

Thermal photons and dileptons Theory

Hendrik van Hees

Goethe University Frankfurt

August 20, 2014



- 1 Electromagnetic probes in heavy-ion collisions
 - Em. current correlation function and electromagnetic probes
 - Sources of dilepton emission in heavy-ion collisions
 - Sources of thermal photons in heavy-ion collisions
- 2 Application to heavy-ion collisions
 - Models for bulk-medium evolution
 - Dielectrons (SIS/HADES) with S. Endres, M. Bleicher, R. Rapp
 - Dimuons (SPS/NA60) with S. Endres, M. Bleicher, R. Rapp
 - Direct Photons at RHIC and LHC with M. He, R. Rapp
- 3 Conclusions and Outlook

Em. current correlator $l^+ l^-$ and γ rates

Electromagnetic probes in heavy-ion collisions

- γ, l^\pm : no strong interactions
- reflect whole “history” of collision:
 - from **pre-equilibrium phase**
 - from thermalized medium
QGP and hot hadron gas
 - from VM decays **after thermal freezeout**

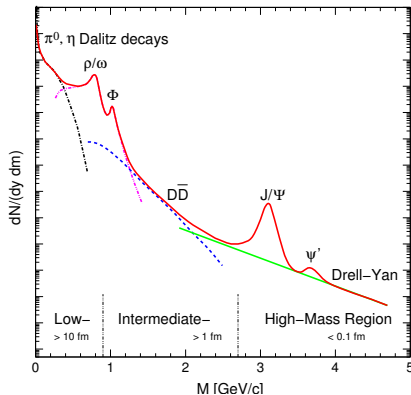
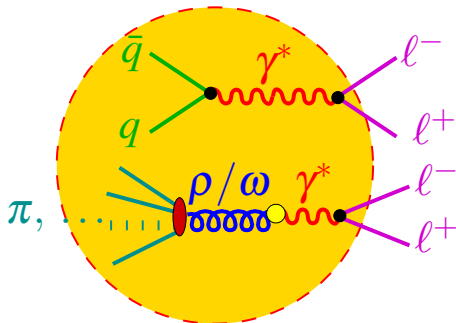


Fig. by A. Drees

- **photon** and **dilepton** thermal emission rates given by **same electromagnetic-current-correlation function** ($J_\mu = \sum_f Q_f \bar{\Psi}_f \gamma_\mu \Psi_f$)

[L. McLerran, T. Toimela 85, H. A. Weldon 90, C. Gale, J.I. Kapusta 91]

$$\Pi_{\mu\nu}^<(q) = \int d^4x \exp(iq \cdot x) \langle J_\mu(0) J_\nu(x) \rangle_T = -2f_B(q \cdot u) \text{Im} \Pi_{\mu\nu}^{(\text{ret})}(q)$$

$$q_0 \frac{dN_\gamma}{d^4x d^3\vec{q}} = \frac{\alpha}{2\pi^2} g^{\mu\nu} \text{Im} \Pi_{\mu\nu}^{(\text{ret})}(q) \Big|_{q_0=|\vec{q}|} f_B(q \cdot u)$$

$$\frac{dN_{e^+e^-}}{d^4x d^4q} = -g^{\mu\nu} \frac{\alpha^2}{3q^2 \pi^3} \text{Im} \Pi_{\mu\nu}^{(\text{ret})}(q) \Big|_{q^2=M_{e^+e^-}^2} f_B(q \cdot u)$$

- u : four-velocity of the fluid cell; $p \cdot u = p_0^{\text{hb}}$ energy in “heat-bath frame”
- to lowest order in α : $e^2 \Pi_{\mu\nu} \simeq \Sigma_{\mu\nu}^{(\gamma)}$
- **vector-meson dominance** model:

$$\Sigma_{\mu\nu}^\gamma = \text{---} \overset{G_\rho}{\text{---}} \text{---}$$

Sources of dilepton emission in heavy-ion collisions

- 1 initial hard processes: Drell Yan
- 2 “core” \Leftrightarrow emission from thermal source [McLerran, Toimela 1985]

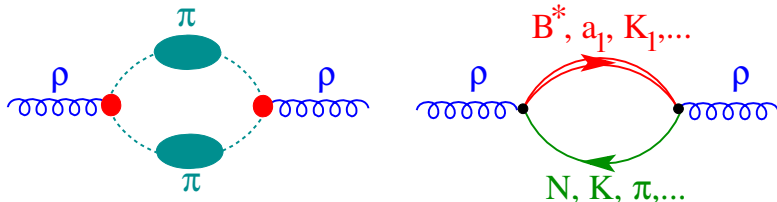
$$\frac{1}{q_T} \frac{dN^{(\text{thermal})}}{dM dq_T} = \int d^4x \int dy \int M d\phi \frac{dN^{(\text{thermal})}}{d^4x d^4q}$$

- 3 “corona” \Leftrightarrow emission from “primordial” mesons (jet-quenching)
- 4 after thermal freeze-out \Leftrightarrow emission from “freeze-out” mesons
[Cooper, Frye 1975]

$$N^{(\text{fo})} = \int \frac{d^3q}{q_0} \int q_\mu d\sigma^\mu f_B(u_\mu q^\mu / T) \frac{\Gamma_{\text{meson} \rightarrow \ell^+ \ell^-}}{\Gamma_{\text{meson}}}$$

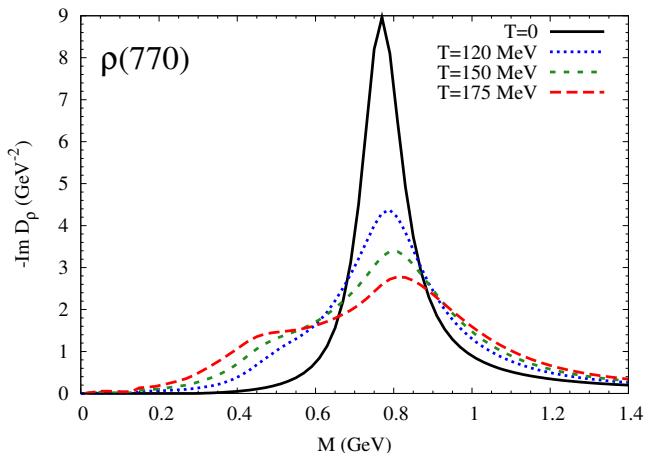
Hadronic many-body theory

- HMBT for vector mesons [Ko et al, Chanfray et al, Herrmann et al, Rapp et al, ...]
- $\pi\pi$ interactions and **baryonic excitations**



- +corresponding vertex corrections \Leftrightarrow gauge invariance
- **Baryon (resonances)** important, even at RHIC with low **net** baryon density $n_B - n_{\bar{B}}$
- reason: $n_B + n_{\bar{B}}$ relevant (CP inv. of strong interactions)

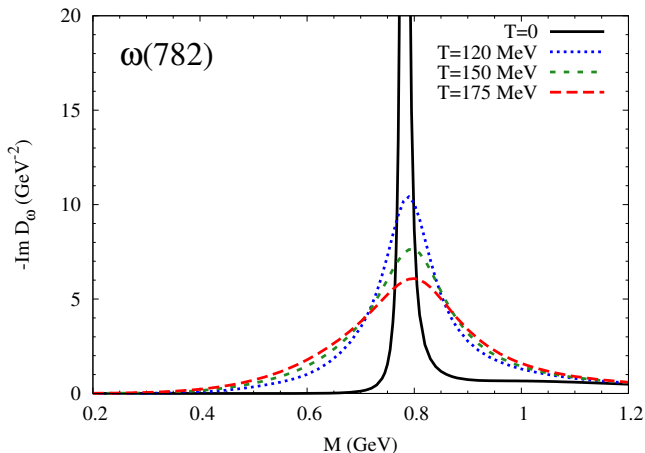
In-medium spectral functions and baryon effects



