

Anisotropy of photon production in magnetic field

Vladimir Skokov



G. Basar, D. Kharzeev, V.S., arXiv:1206.1334
A. Bzdak, V.S., arXiv:1208.5502

Outline

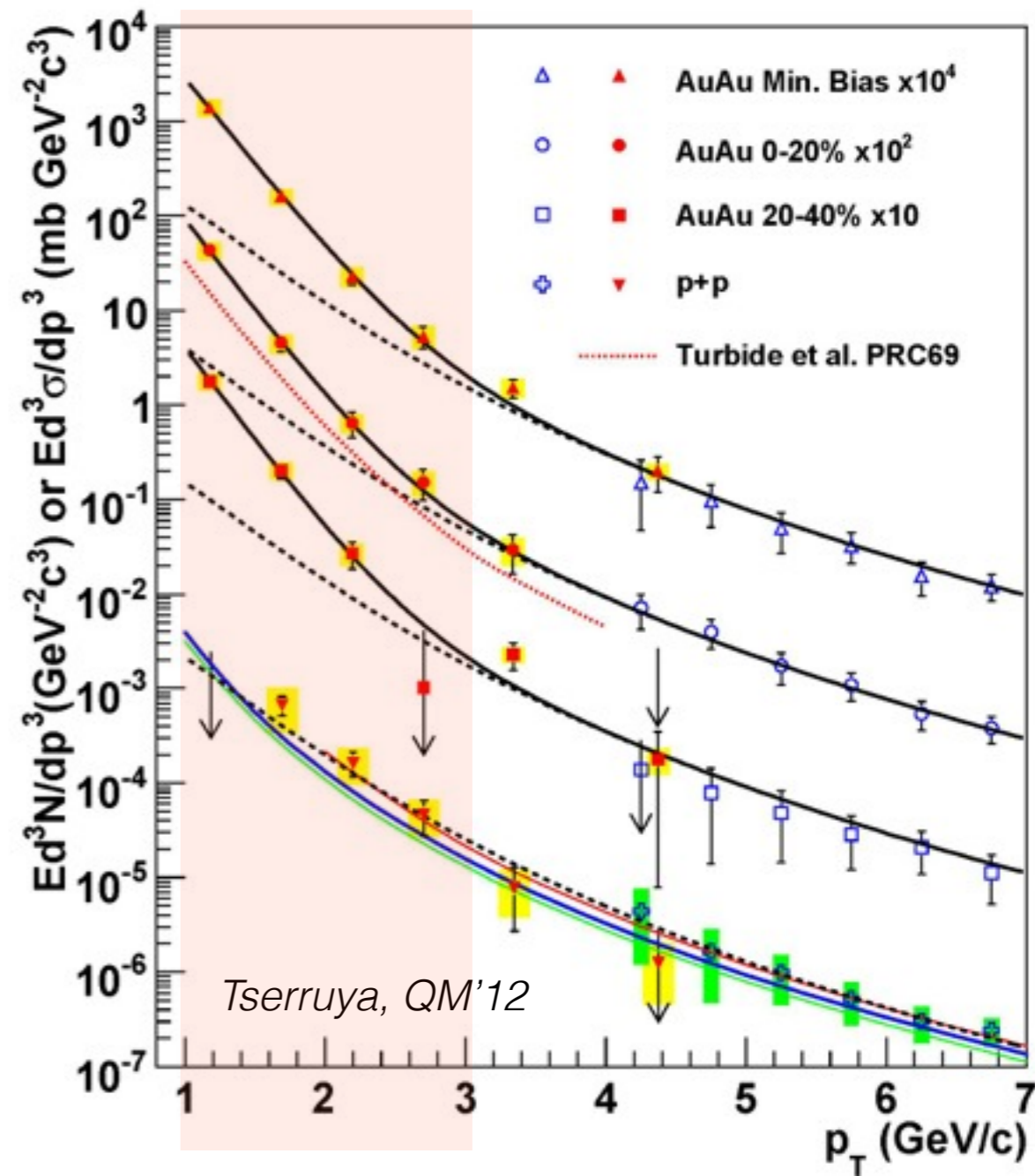
- motivation: photon puzzle
measurements vs expectations
- possible solutions:
“hadronic”
“partonic” + magnetic field
- magnetic field in heavy-ion collisions:
essential properties: magnitude, lifetime, b-dependence
natural source of anisotropy
- photon production and magnetic field:
results & possible experimental signatures

Experimental facts about γ

- transverse momentum spectrum

$$T_{ave} = 221 \text{ MeV} \rightarrow$$

$$T_{ini} = 300 \text{ to } 600 \text{ MeV} \quad \tau_0 = 0.15 \text{ to } 0.6 \text{ fm}/c$$

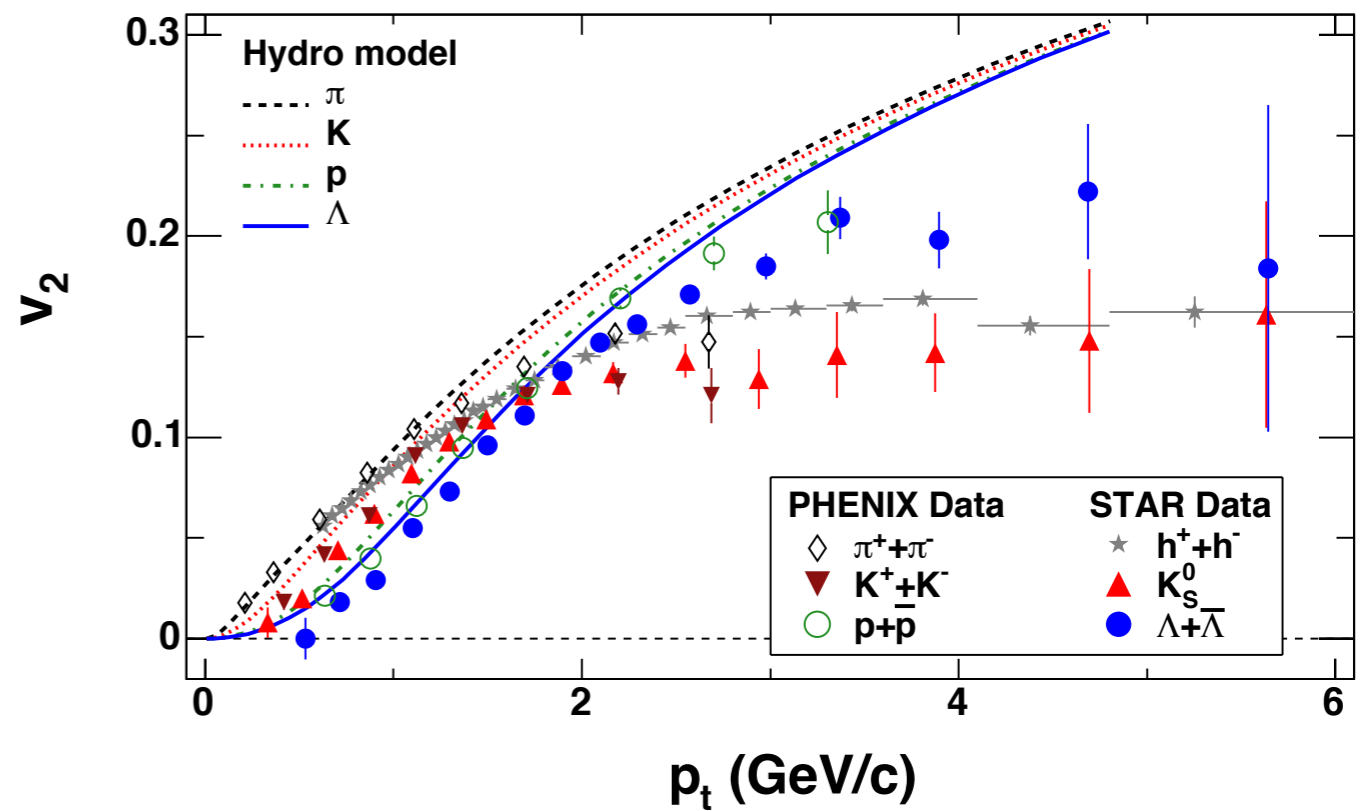


Azimuthal anisotropy

- Direct photons:

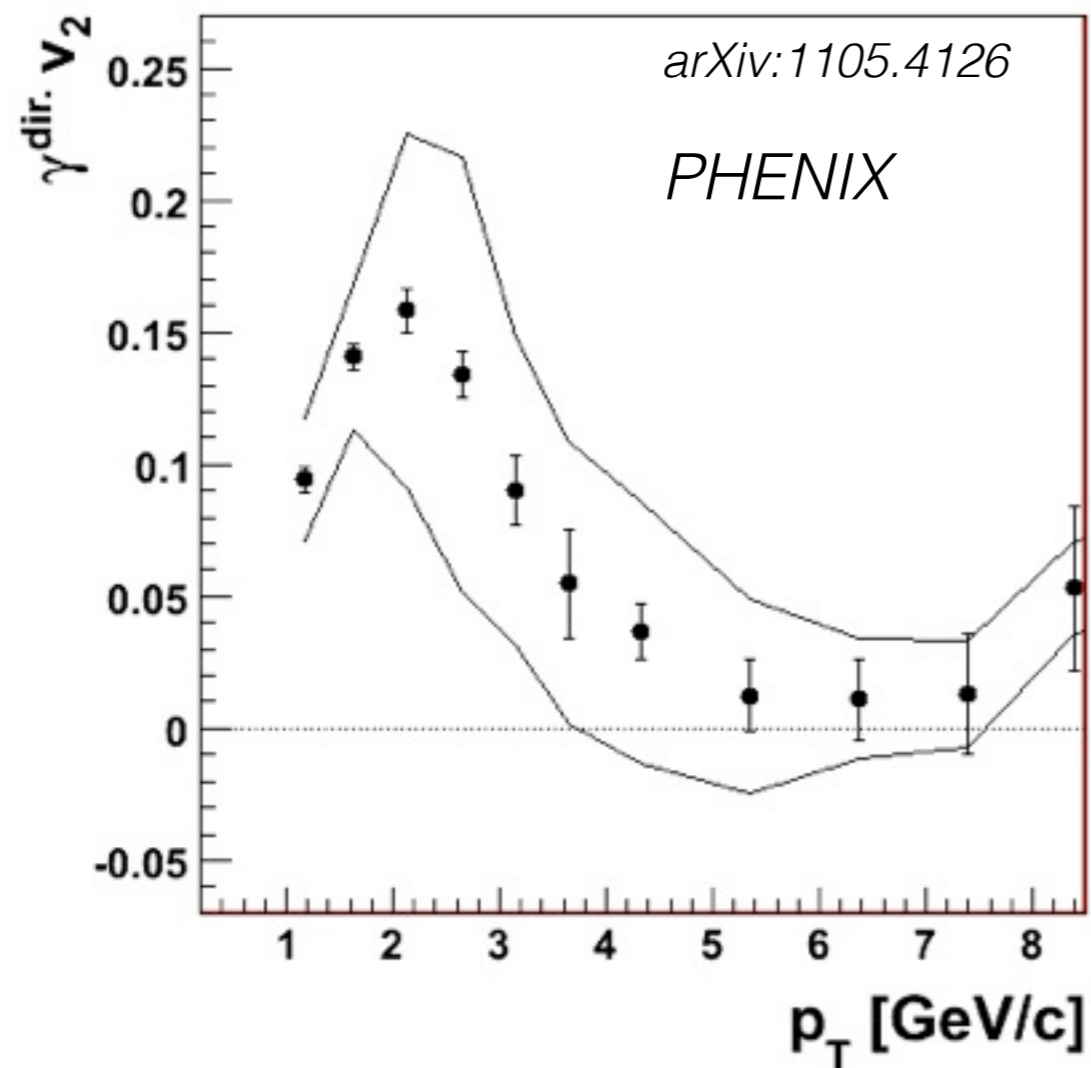
PHENIX

- Hadrons:

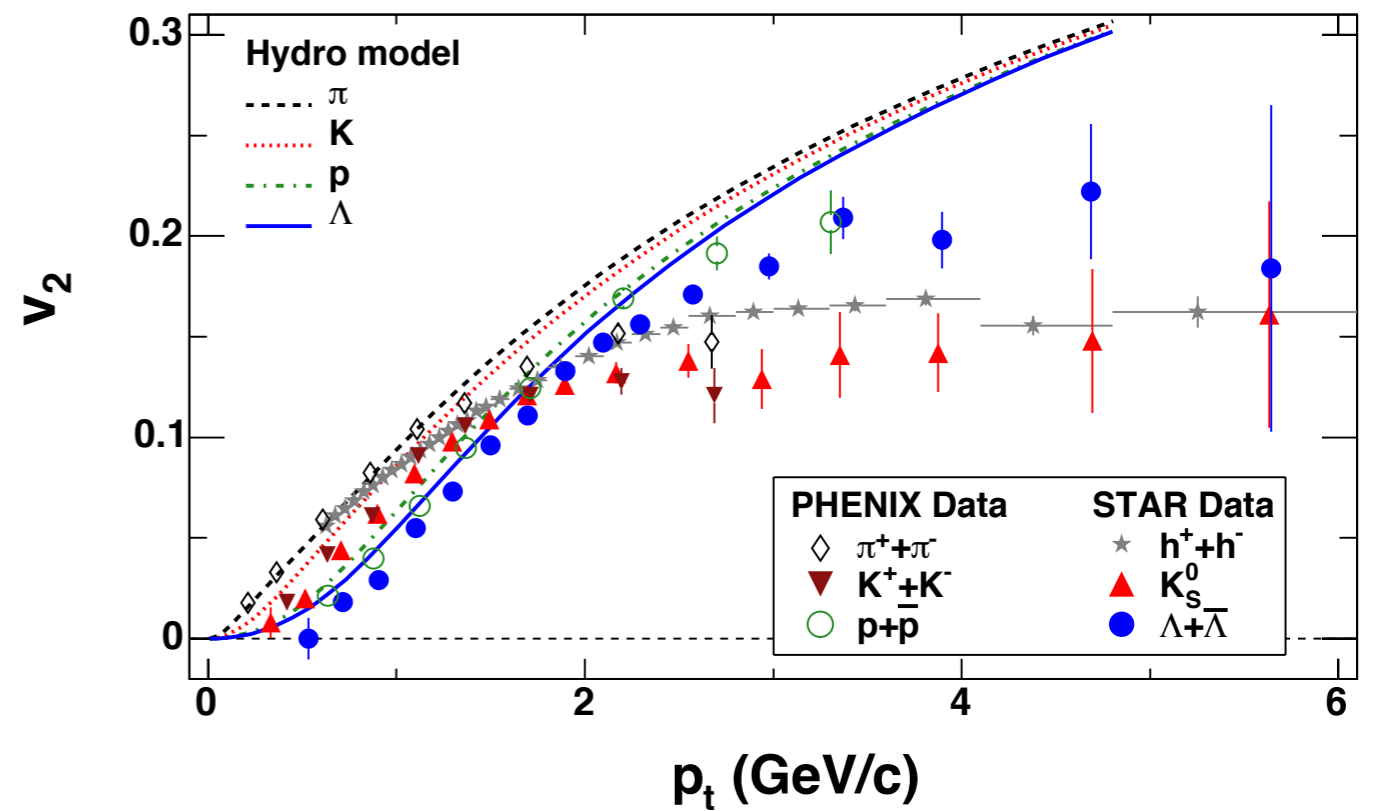


Azimuthal anisotropy

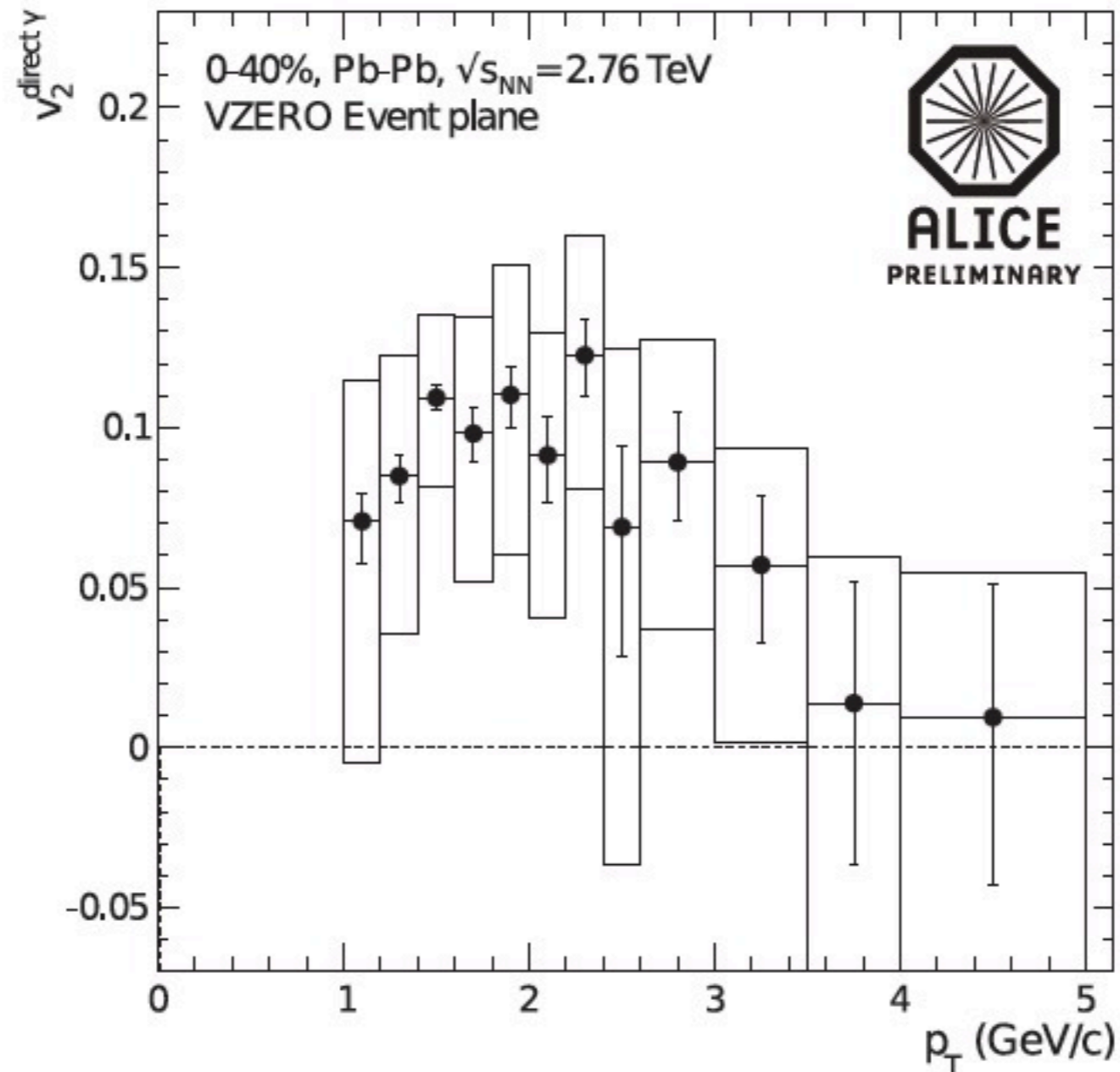
- Direct photons:



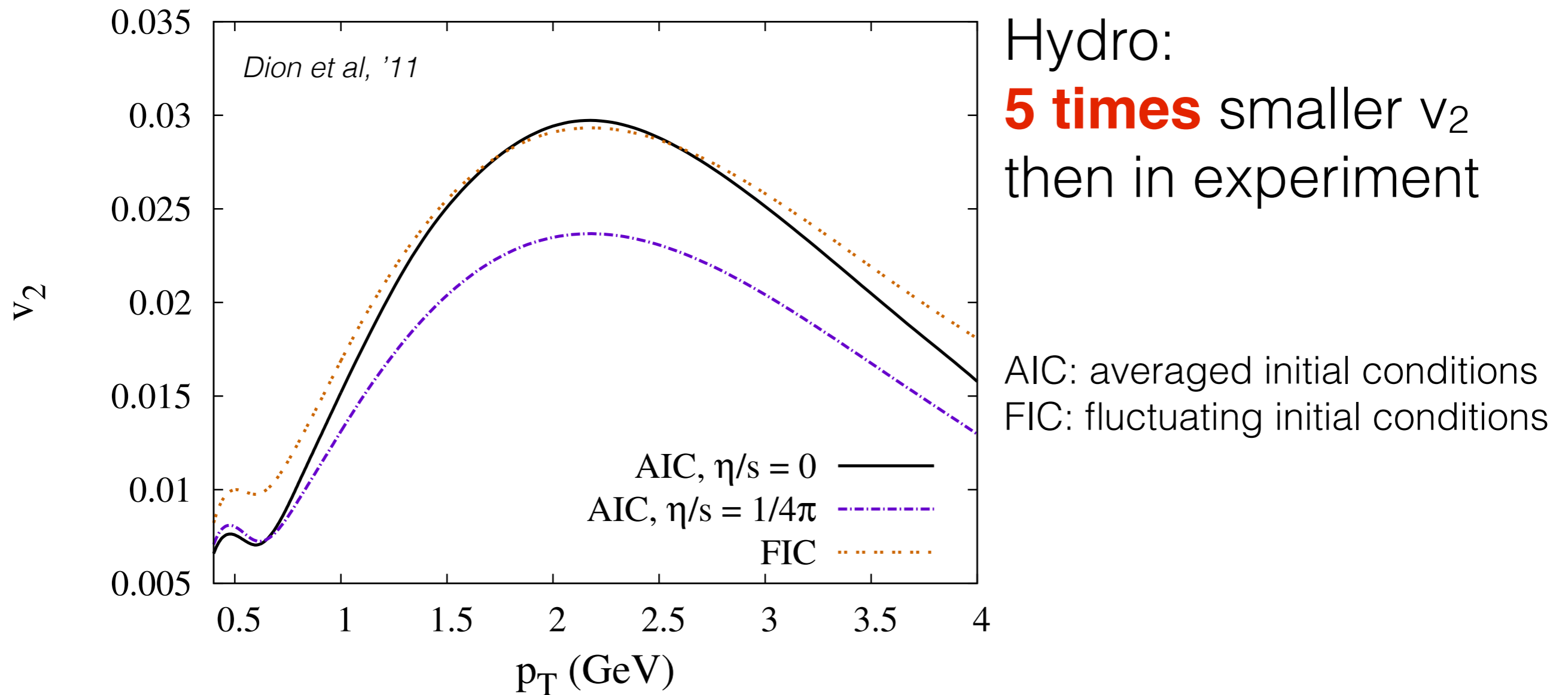
- Hadrons:



Azimuthal anisotropy: LHC



Theory: Hydrodynamics



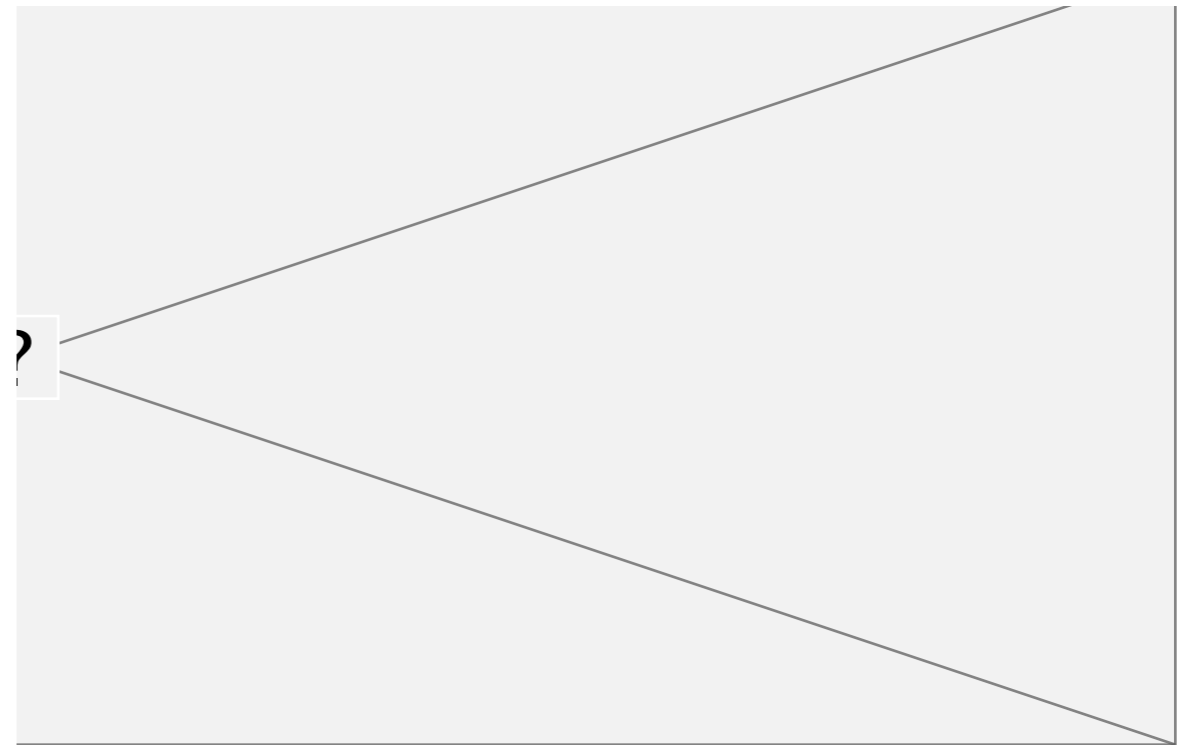
Large photon v_2 is difficult to explain with dominant QGP source

Another source of anisotropy?

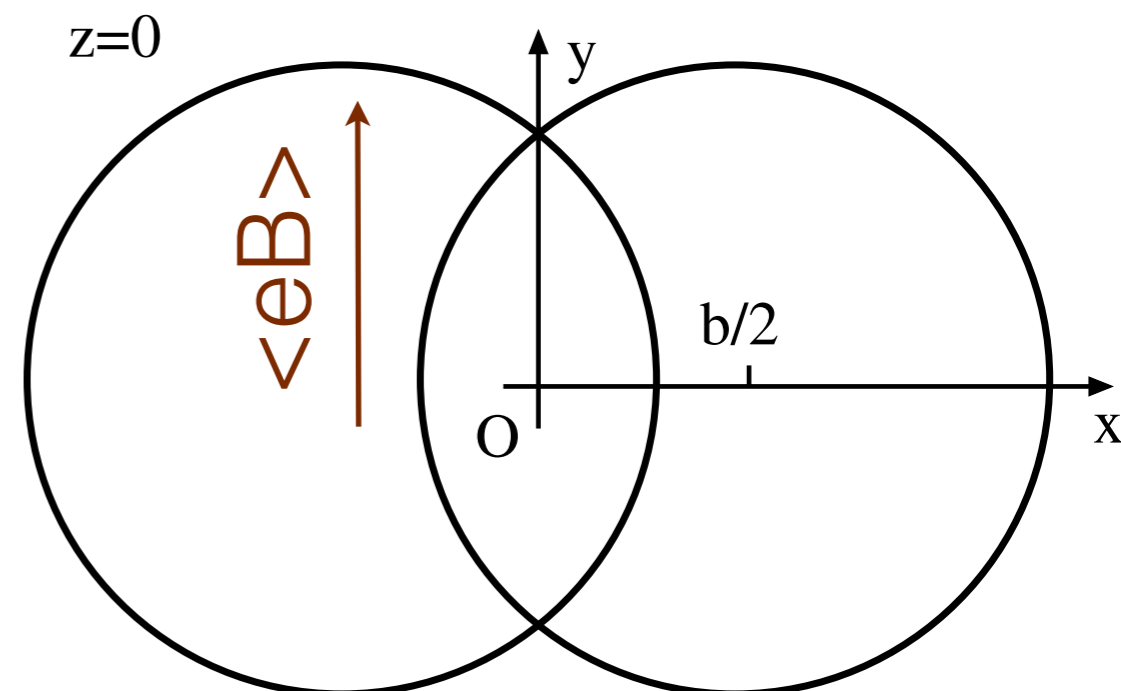
- anisotropy \neq flow!
- other sources for anisotropy not related to flow?!
- magnetic field?! Perfect candidate for anisotropic photon production.

