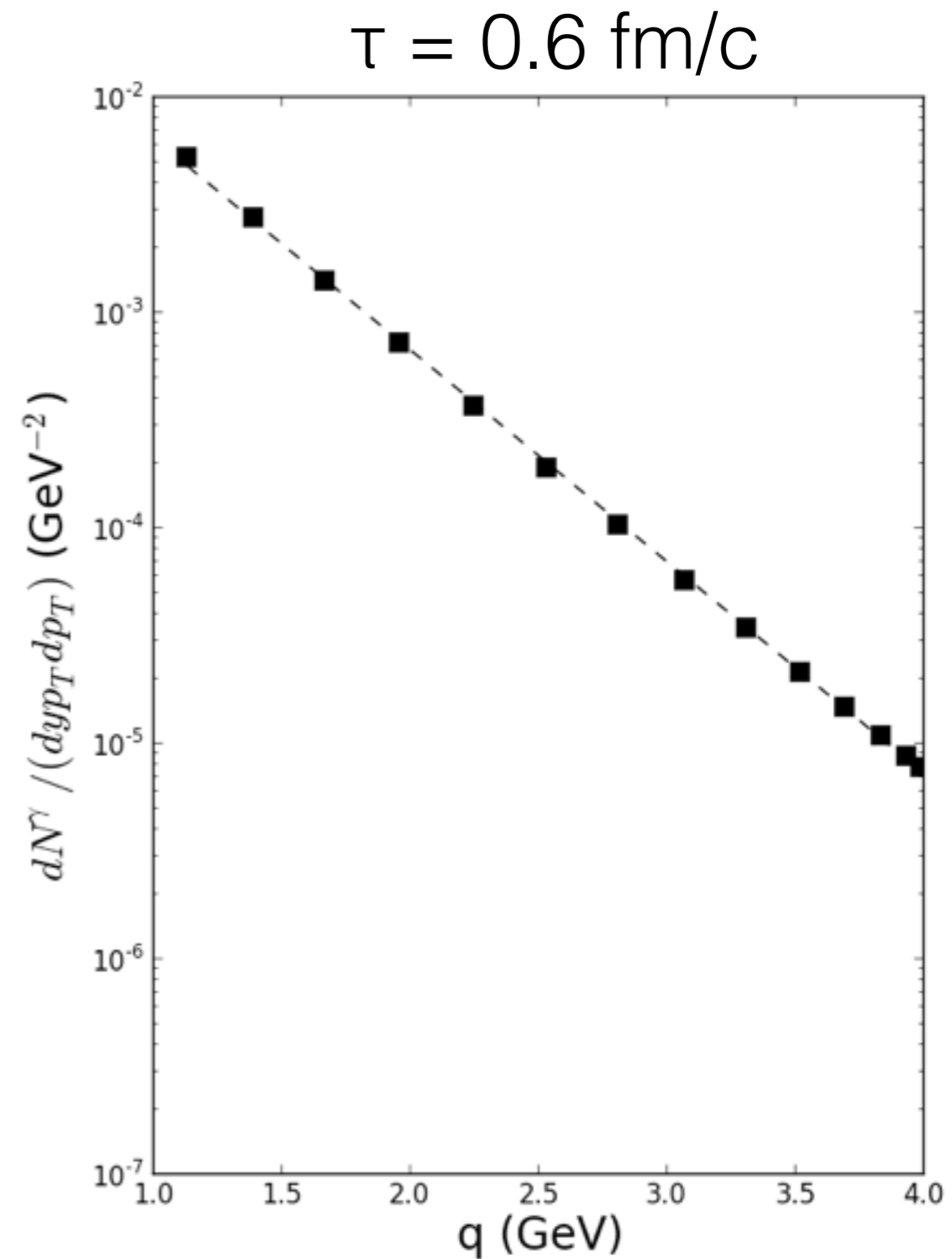
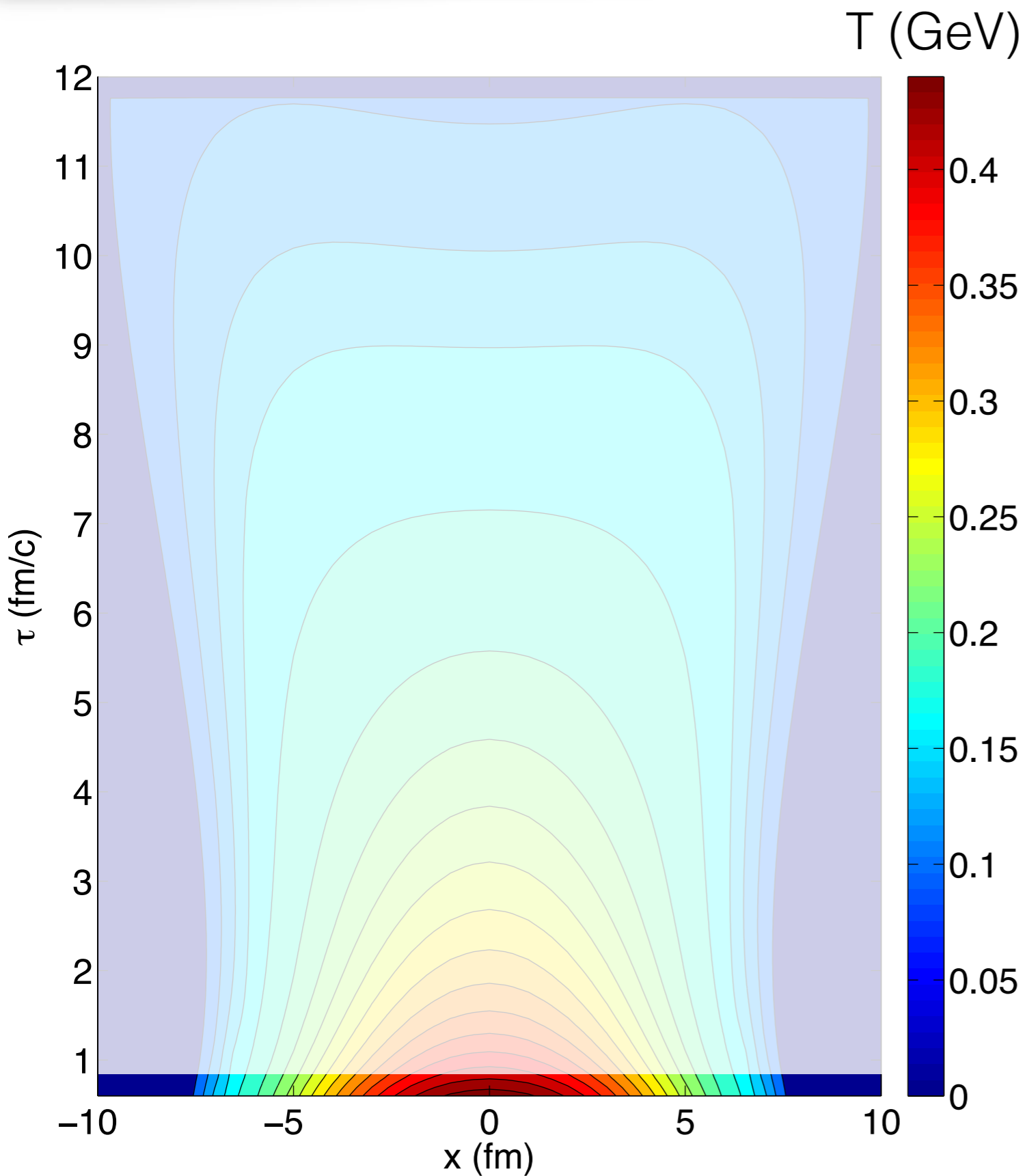
C. Shen, U. Heinz, J.-F. Paquet and C. Gale, Phys. Rev. C 89, 044910 (2014)



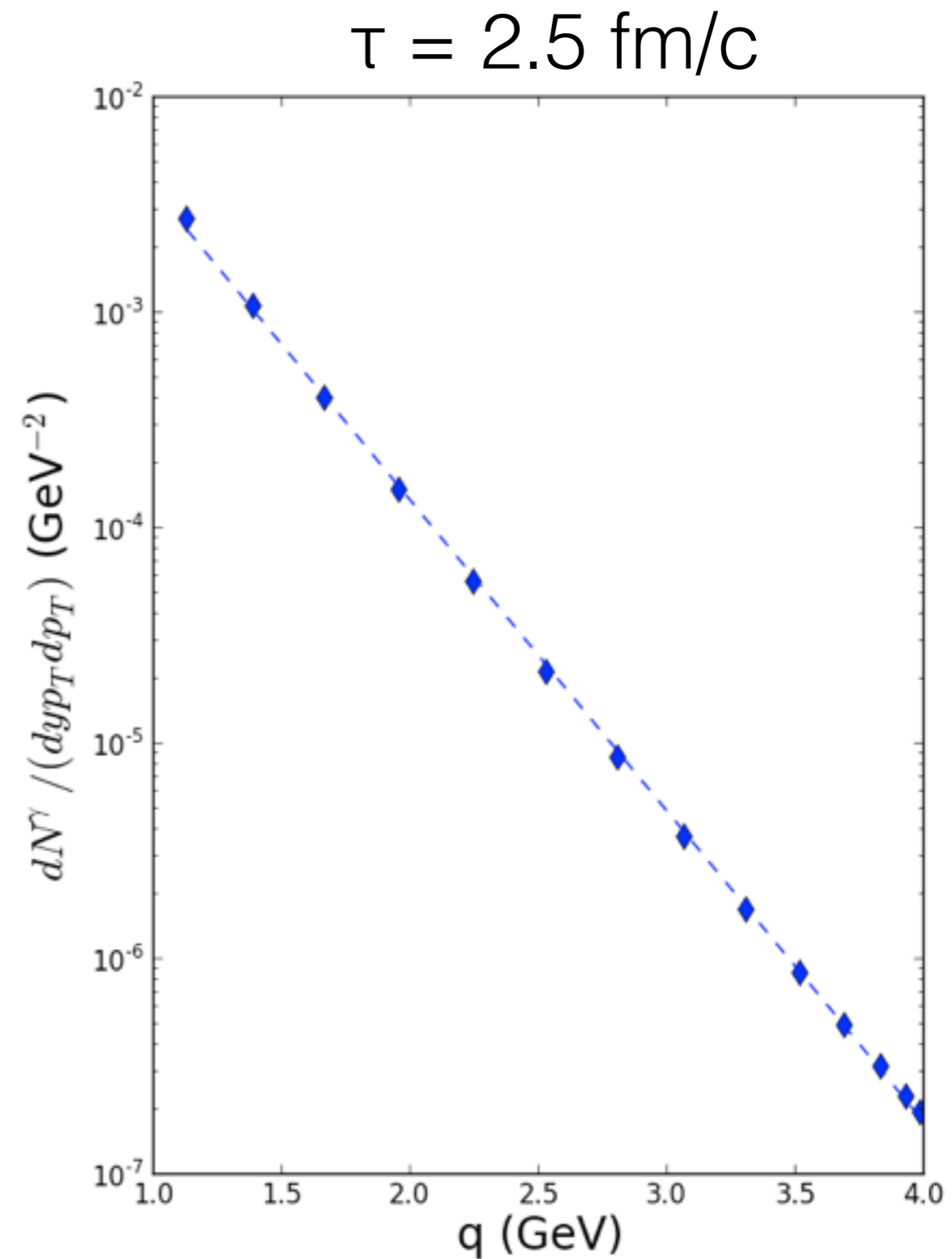
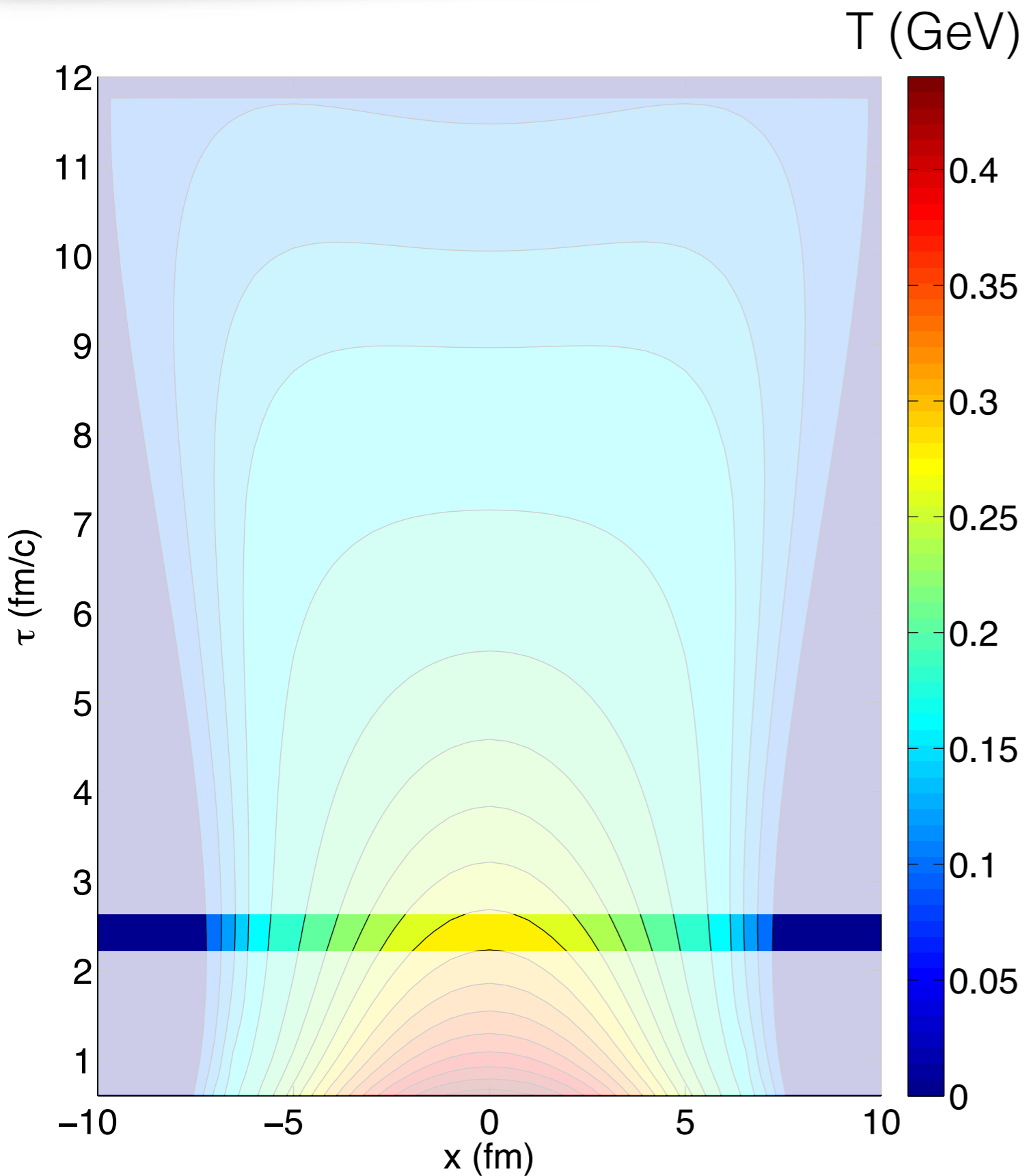
- **All** photons with $T < 250$ MeV at RHIC and < 300 MeV at LHC carries T_{eff} within the experimental fitted region
- About **50-60%** of photons are emitted from $T = 165 \sim 250$ MeV, they are strongly blue shifted by radial flow

$$T_{\text{eff}} = T \sqrt{\frac{1+v}{1-v}}$$

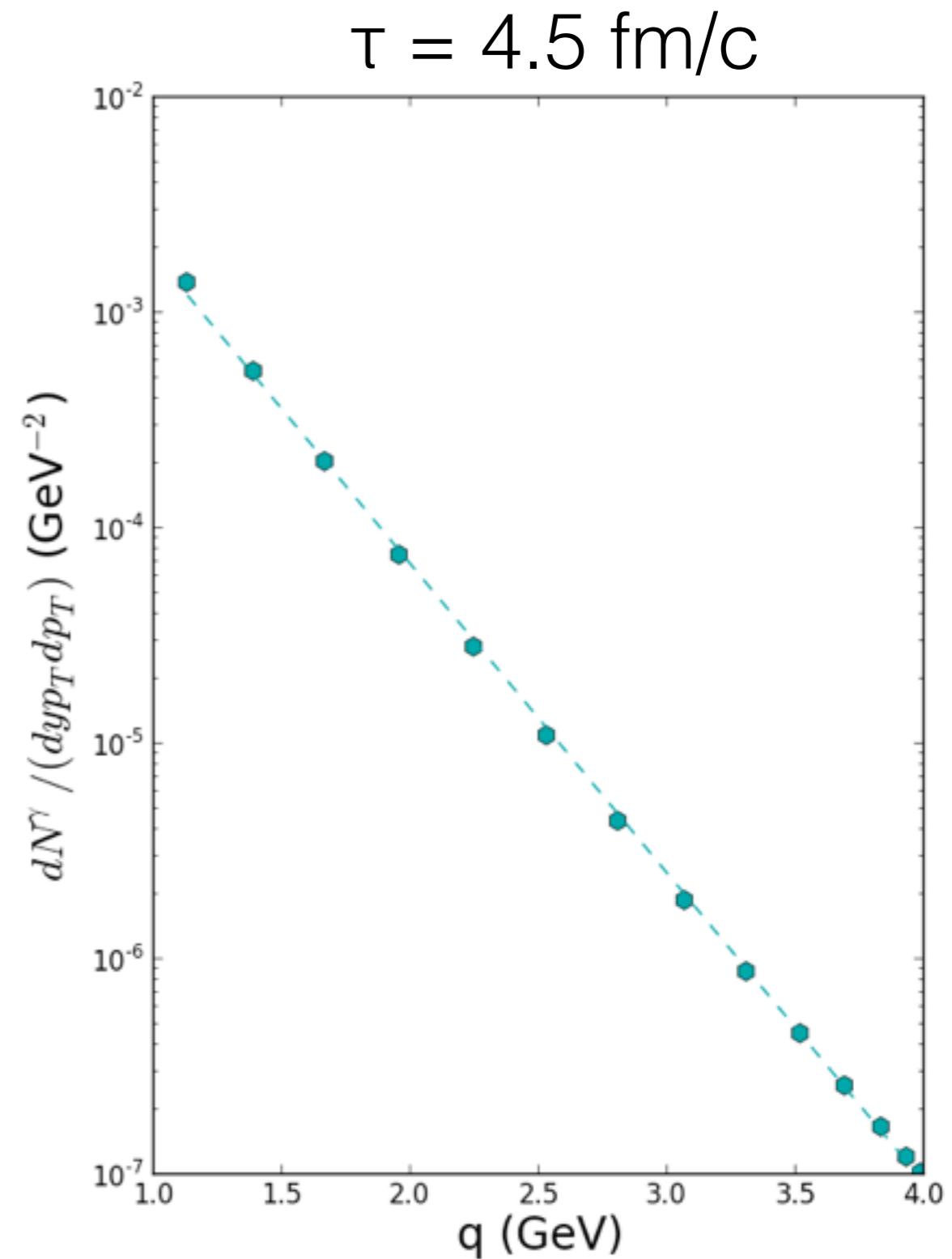
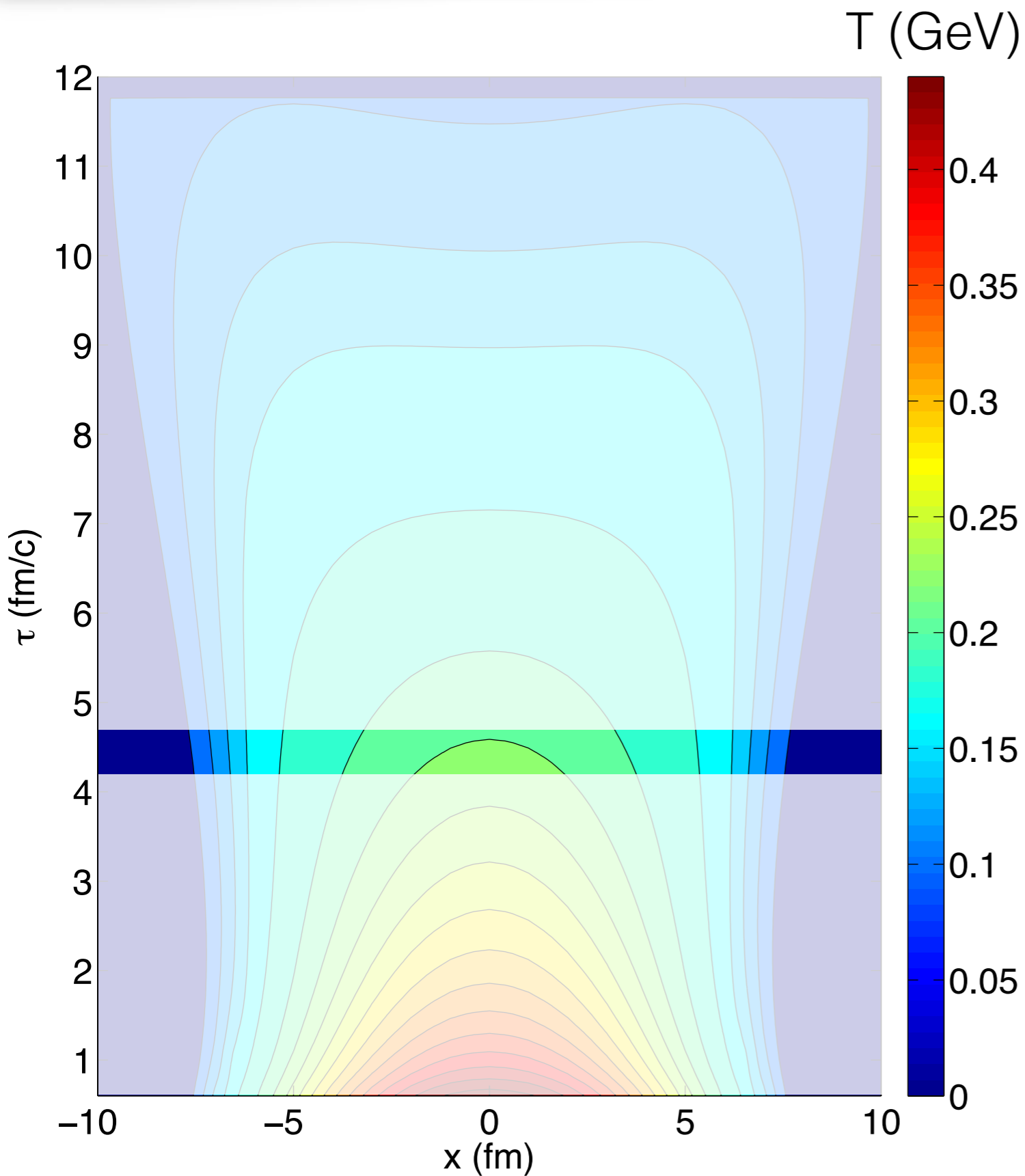
Slope of Photon Spectrum



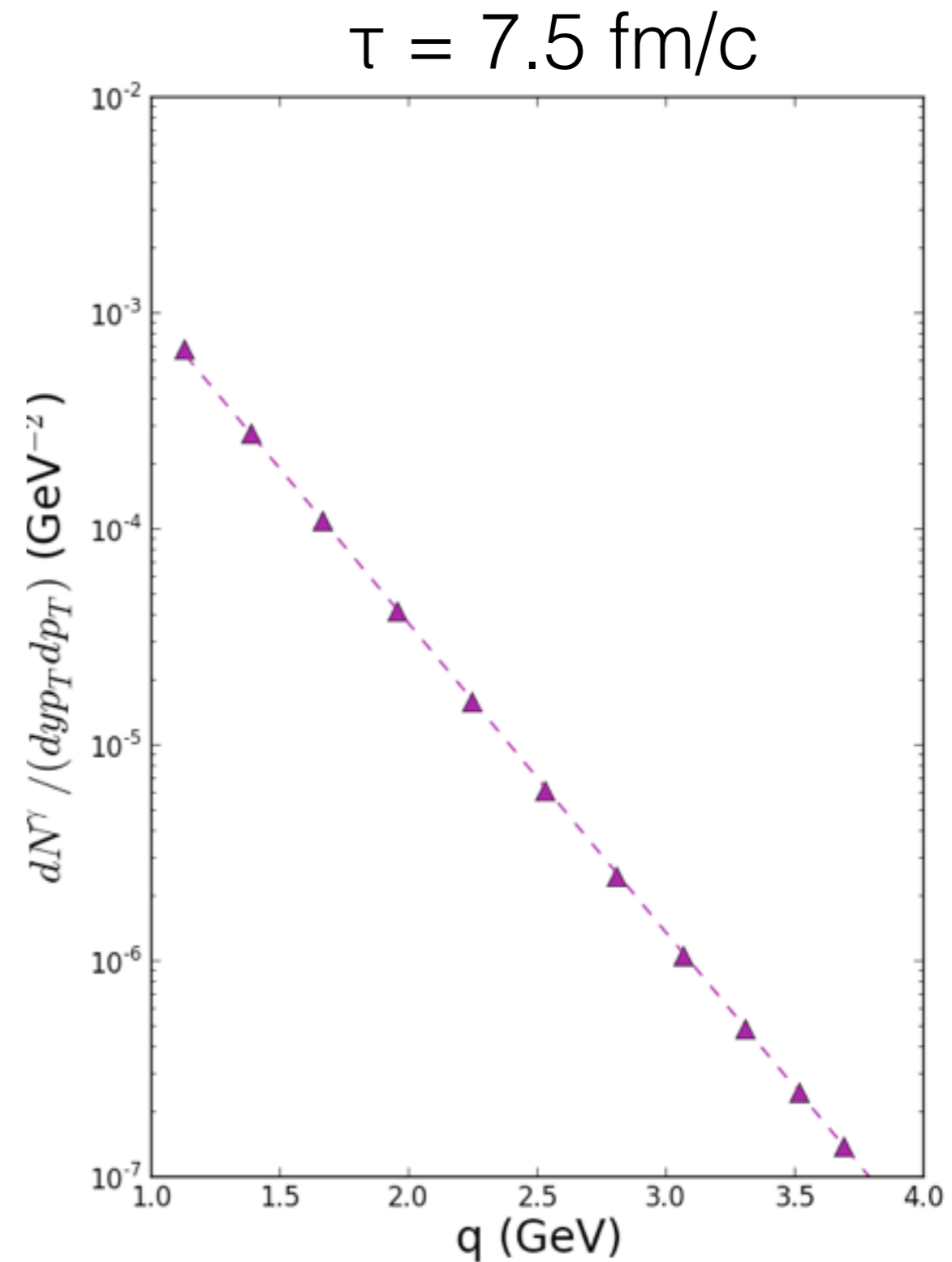
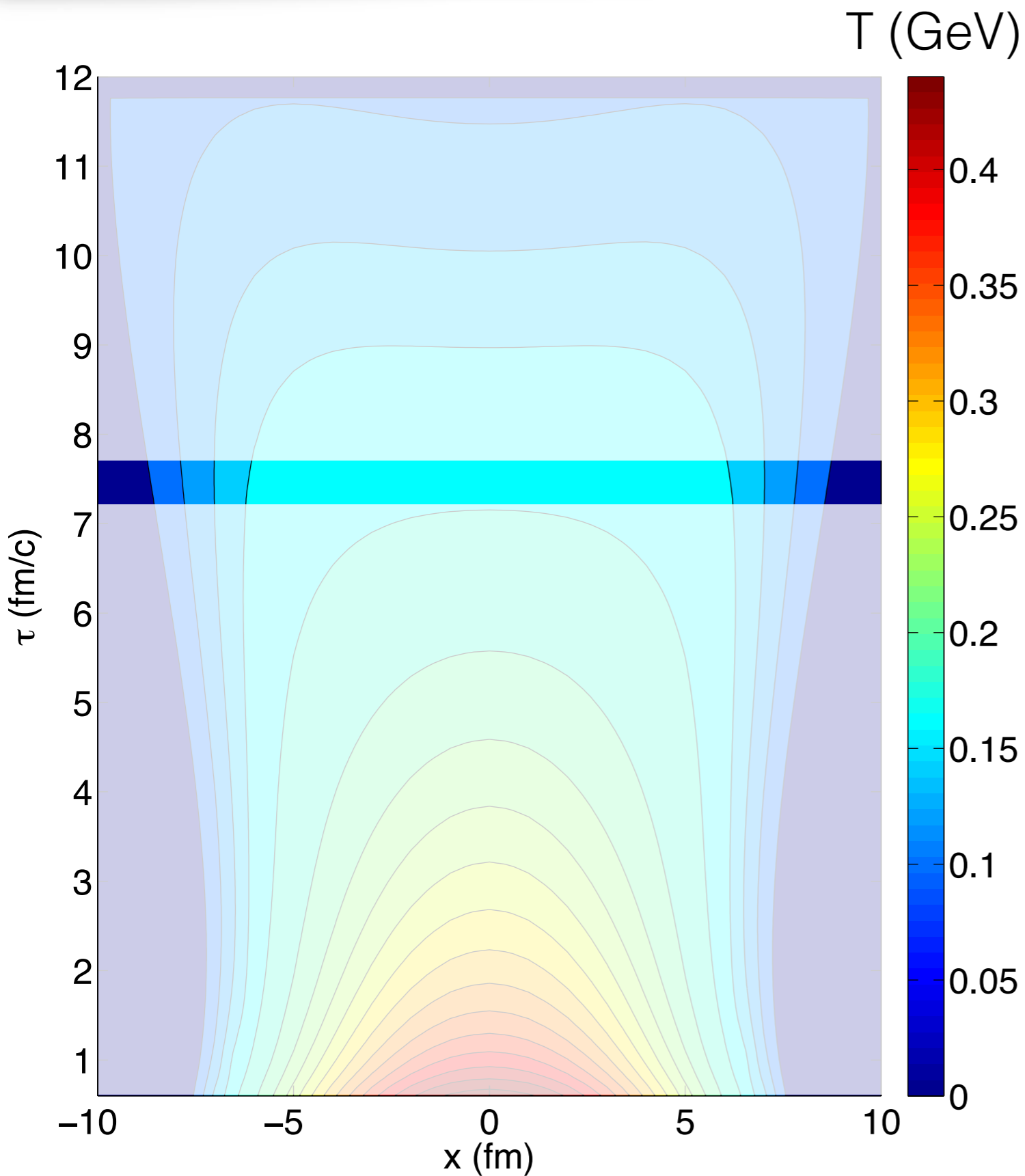
Slope of Photon Spectrum