[R. Rapp, J. Wambach 99]

- **baryon effects** important
 - large contribution to broadening of the peak
 - responsible for most of the strength at small M

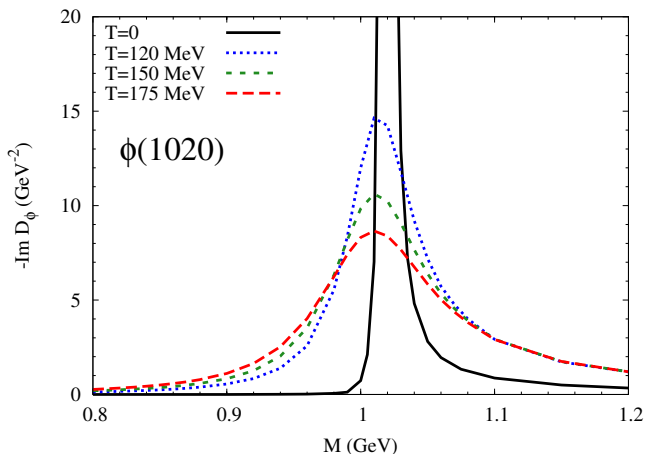
In-medium spectral functions and baryon effects



[R. Rapp, J. Wambach 99]

- **baryon effects** important
 - large contribution to broadening of the peak
 - responsible for most of the strength at small M

In-medium spectral functions and baryon effects

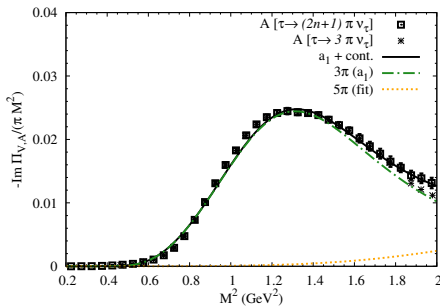
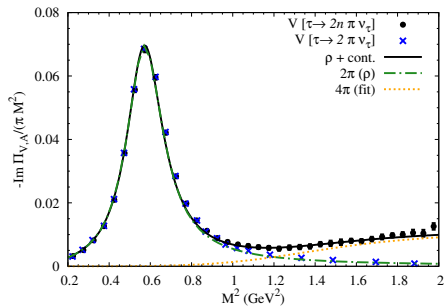


[R. Rapp, J. Wambach 99]

- **baryon effects** important
 - large contribution to broadening of the peak
 - responsible for most of the strength at small M

Intermediate masses: hadronic “ 4π contributions”

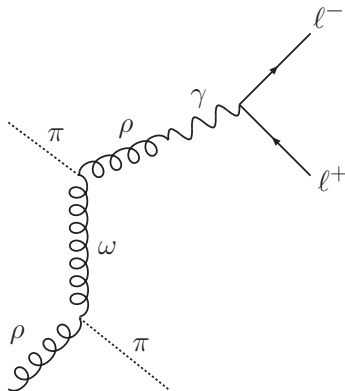
- e.m. current-current correlator $\Leftrightarrow \tau \rightarrow 2n\pi$



- “ 4π contributions”: $\pi + \omega, a_1 \rightarrow \mu^+ + \mu^-$
- leading-order virial expansion for “four-pion piece”
- additional strength through “chiral mixing”

Radiation from thermal sources: Meson t-channel exchange

- motivation: q_T spectra too soft compared to NA60 data
- **thermal contributions** not included in models so far



- also for π, a_1

Dileptons from thermal QGP

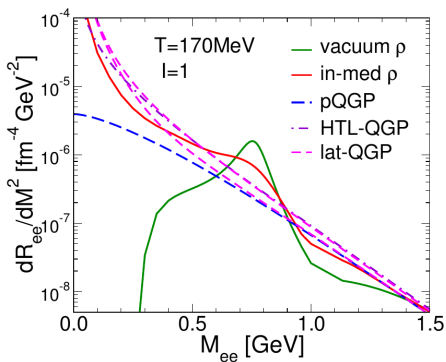
- in QGP phase: $q\bar{q}$ annihilation
- HTL improved electromagnetic current correlator

$$-i\Pi_{\text{em,QGP}} = \text{Diagram}$$

- or electromagnetic current correlator from the **lattice** [H.-T. Ding, A. Francis et al (Bielefeld) 2011] (extrapolated to finite q)
- “quark-hadron duality” around T_c

Dilepton rates: Hadron gas \leftrightarrow QGP

- in-medium **hadron gas** matches with **QGP**
- similar results also for γ rates
- “quark-hadron duality”?

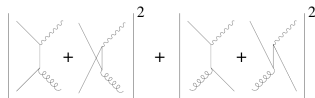


[R. Rapp, arXiv: 1304.2309 [hep-ph]]

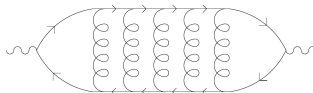
Sources of thermal photons in heavy-ion collisions

- **QGP:** rates from [Arnold, Moore, Yaffe, JHEP **12**, 009 (2001)]

- $q\bar{q} \rightarrow \gamma g, qg \rightarrow \gamma q$

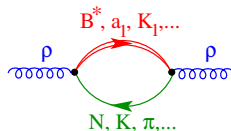
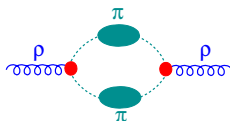


- resummation of soft-gluon bremsstrahlung contributions
- Landau-Pomeranchuk-Migdal effect



- **hadronic matter** from [Turbide, Rapp, Gale, PRC **69**, 014903 (2004); Rapp, Wambach EPJ A **6**, 415 (1999)]

- pion-cloud dressing + vector meson-baryon/meson interactions

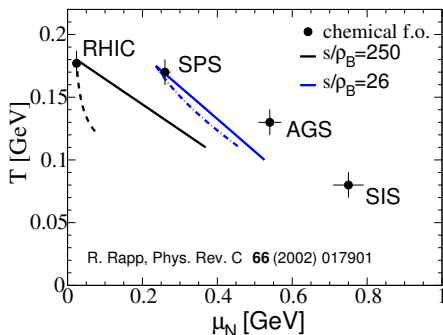


- $\pi\rho a_1, \omega$ -t-channel exchange

Medium evolution

Thermal fireball

- cylindrical fireball model: $V_{\text{FB}} = \pi(z_0 + v_{z0}t + \frac{a_z}{2}t^2) (\frac{a_{\perp}}{2}t^2 + r_0)^2$
- thermodynamics:
 - isentropic expansion; S_{tot} fixed by N_{ch} ; $T_c = T_{\text{chem}} = 175$ MeV
 - $T > T_c$: QGP; lattice equation of state
 - continuous cross-over (no 1st-order mixed state!)
 - $T < T_c$: hadron-resonance gas
- $\Rightarrow T(t), \mu_{\text{baryon,meson}}(t)$
- chemical freezeout:
 - $\mu_N^{\text{chem}} = 232$ MeV
 - hadron ratios fixed
 $\Rightarrow \mu_N, \mu_{\pi}, \mu_K, \mu_{\eta}$ at fixed
 $s/\rho_B = 27$
- thermal freezeout:
 $(T_{\text{fo}}, \mu_{\pi}^{\text{fo}}) \simeq (120, 80)$ MeV