Magnetic field in HIC I

- spectators form two currents

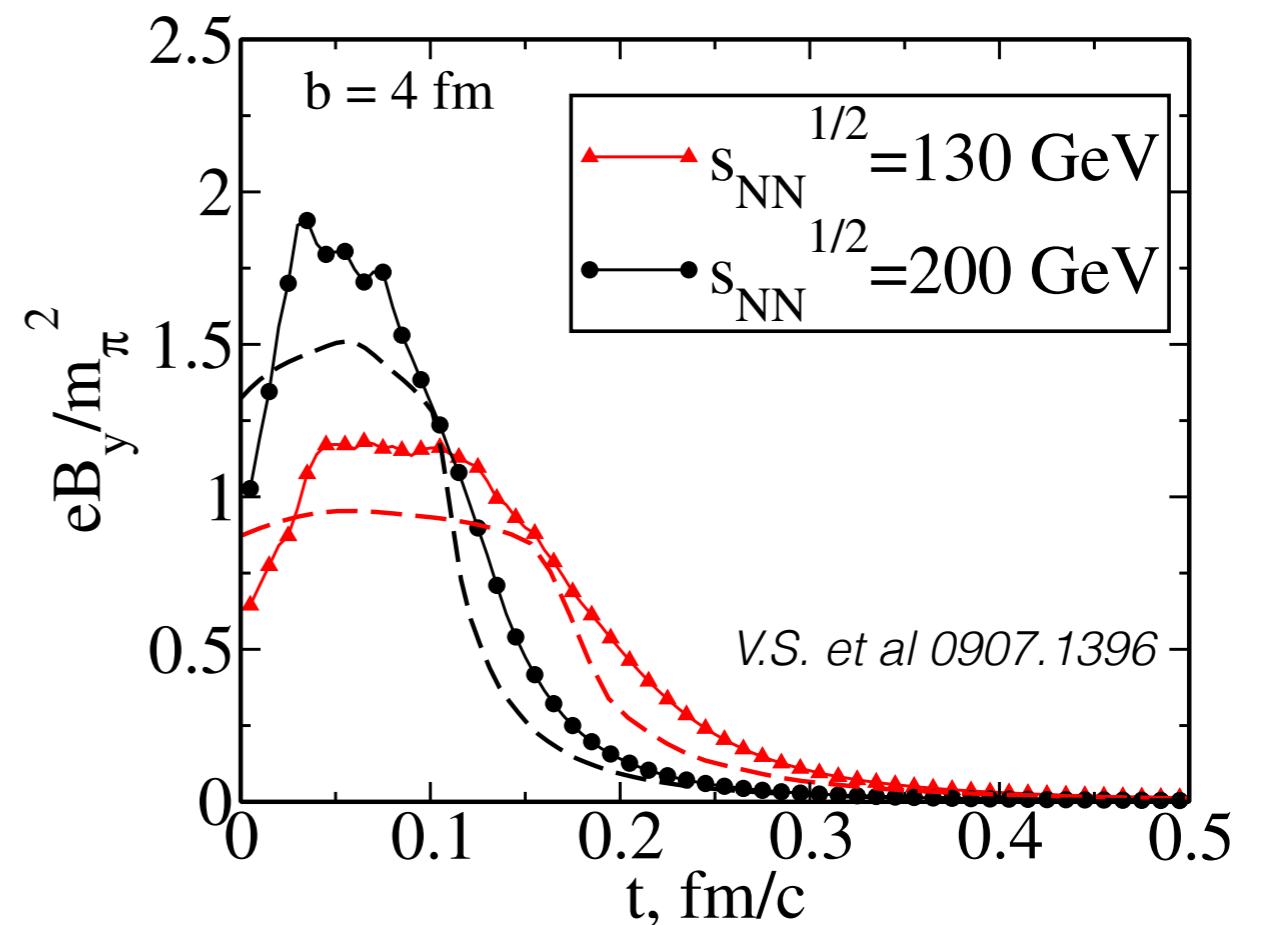


- resulting event
average magnetic field
 $\langle eB_y \rangle \sim m_\pi^2$ (out-plane)
 $\langle eB_x \rangle \sim 0$ (in-plane)

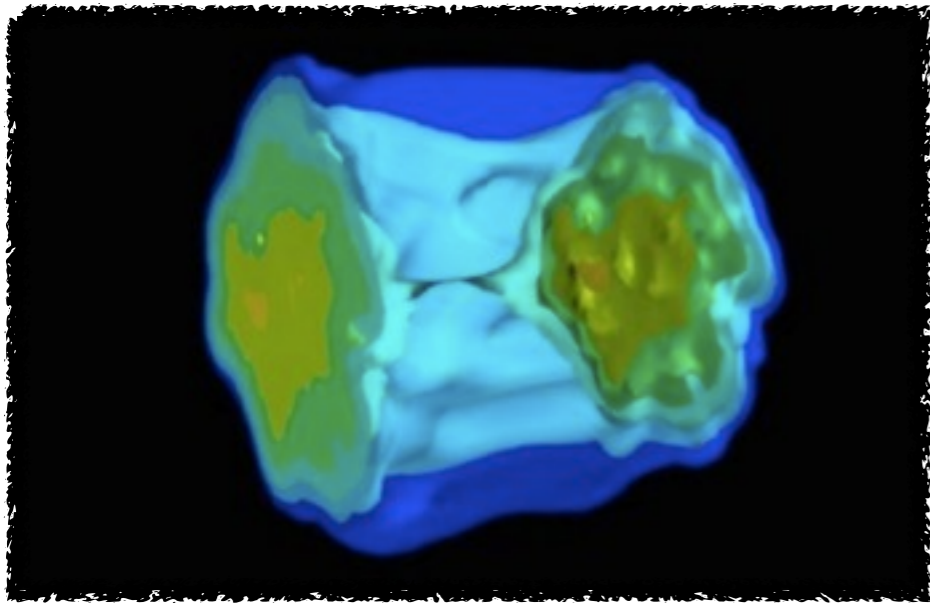


Magnetic field in HIC II

- maximal $eB \sim \sqrt{s}$
- maximum at $t_M \sim 1/\sqrt{s}$
- life time $t_{lt} \sim 1/\sqrt{s}$
- integral $\sim \text{const}$

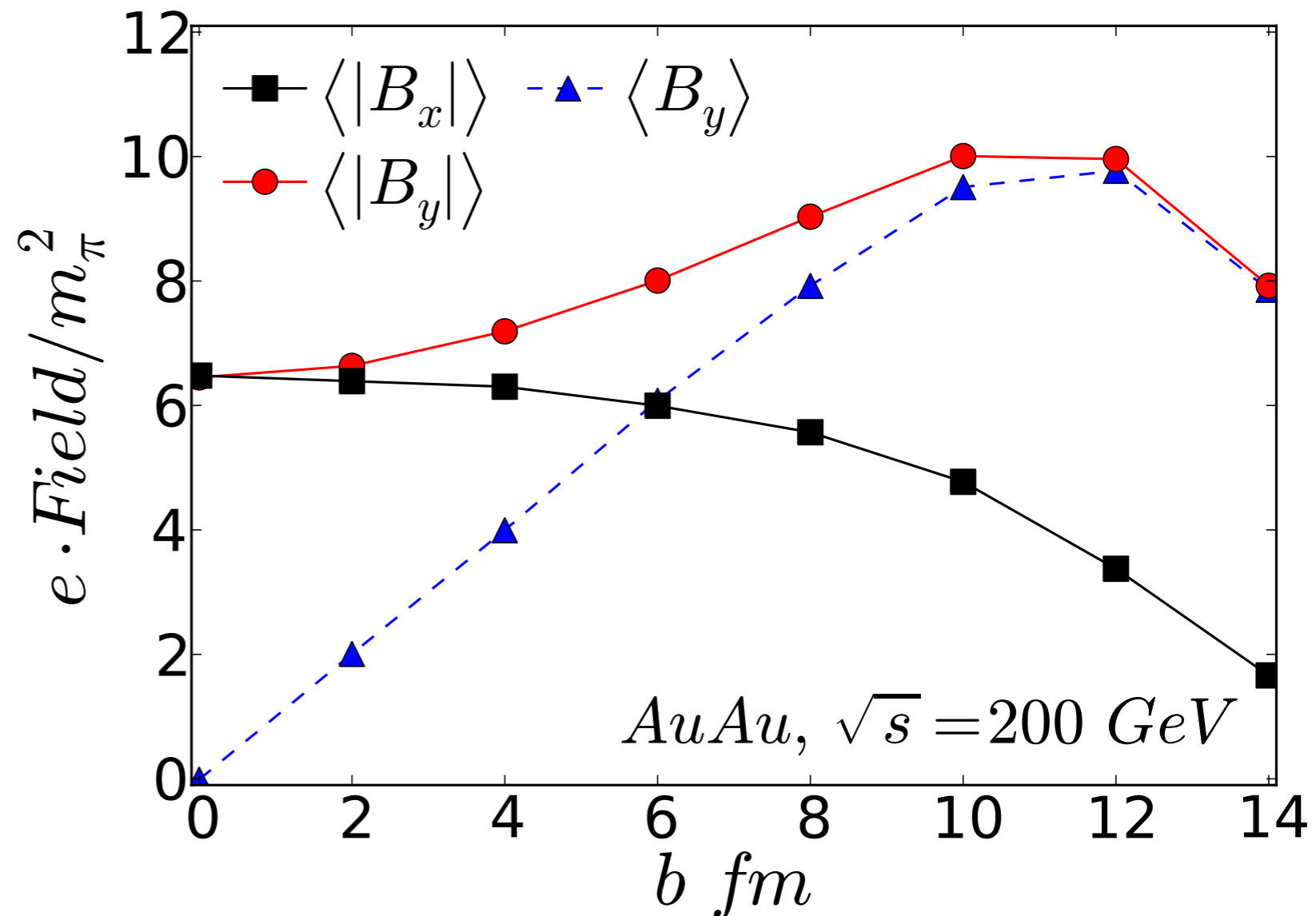


Magnetic field in HIC III



- lumpy distribution of electric charge in colliding nuclei results in nonzero randomly oriented magnetic field even in central collisions