Slope of Photon Spectrum

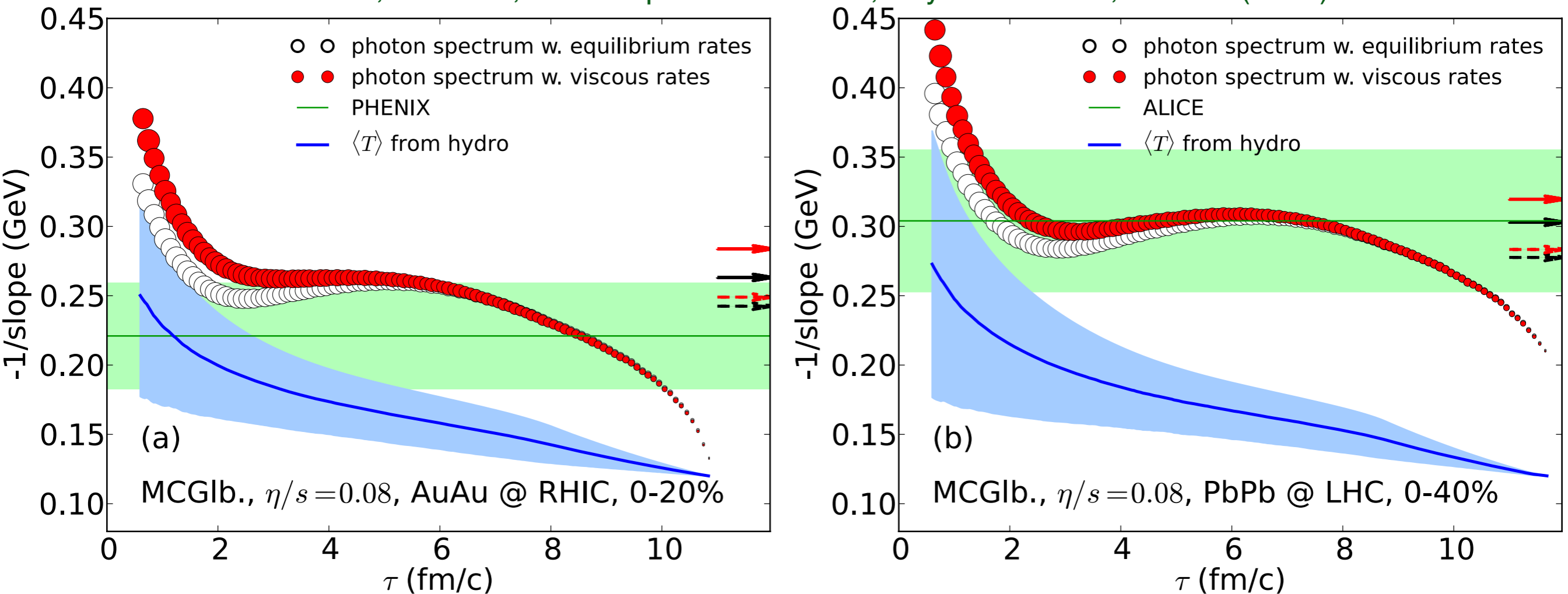


Slope of Photon Spectrum



Fitted T_{eff} vs. Emission Time

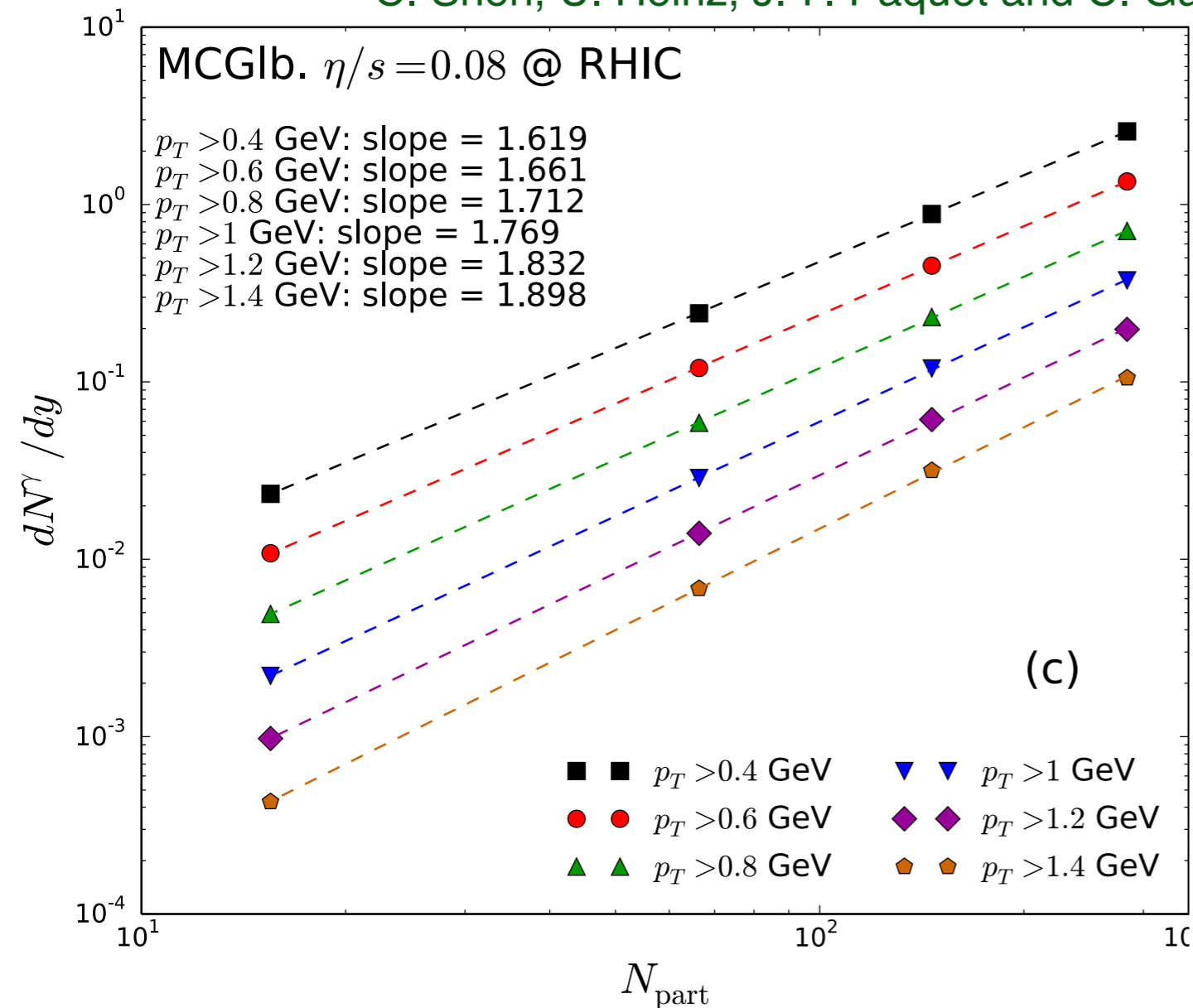
C. Shen, U. Heinz, J.-F. Paquet and C. Gale, Phys. Rev. C 89, 044910 (2014)



- About 25% of thermal photons are emitted in the first 2 fm/c
- After 2 fm/c, thermal photons are significantly blue shifted by radial flow
- Viscous corrections to the slope of photon spectra are stronger during the early part of the evolution

Centrality dependence of photon yield

C. Shen, U. Heinz, J.-F. Paquet and C. Gale, Phys. Rev. C 89, 044910 (2014)

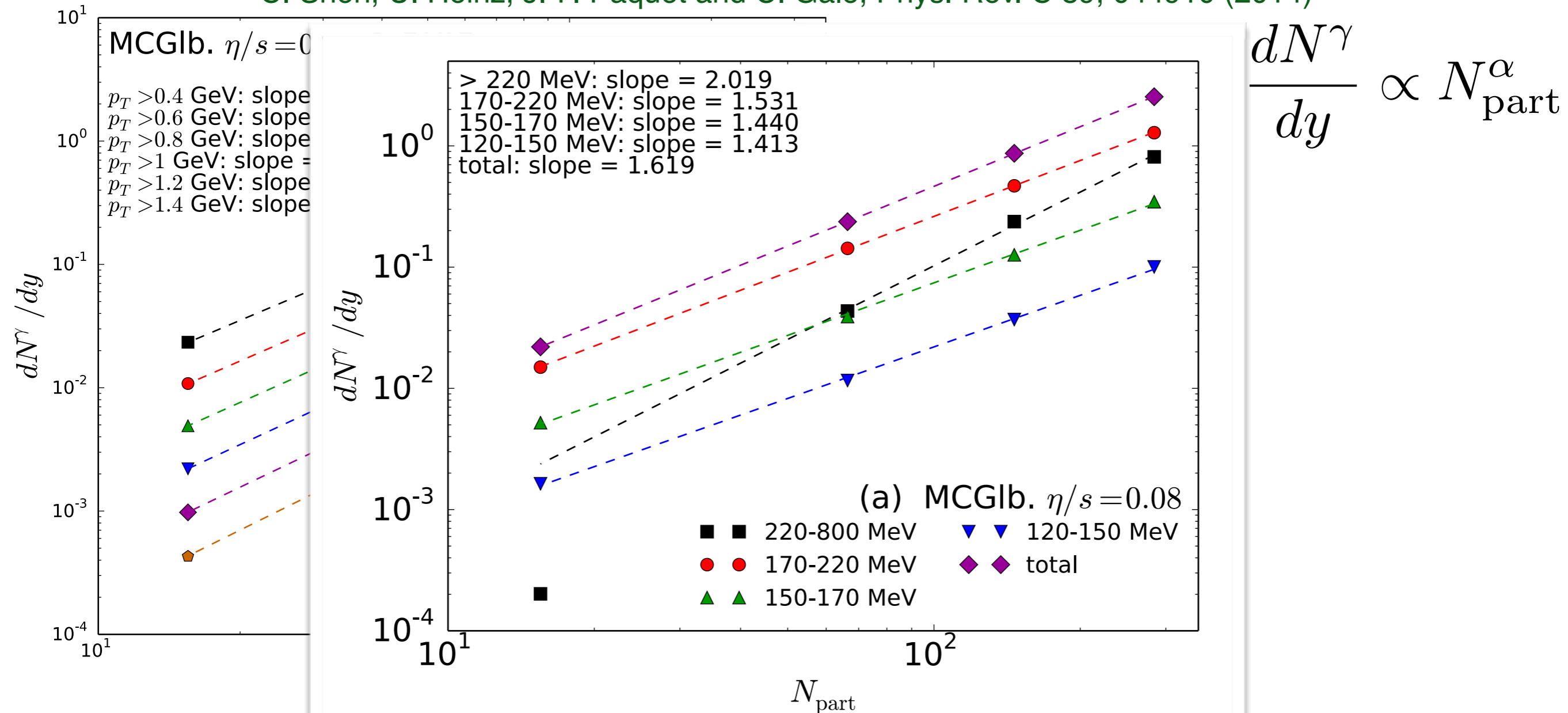


$$\frac{dN^\gamma}{dy} \propto N_{\text{part}}^\alpha$$

- Thermal photons from hydrodynamic medium qualitatively reproduce the centrality dependence of the direct excess photon yield at the top RHIC energy

Centrality dependence of photon yield

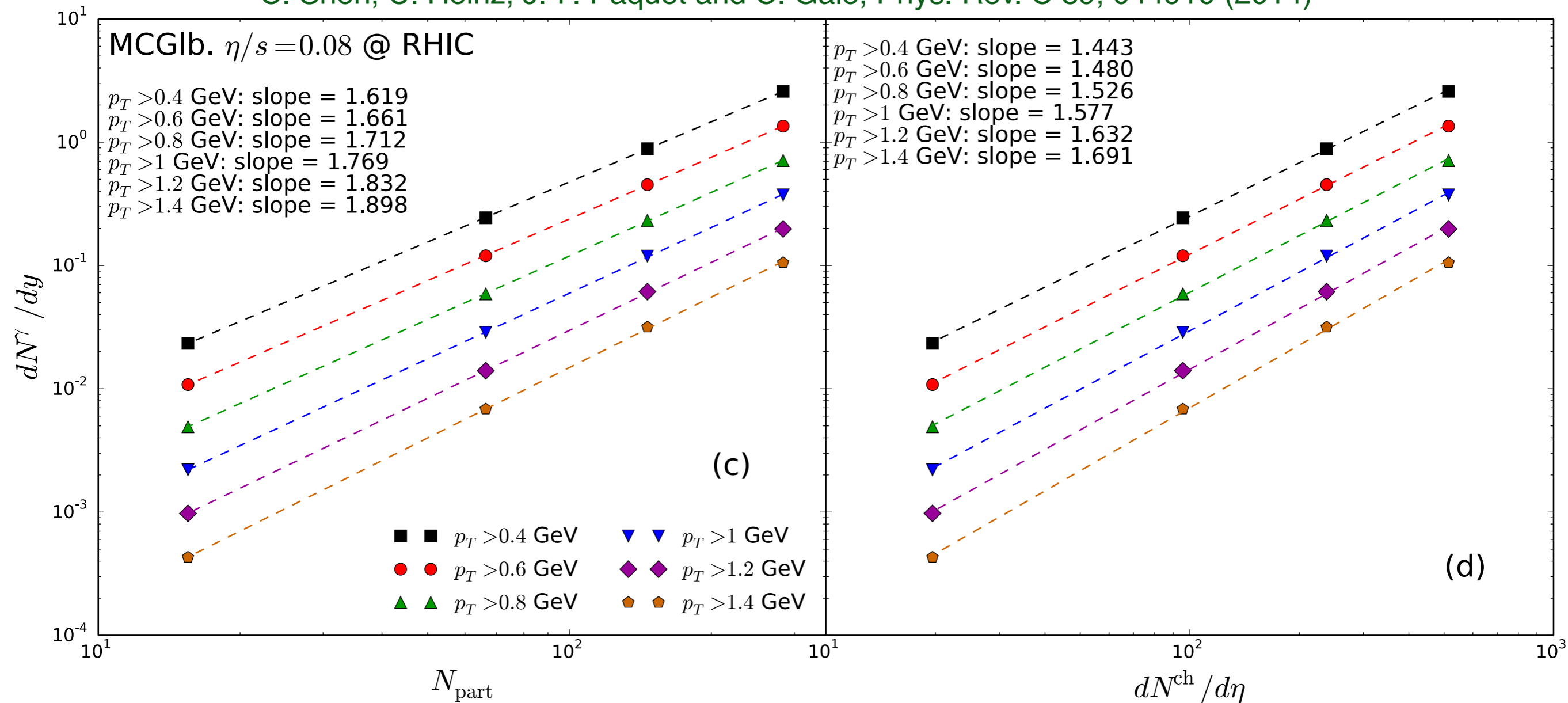
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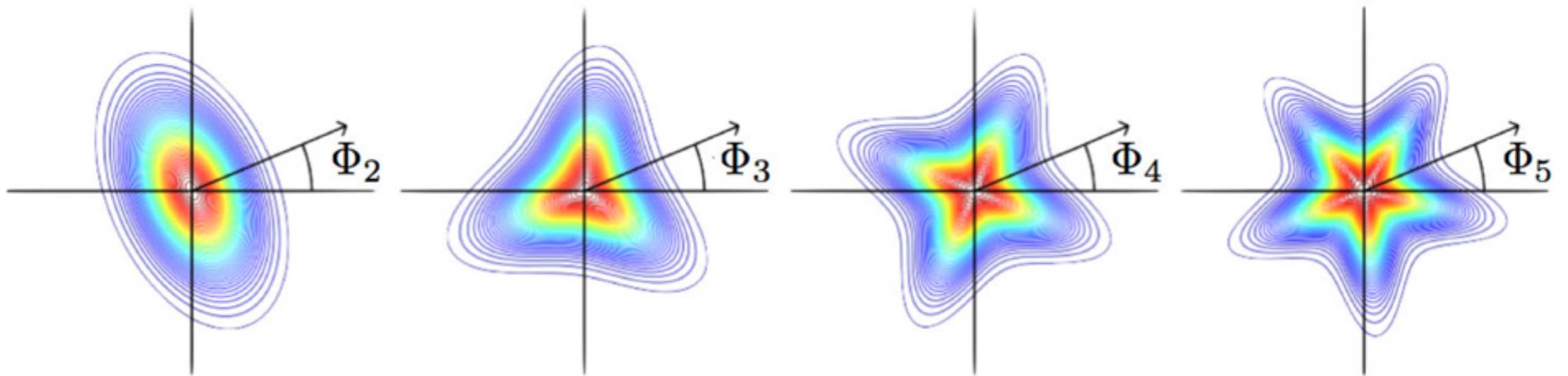
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dN^γ/dy vs. $dN^{\text{ch}}/d\eta$

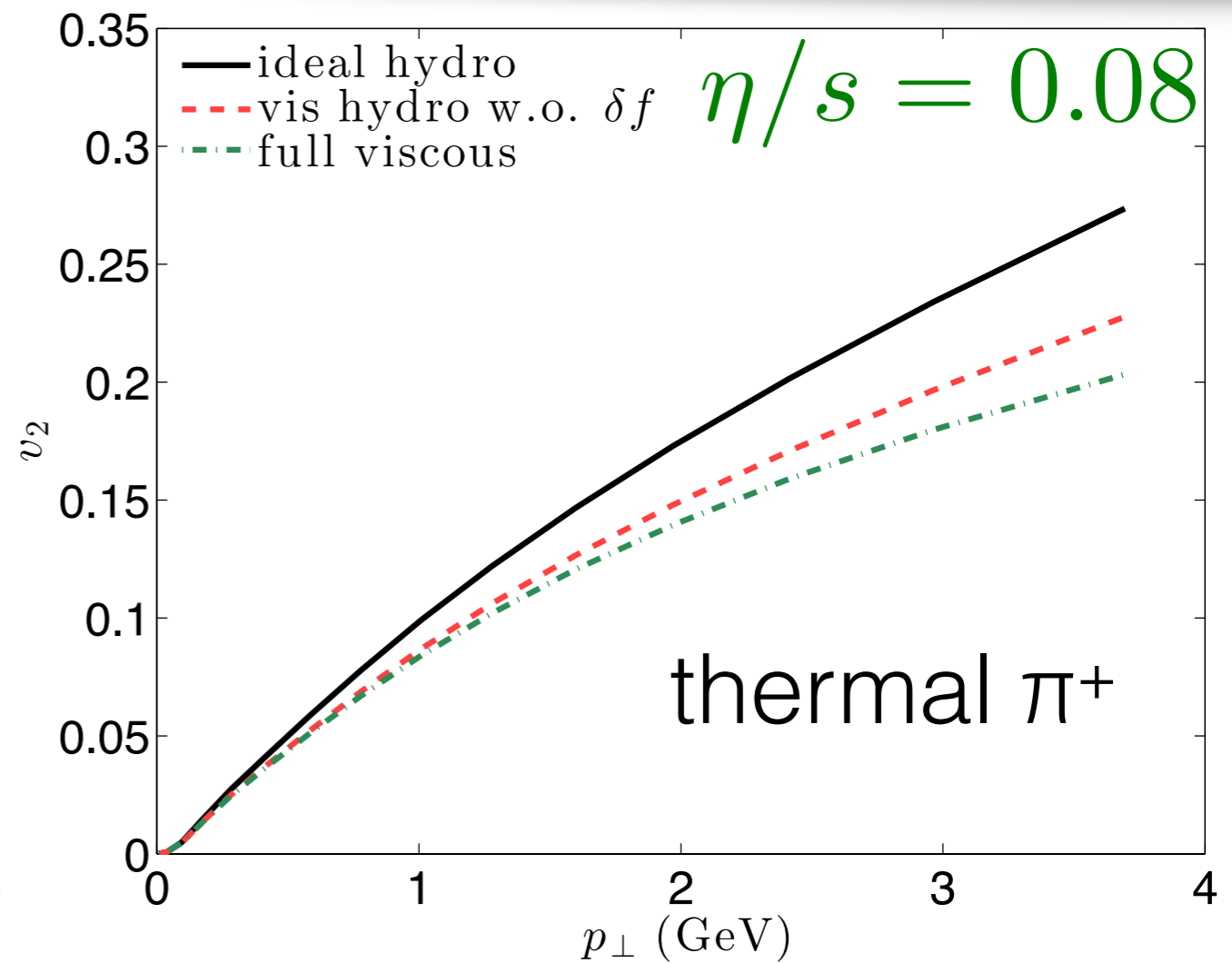
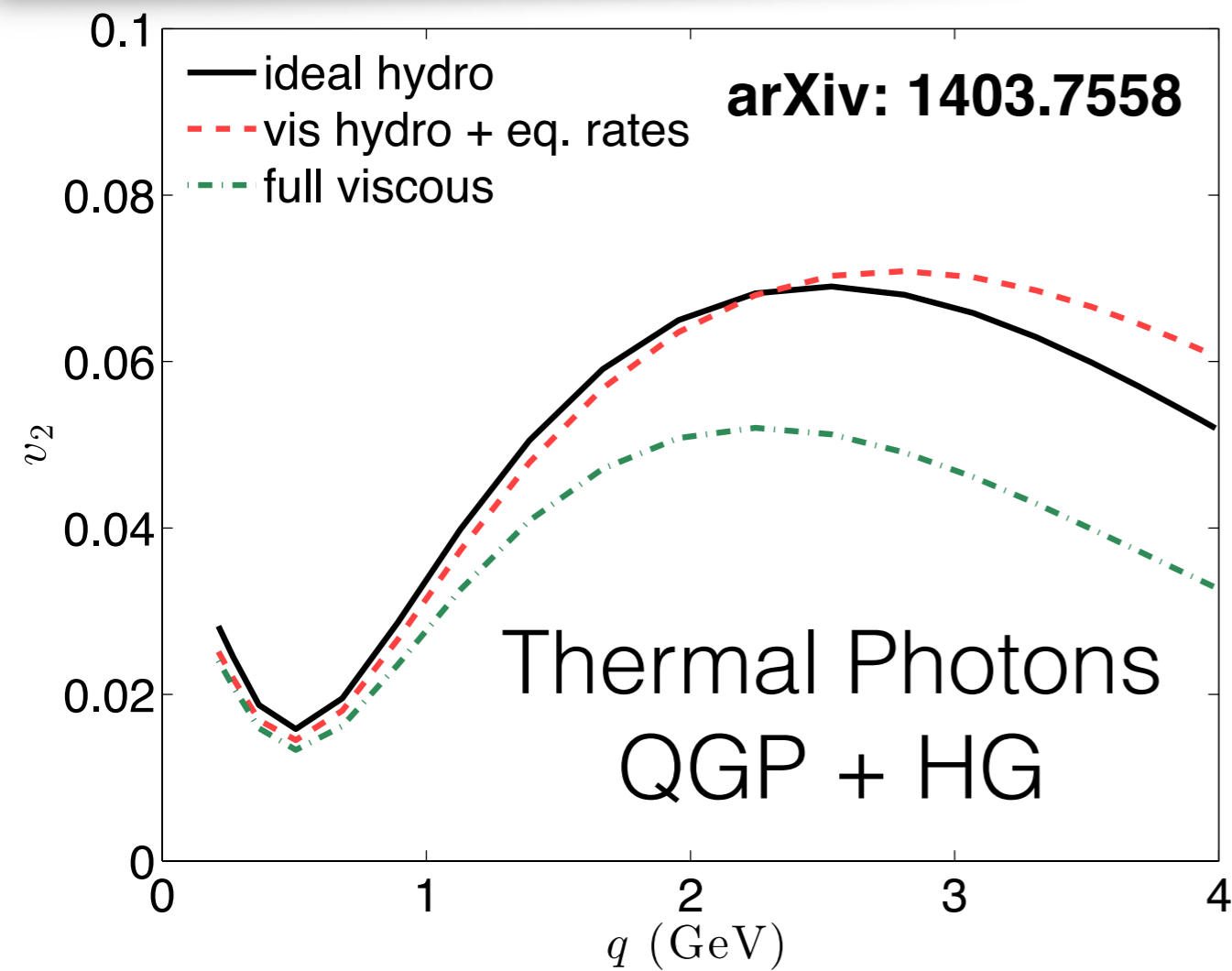


**less model dependent
comparison**

Photon anisotropic flow

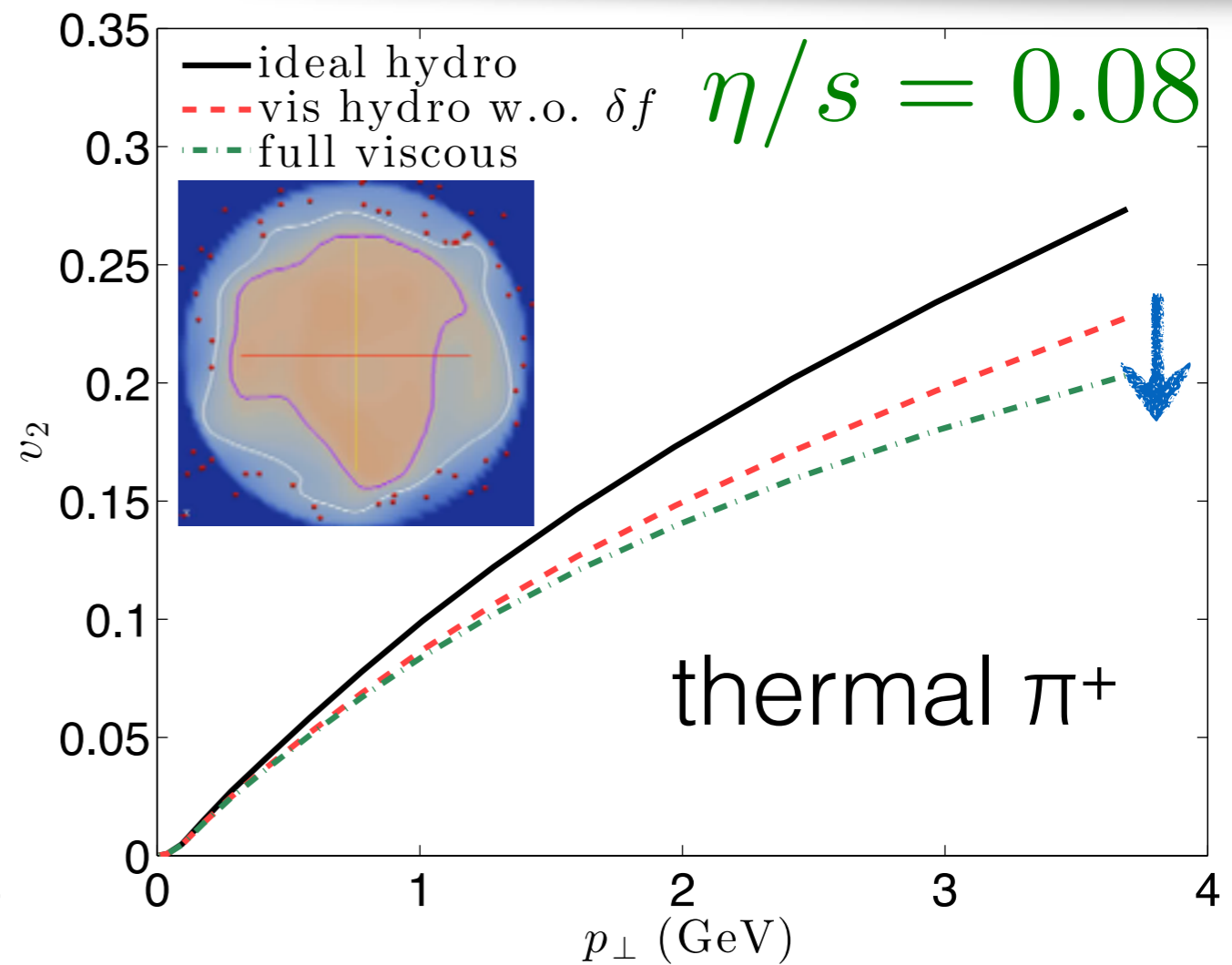
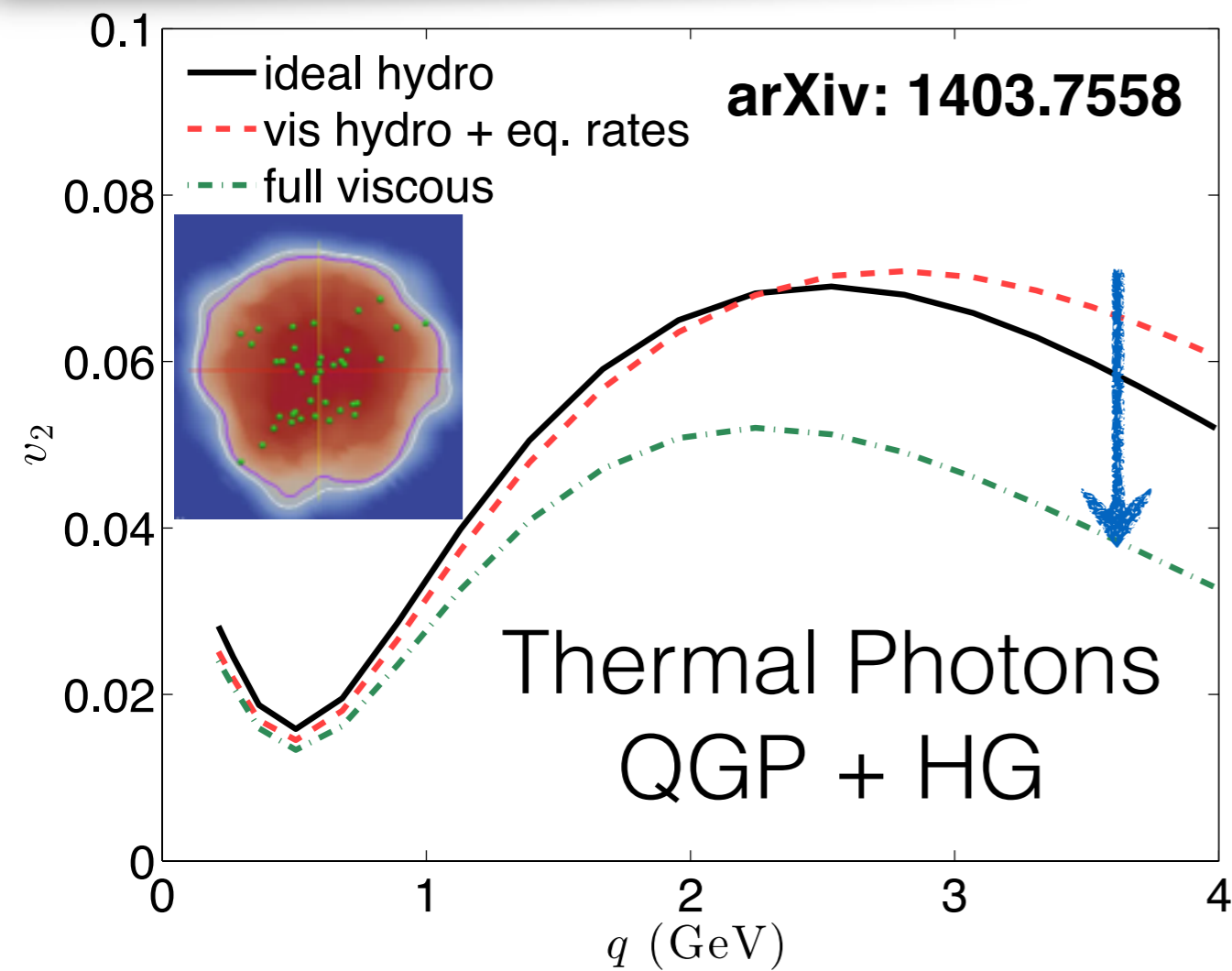