Coarse-grained transport (UrQMD)

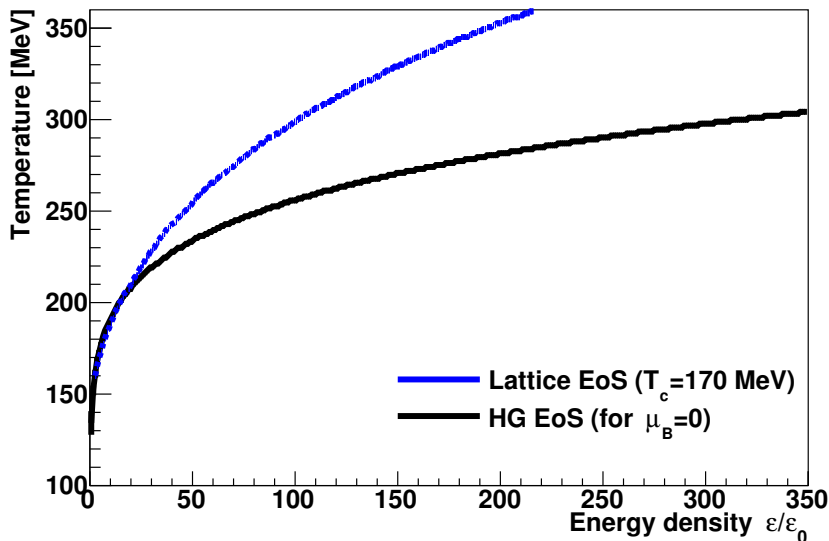
- Use **ensemble of UrQMD** runs with an **equation of state**
- map evolution of medium to **locally thermalized fluid cells**
- fit **temperature, chemical potentials, flow-velocity field**
from anisotropic energy-momentum tensor [W. Florkowski et al, NPA **904-905**, 803c (2013)]

$$T^{\mu\nu} = (\varepsilon + P_{\perp})u^{\mu}u^{\nu} - P_{\perp}g^{\mu\nu} - (P_{\parallel} - P_{\perp})V^{\mu}V^{\nu}$$

- thermal rates from **partonic/hadronic QFT** become **applicable**
- here: **extrapolated lattice QGP** and **Rapp-Wambach hadronic many-body theory**
- caveat: **consistency between EoS, matter content of QFT model/UrQMD!**

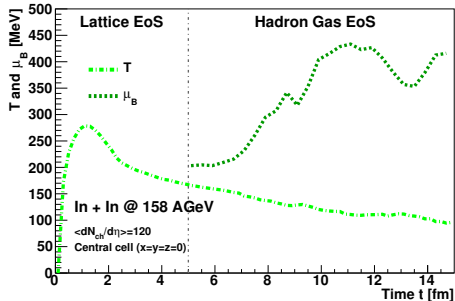
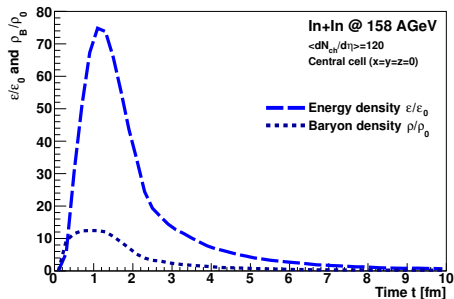
Coarse-grained transport (UrQMD)

- $T_c = 170$ MeV; $T > T_c \Rightarrow$ lattice EoS; $T < T_c \Rightarrow$ HRG EoS

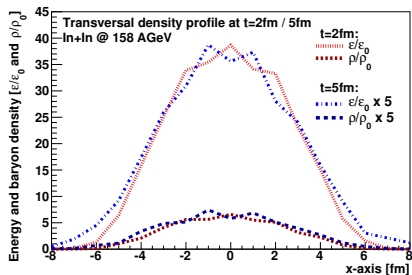
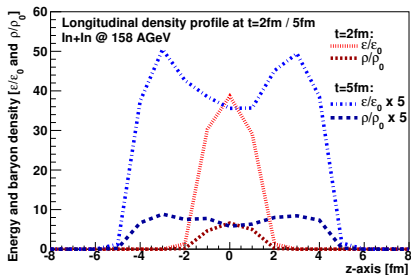


Coarse-grained transport (UrQMD)

- energy/baryon density $\Rightarrow T, \mu_B$ (for In+In @ SPS; NA60)
- central “fluid” cell!

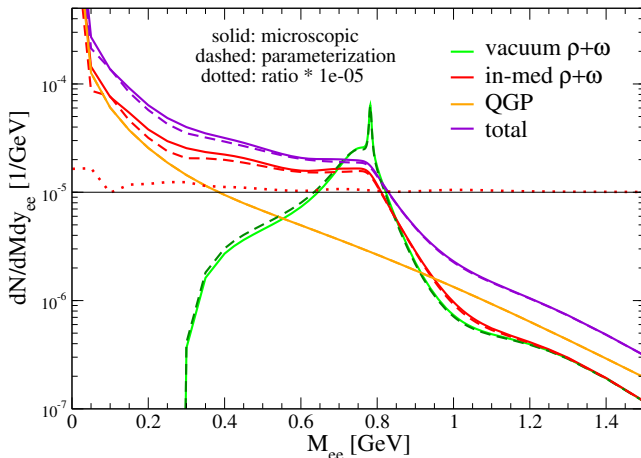


- temperature/density profiles (for In+In@SPS; NA60)



Parametrized Rapp-Wambach rates

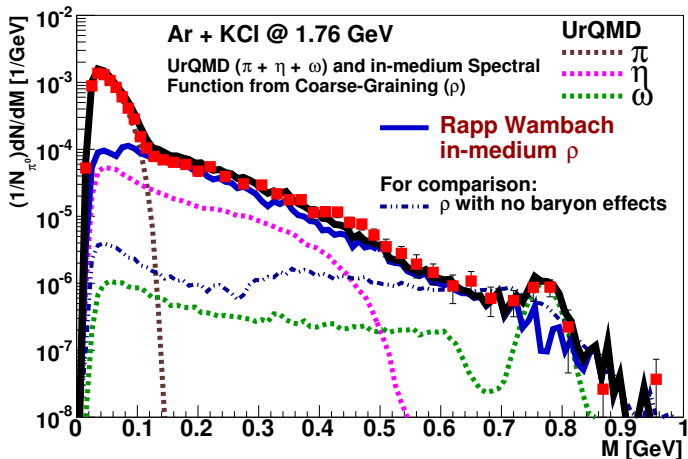
- need rates as function of T , μ_B , μ_π , μ_K
- parametrization of the **microscopic rates** necessary
- comparison for 20 AGeV Au Au collisions (min bias) [R. Rapp private commun.]
- pion-cloud effects not fully implemented \Rightarrow some deviations in LMR



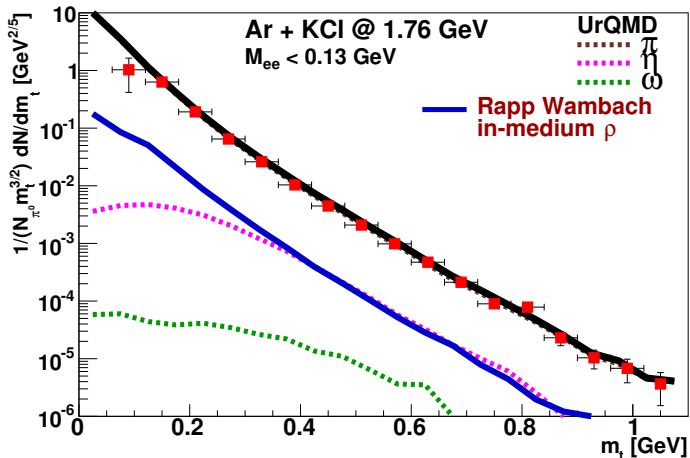
Dielectrons (SIS/HADES)