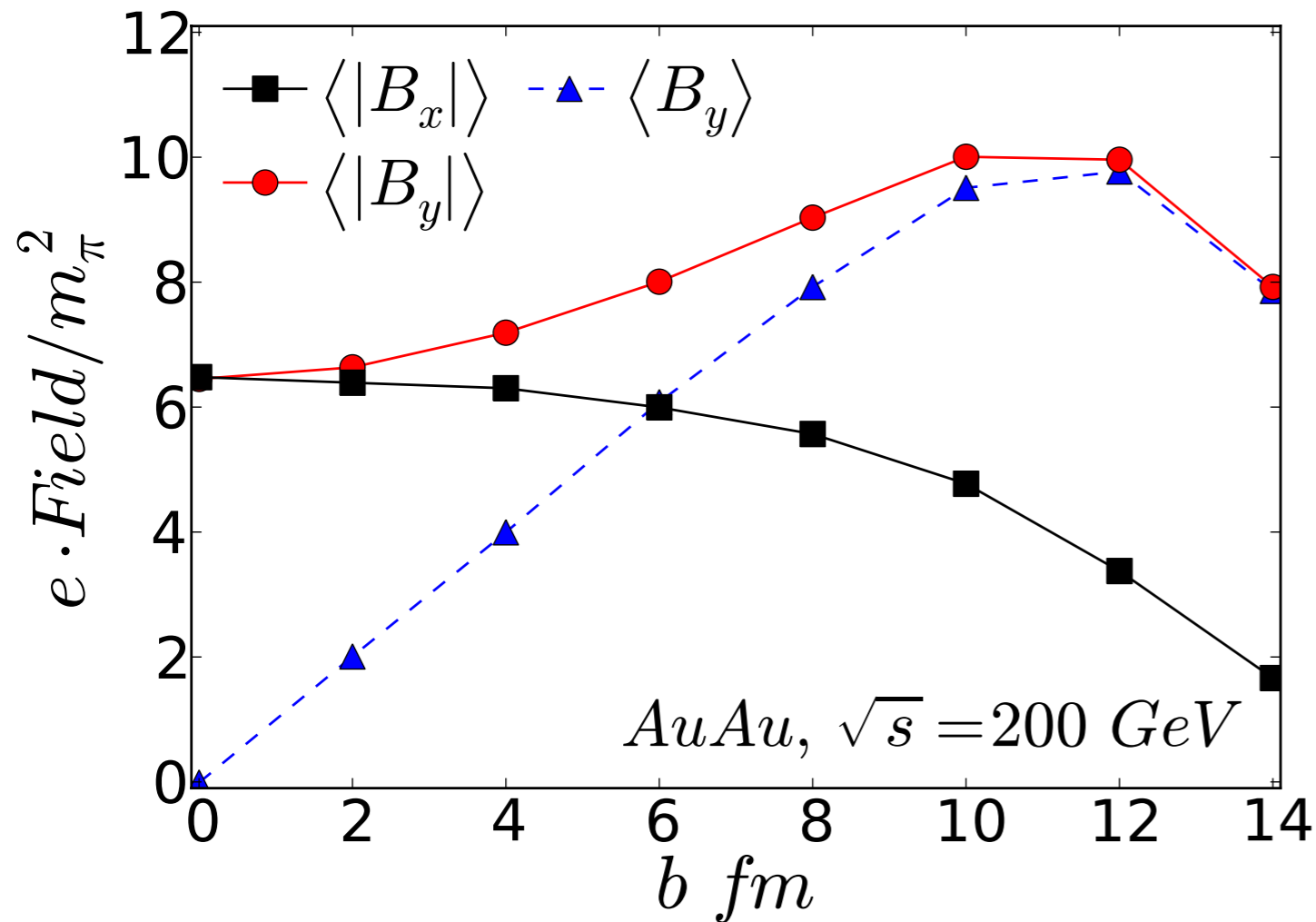
- fluctuations can play important role



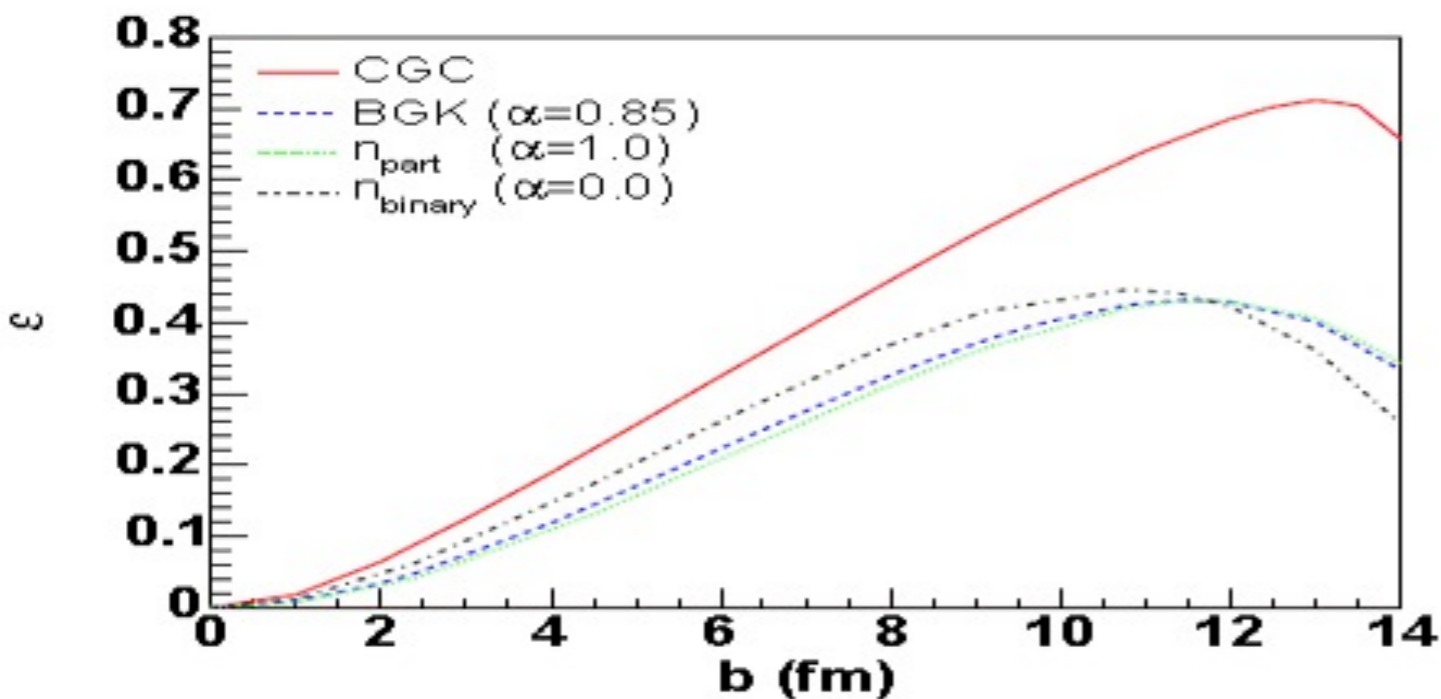
V.S. et al, 0907.1396;

A. Bzdak and V.S., 1111.1949

Magnetic field in HIC IV



- $\langle eB_y \rangle$ is linear as a function of impact parameter
- this is common feature of $\langle eB \rangle$ and eccentricity ϵ_2



Magnetic field in HIC V

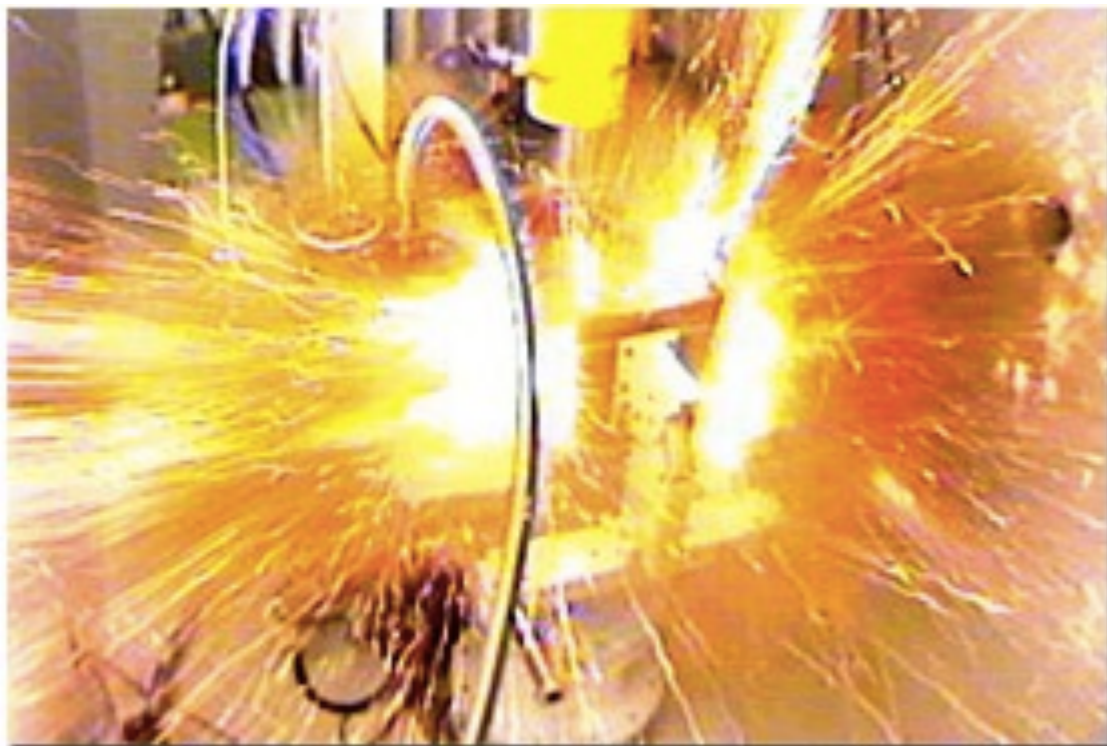
eB in HIC compared to

- Hybrid magnet at National High Magnetic field Lab
45 Tesla $\sim 4.5 \times 10^{-13} m_{\pi}^2$
- Pulsed magnets:
100 Tesla $\sim 10^{-12} m_{\pi}^2$

Vital Statistics	
Strength	45 tesla
Type	Hybrid
Bore size	32 mm (~1.25 inches)
Online since	December 1999
Cost	\$14.4 million
Weight	31,752 kg (35 tons)
Height	6.7 meters (22 feet)
Operating temperature	-271 ° C (-456 ° F)
Water used per minute	15,142 liters (4,000 gallons)
Power required	33 MW



Photo Credit: Larry Gordon



Watch an exploding pulsed magnet at work.

- Radio pulsars:
 10^{-6} - $10^{-5} m_{\pi}^2$
- Magnetars:
 10^{-4} - $10^{-3} m_{\pi}^2$



High eB ... So what?