Shear viscous effects on photon elliptic flow



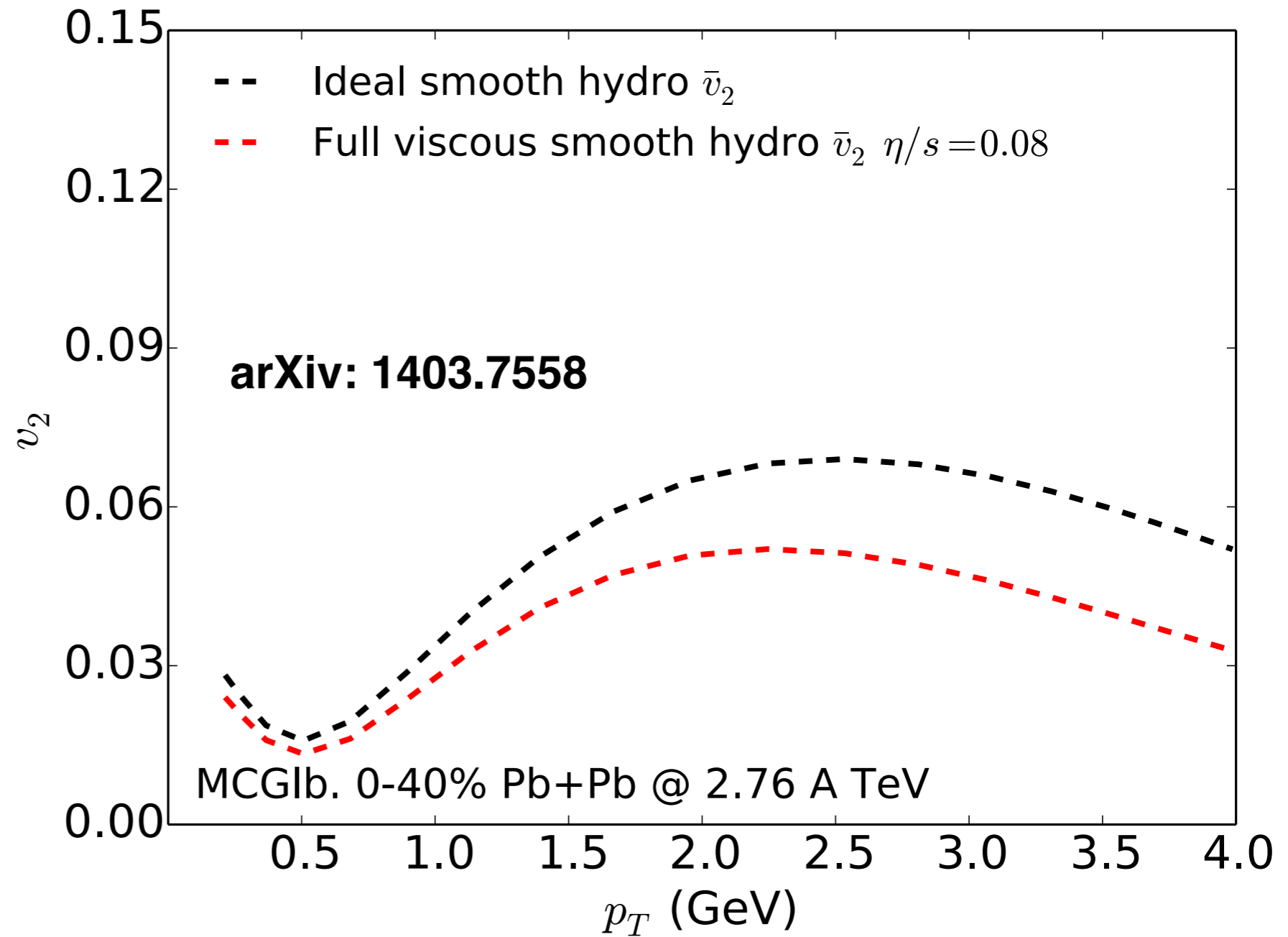
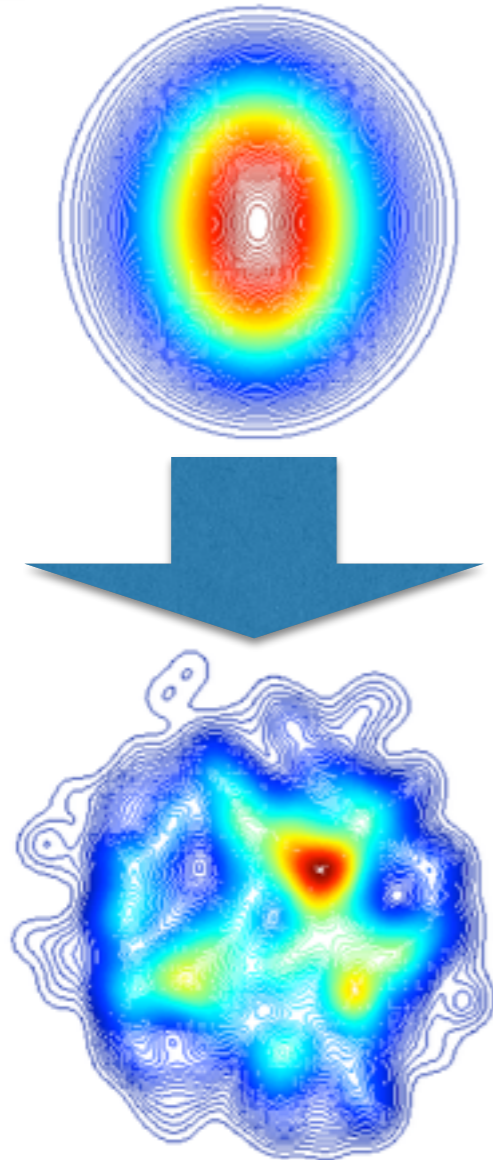
- Shear viscous suppression of photon v_2 is dominated by the viscous corrections to the photon emission rate
- Photon elliptic flow is sensitive to the larger shear stress tensor at early times

Shear viscous effects on photon elliptic flow

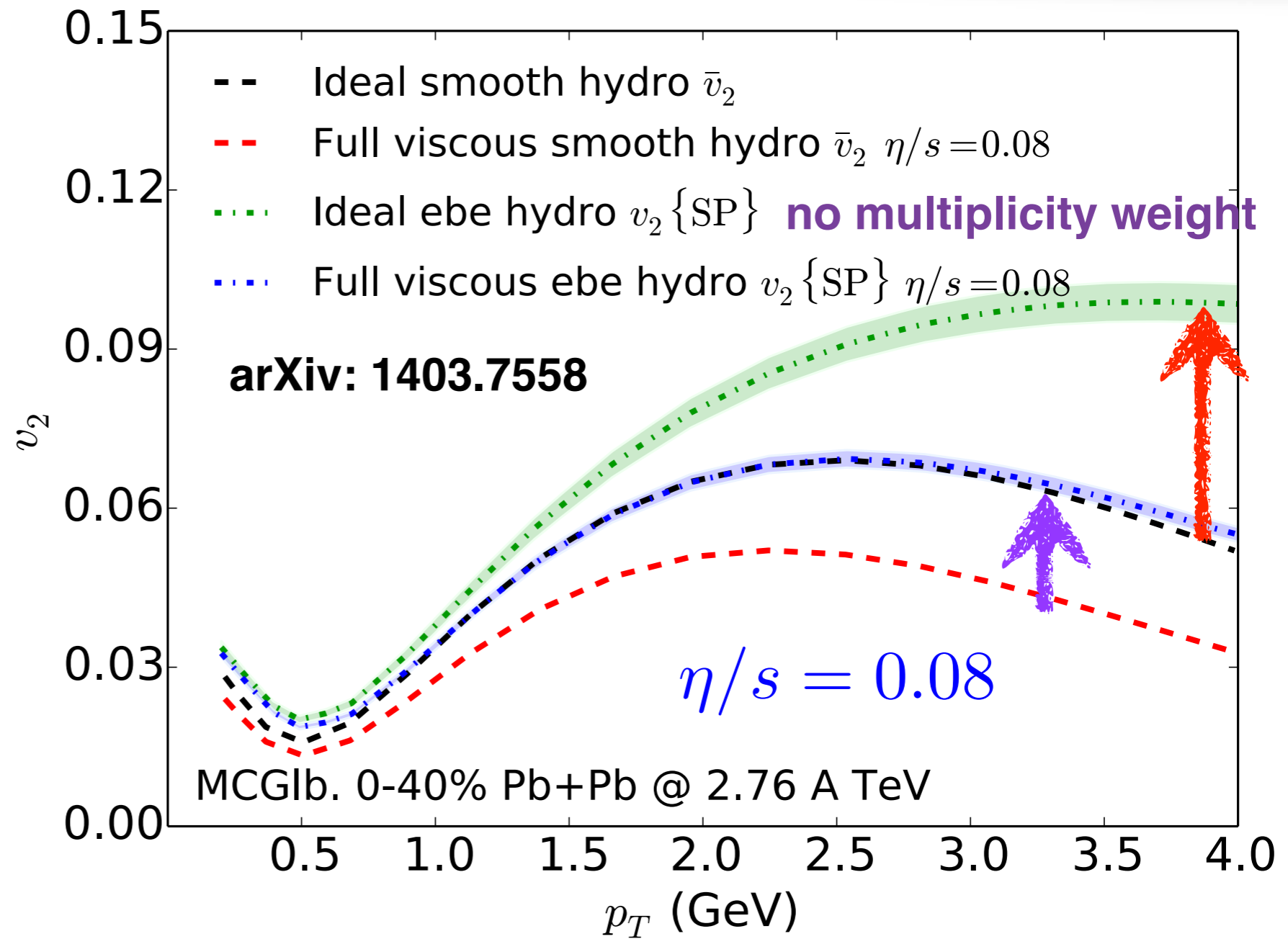
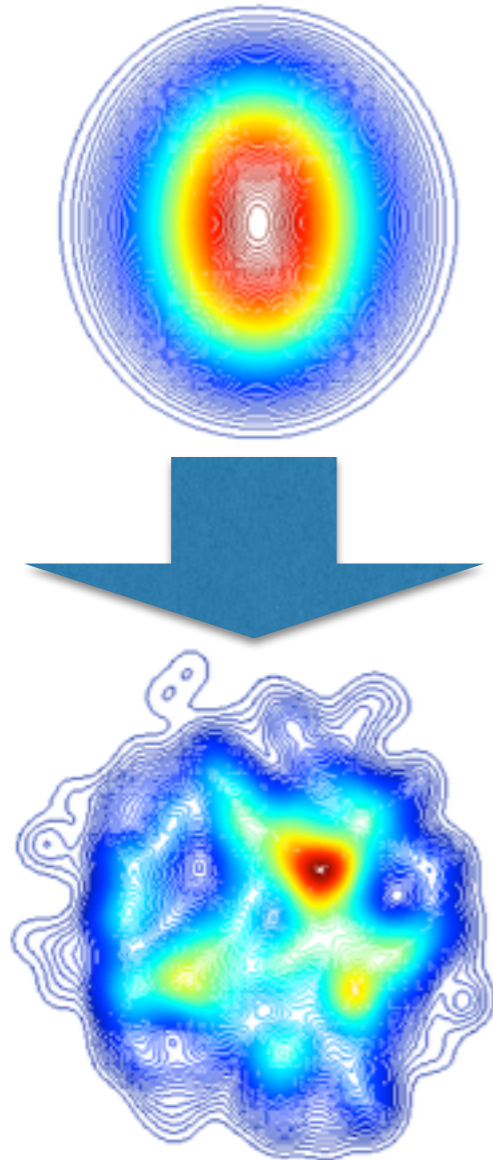


- Shear viscous suppression of photon v_2 is dominated by the viscous corrections to the photon emission rate
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Fluctuation effects on photon elliptic flow

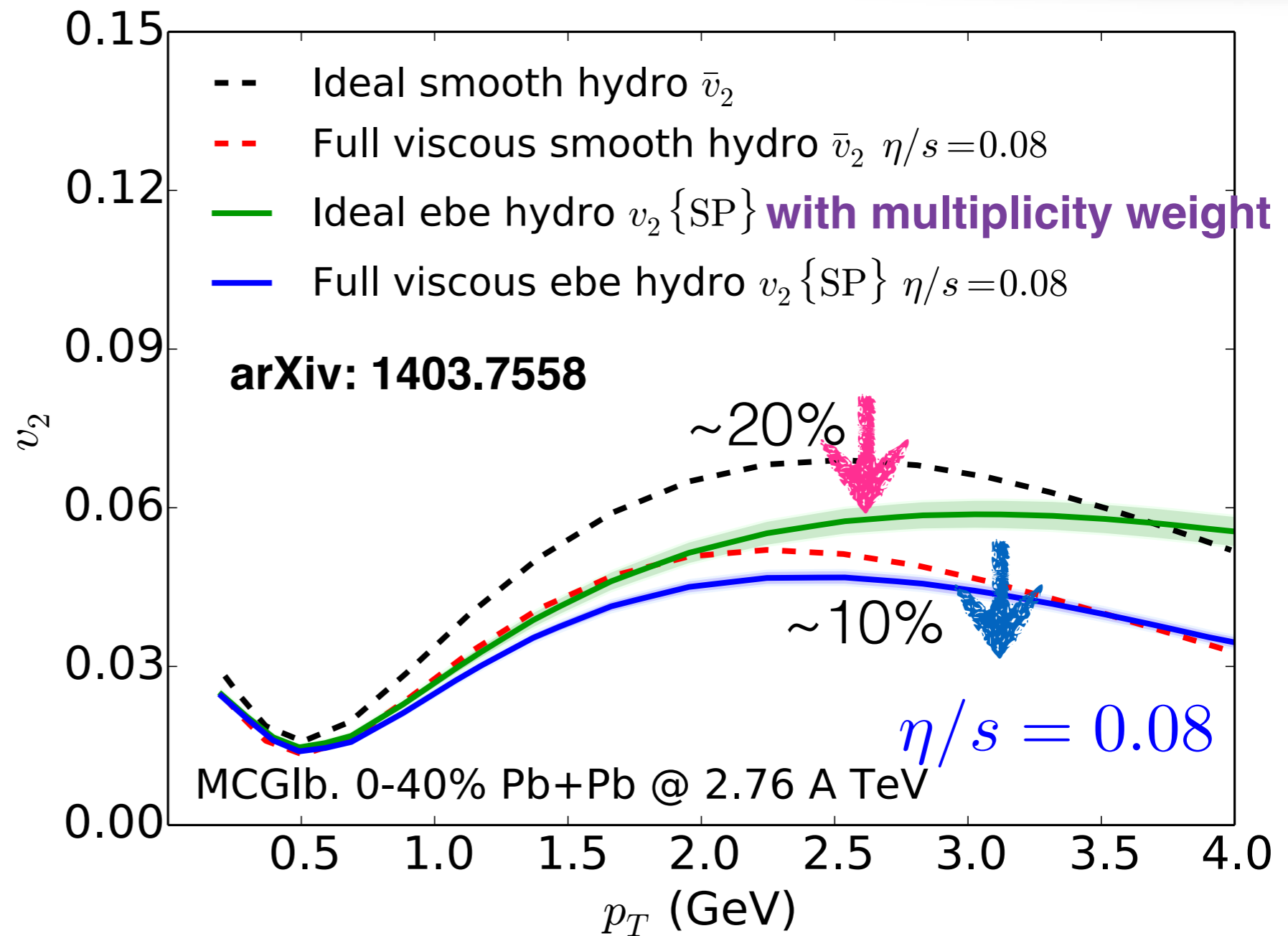
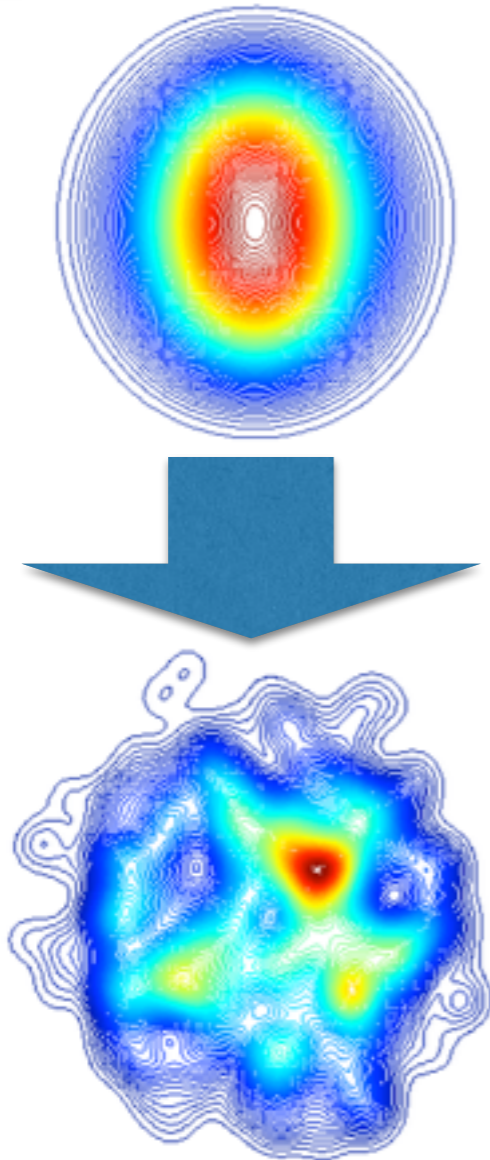


Fluctuation effects on photon elliptic flow



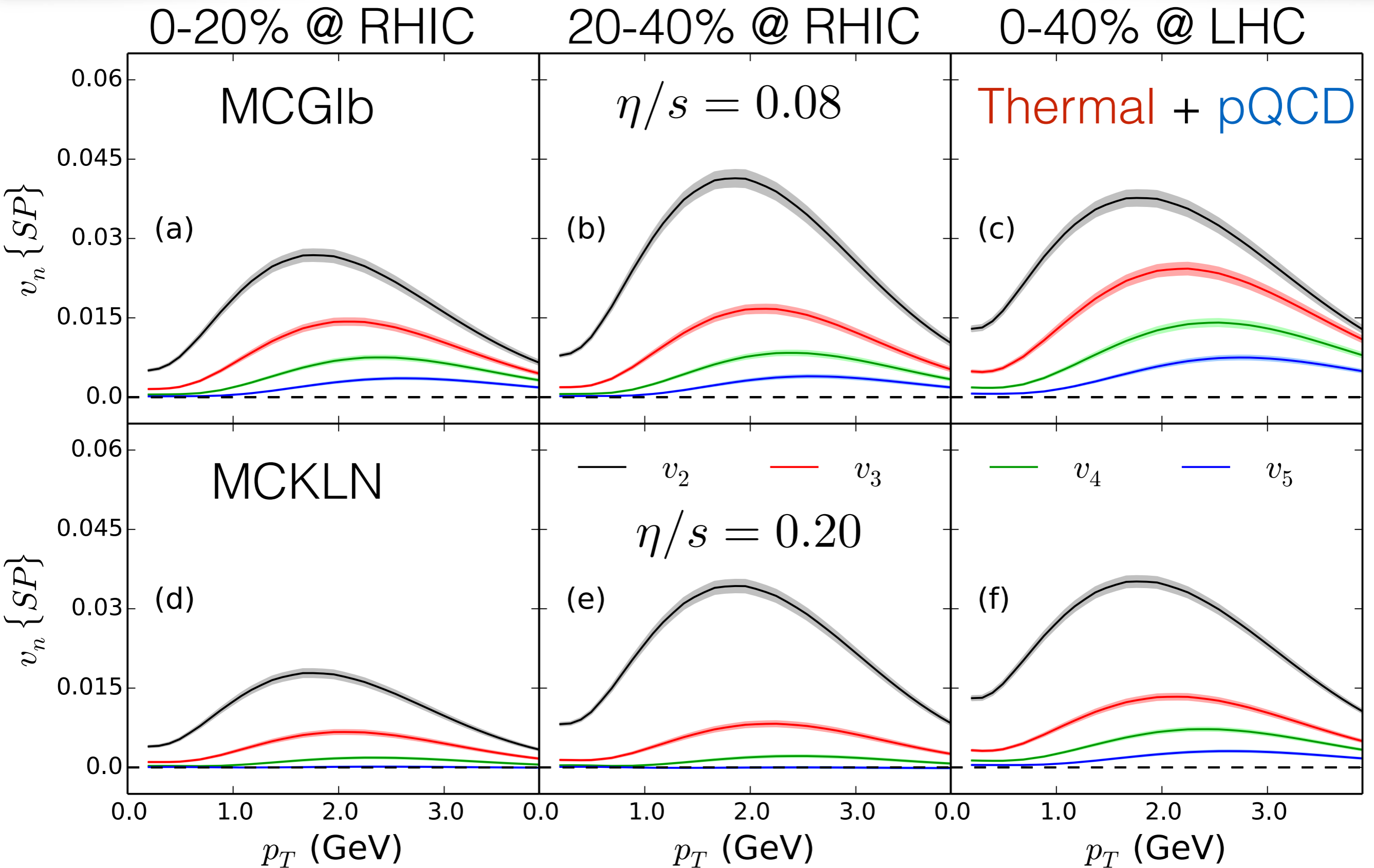
- ▶ Initial fluctuations increase photons' elliptic flow

Fluctuation effects on photon elliptic flow



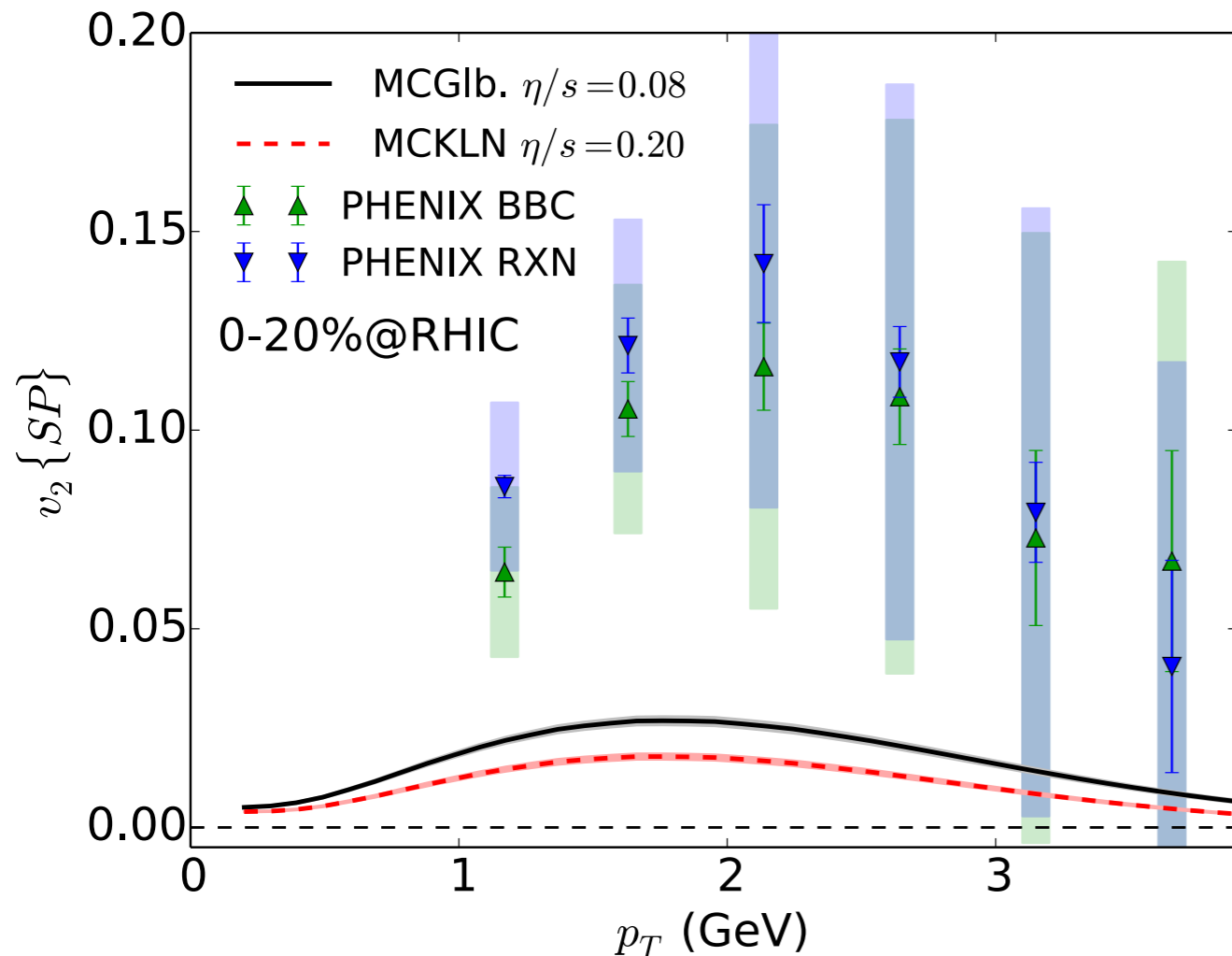
- ▶ Initial fluctuations increase photons' elliptic flow
- ▶ The additional photon multiplicity weighting biases e-b-e v_2 towards central collisions, resulting in ~10-20% smaller v_2 compared to smooth hydro

Event-by-Event Full Viscous Photon v_n

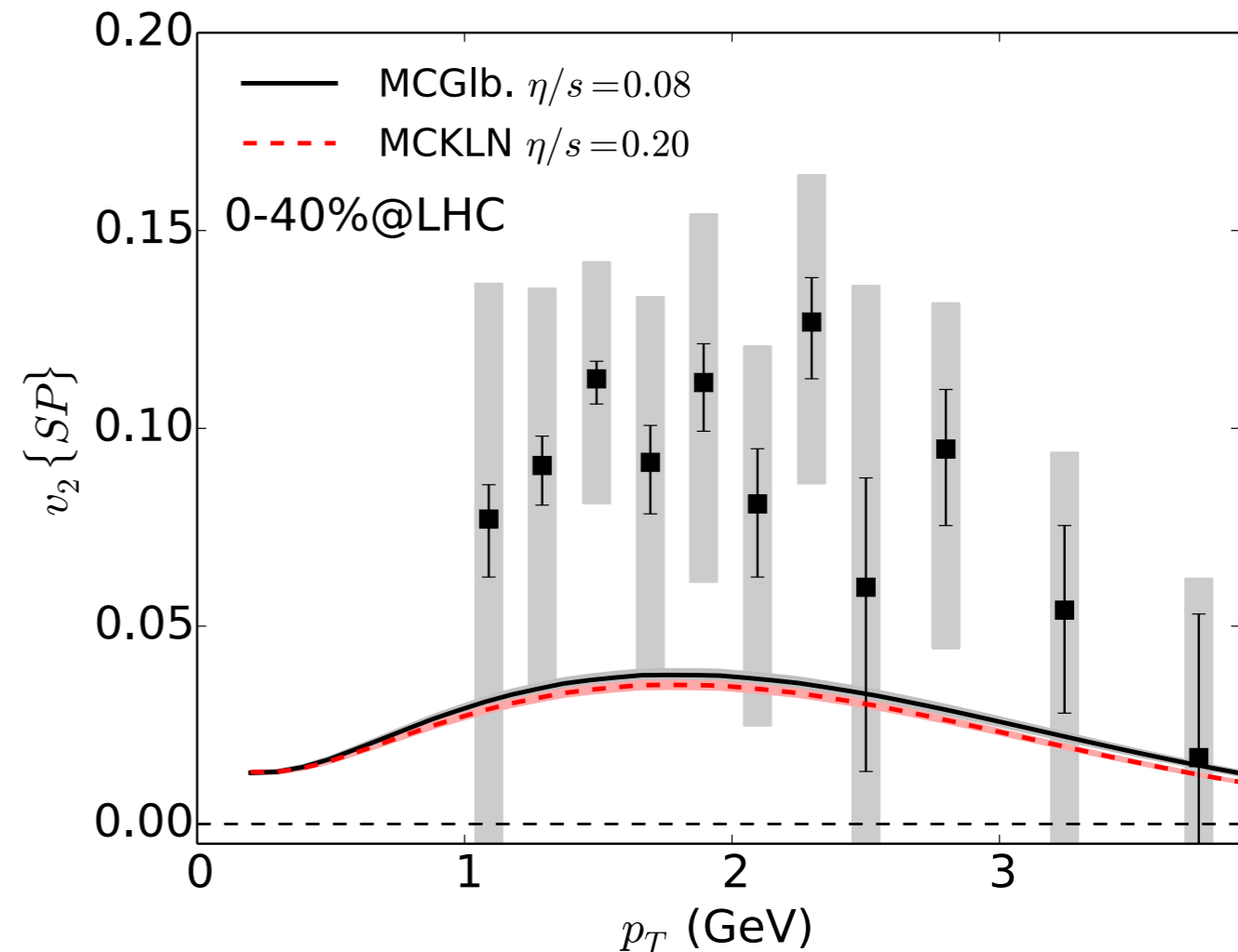


The sky falls ...

RHIC 0-20%



LHC 0-40%

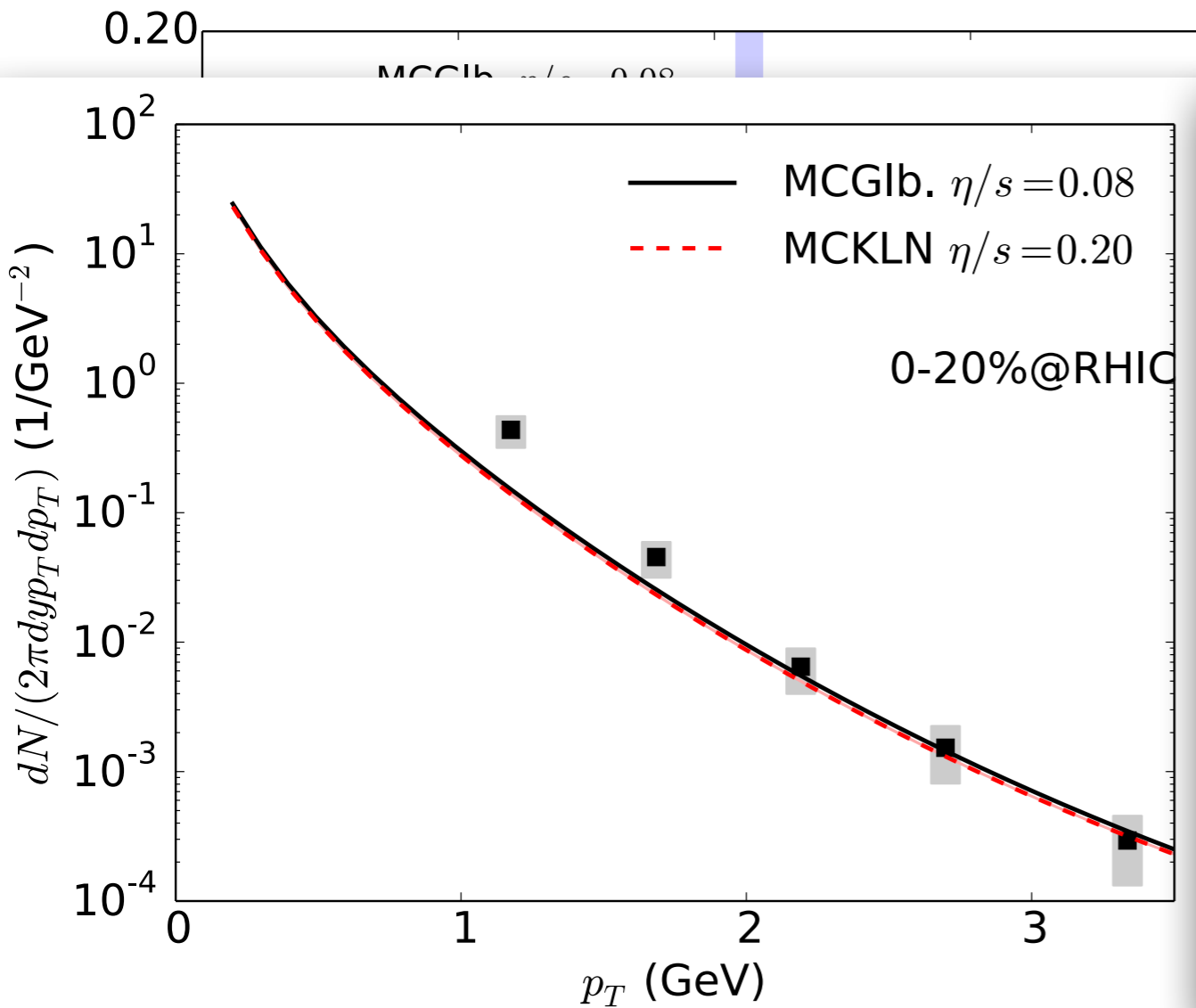


- Current calculations still underestimate the experimental data by a factor of 3!

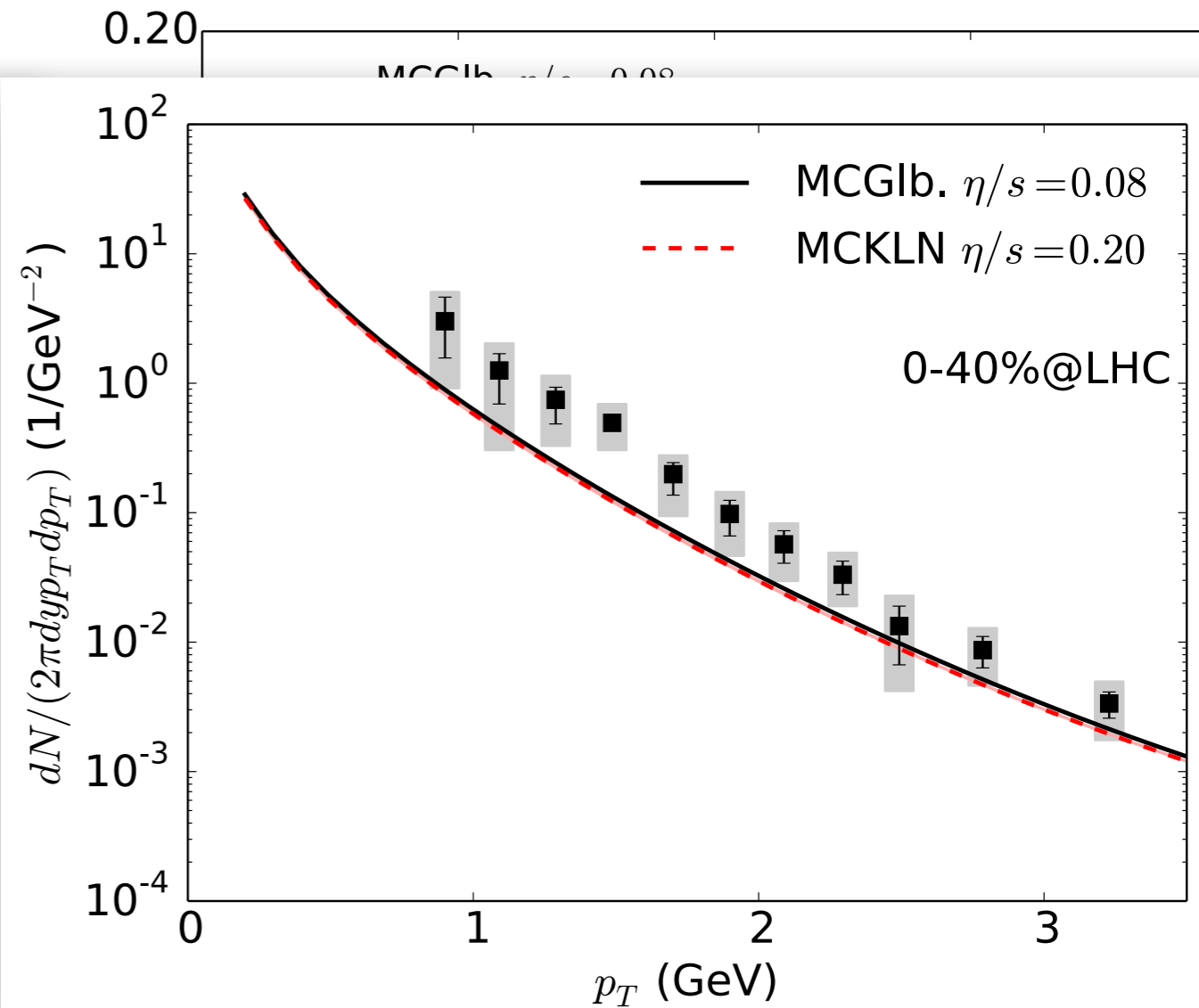


The sky falls ...