$e^+e^- M$ spectrum (SIS/HADES)

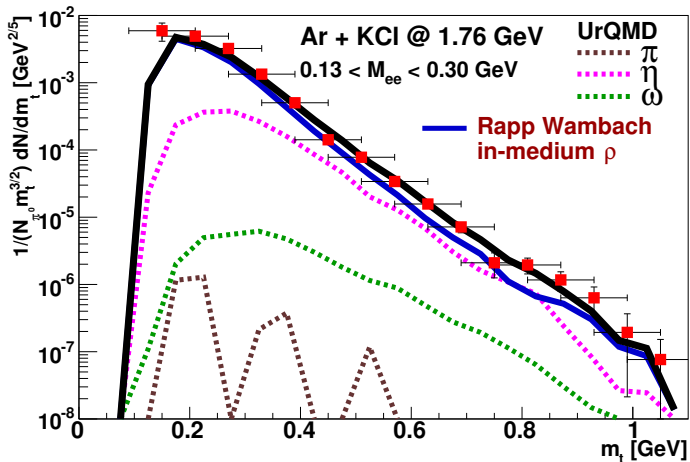
- coarse-graining method works at low energies!
- UrQMD-medium evolution + RW-QFT rates



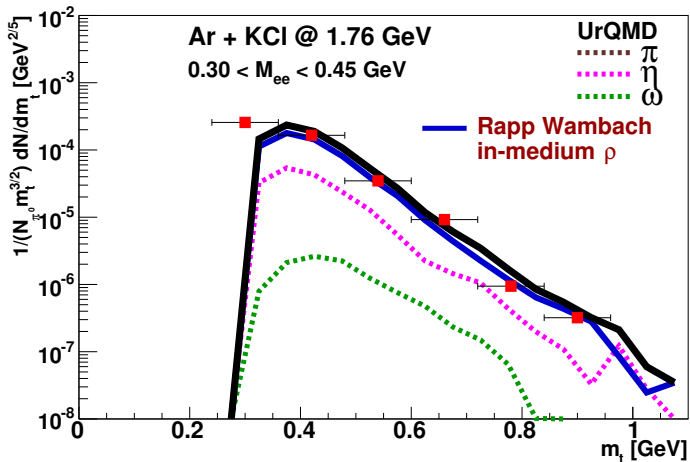
- dielectron spectra from $\text{Ar} + \text{KCl}(1.76 \text{ AGeV}) \rightarrow e^+e^-$ (SIS/HADES)



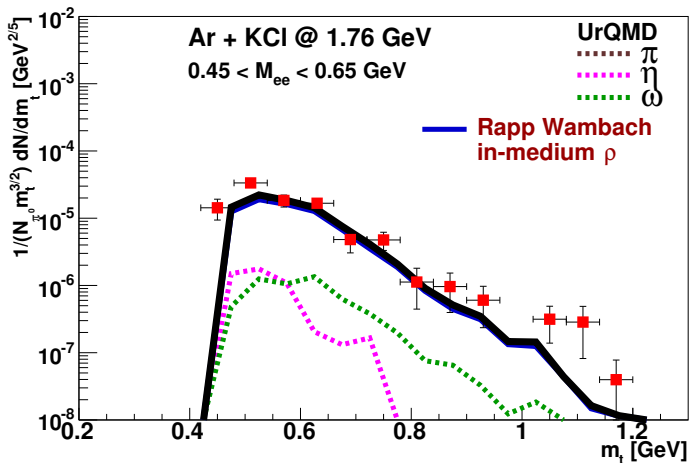
- dielectron spectra from $\text{Ar} + \text{KCl} (1.76 \text{ AGeV}) \rightarrow e^+e^-$ (SIS/HADES)



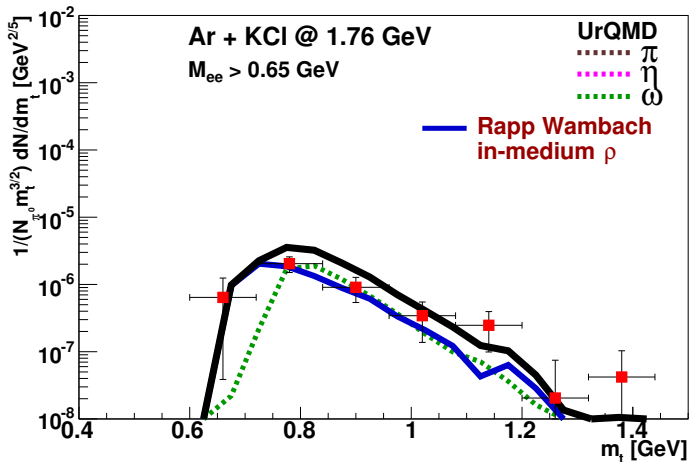
- dielectron spectra from $\text{Ar} + \text{KCl} (1.76 \text{ AGeV}) \rightarrow e^+e^-$ (SIS/HADES)



- dielectron spectra from $\text{Ar} + \text{KCl} (1.76 \text{ AGeV}) \rightarrow e^+e^-$ (SIS/HADES)



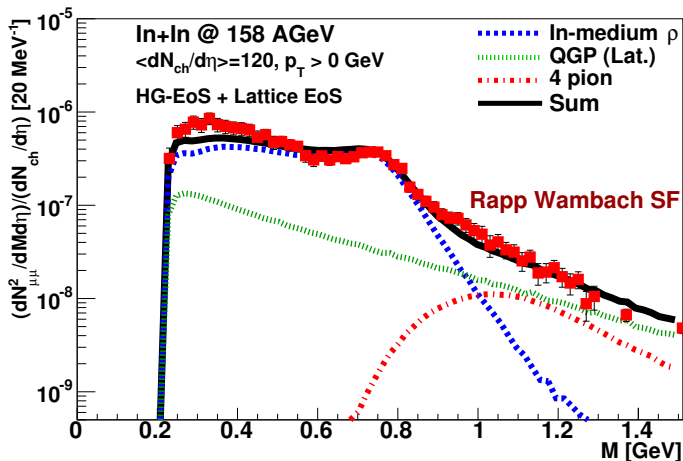
- dielectron spectra from $\text{Ar} + \text{KCl} (1.76 \text{ AGeV}) \rightarrow e^+e^-$ (SIS/HADES)



Dimuons (SPS/NA60)

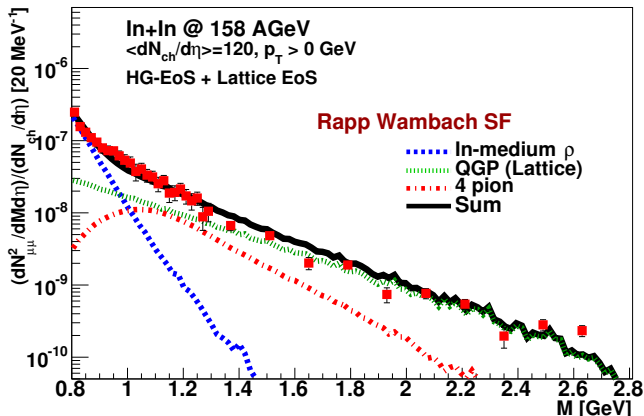
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