Effects, that can be potentially observed:

- modification of QCD phase diagram
(not really, short lifetime of B)
- chiral magnetic effect
(sphaleron transition rate?!)
- chiral magnetic wave
(life time for magnetic field $\sim 4 \text{ fm}/c$)
- Photon splitting, and many other in the next talk
- **photon production!**

Photon production from eB

Several mechanisms:

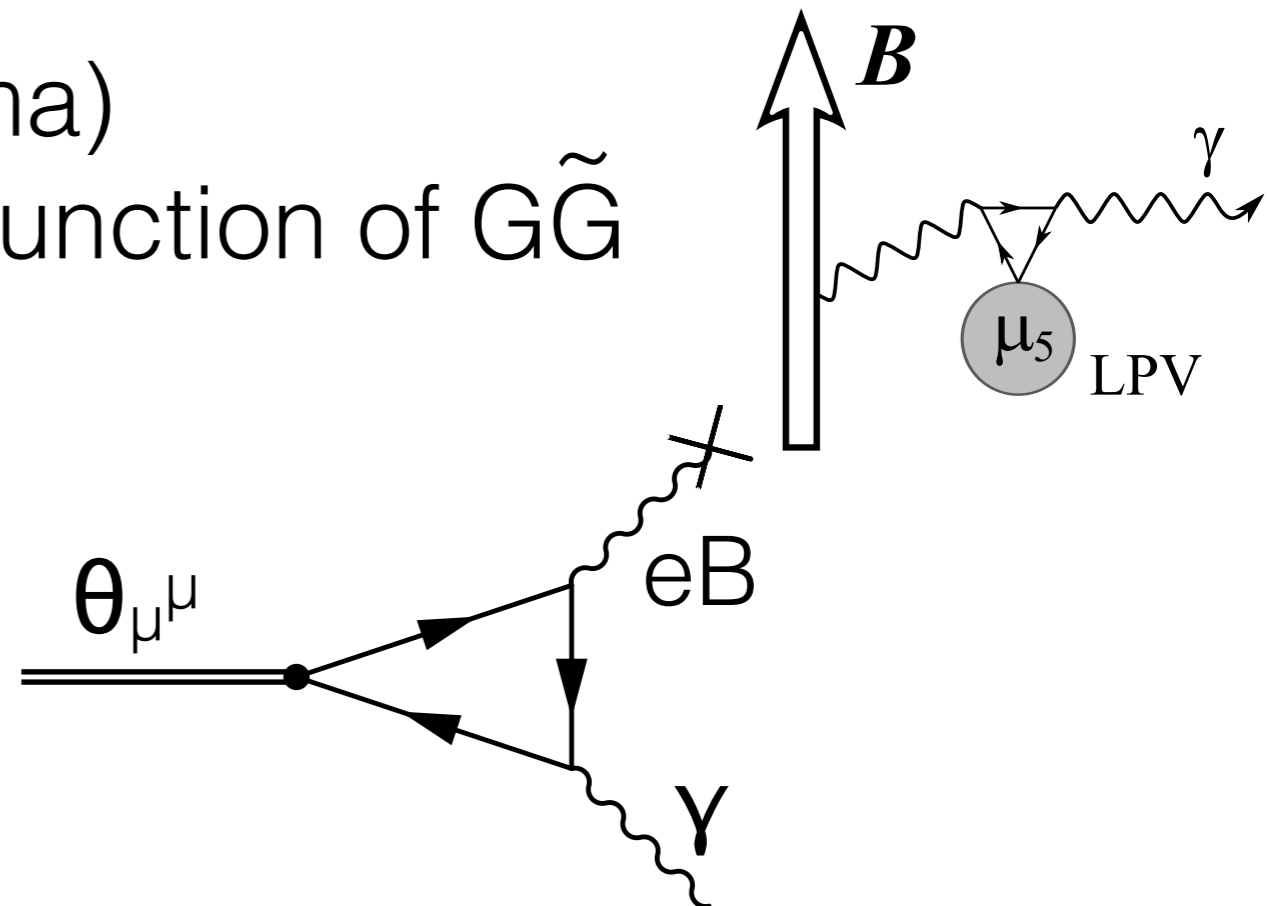
- synchrotron radiation of quarks in eB (K. Tuchin)
unknown: density and distribution function of quarks in early stage

R. Venugopalan and V.S. Quark production in Glasma

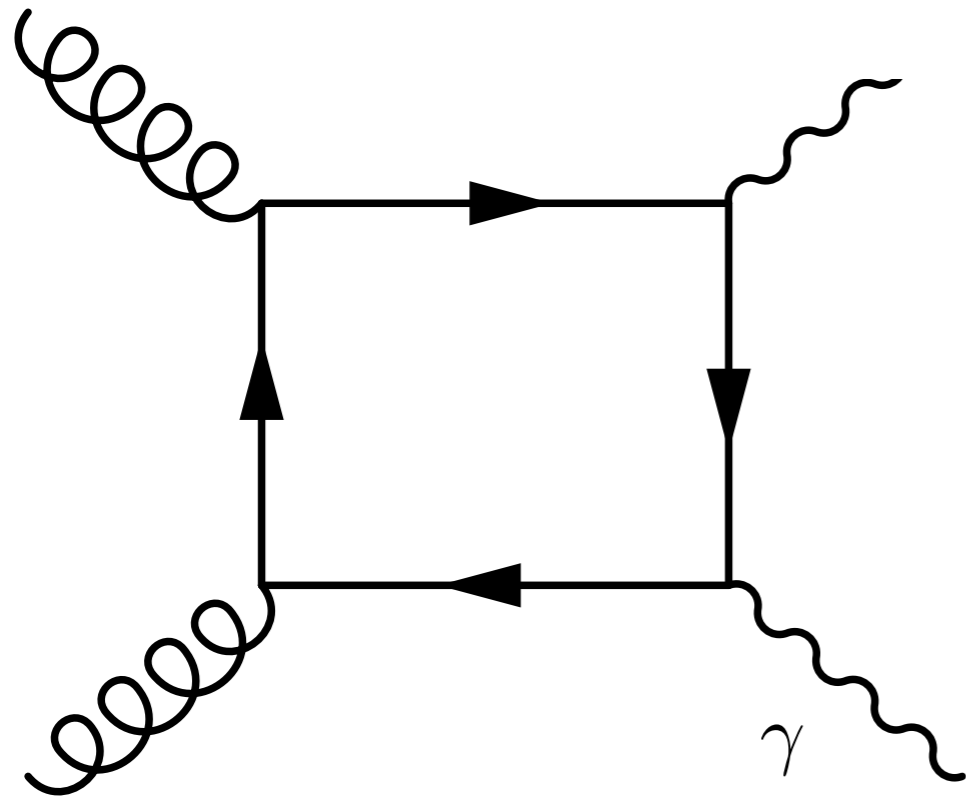
- axial anomaly (K. Fukushima)
unknown: μ_5 and spectral function of $G\tilde{G}$

G. Basar and D. Kharzeev

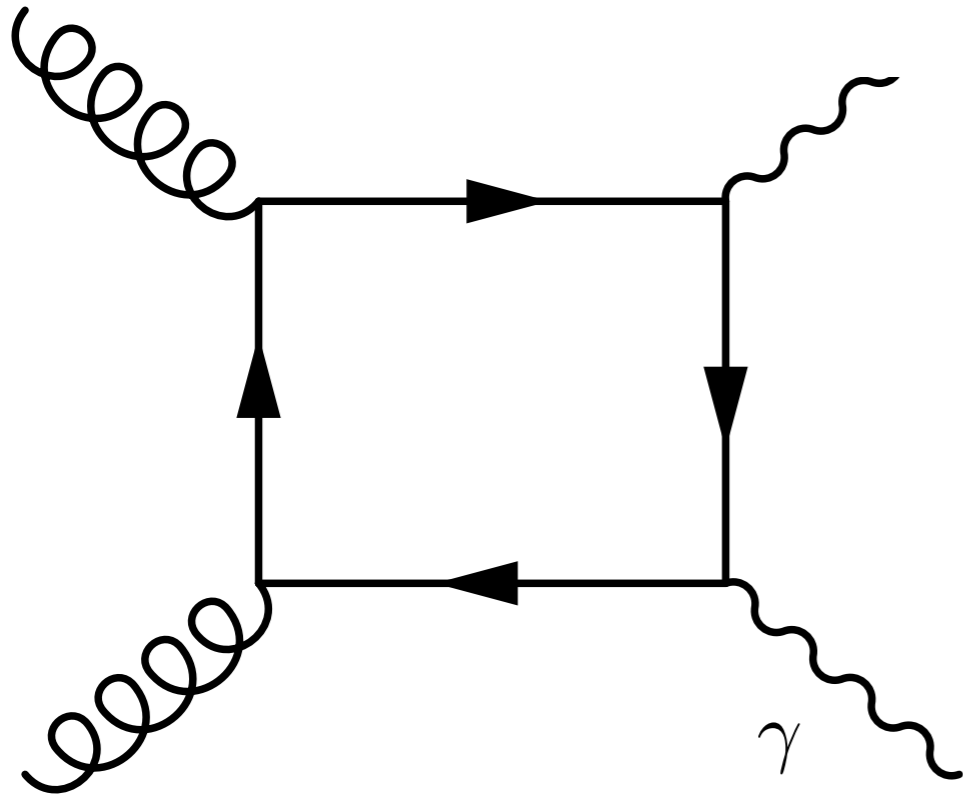
- **conformal anomaly**
unknown?!: bulk viscosity