RHIC 0-20%



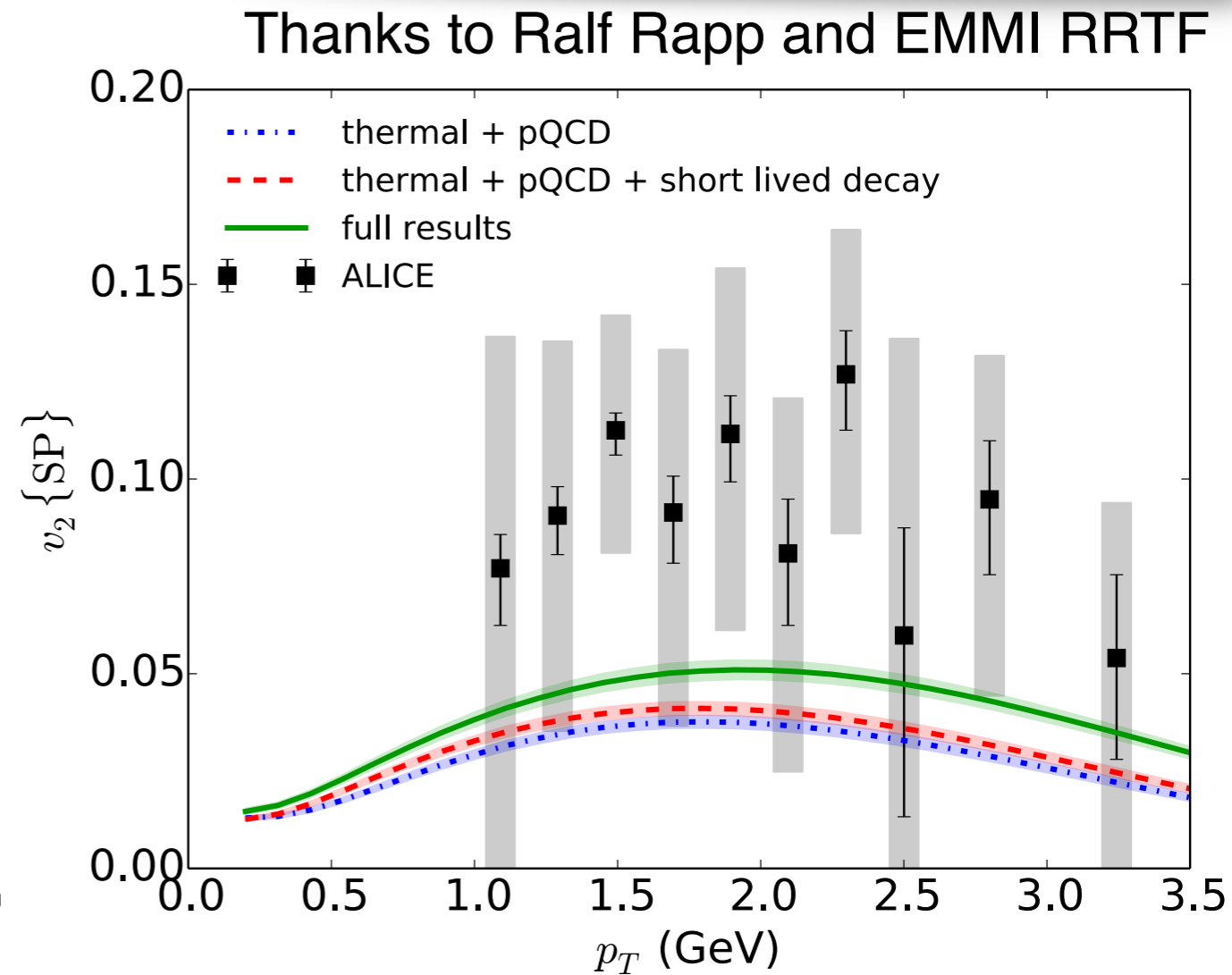
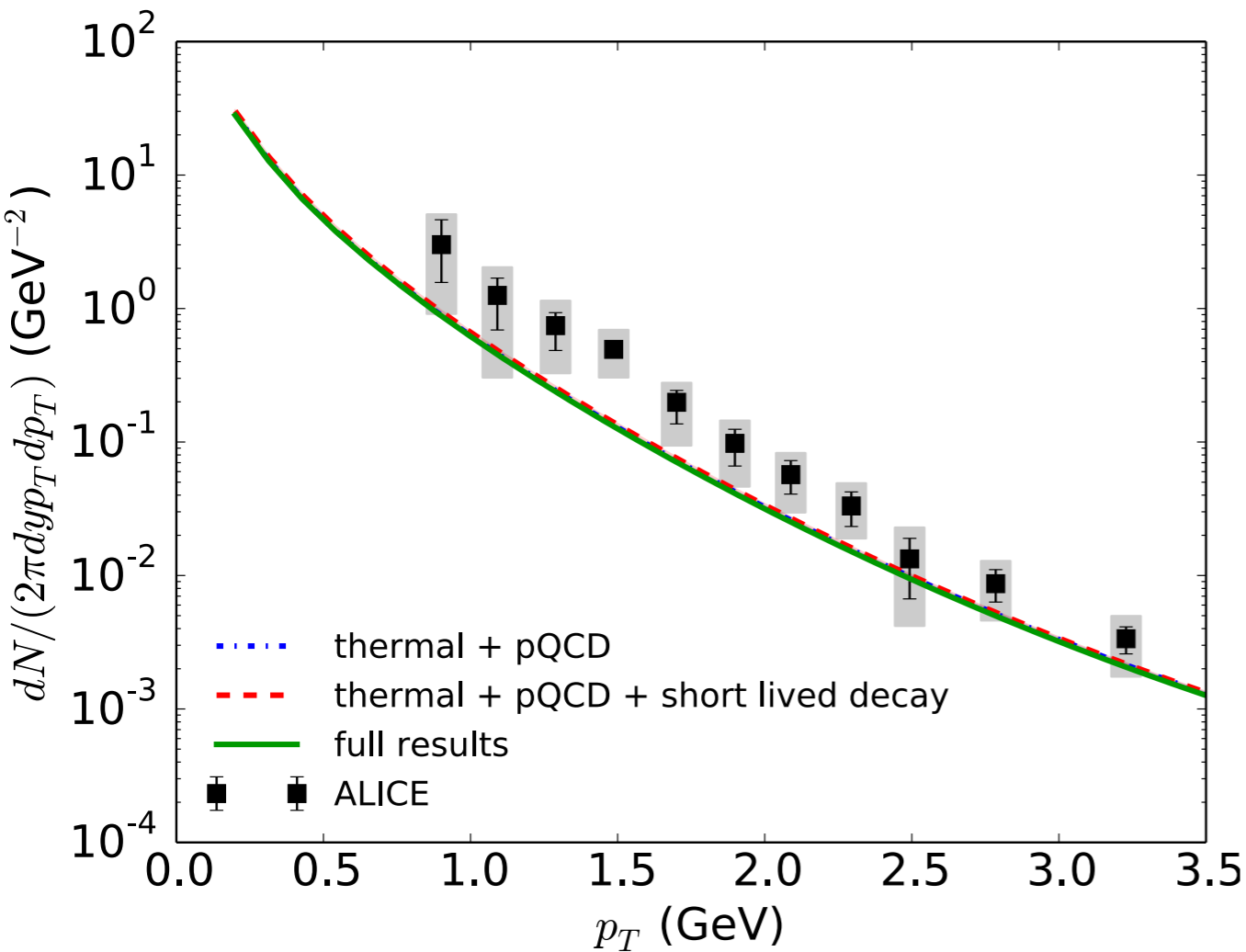
LHC 0-40%



- Current calculations still underestimate the experimental data by a factor of 3!
- Thermal yield is also missing in the azimuthally integrated photon spectra at low



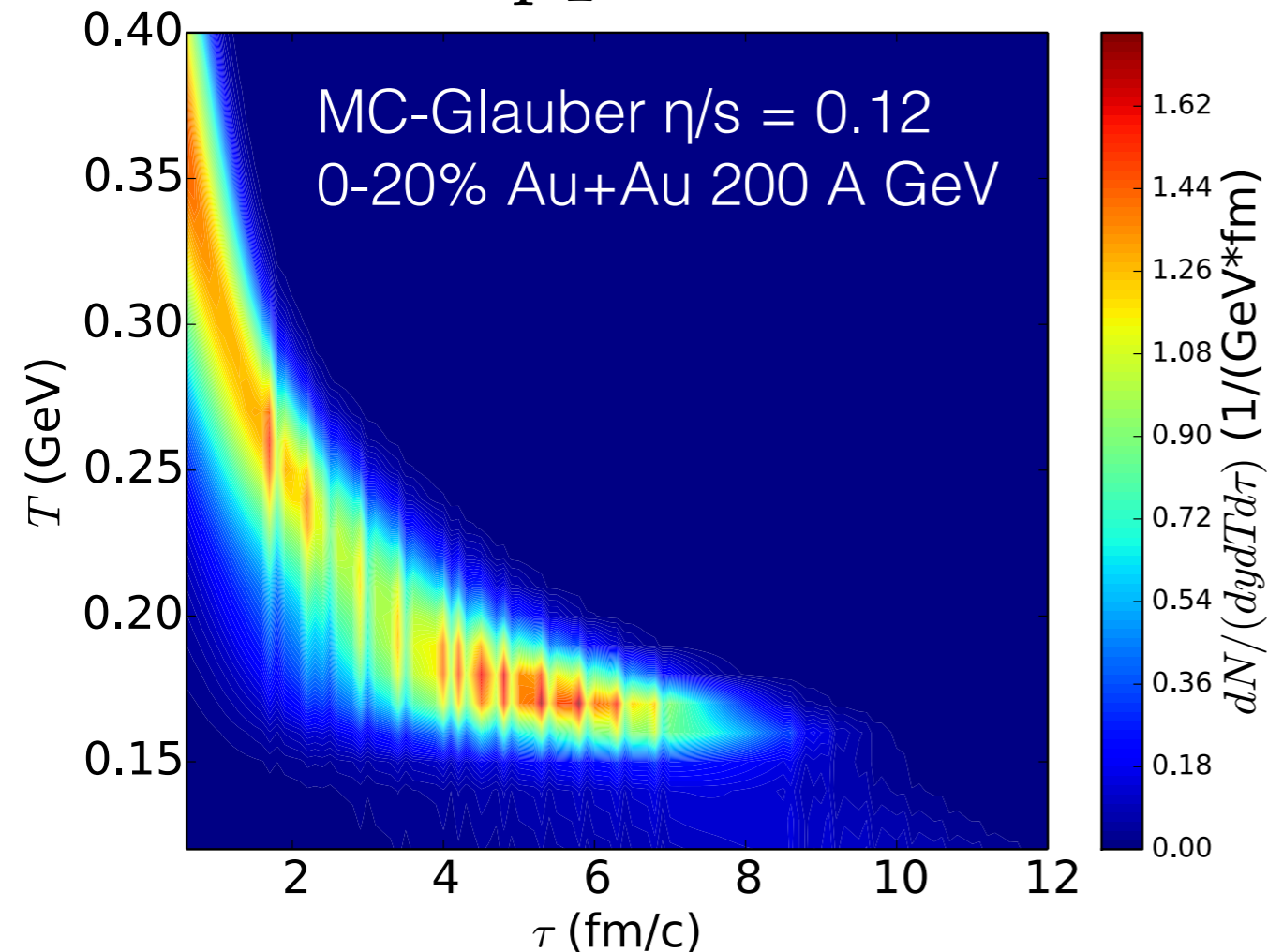
Efforts to resolve the photon flow puzzle



- The post freeze-out **short-lived resonances** give small but positive contributions
- **Pre-equilibrium flow** helps the fireball to develop the flow anisotropy more quickly and improves the theoretical calculations

Thermal photon tomography

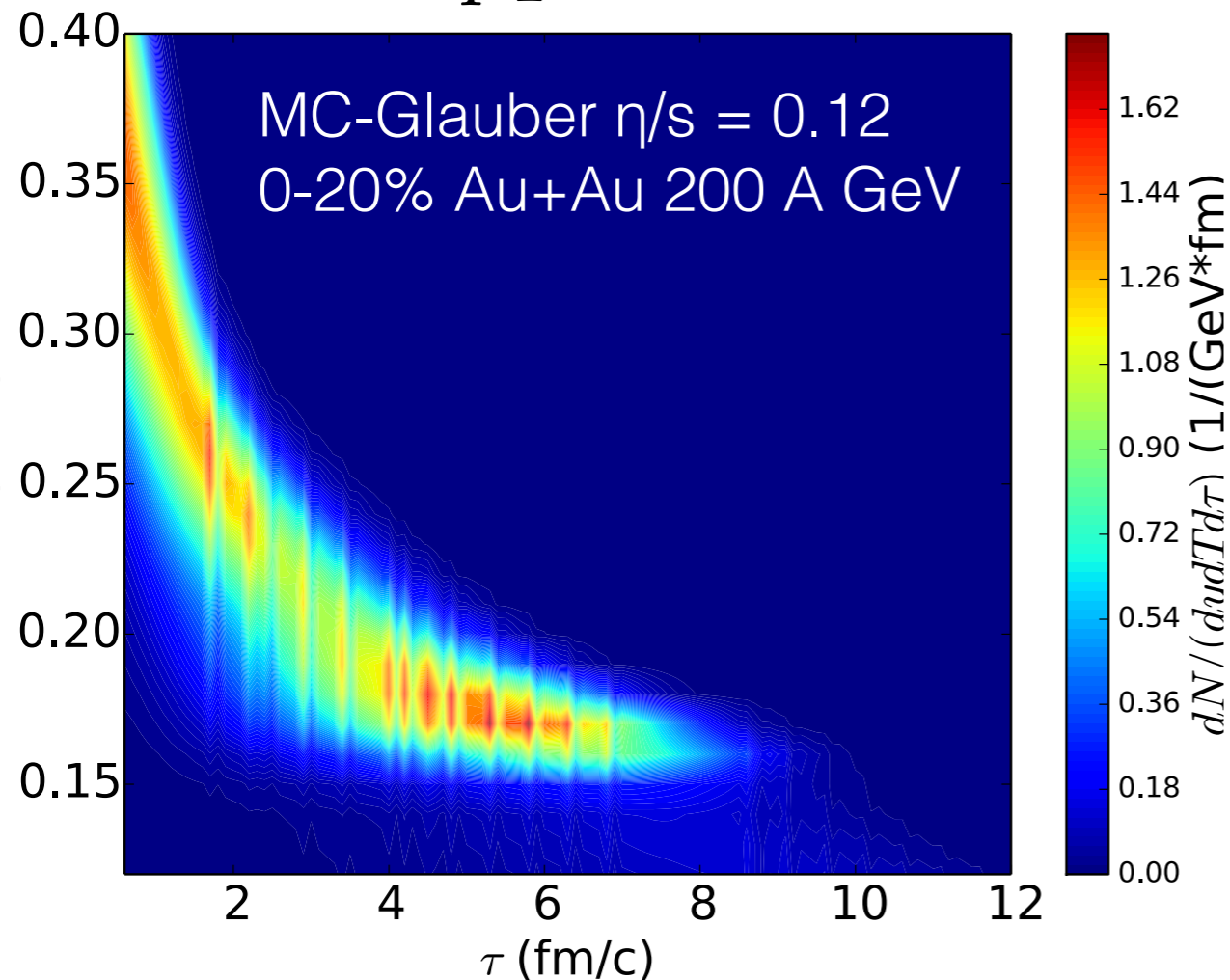
$$1 \leq p_T \leq 4 \text{ GeV}$$



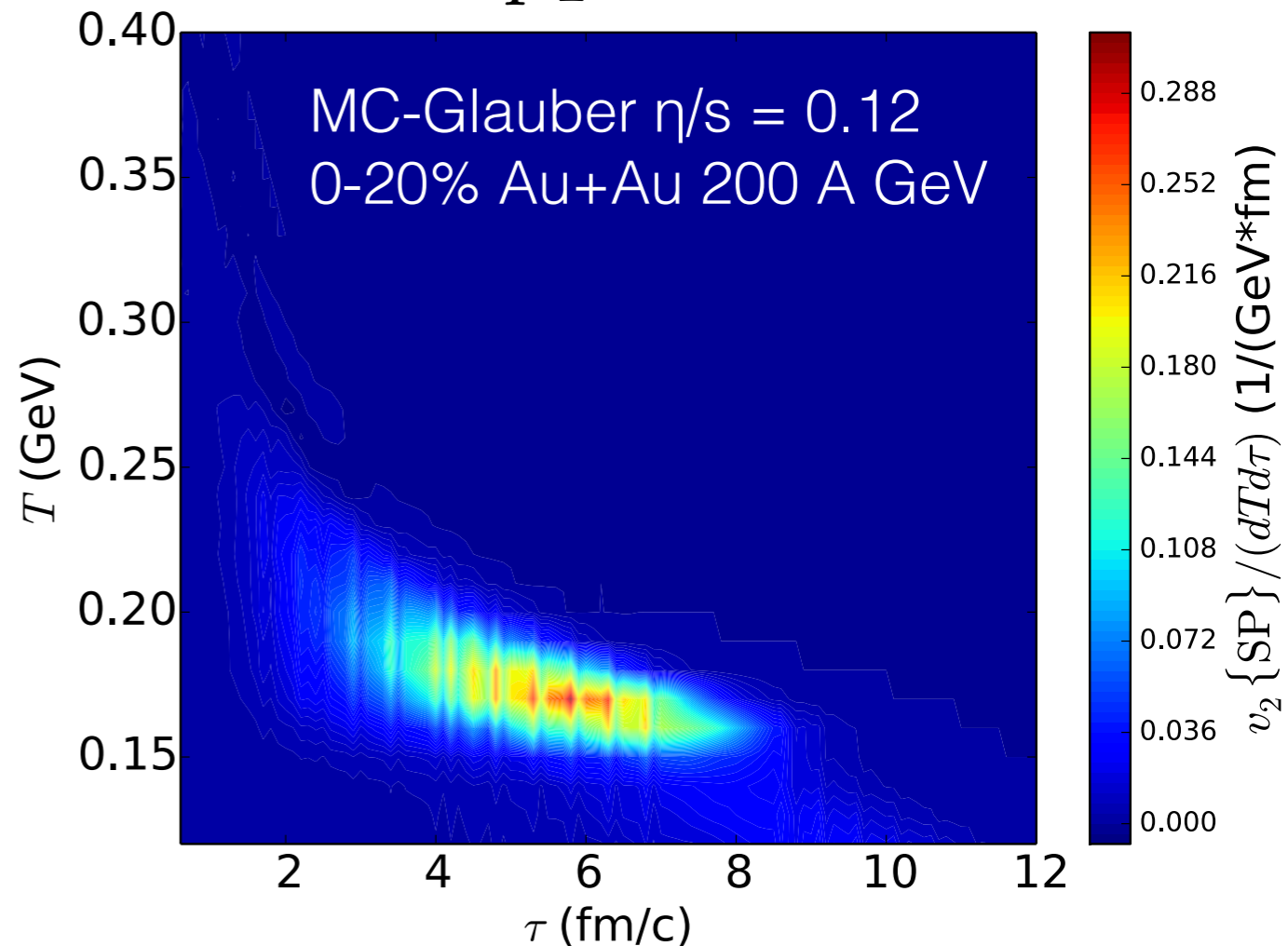
- By cutting hydro medium both in T and τ , we observe a **two-wave** thermal photon production
 - early time production — high rates at high temperatures
 - near transition region — growing of space-time volume

Thermal photon tomography

$$1 \leq p_T \leq 4 \text{ GeV}$$



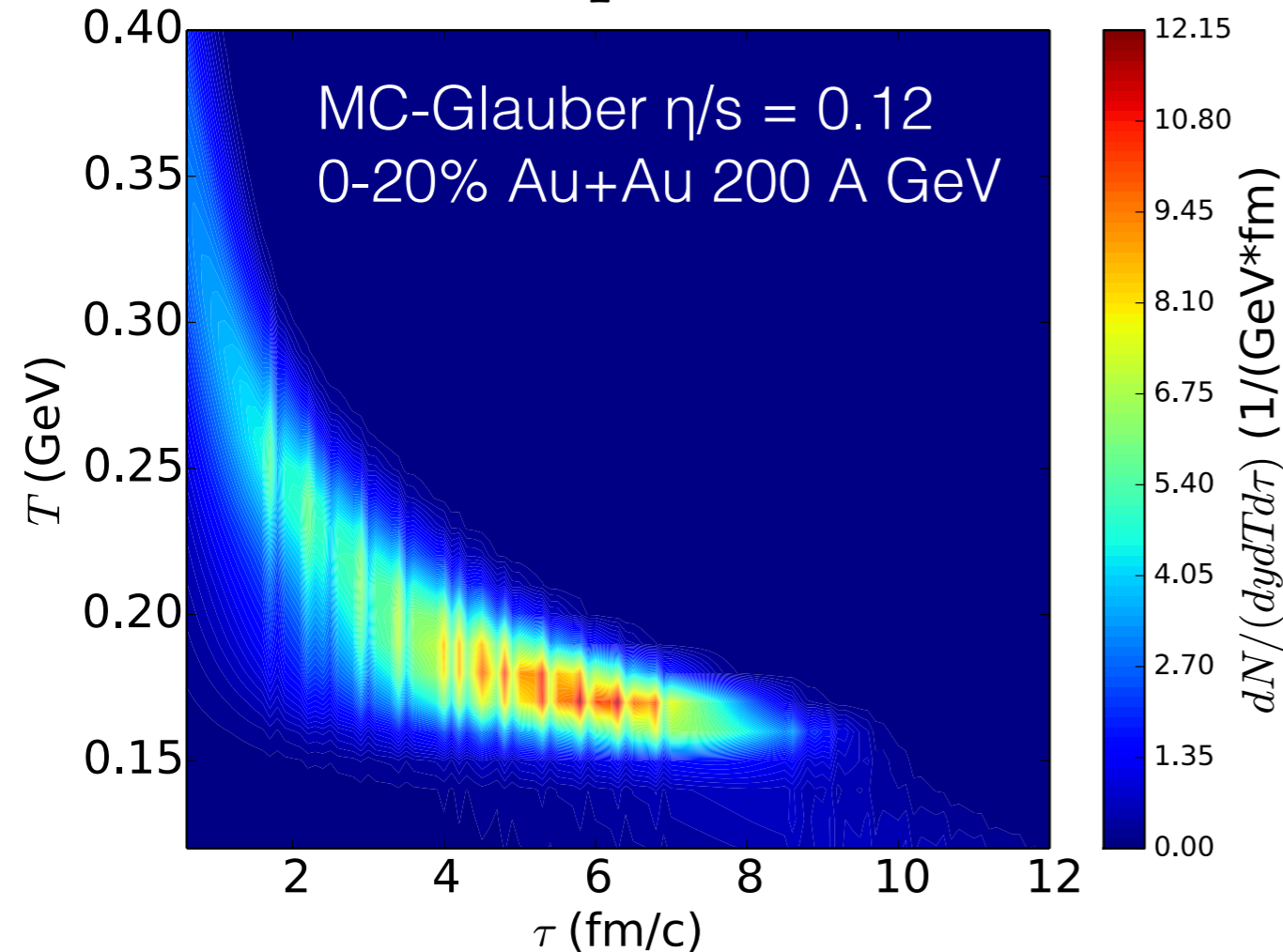
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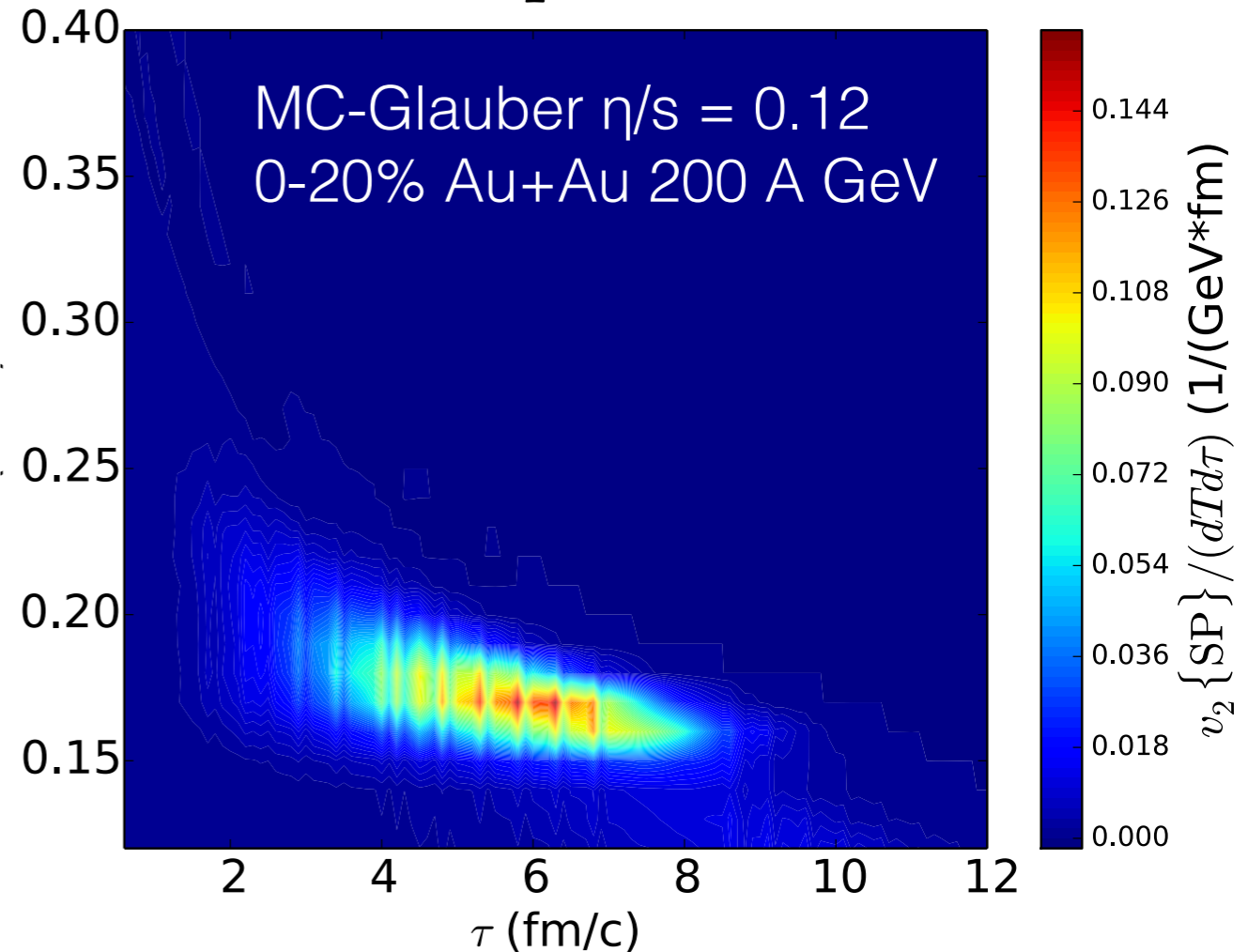
- By cutting hydro medium both in T and τ , we observe a **two-wave** thermal photon production
- Thermal photon v_2 is mostly coming from the transition region, $T = 150 \sim 200 \text{ MeV}$, $\tau = 3 \sim 8 \text{ fm}$ @ RHIC

Thermal photon tomography

$$0.4 \leq p_T \leq 1.0 \text{ GeV}$$



$$0.4 \leq p_T \leq 1.0 \text{ GeV}$$

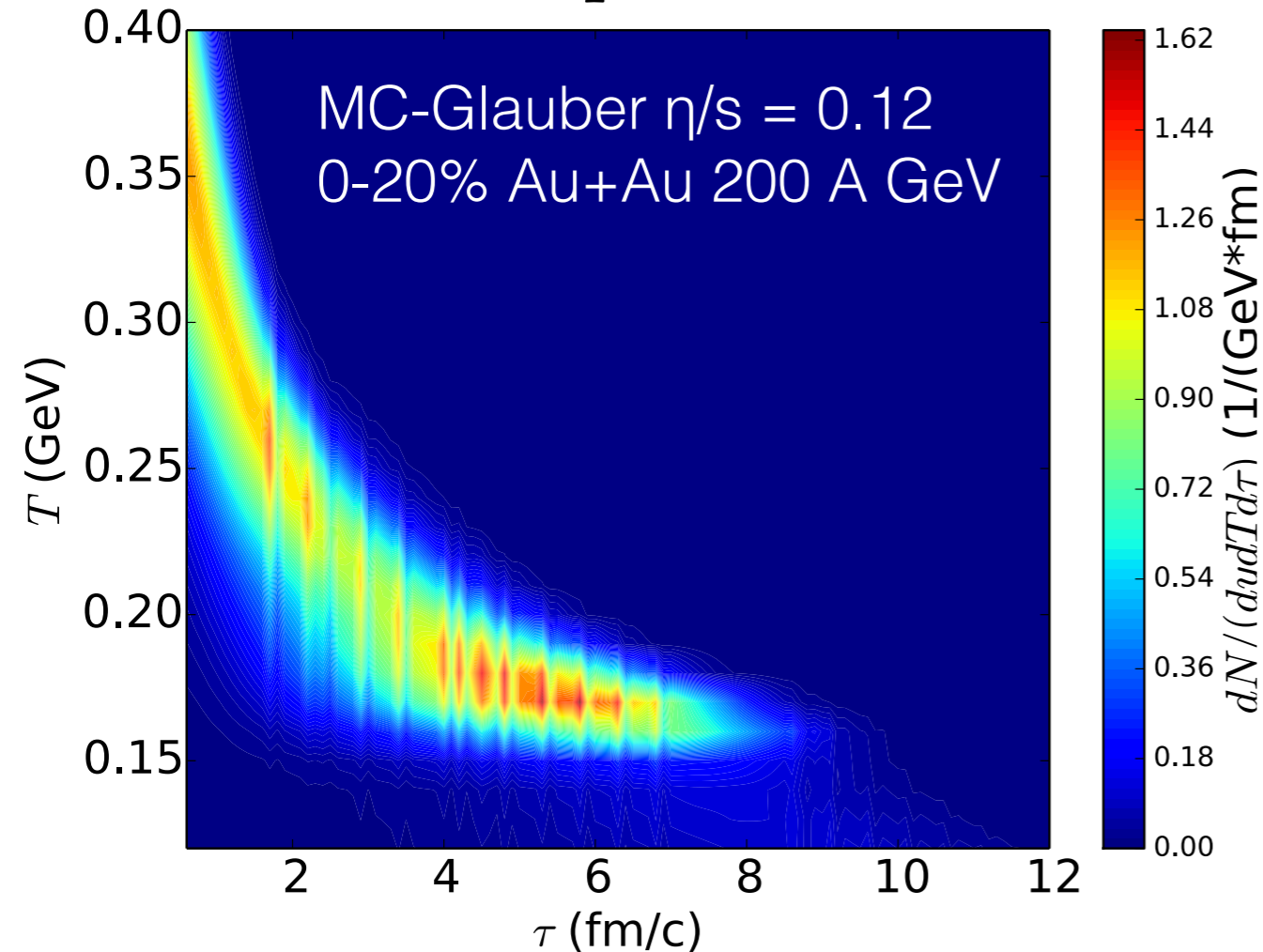


- Thermal photons with $p_T = 0.4 \sim 1.0 \text{ GeV}$ are mostly produced around transition region, their v_2 also reflect the flow anisotropy in this region.