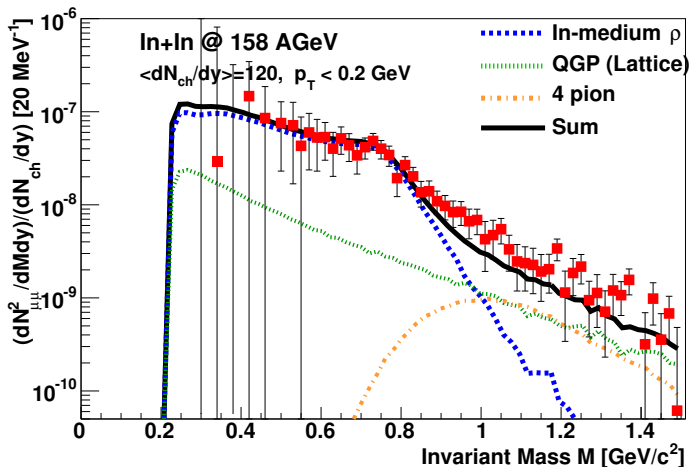
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)
- higher IMR: provides **averaged true temperature** (no blueshifts in the **invariant-mass** spectra!)



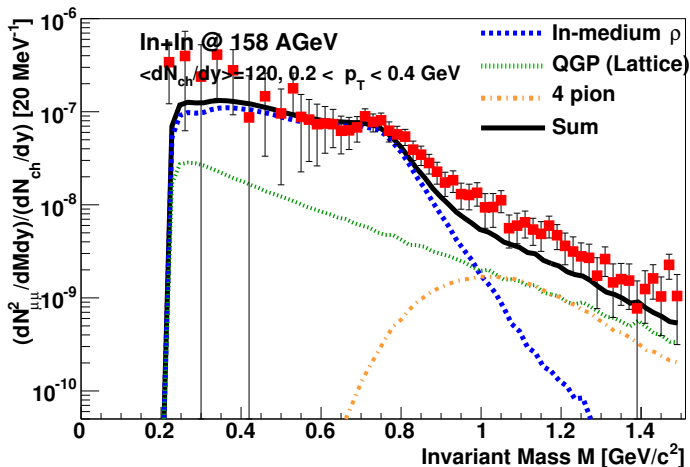
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



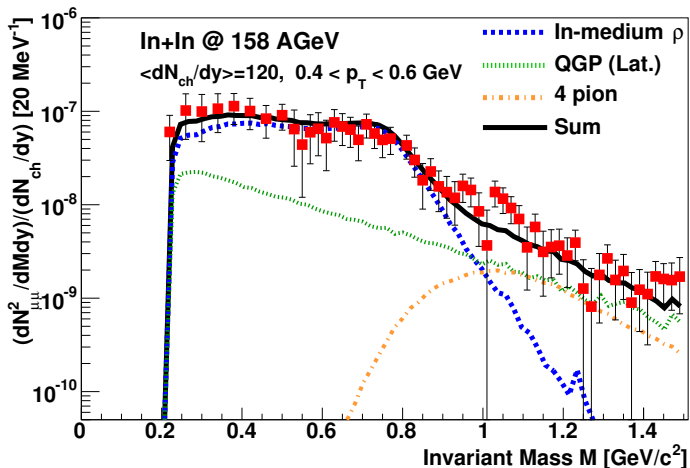
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In (158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



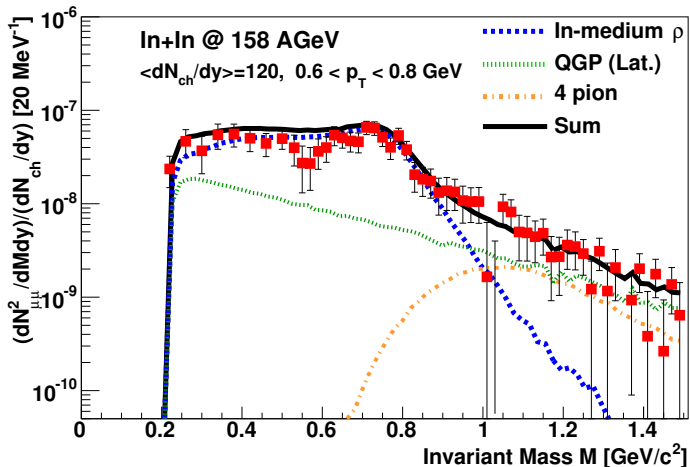
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



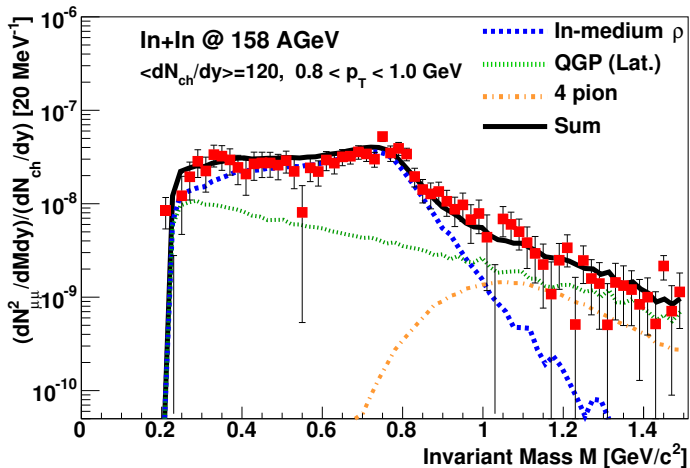
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



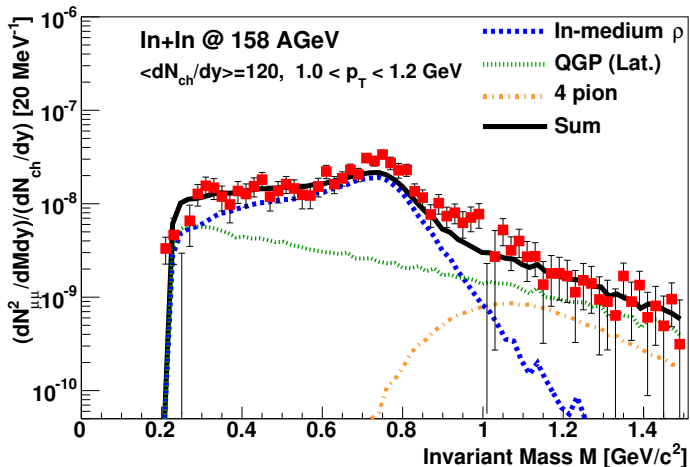
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



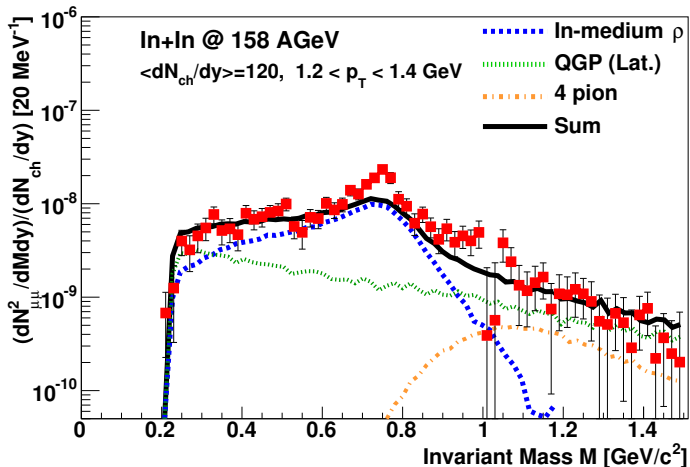
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