Simple explanation

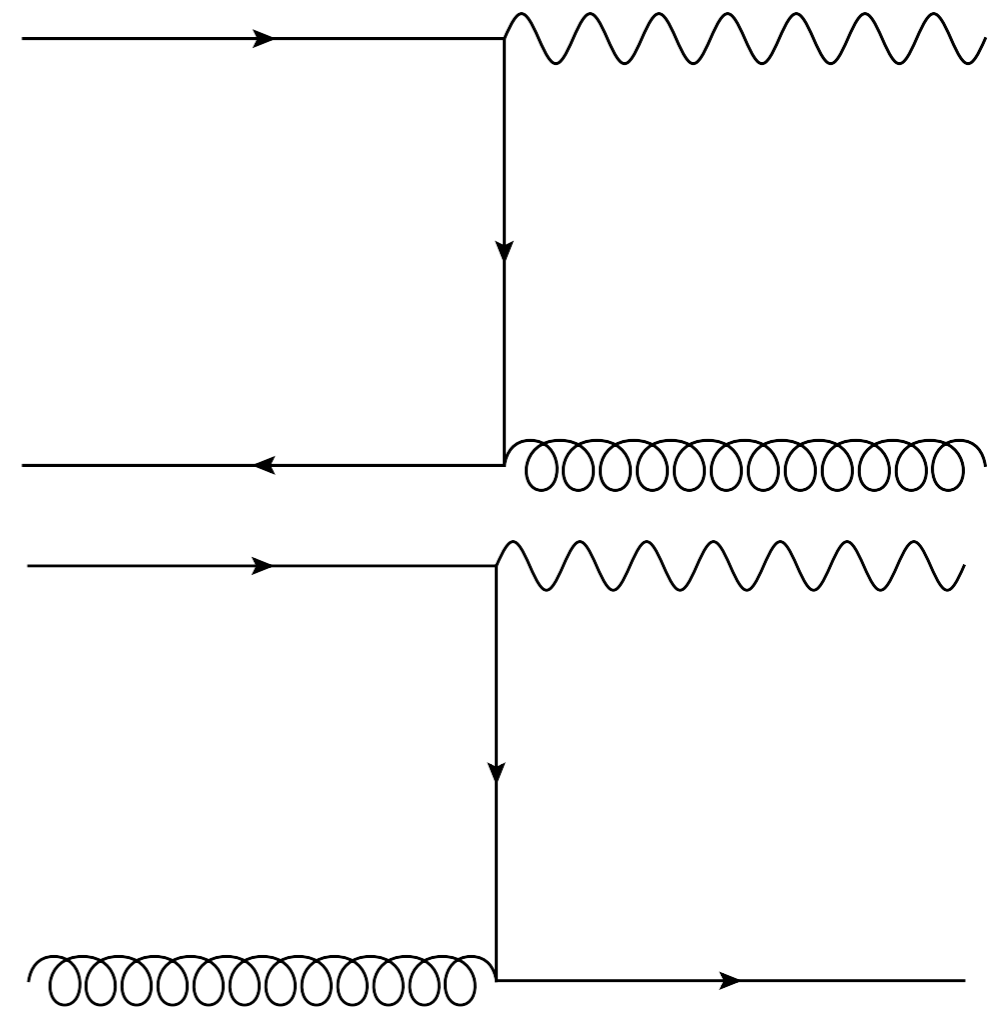
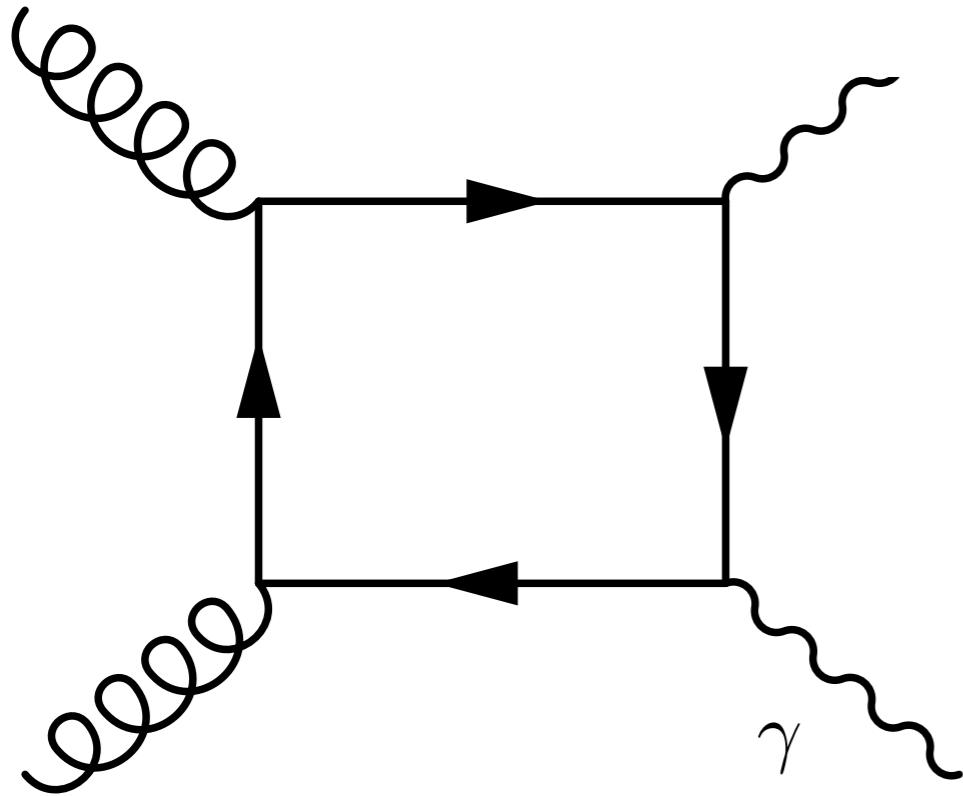


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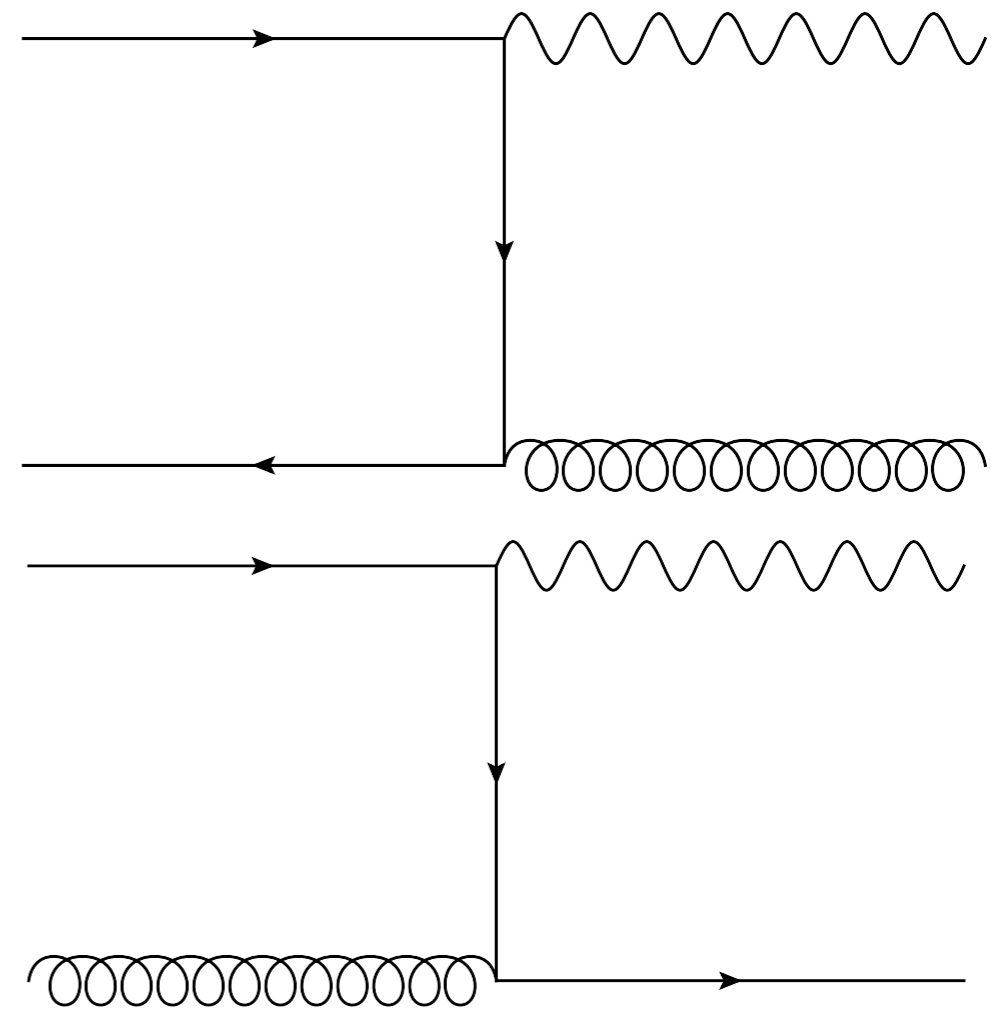
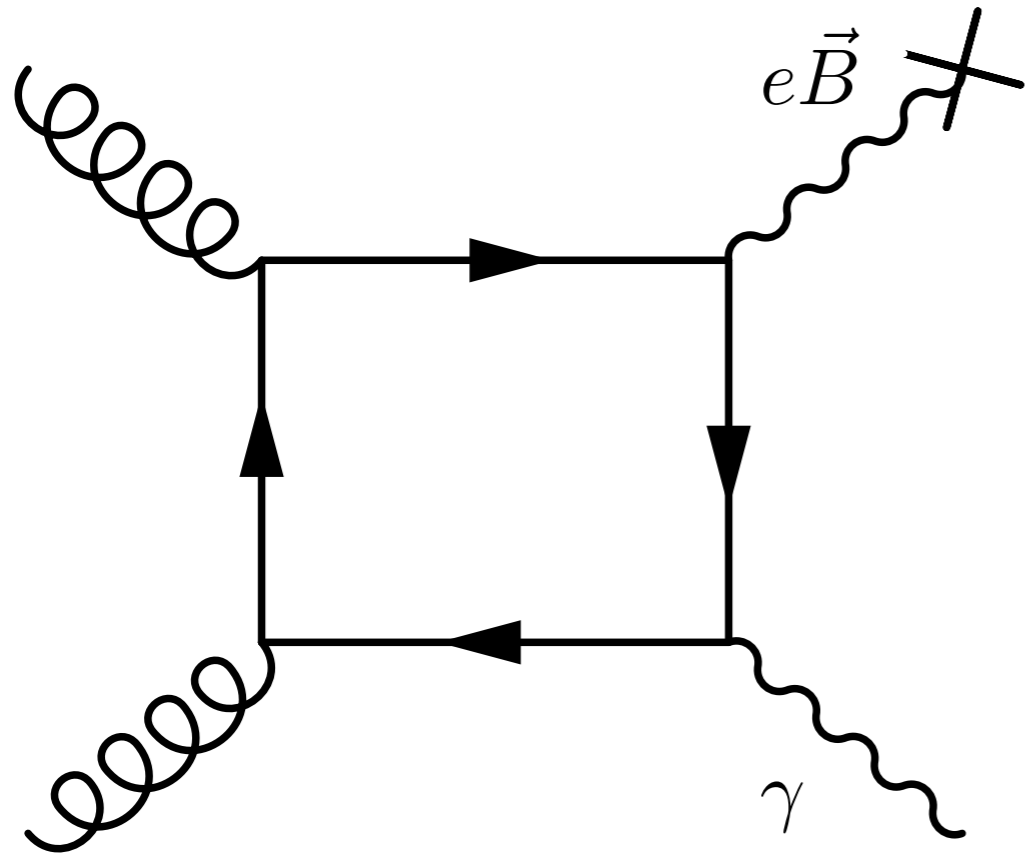
- Two photon production: $\alpha_s \alpha G^2 F^2$, $F^2 = F_{\mu\nu} F^{\mu\nu}$
thus rate $\sim \alpha^2$
- Replace one photon with eB
rate $\sim \alpha$