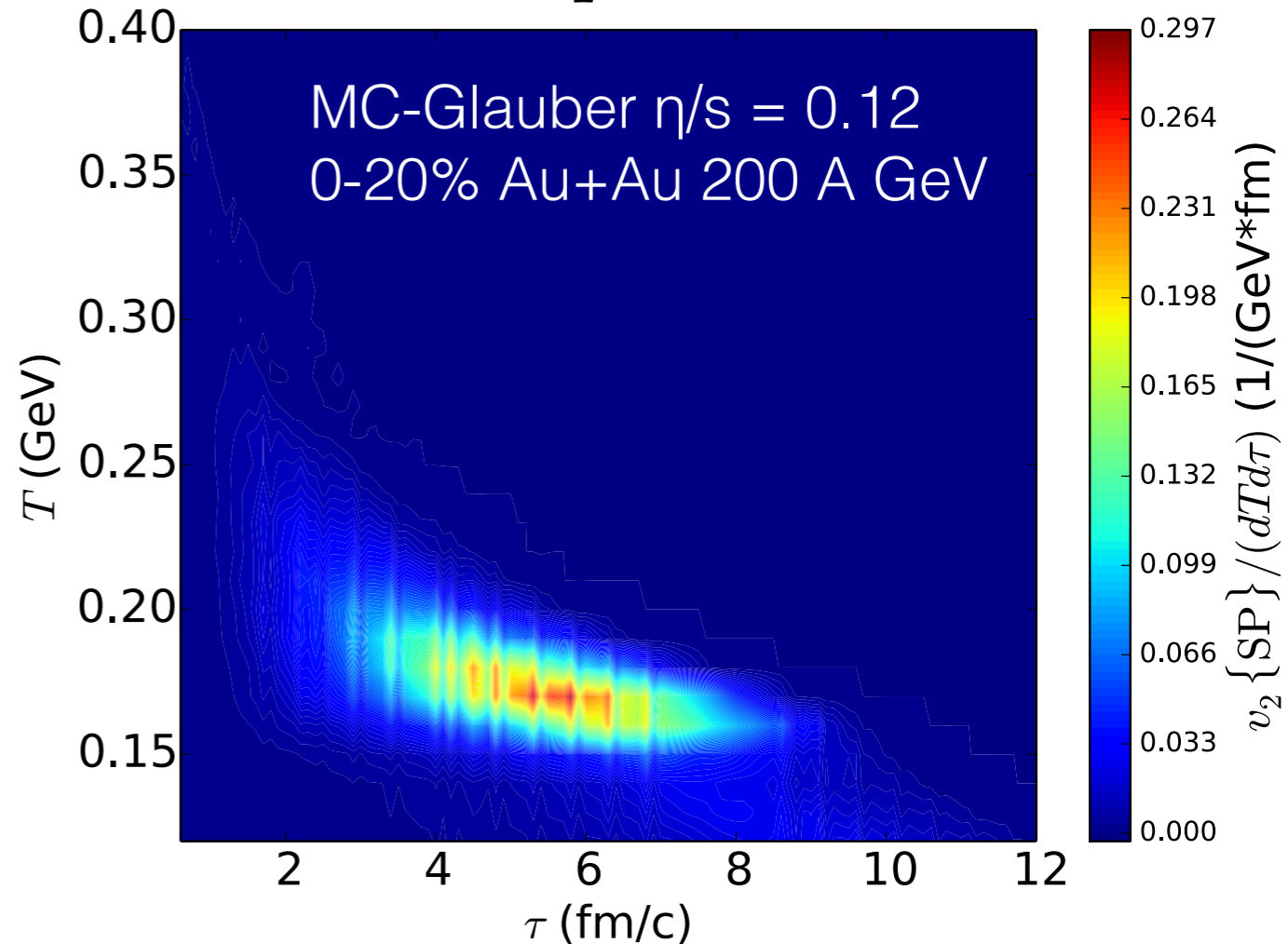
T = 150~200 MeV @ RHIC

Thermal photon tomography

$$1.0 \leq p_T \leq 2.0 \text{ GeV}$$



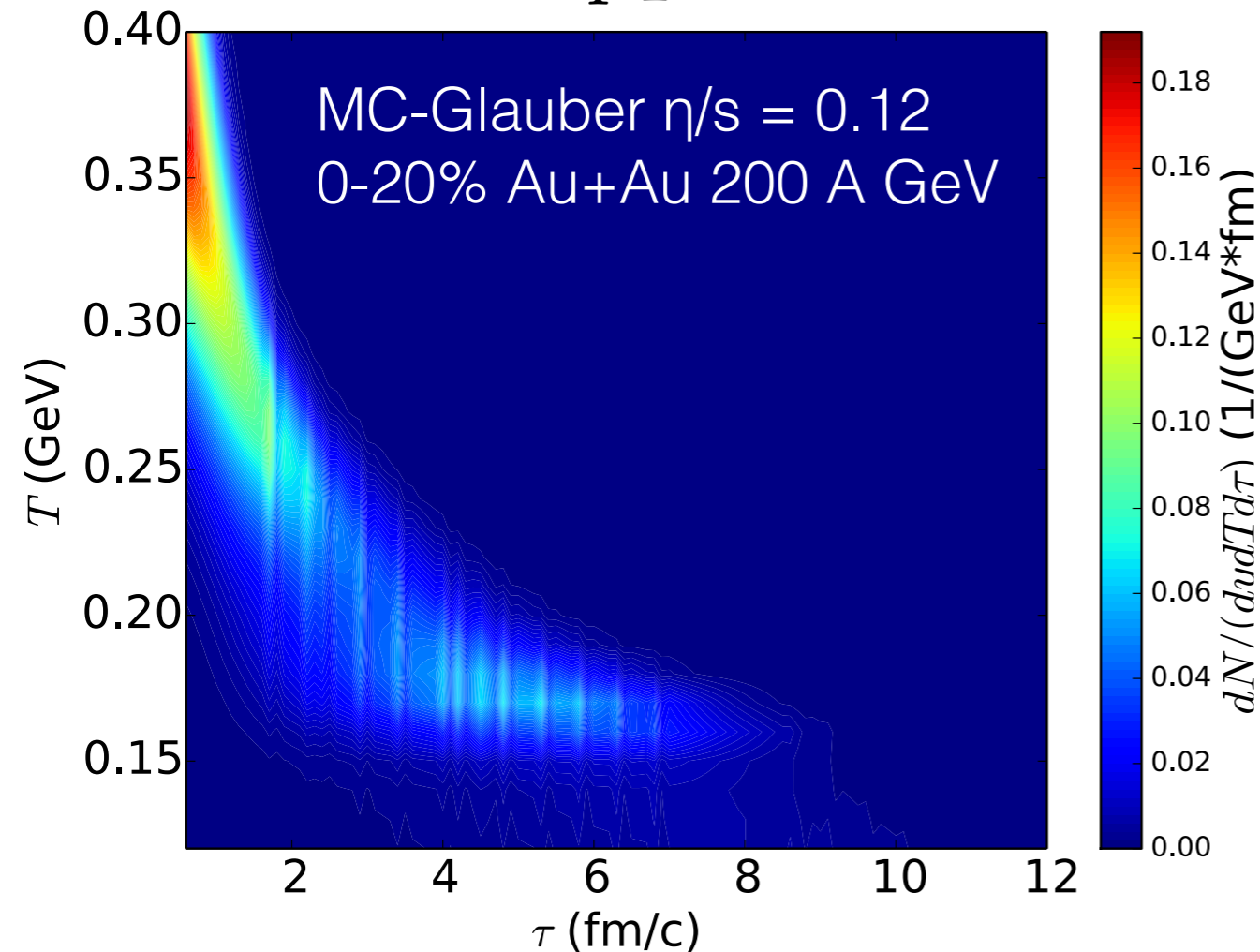
$$1.0 \leq p_T \leq 2.0 \text{ GeV}$$



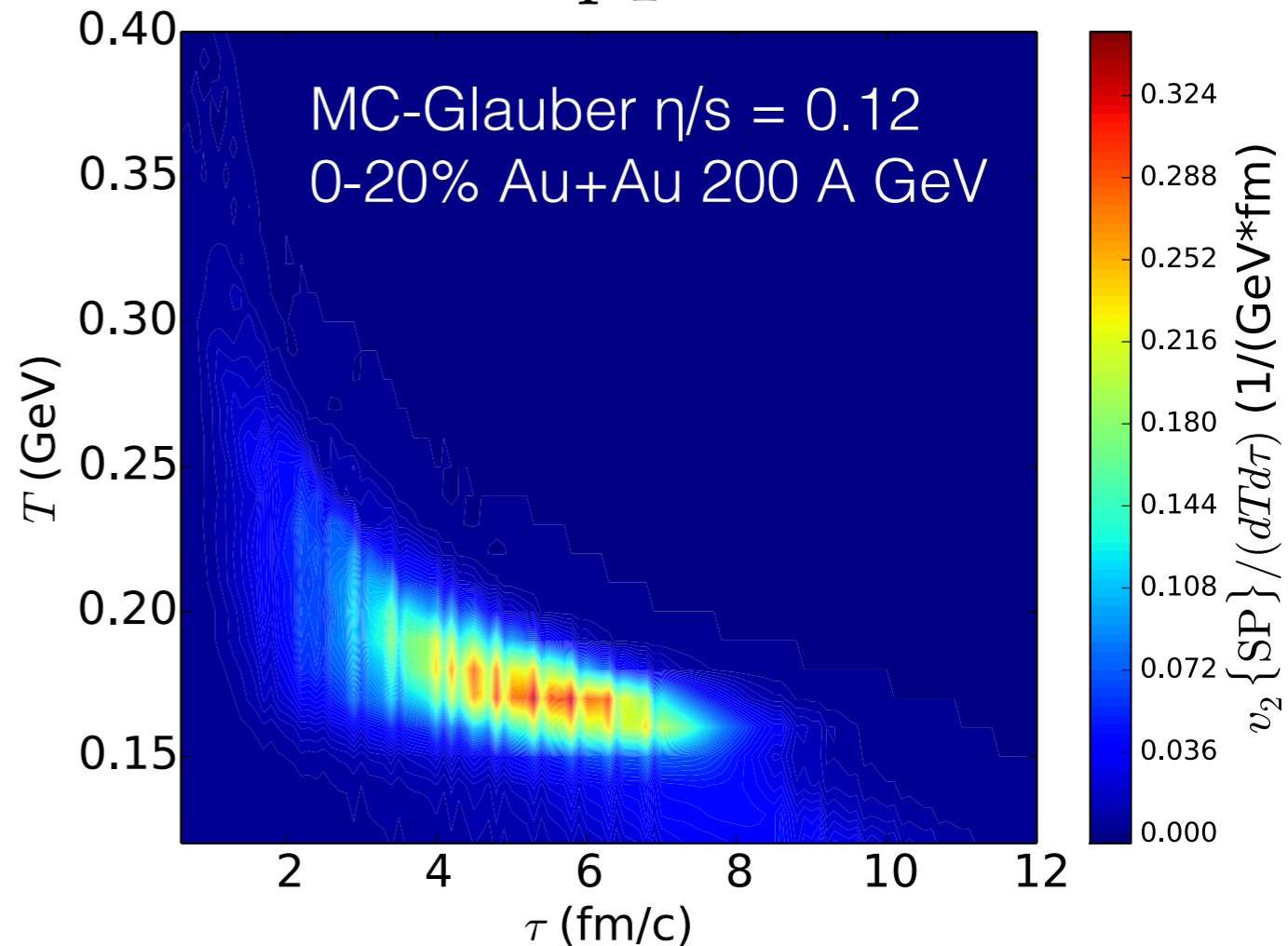
- Thermal photons with $p_T = 1.0 \sim 2.0 \text{ GeV}$ are produced in two waves, their v_2 reflect the flow anisotropy around the transition region. $T = 150 \sim 200 \text{ MeV @ RHIC}$

Thermal photon tomography

$$2.0 \leq p_T \leq 3.0 \text{ GeV}$$



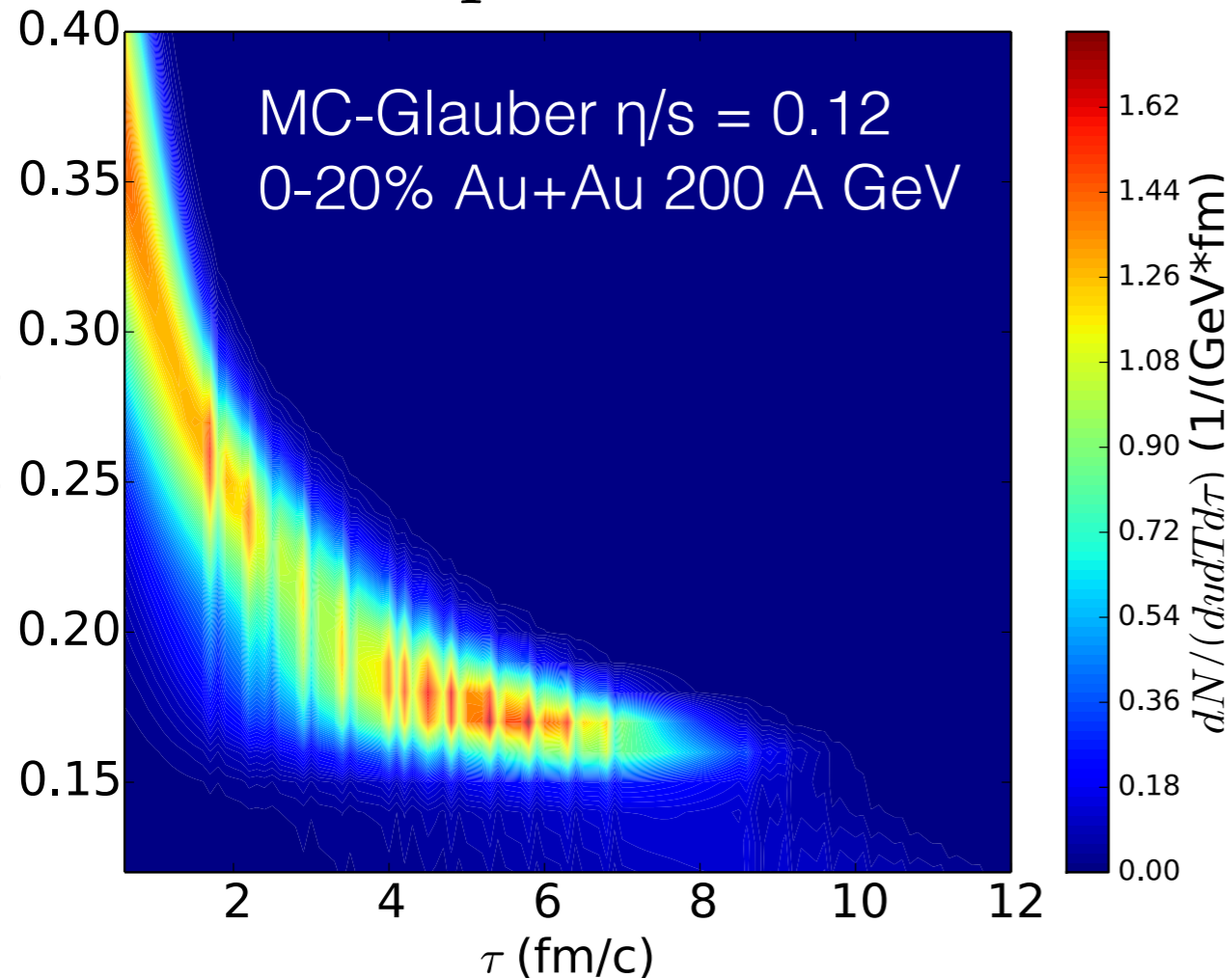
$$2.0 \leq p_T \leq 3.0 \text{ GeV}$$



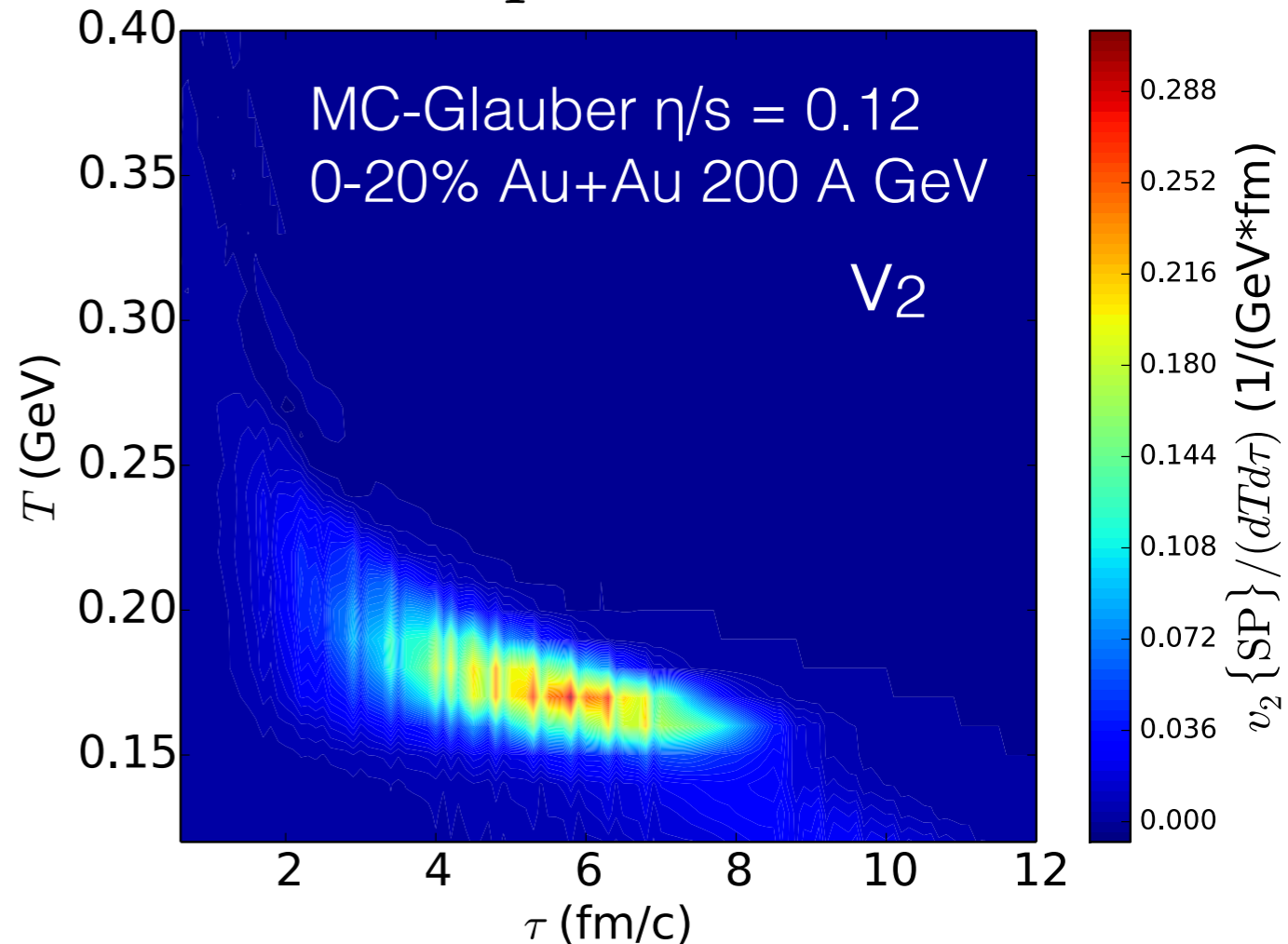
- Thermal photons with $p_T = 2.0 \sim 3.0 \text{ GeV}$ are produced very early, however their v_2 still probes the transition region $T = 150 \sim 200 \text{ MeV}$ @ RHIC

Thermal photon tomography

$$1 \leq p_T \leq 4 \text{ GeV}$$



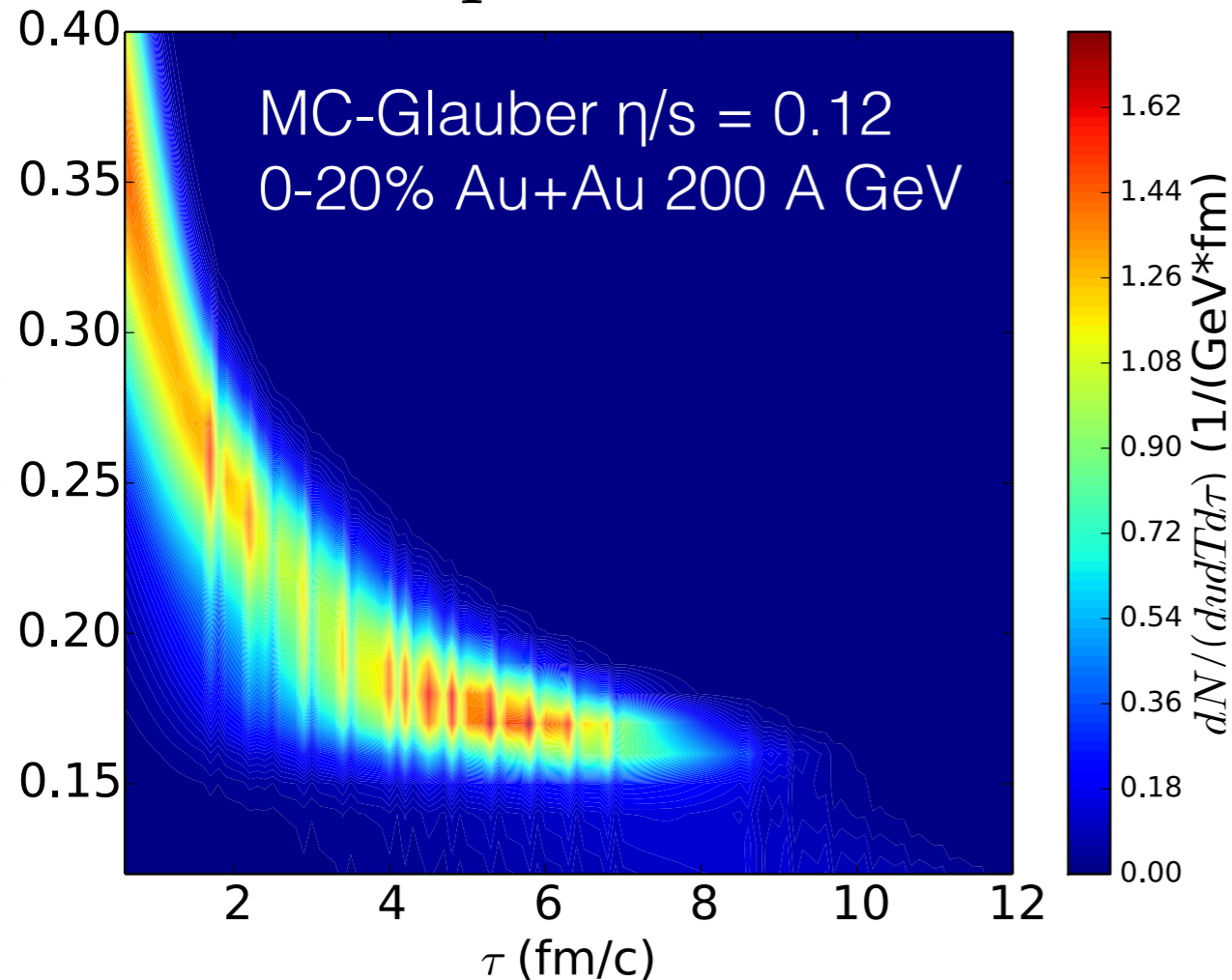
$$1 \leq p_T \leq 4 \text{ GeV}$$



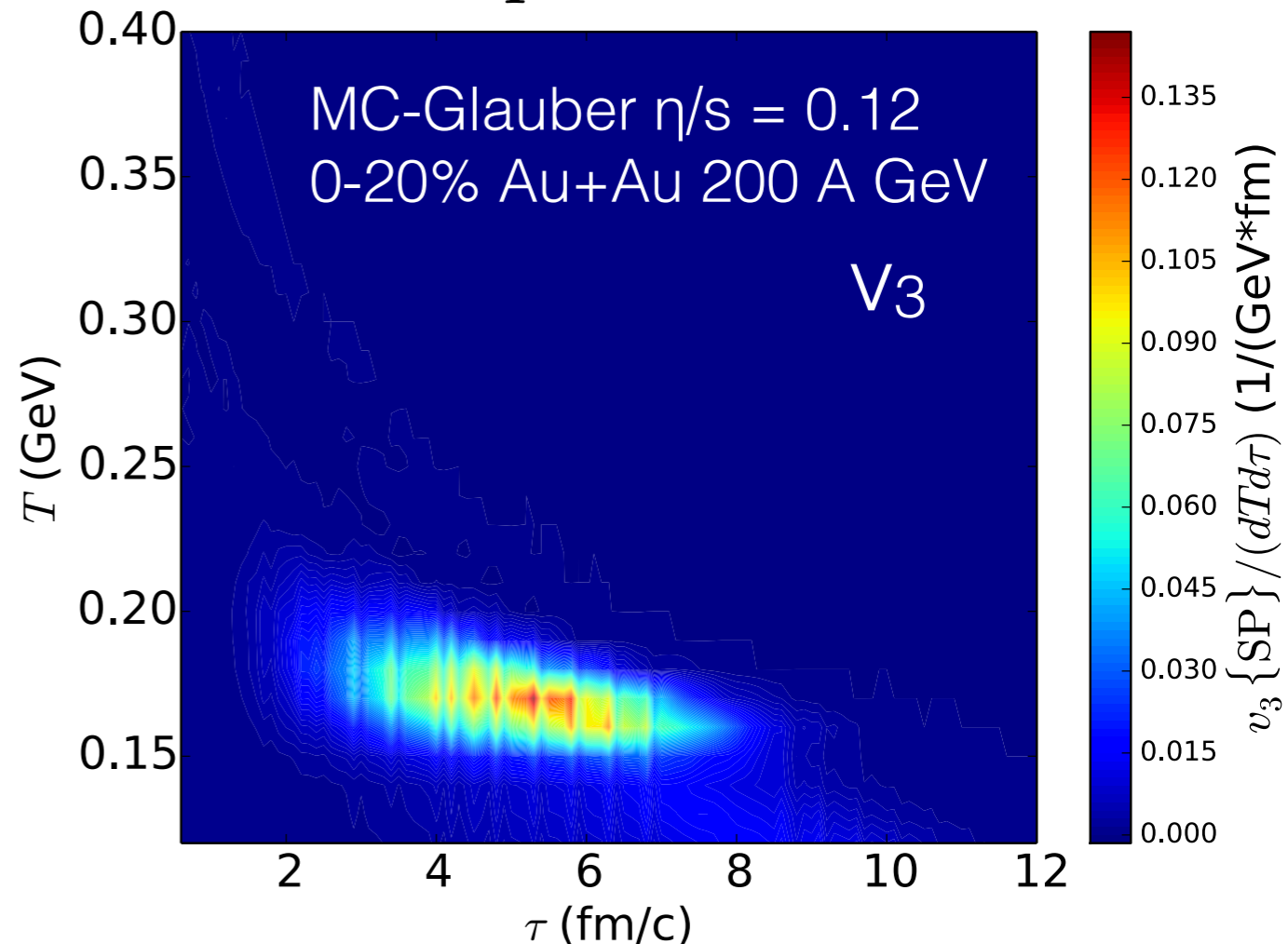
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Thermal photon tomography

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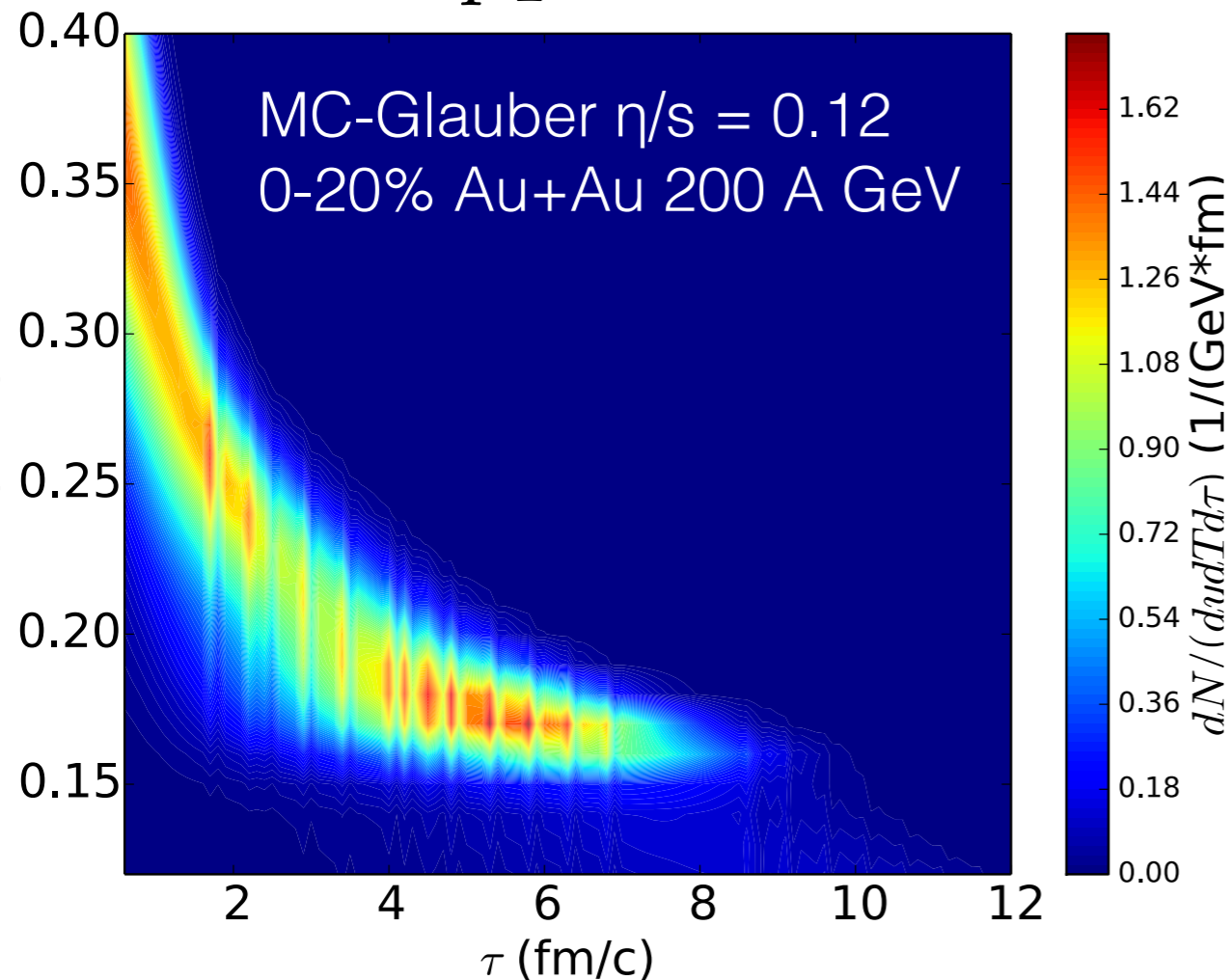
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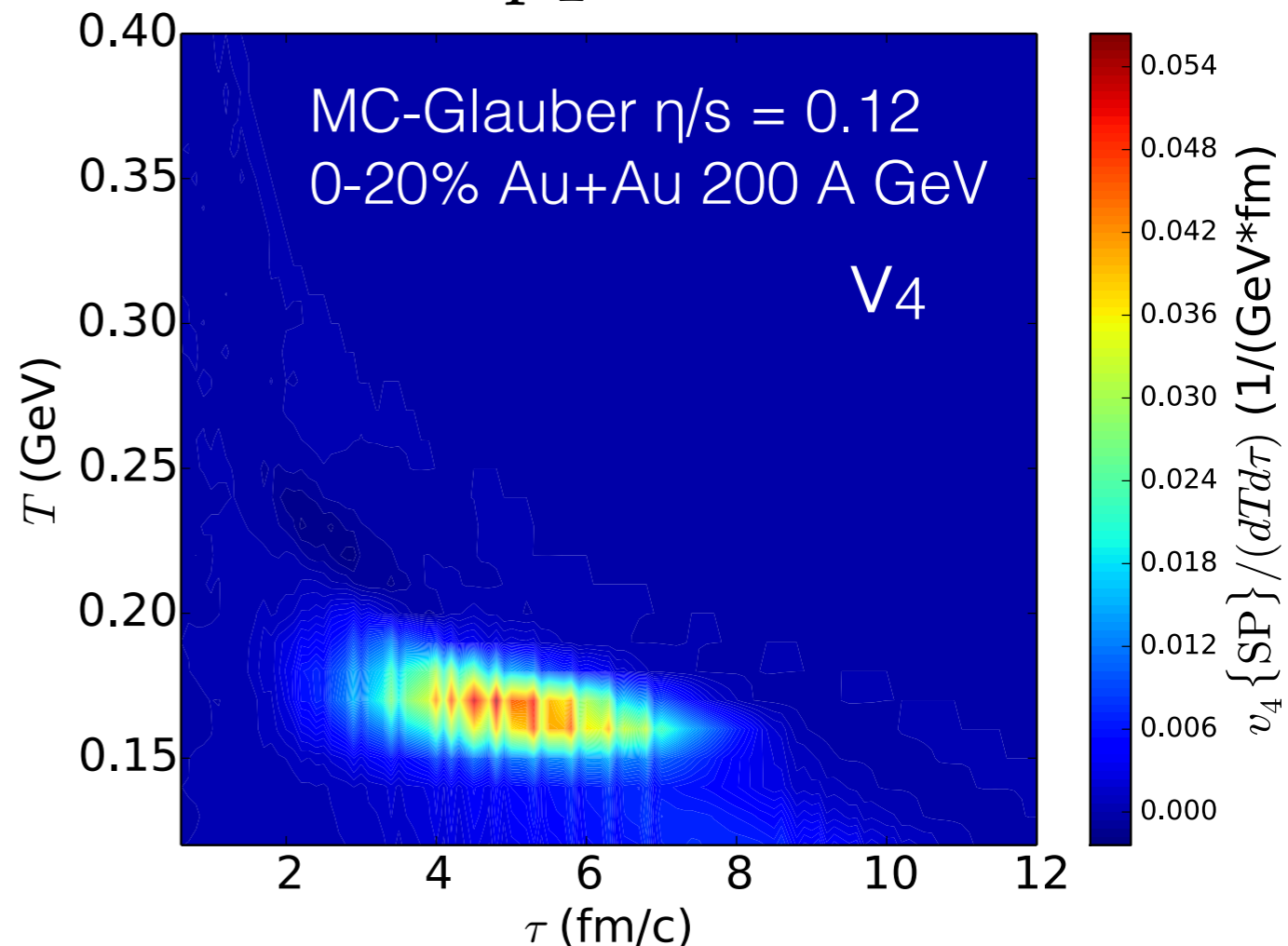
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Thermal photon tomography