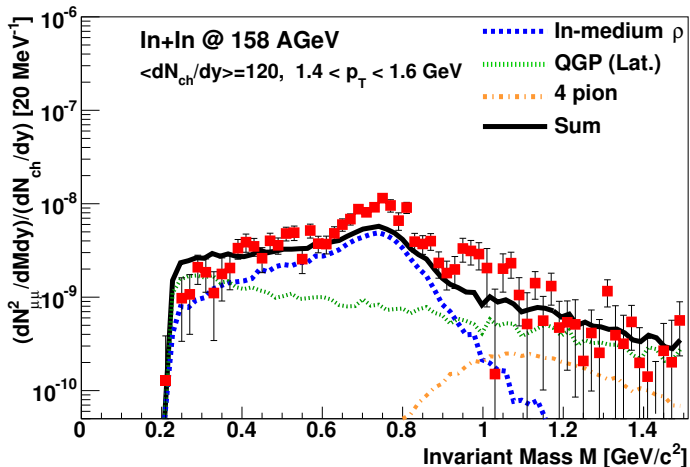
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



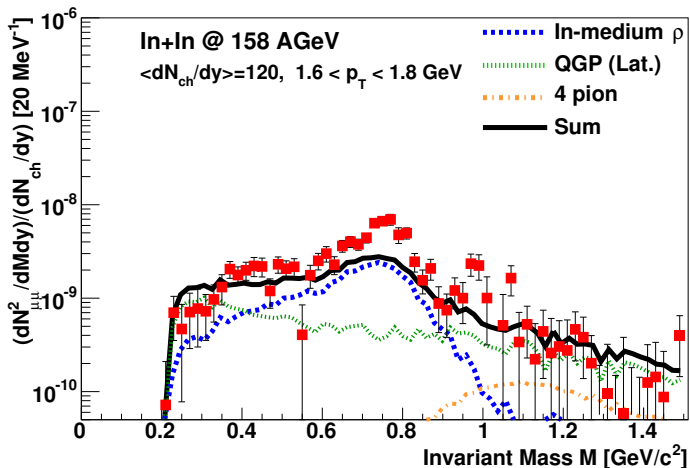
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



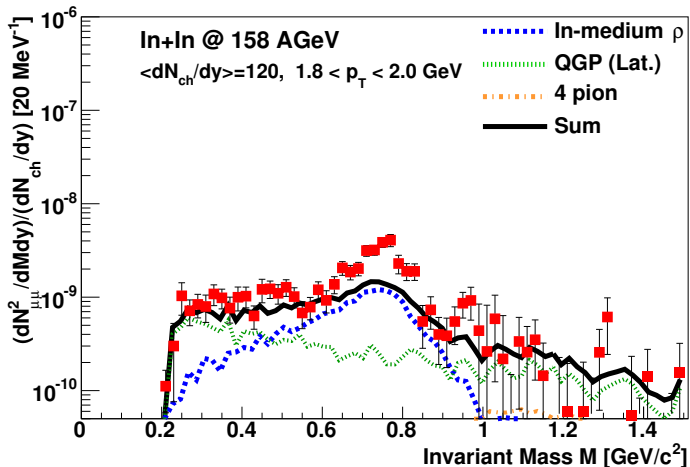
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In (158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



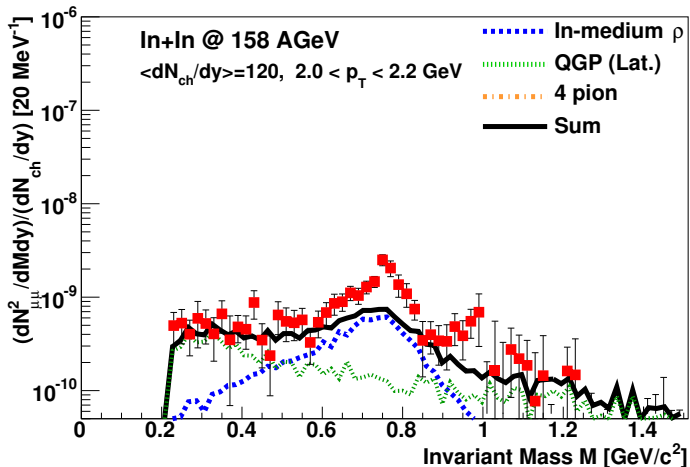
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



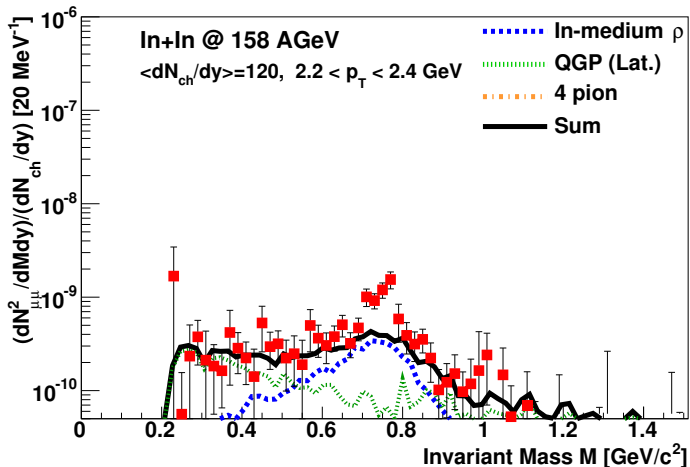
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



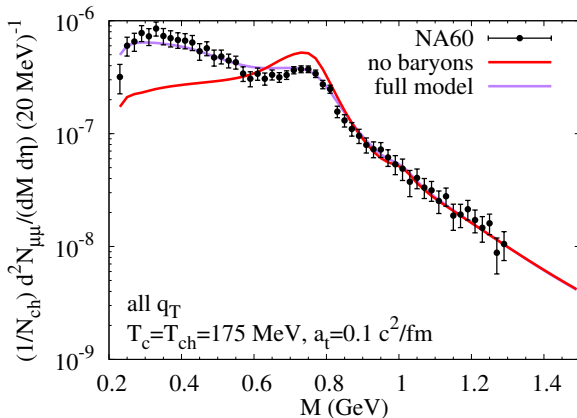
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



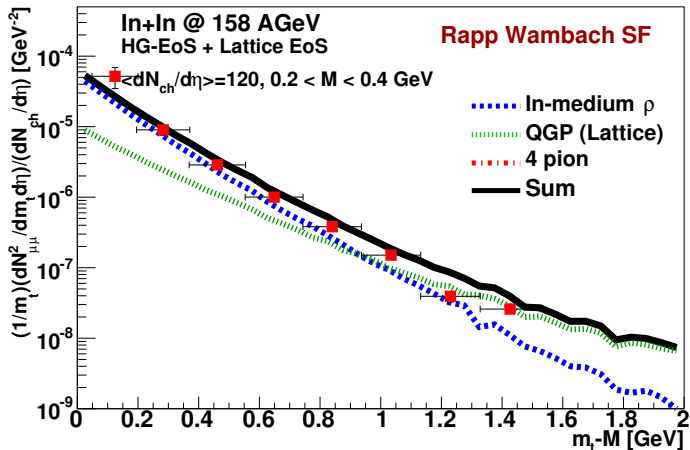
$\mu^+\mu^-$ M spectra (SPS/NA60)

- dimuon spectra from $\text{In} + \text{In}(158\text{A GeV}) \rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)
- influence of baryon interactions in spectral function
- from previous calculation with thermal-fireball parametrization (compatible with course-grained UrQMD)



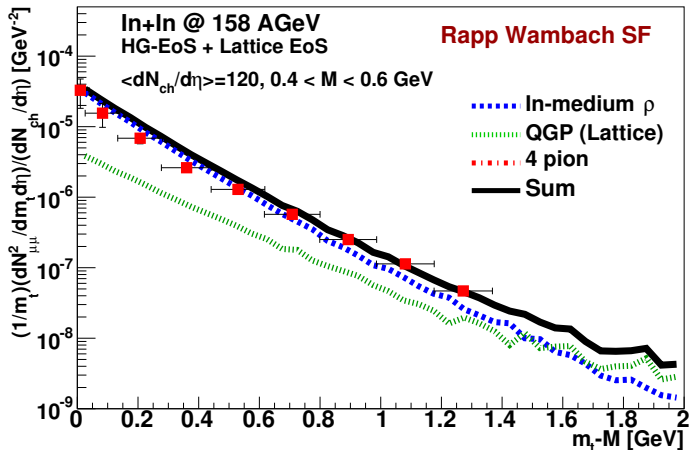
$\mu^+\mu^- m_T$ spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



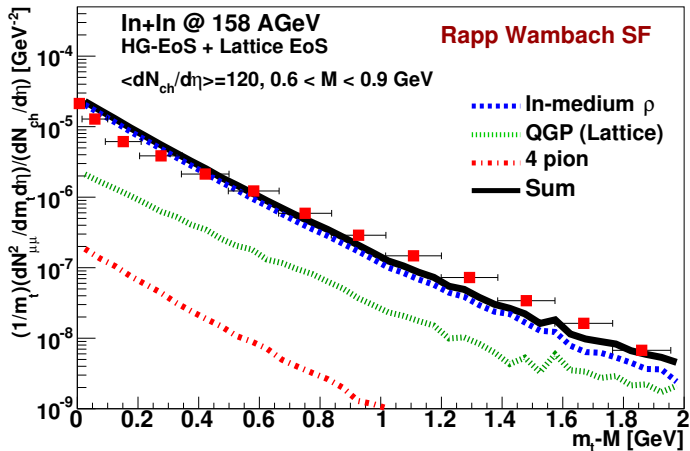
$\mu^+\mu^- m_T$ spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{ch}/dy = 120$)



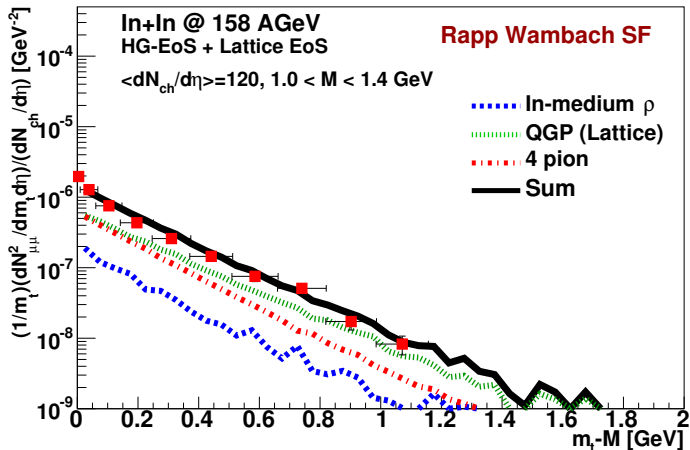
$\mu^+\mu^- m_T$ spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+\mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)



$\mu^+ \mu^- m_T$ spectra (SPS/NA60)

- dimuon spectra from In + In(158 AGeV) $\rightarrow \mu^+ \mu^-$ (NA60)
- min-bias data ($dN_{\text{ch}}/dy = 120$)

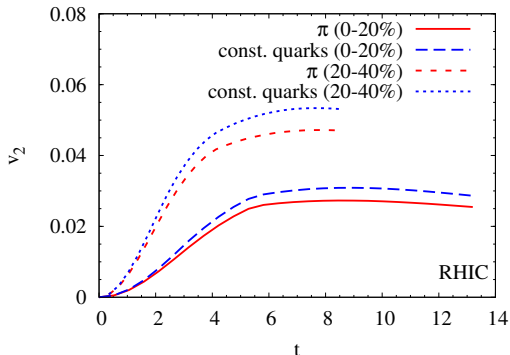


Direct Photons at RHIC and LHC

Direct Photons at RHIC and LHC: elliptic fireball/hydro

- fitted to **measured p_T spectra and v_2** ; multi-strange hadrons: fo at T_c !
- can be achieved with (ideal) hydro

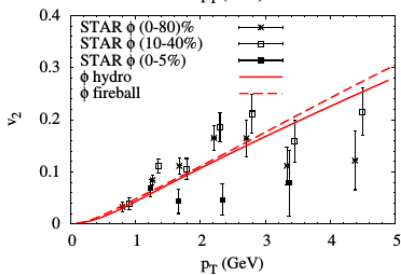
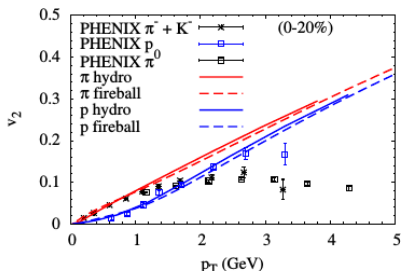
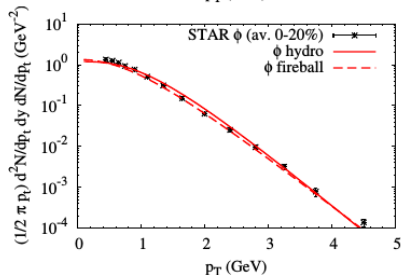
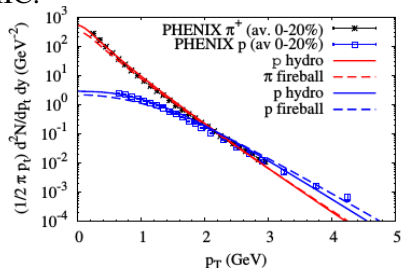
[He, Fries, Rapp, PRC 85, 044911 (2012); HvH, He, Rapp arXiv: 1404.2846 [nucl-th]]



- important for “sufficient” photon v_2 :
 - rapid buildup of v_2
 - (nearly) full v_2 at end of mixed phase
 - consistent with **CQN scaling** for multi-strange and other hadrons!