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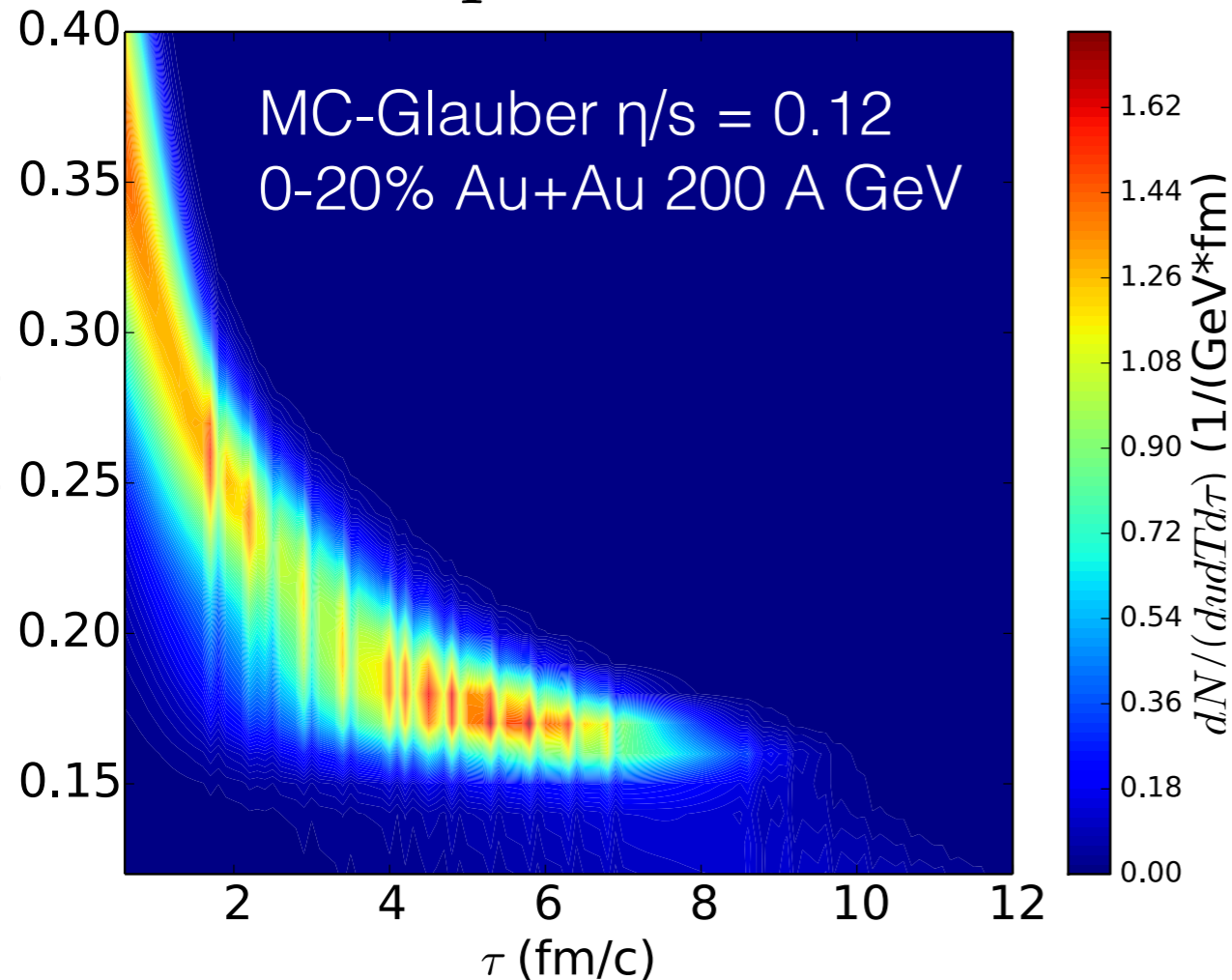
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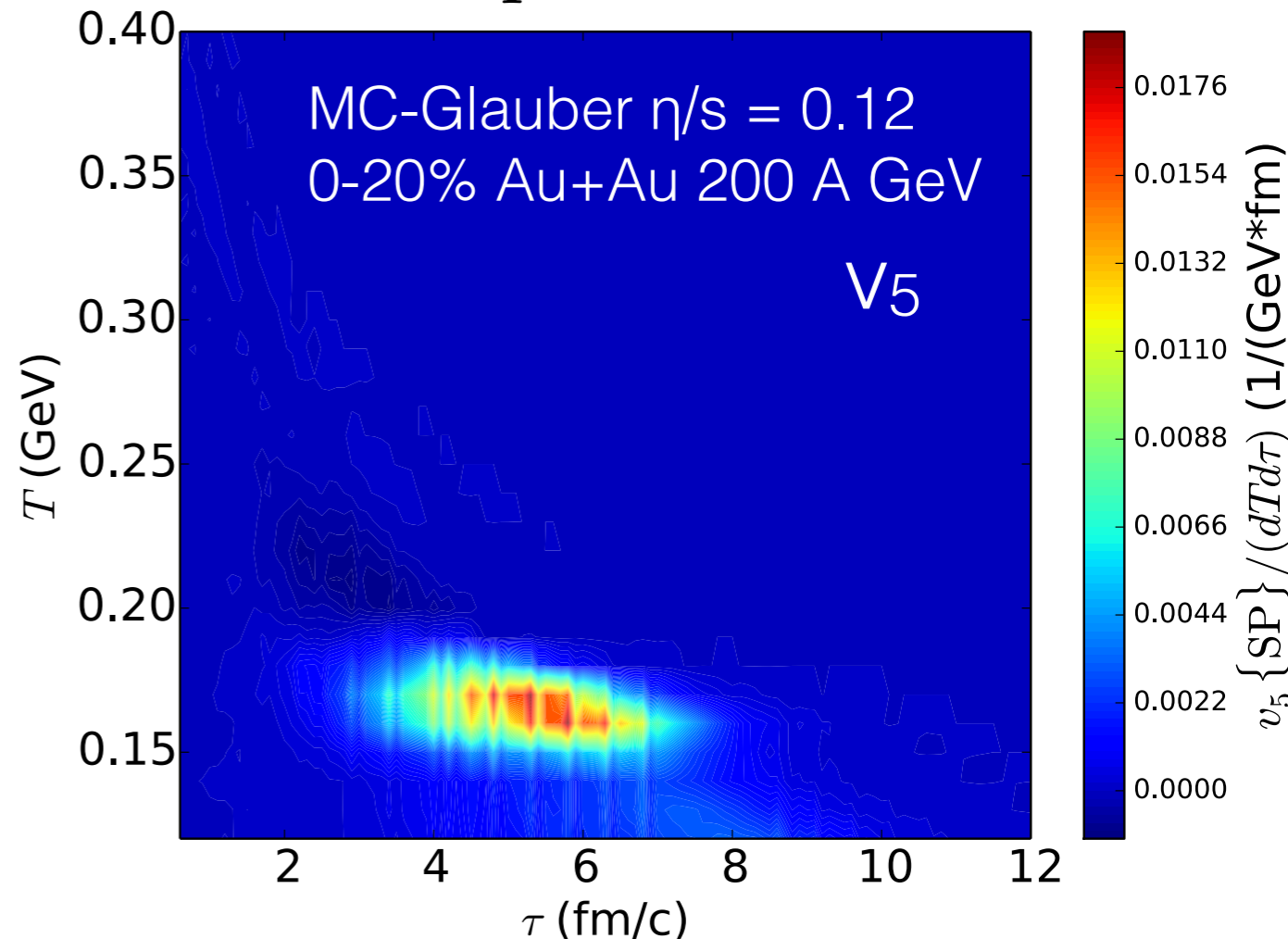
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Thermal photon tomography

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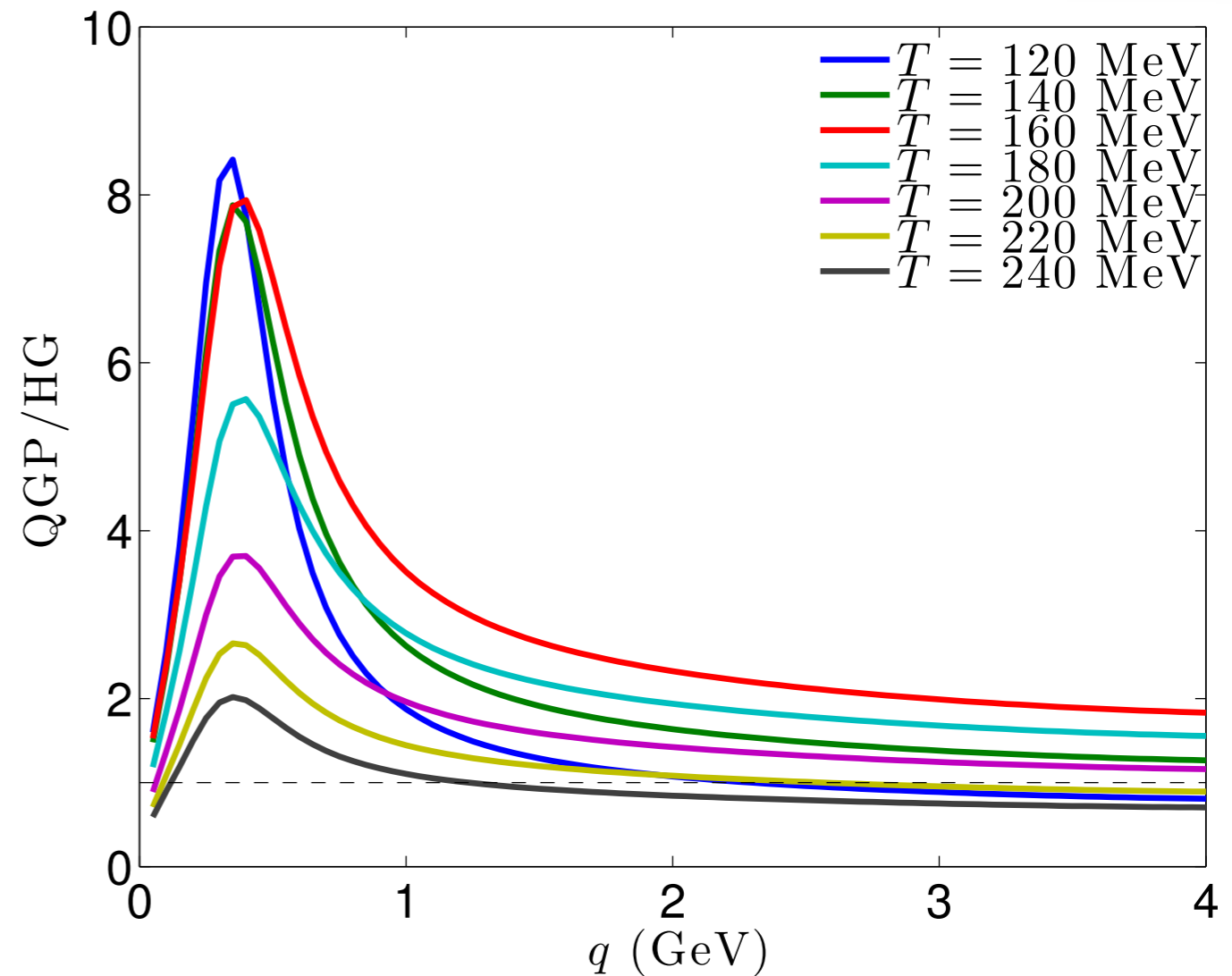
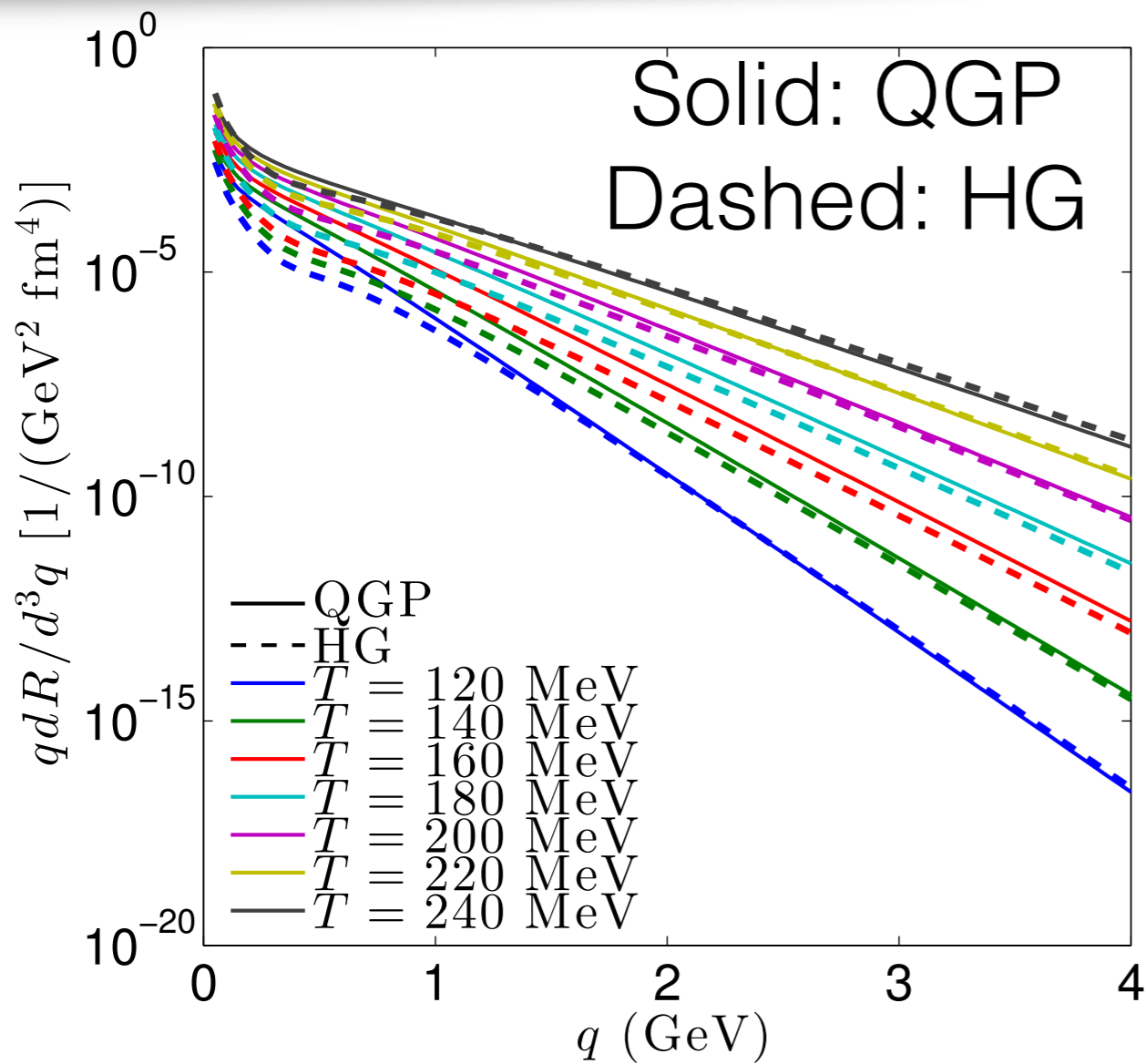


$$1 \leq p_T \leq 4 \text{ GeV}$$



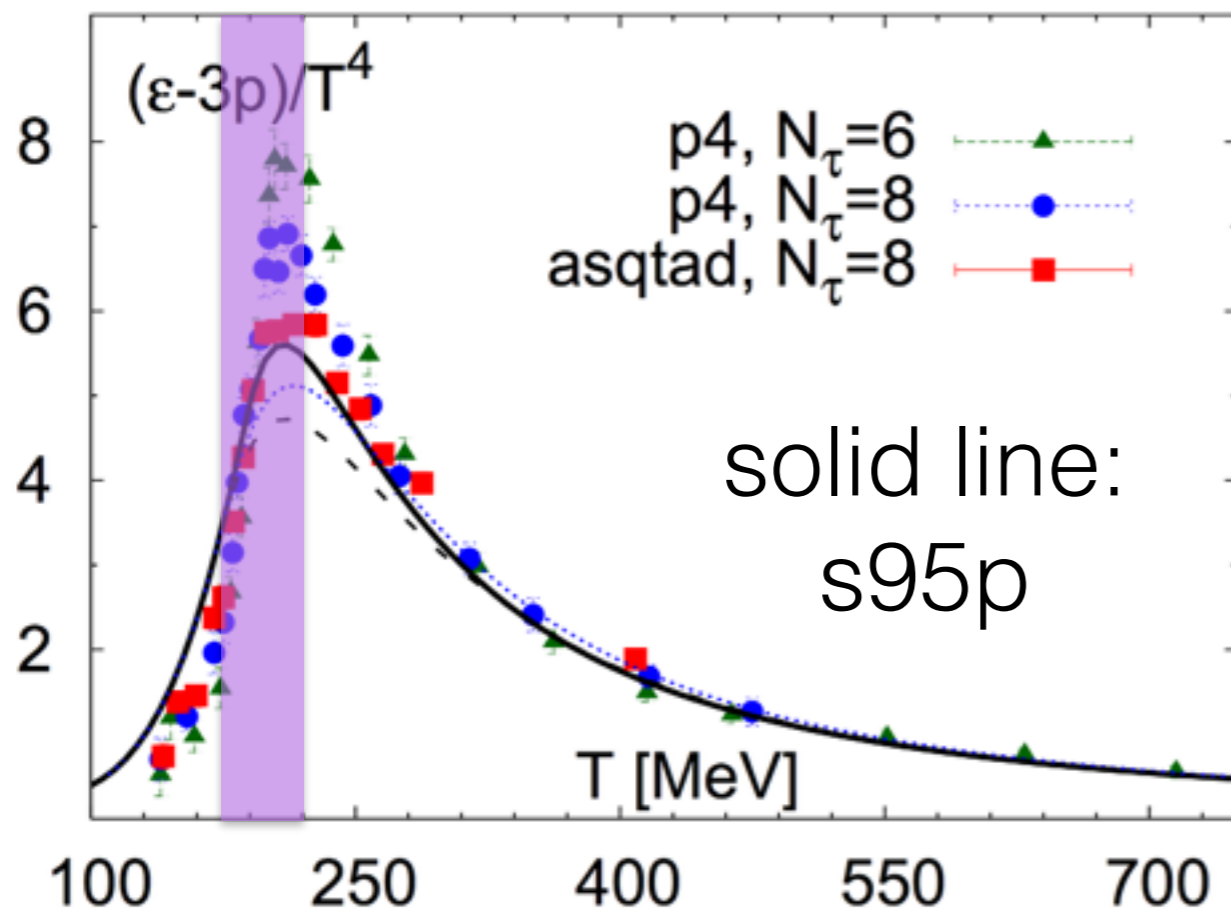
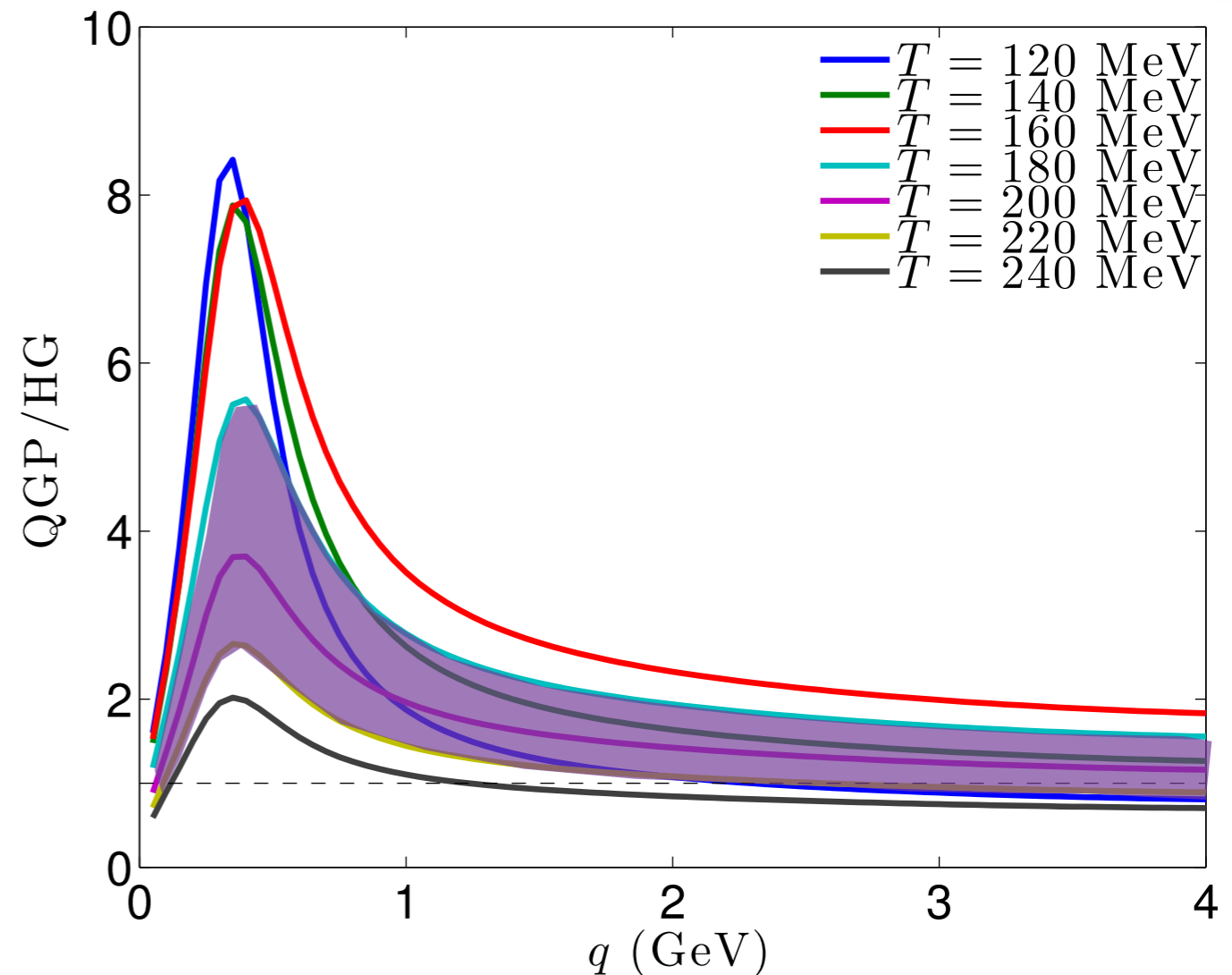
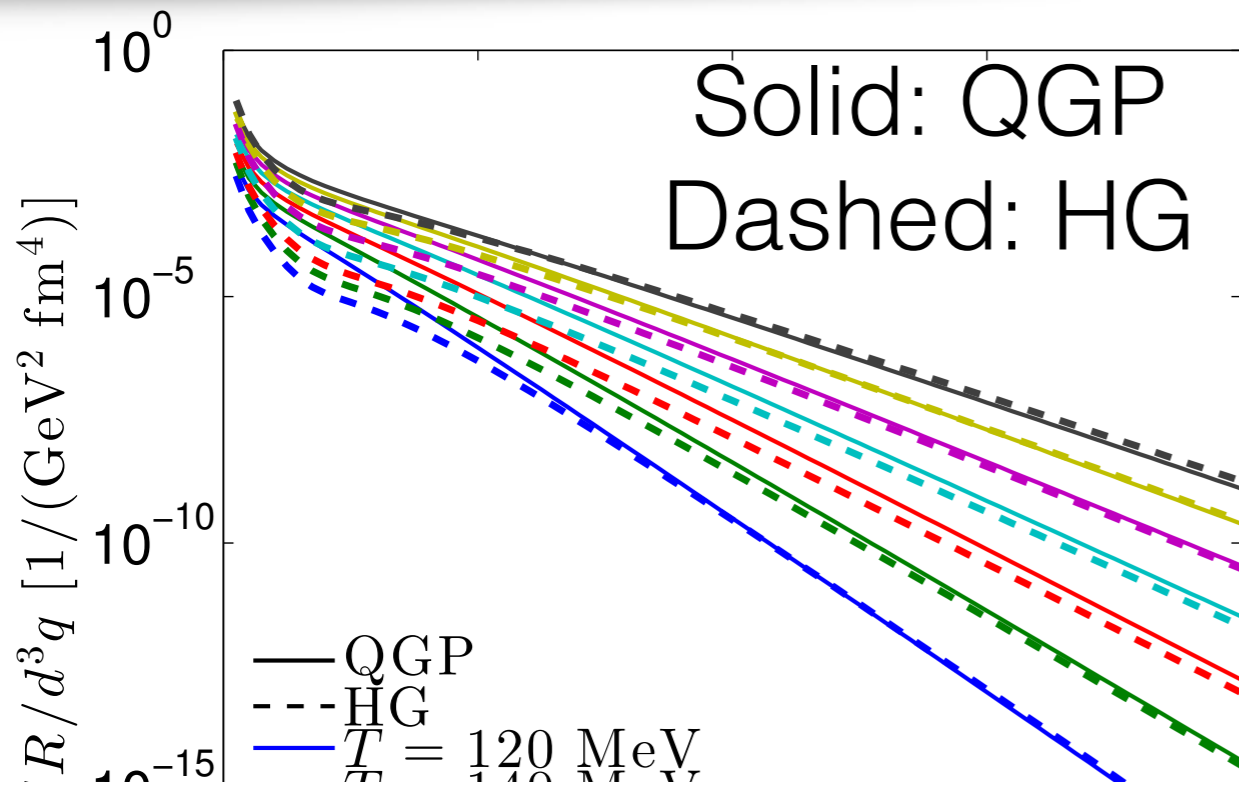
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Photon Emission Rates in the transition region