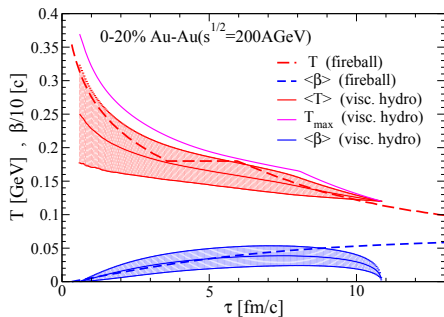
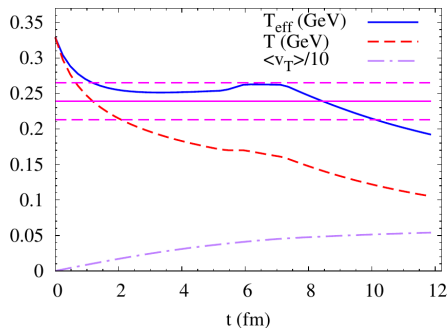
Direct Photons at RHIC and LHC: elliptic fireball/hydro

RHIC:



[HvH, Gale, Rapp, Phys. Rev. C **84**, 054906 (2011); HvH, He, Rapp arXiv: 1404.2846 [nucl-th]]

Temperature vs. effective Slope



[Rapp, Hvh, He, arXiv: 1408.0612 [nucl-th]]

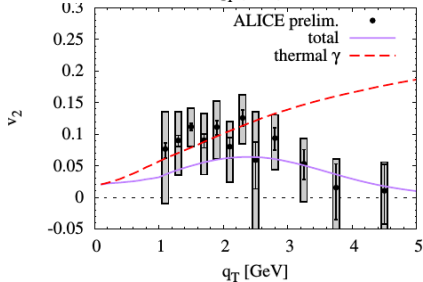
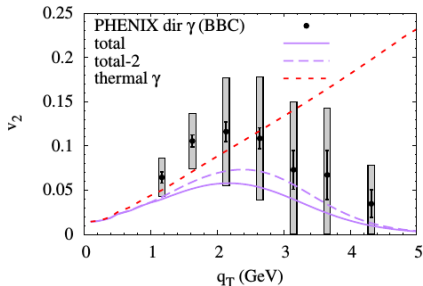
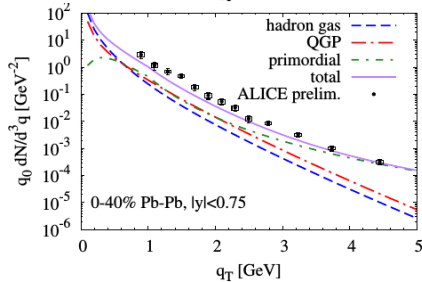
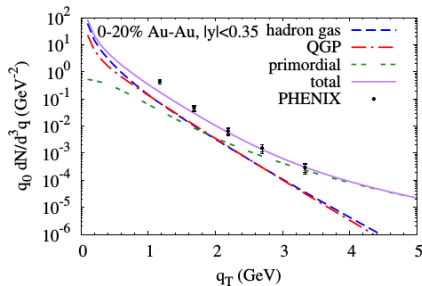
[C. Shen, U. W. Heinz, J.-F. Paquet, C. Gale]
[arXiv:1308.2440 [nucl-th]]

- blue-shift formula (Doppler effect) translates into

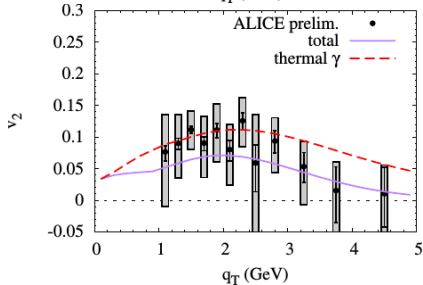
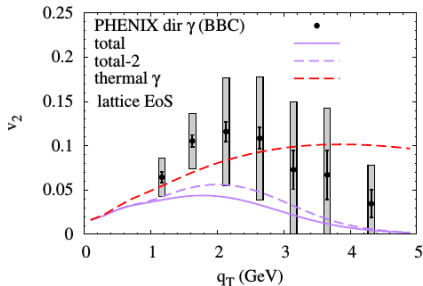
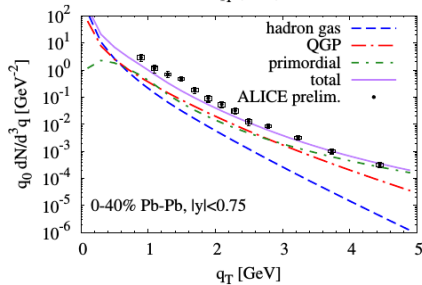
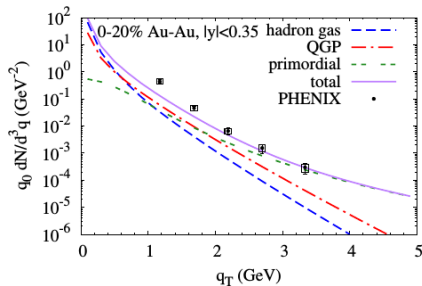
$$T_{\text{eff}} \simeq T \sqrt{\frac{1 + \langle v_T \rangle}{1 - \langle v_T \rangle}}, \quad v_T : \text{ transverse fluid flow}$$

- measured slope indicates **emission from source around T_c**

Direct photons: fireball

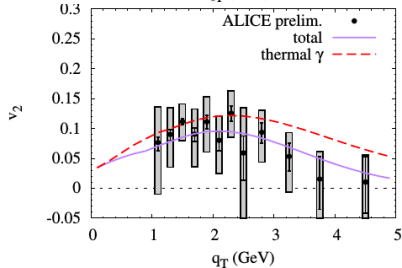
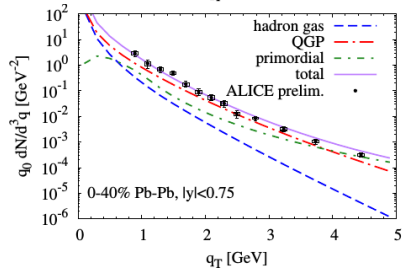
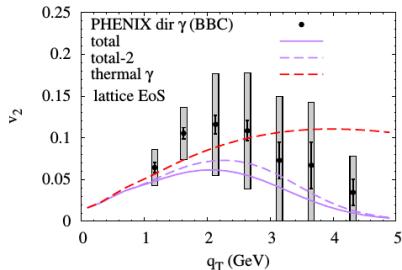
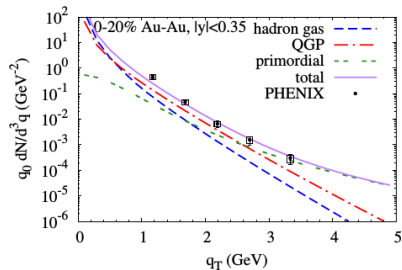


Direct photons: ideal hydro



Direct photons: enhanced rates

- assume enhancement of baseline rates by factor of 2
- augmented up to factor of 3 for $140 \text{ MeV} < T < 200 \text{ MeV}$



Conclusions and Outlook

● General ideas

- em. probes \Leftrightarrow in-medium em. current-correlation function
- dual rates around T_c (compatible with χ symmetry restoration)
 \Rightarrow see Paul Hohler's talk
- medium modifications of ρ , ω , ϕ
- importance of baryon-resonance interactions

● Application to dileptons in HICs

- need realistic bulk-medium evolution
- thermal fireball, (ideal) hydrodynamics
- new: coarse-grained transport
- applicable also at low collision energies
- allows use of thermal-QFT models for em. current-correlation functions
- successful description at HADES, SPS, and RHIC (STAR)
- consistent description of M and m_T spectra!
- Outlook: effective slope of M spectra in higher IMR ($1.5 \text{ GeV} < M < M_{J/\psi}$) provides $\langle T \rangle$
- applied in beam-energy scan at RHIC and FAIR \Rightarrow signature of phase transition?
- signature of cross-over vs. 1st order (or even critical endpoint)?

- Application to photons in HICs
 - so far: bulk evolution with **elliptic thermal fireball and hydro**
 - direct-photon “ **v_2 puzzle**”
 - dominated from **fireball temperatures around T_c (remnant of latent heat)**
- ⇒ Early build-up of elliptic flow
 - compatible with early freeze-out of **multi-strange hadrons**
 - can be achieved with fireball parametrization or choice of appropriate hydro-initial conditions (initial flow)
 - still yield missing ⇒ probable **enhancement of rates** due to non-perturbative **enhanced cross sections around T_c ???**