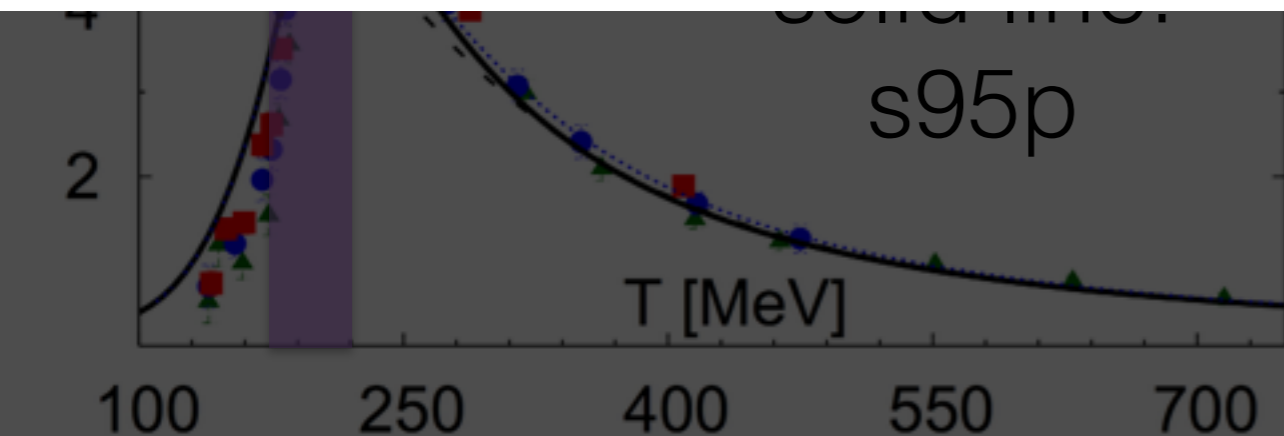
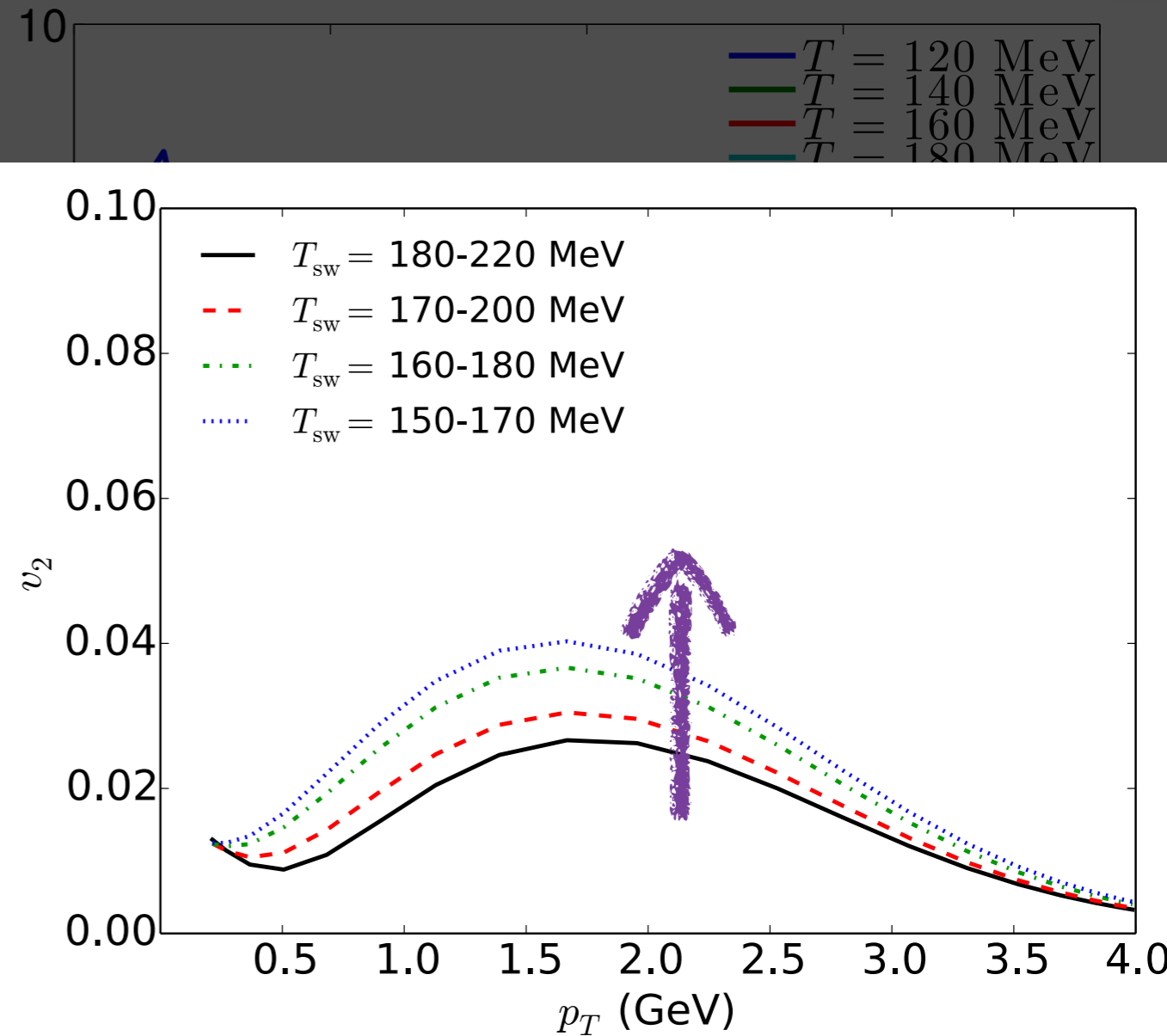
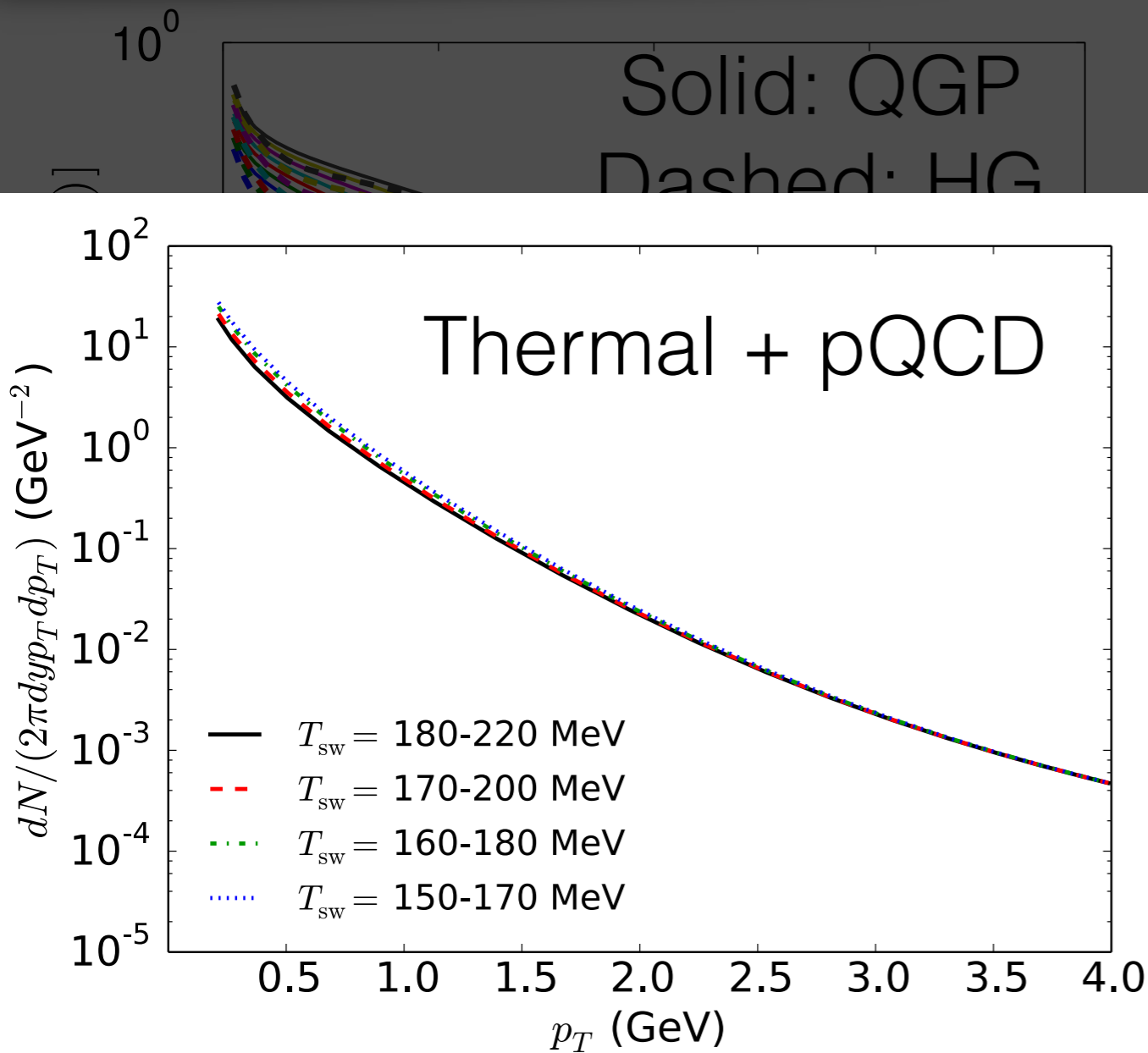
- QGP rates have very different p_T dependence compared to HG rates

Photon Emission Rates in the transition region



- QGP rates have very different p_T dependence compared to HG rates
- Estimated transition region for production rates, $T \sim \mathbf{184 - 220 \text{ MeV}}$

Photon Emission Rates in the transition region



- p_T dependence compared to HG rates
- Estimated transition region for production rates, $T \sim \mathbf{184 - 220}$ MeV

Bulk viscous corrections to photon emission rates

Thermal photon emission rates can be calculated by

$$E_q \frac{dR}{d^3q} = \int \frac{d^3p_1}{2E_1(2\pi)^3} \frac{d^3p_2}{2E_2(2\pi)^3} \frac{d^3p_3}{2E_3(2\pi)^3} \frac{1}{2(2\pi)^3} |\mathcal{M}|^2 \\ \times f_1(p_1^\mu) f_2(p_2^\mu) (1 \pm f_3(p_3^\mu)) (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - q)$$

With

$$f^i(p^\mu) = f_0^i(p \cdot u) + f_0^i(p \cdot u) (1 \pm f_0^i(p \cdot u)) \frac{\pi^{\mu\nu} \hat{p}_\mu \hat{p}_\nu}{2(e + \mathcal{P})} \chi \left(\frac{p \cdot u}{T} \right) \\ + f_0^i(p \cdot u) (1 \pm f_0^i(p \cdot u)) \Pi(B^i(T) + D^i(T)(p \cdot u) + E^i(T)(p \cdot u)^2)$$

We can expand photon emission rates around the thermal equilibrium:

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e + \mathcal{P})} a_{\alpha\beta} \Gamma^{\alpha\beta}(q, T) + \frac{\Pi}{\mathcal{P}} \Gamma_\Pi(q, T)$$

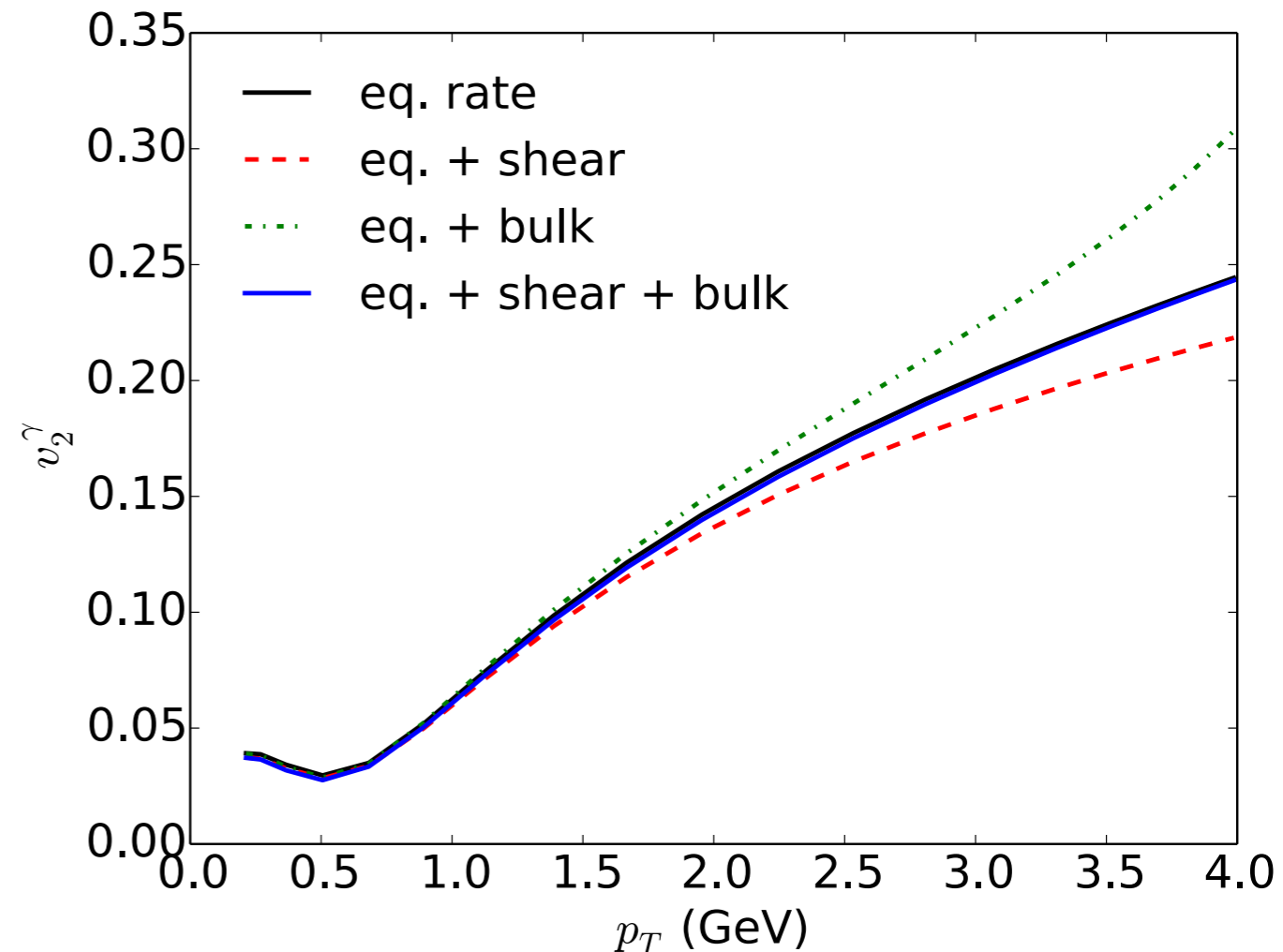
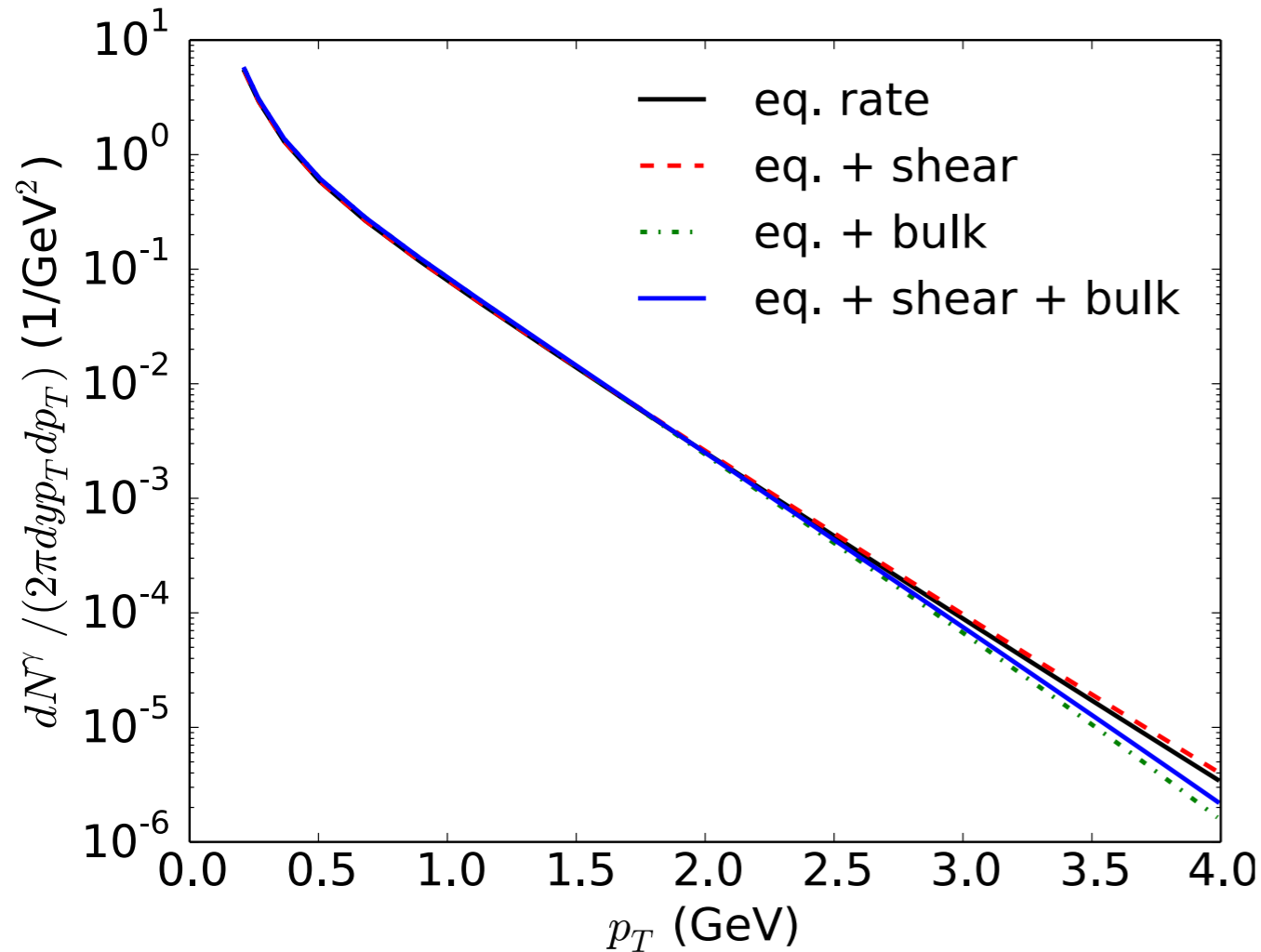
$$a_{\mu\nu} = \frac{3}{2(u \cdot \hat{q})^4} \hat{q}_\mu \hat{q}_\nu + \frac{1}{(u \cdot \hat{q})^2} u_\mu u_\nu + \frac{1}{2(u \cdot \hat{q})^2} g_{\mu\nu} - \frac{3}{2(u \cdot \hat{q})^3} (\hat{q}_\mu u_\nu + \hat{q}_\nu u_\mu).$$

Peek of bulk viscous effects on thermal photon observables

$$\frac{\zeta}{s} = \frac{1}{2} \frac{\eta}{s} \left(\frac{1}{3} - c_s^2 \right) \quad \eta/s = 0.08$$

J. Noronha-Hostler, G. S. Denicol, J. Noronha, R. P. G. Andrade and F. Grassi, Phys. Rev. C **88**, 044916 (2013)

Hadronic photons:



- Bulk viscosity **steepens** thermal photon spectrum
- It **increases** thermal photon pT differential elliptic flow
reduces hydrodynamic radial flow

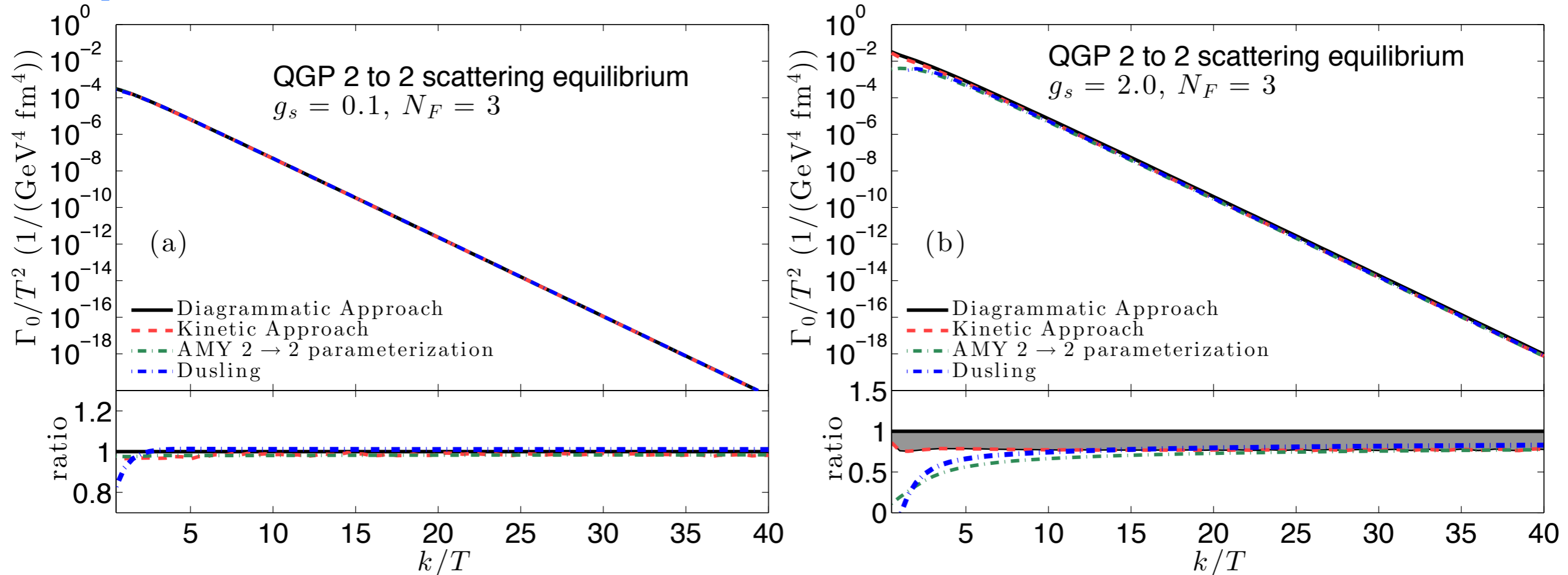
Conclusion

- We study photon spectra and their anisotropic flows \mathbf{v}_n from *event-by-event* viscous hydrodynamic medium
- Thermal photon spectra are strongly **blue shifted** by hydrodynamic radial flow
- Shear viscosity **suppresses** photon v_n . Dominant suppression comes not from flow, but from the viscous correction to the production rates.
- Uncertainty of the photon emission rates in the **transition region** plays a crucial role in the theoretical calculations
- The interplay between **bulk** and **shear** viscous effects need to be carefully studied

Back up

Photon Rates (QGP 2 to 2 processes only)

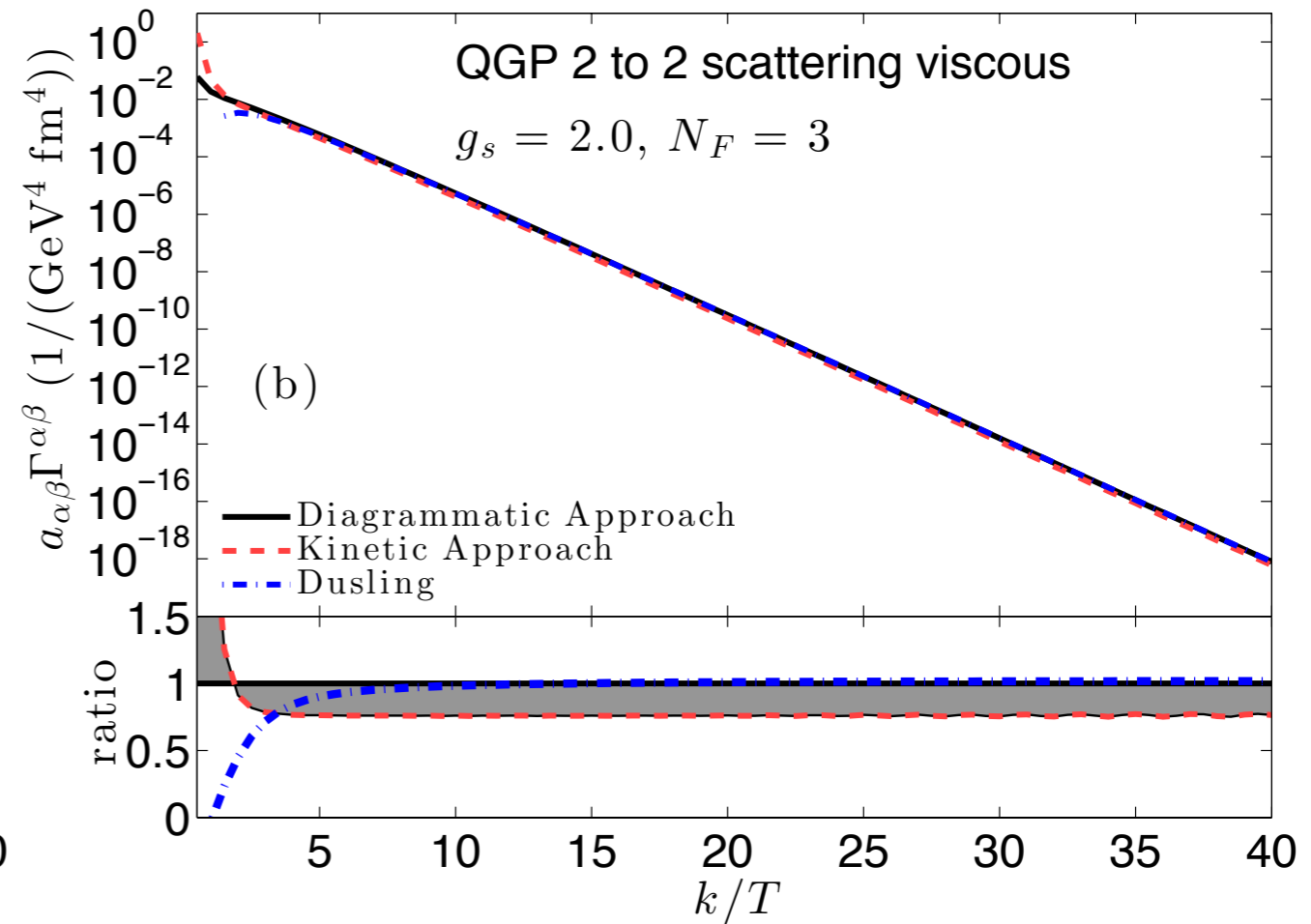
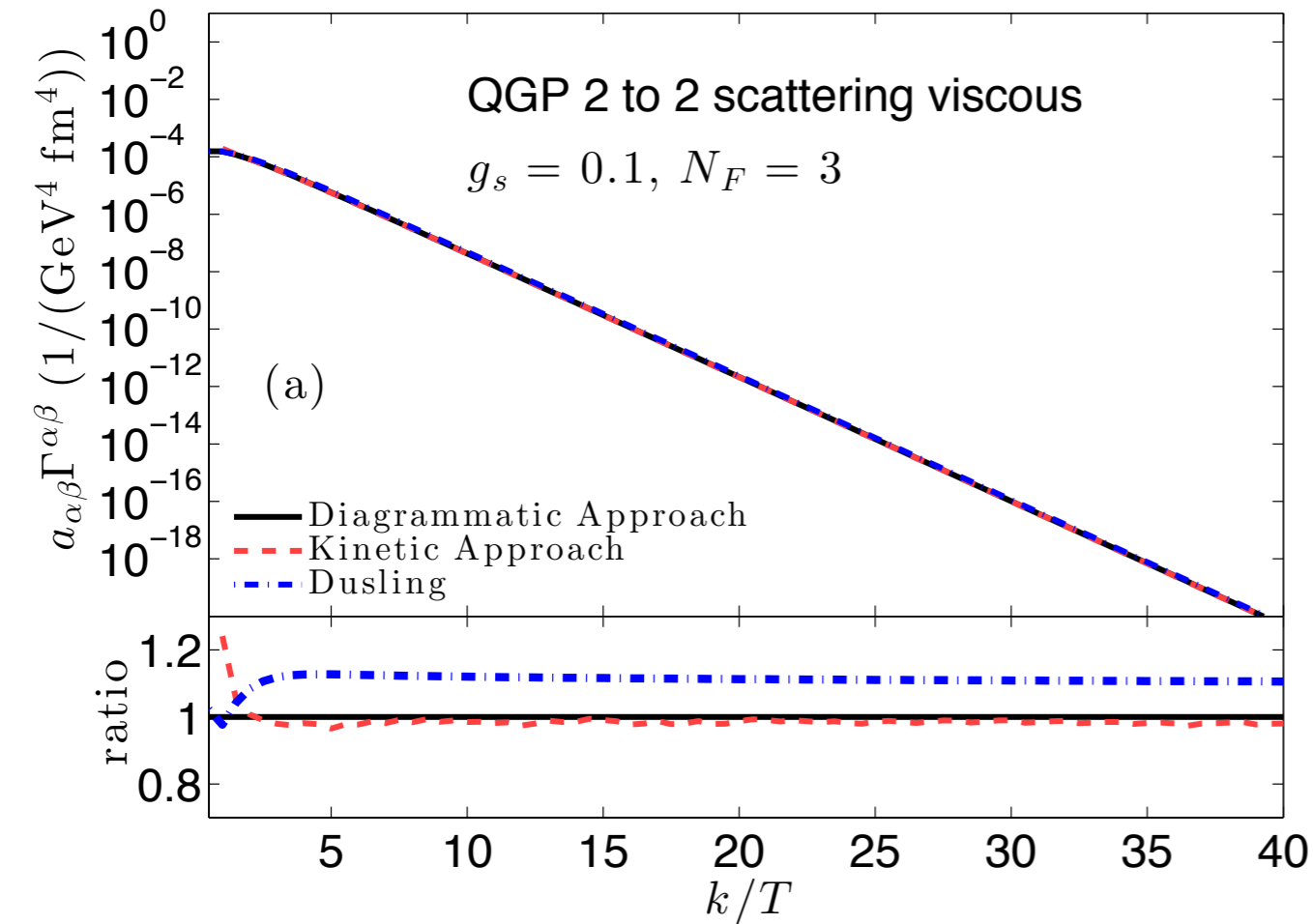
Equilibrium rates:



- For small g , results from diagrammatic approach agree well with kinetic approach and AMY
- For $g = 2.0$, diagrammatic approach gives 25% larger results compared to kinetic approach; difference are due to cut-off dependence.

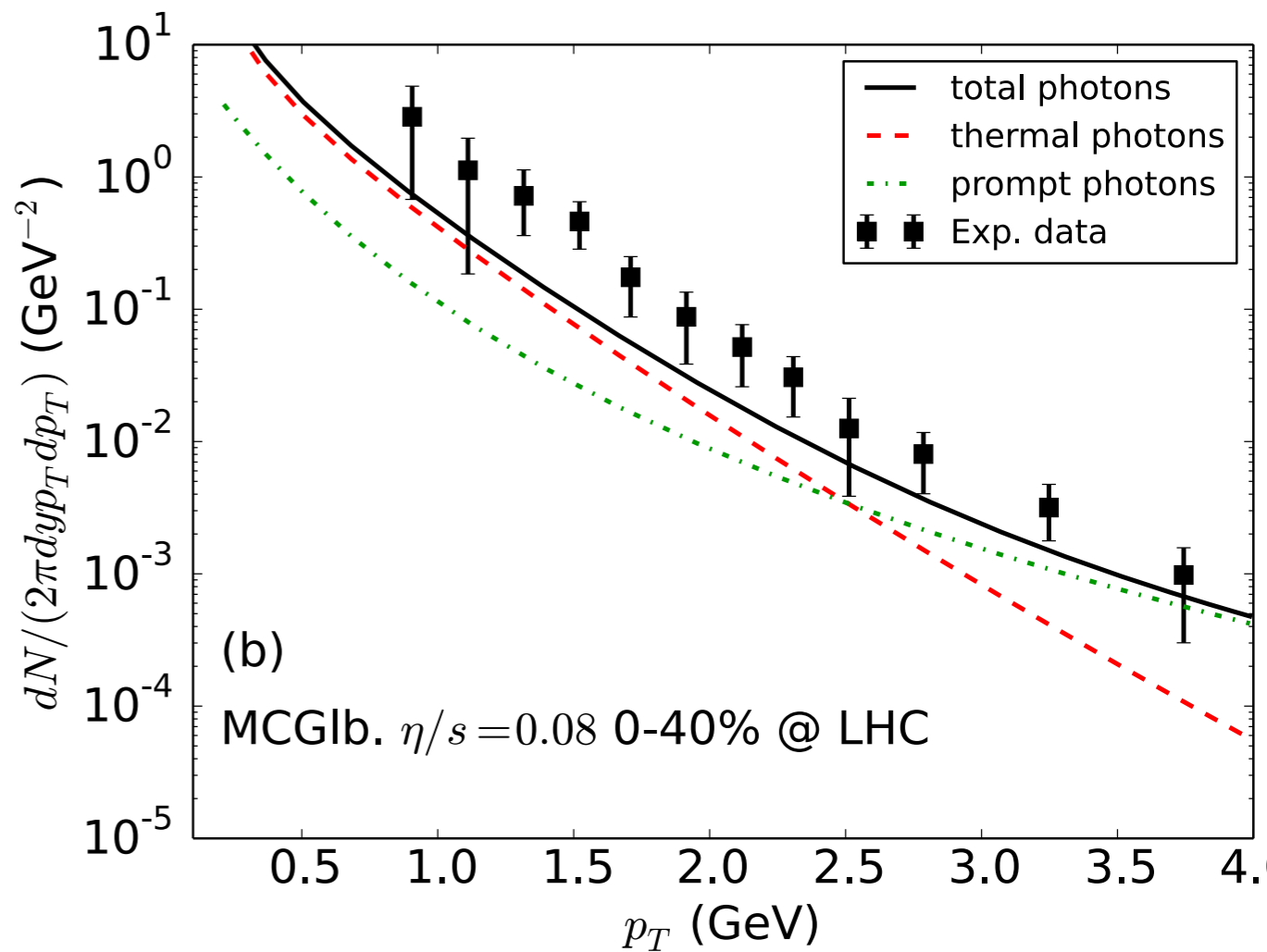
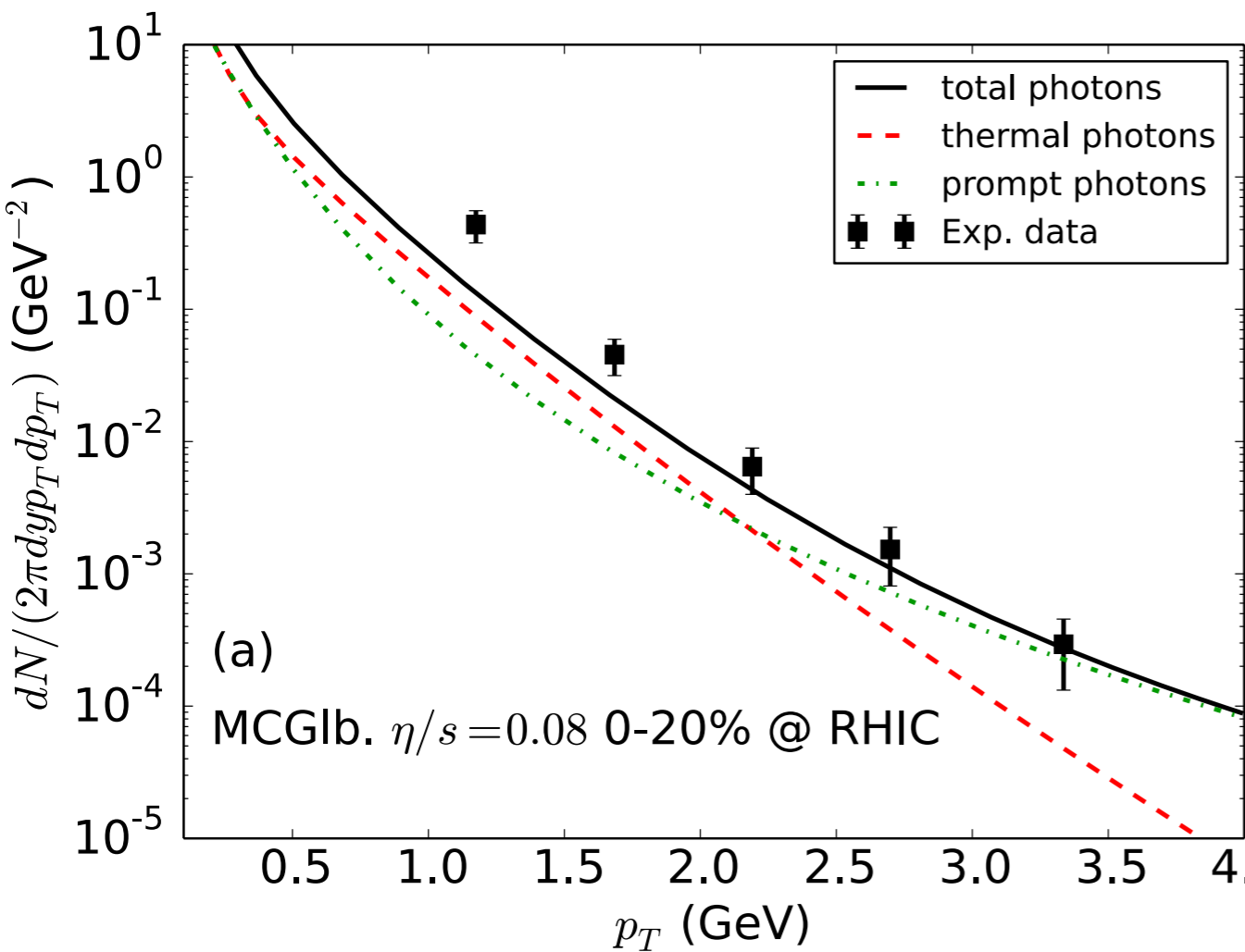
Photon Rates (QGP 2 to 2 processes only)

Viscous corrections:



- For small g , diagrammatic approach agrees with kinetic approach
- For $g = 2$, the deviations at small k/T may originate from different higher order $O(g^2 T)$ contributions

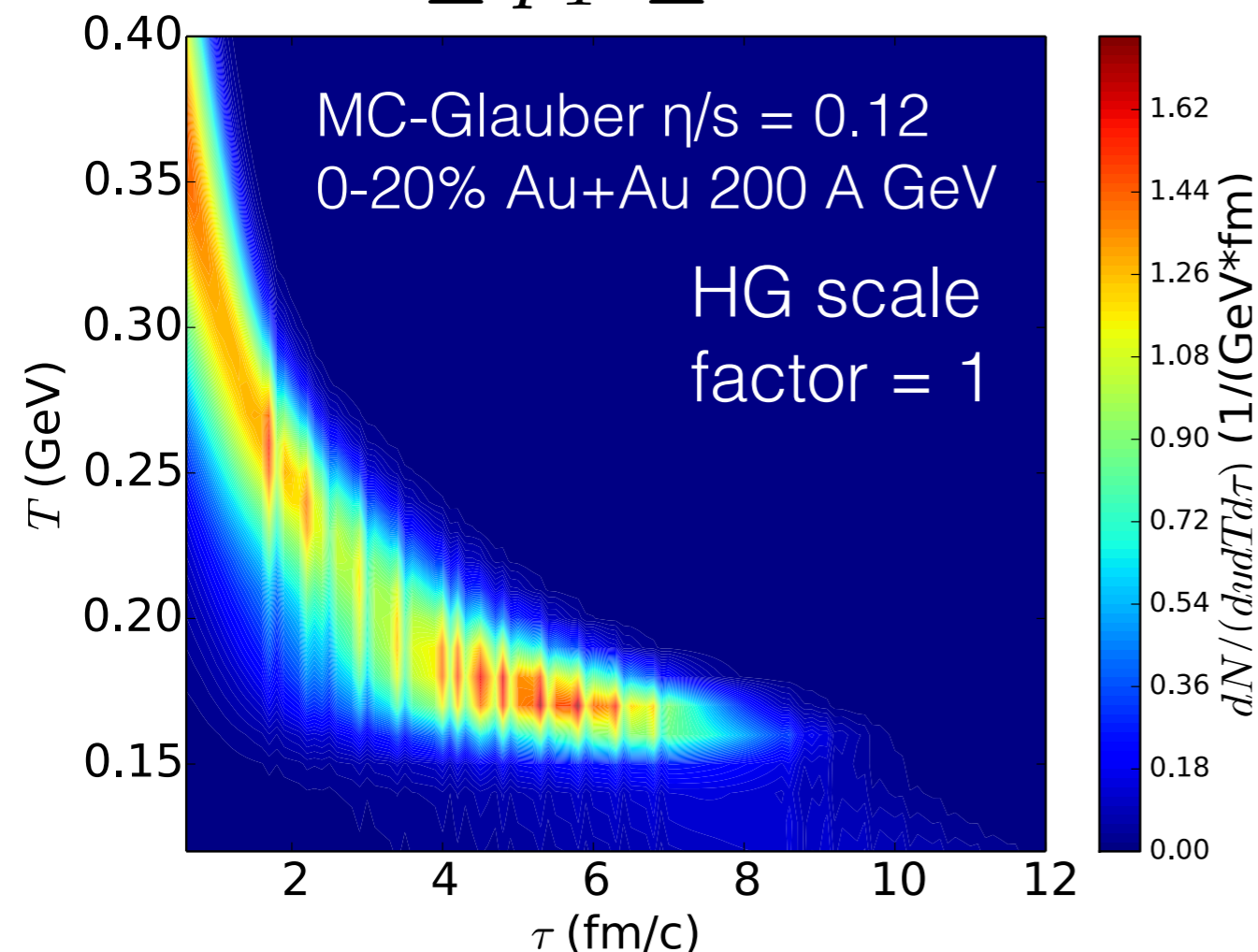
Thermal Photon Spectra



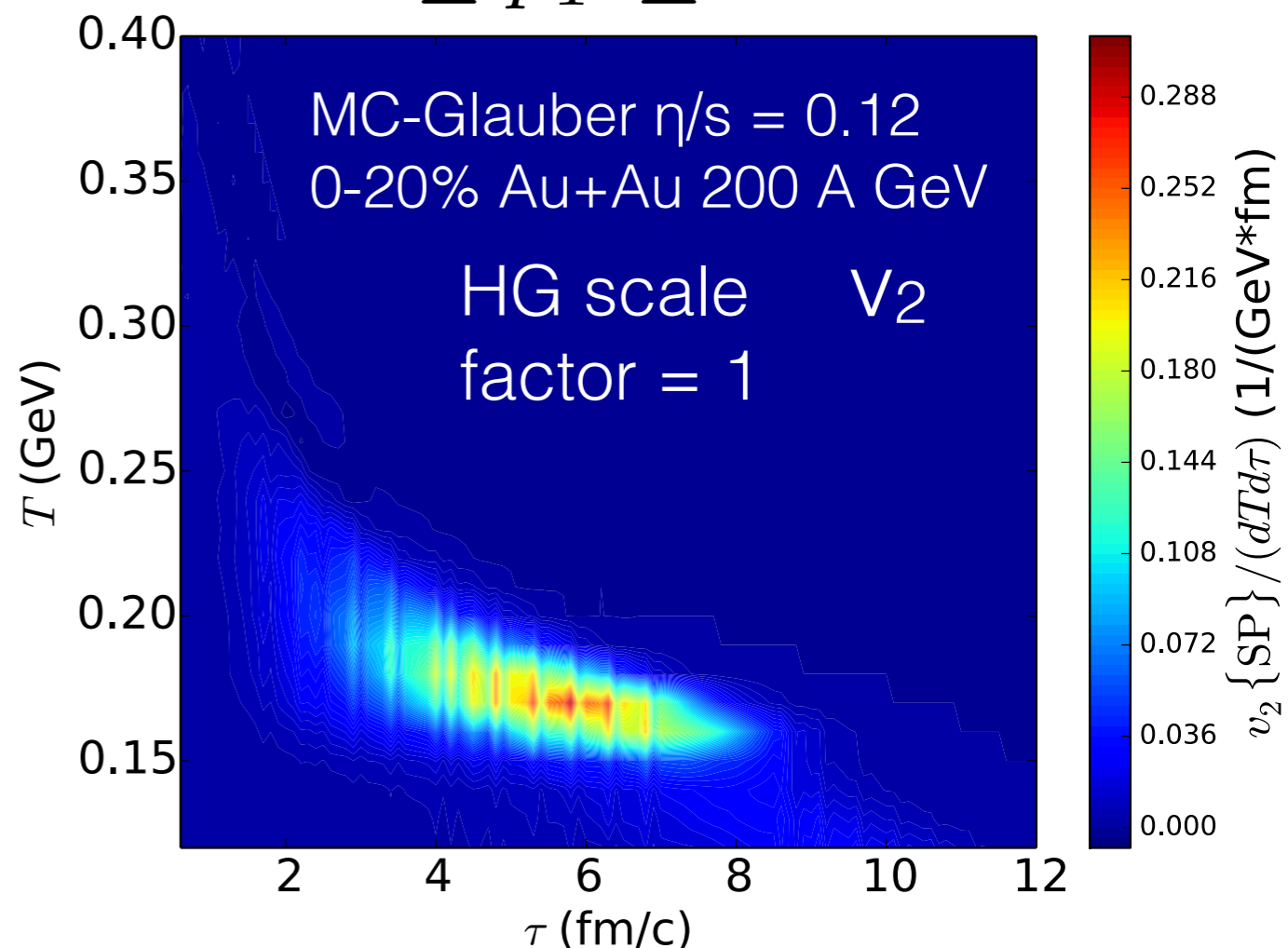
- With all available thermal emission sources, our current calculations still underestimate measured direct photon spectra at low p_T at both RHIC and LHC energies
- Additional emission sources need to be included to improve the agreement between theory and data

Thermal photon tomography

$$1 \leq p_T \leq 4 \text{ GeV}$$



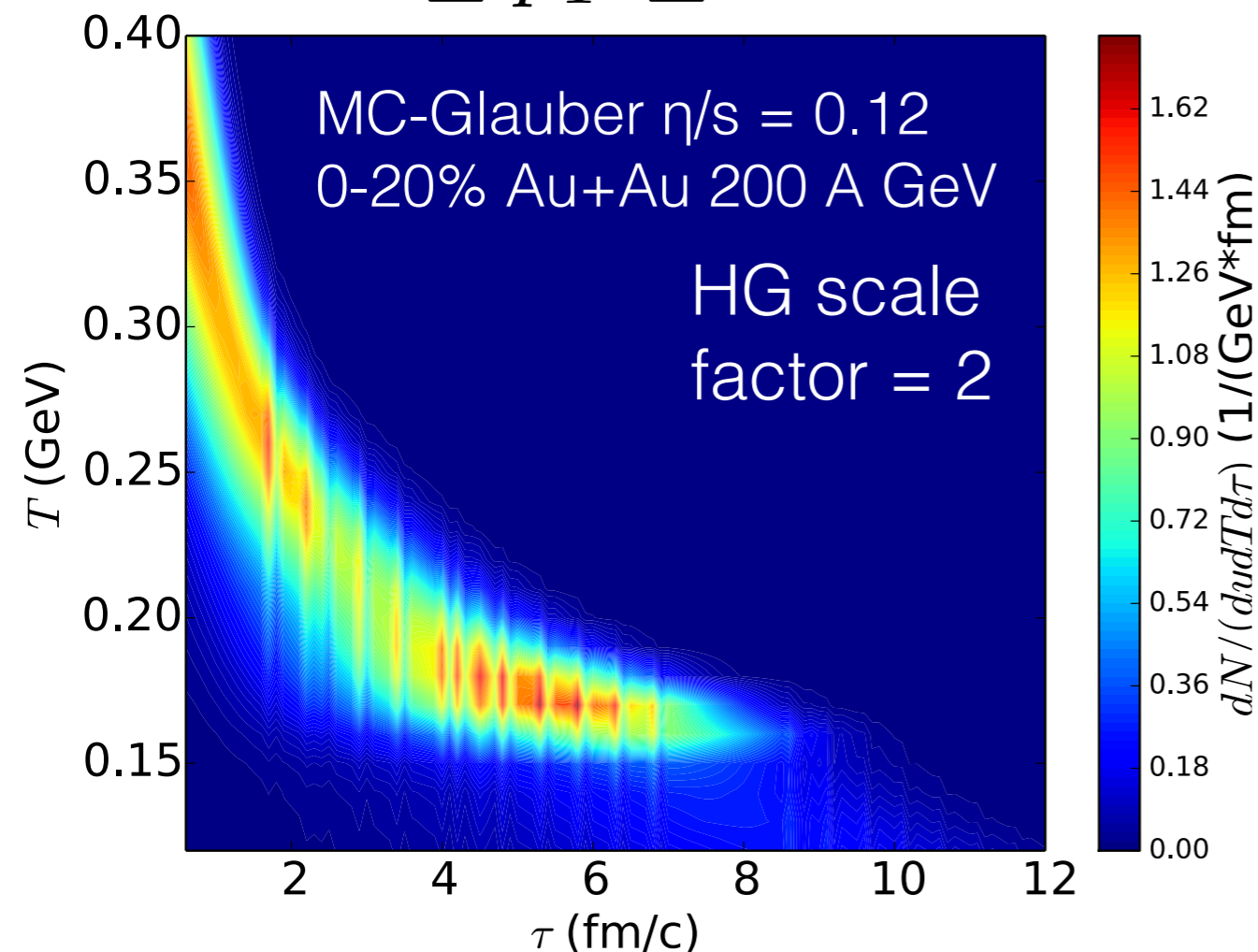
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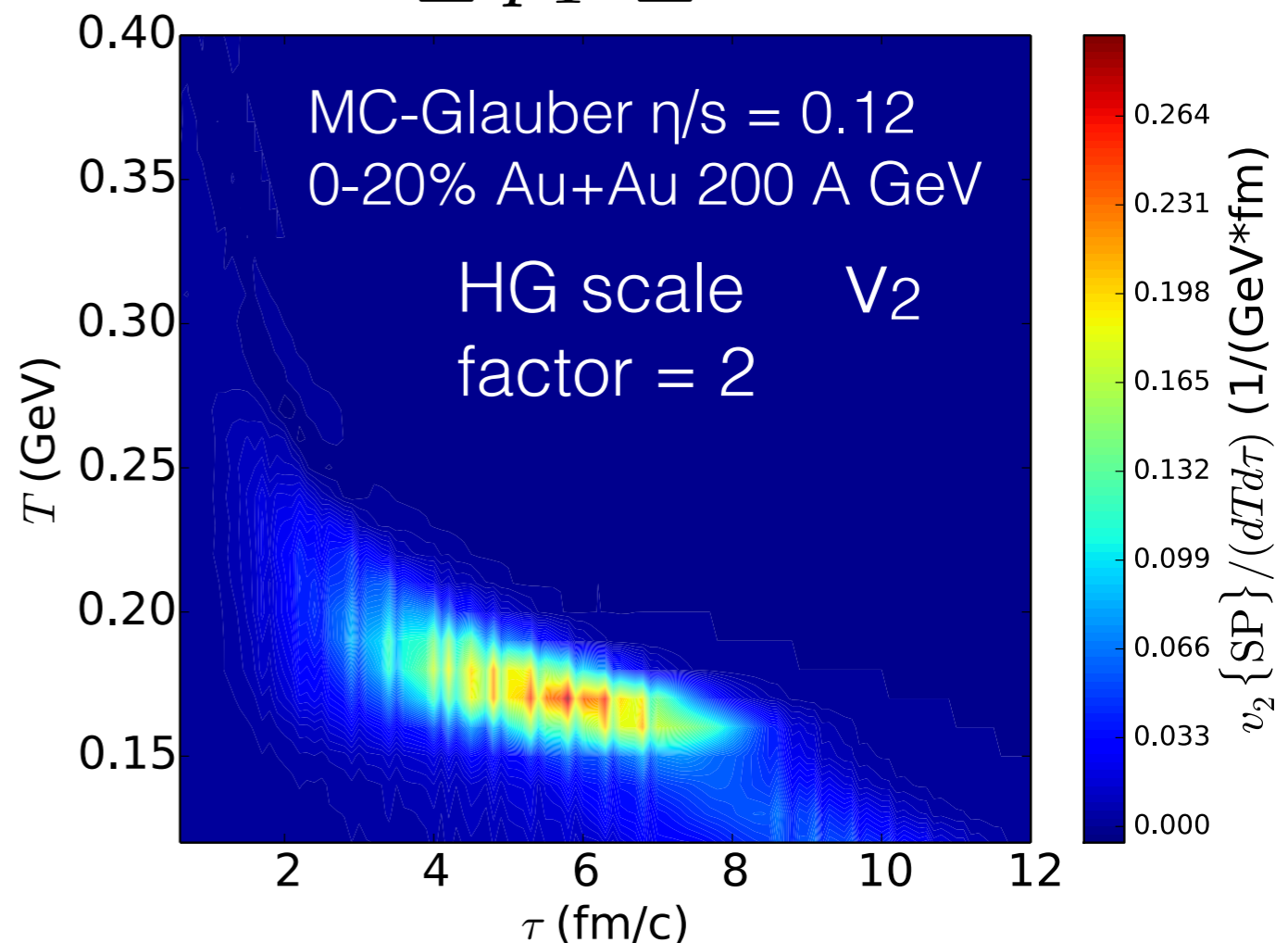
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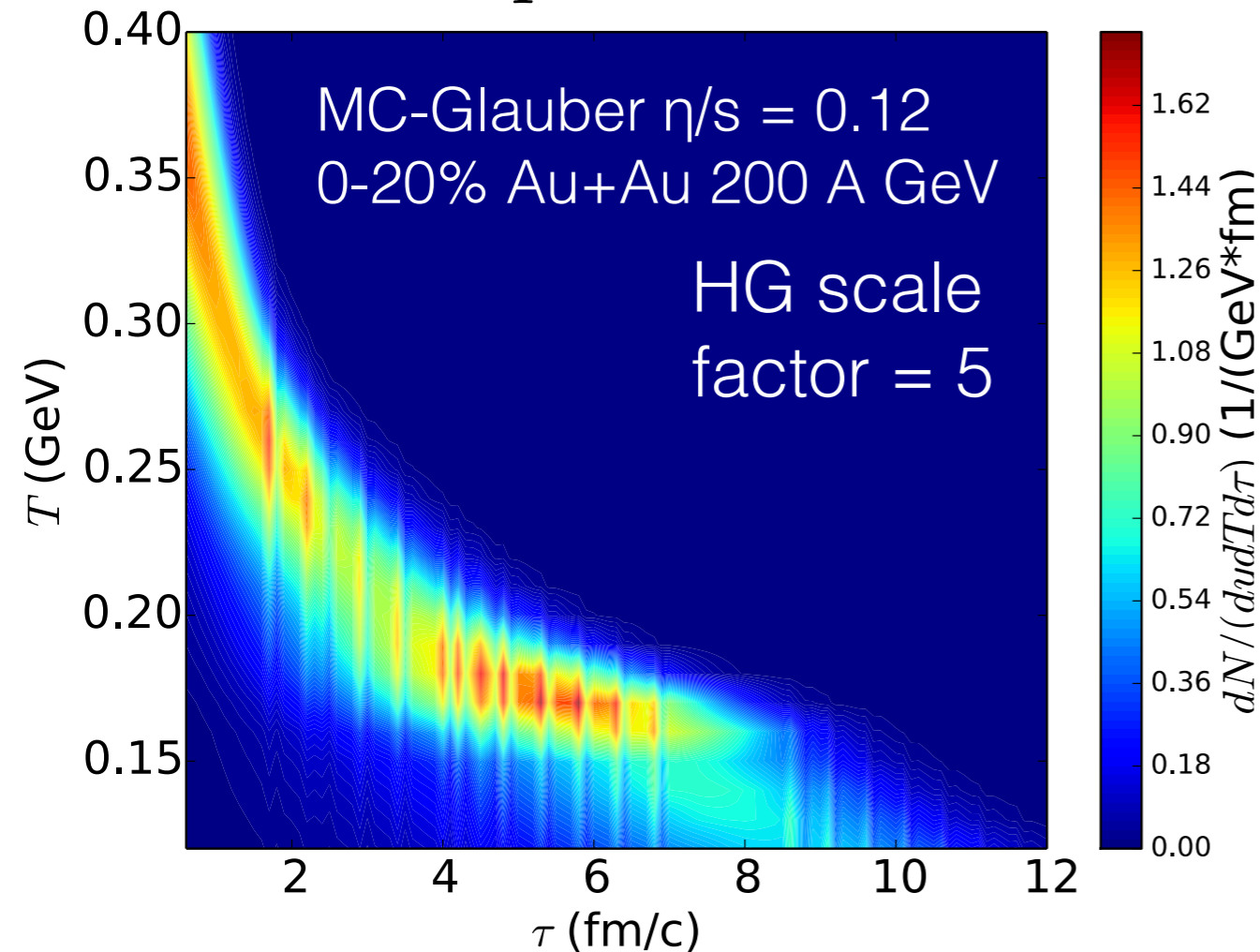
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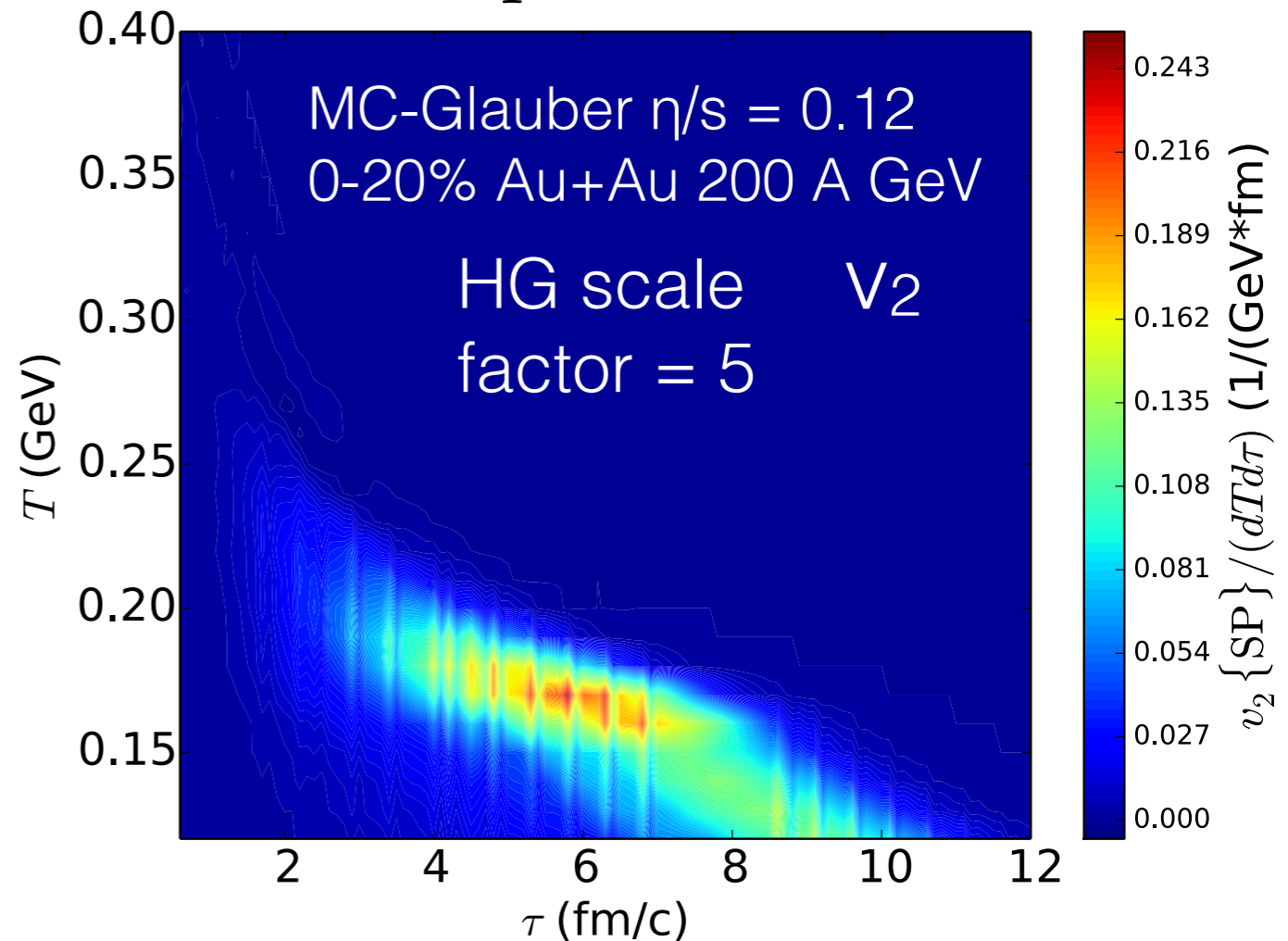
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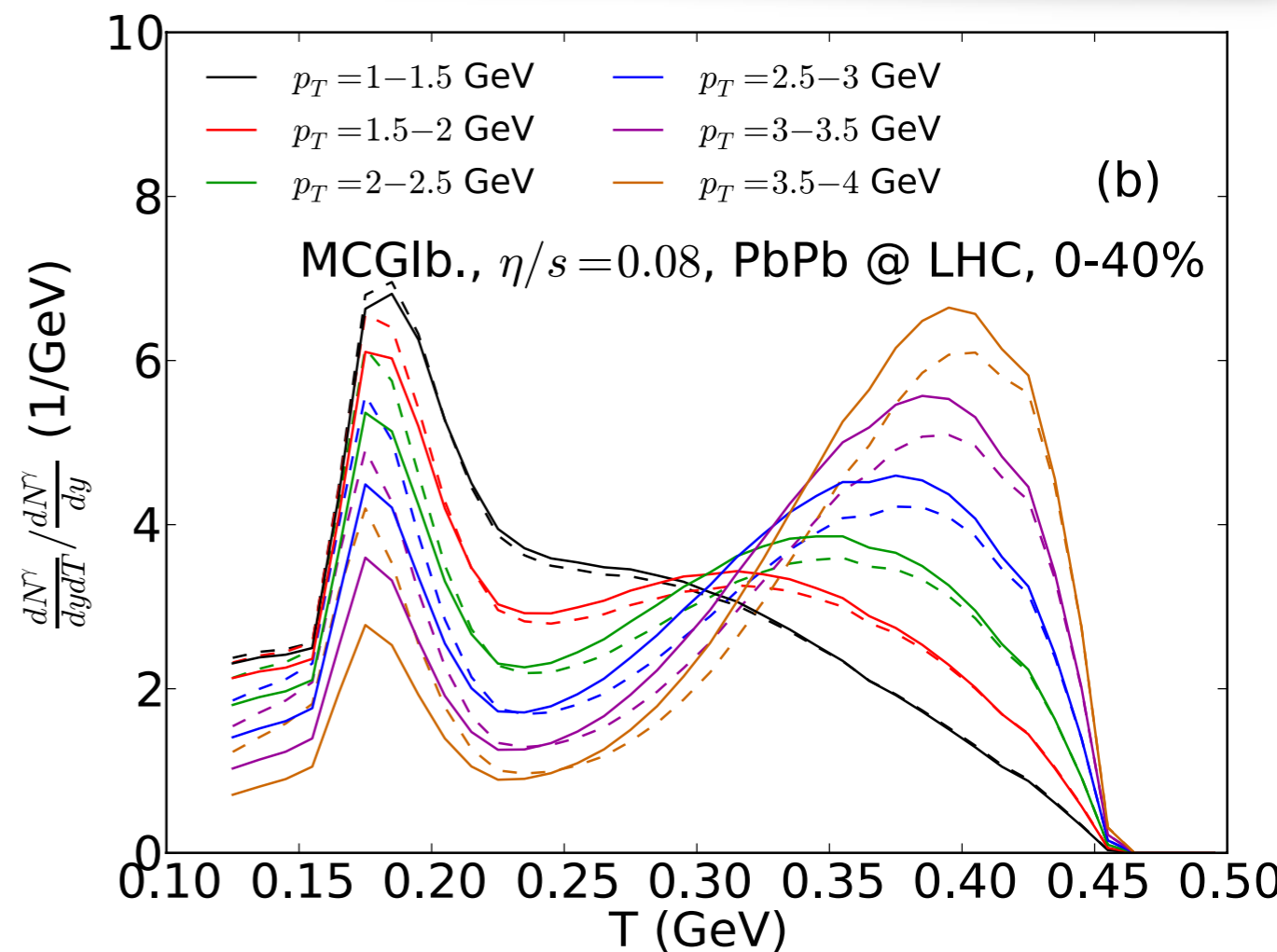
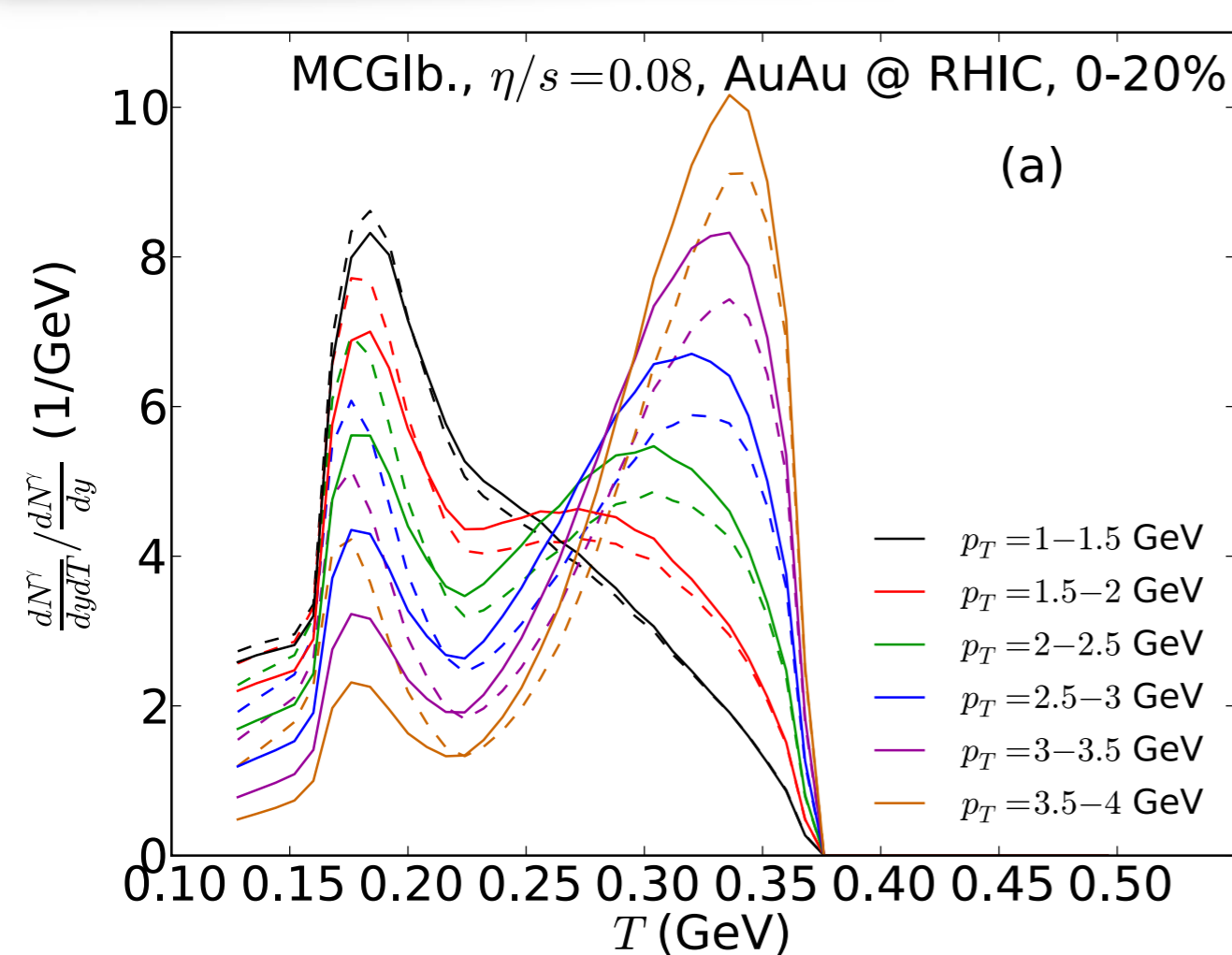


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Emission vs. Temperature



- High p_T photons are mostly emitted from high temperature region
- Peak photon production around $T = 165-200 \text{ MeV}$ due to large hydrodynamic space-time volume