Simple explanation



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Conformal anomaly

- divergence of dilatation current

$$\partial^\mu S_\mu = \theta_\mu^\mu = \frac{\beta(g)}{2g} G^{\mu\nu a} G_{\mu\nu a} + \sum_q m_q [1 + \gamma_m(g)] \bar{q}q$$

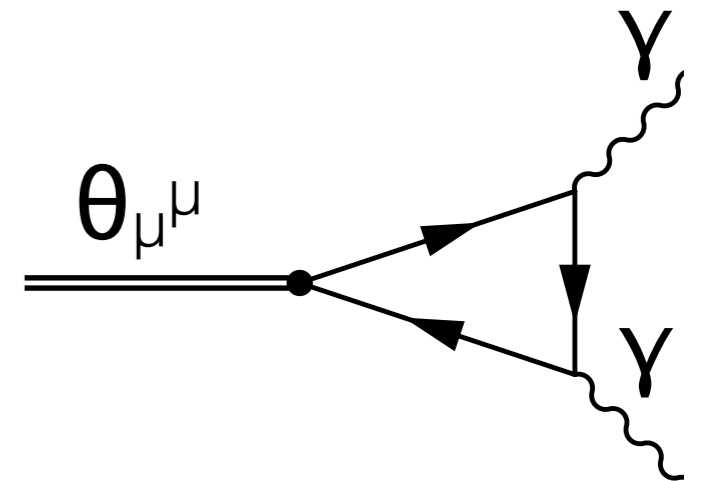
- color singlet states $\sigma \sim \theta_{\mu\mu}$ *Migdal, Shifman*

$$\langle 0 | S^\mu | \sigma \rangle = i q^\mu f_\sigma; \quad \langle 0 | \partial_\mu S^\mu | \sigma \rangle = m_\sigma^2 f_\sigma$$

- effective Lagrangian

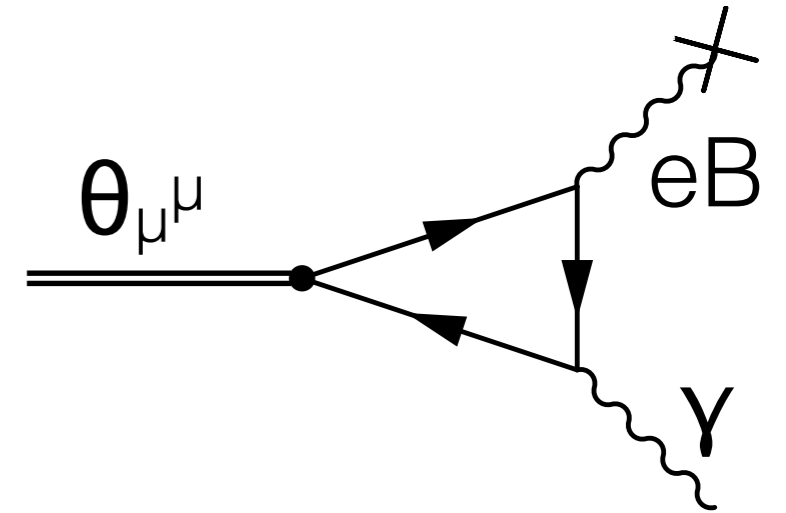
$$\mathcal{L}_{\sigma\gamma\gamma} = g_{\sigma\gamma\gamma} \sigma F_{\mu\nu} F^{\mu\nu}$$

- $g_{\sigma\gamma\gamma} \cong 0.02 \text{ GeV}^{-1}$ *Ellis and Lanik; Crewther; Chanowitz*

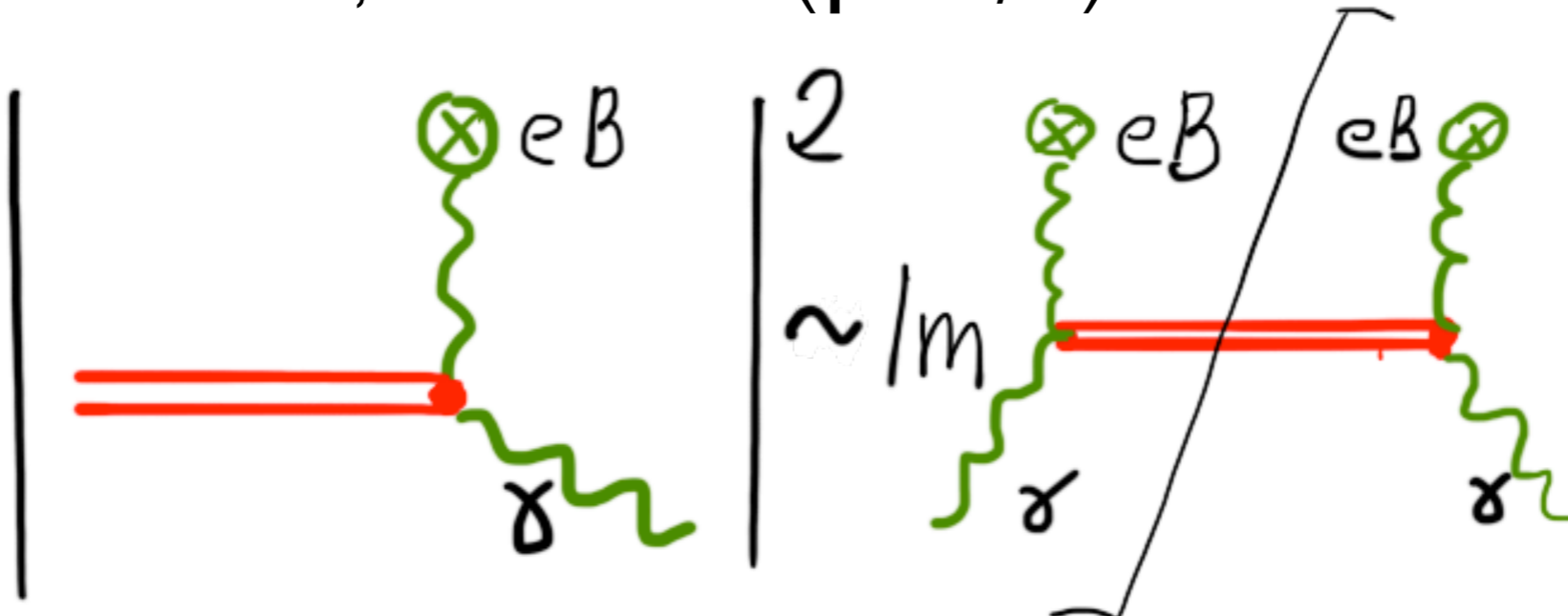


Photon production rate

- one of the photons: classical field eB



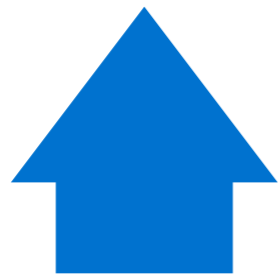
- production rate, as usual ($\beta=1/T$):



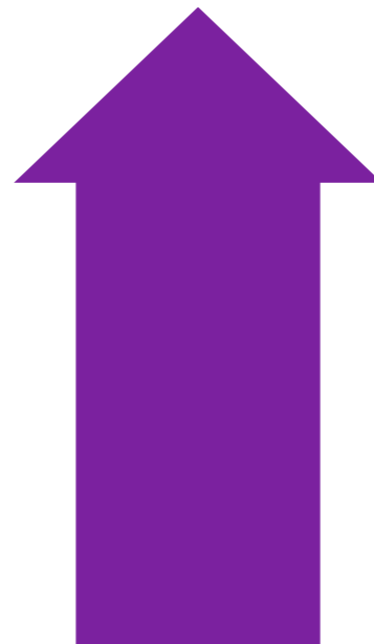
$$q_0 \frac{d\Gamma_B}{d^3q} = 2 \left(\frac{g_{\sigma\gamma\gamma}}{\pi f_\sigma m_\sigma^2} \right)^2 \times \frac{(B_y^2 - B_x^2)q_x^2 + q_\perp^2 B_x^2}{\exp(\beta q_0) - 1} \rho_\theta(q_0 = |\vec{q}|).$$

Final answer

$$q_0 \frac{d\Gamma_B}{d^3q} = 2 \left(\frac{g_{\sigma\gamma\gamma}}{\pi f_\sigma m_\sigma^2} \right)^2 \times \frac{(B_y^2 - B_x^2) q_x^2 + q_\perp^2 B_x^2}{\exp(\beta q_0) - 1} \rho_\theta(q_0 = |\vec{q}|).$$



Numerical coefficient;
constrained by hadronic
observables



Momentum dependence;
 $\beta = 1/T$; if e-b-e fluc. of
magnetic field are
neglected:

$$\frac{B_y^2 q_x^2}{\exp(\beta q_0) - 1}$$



Spectral function
for G^2 , or trace of
energy momentum
tensor

Spectral function of Θ_{μ}^{μ}

- hydrodynamic approximation

$$\rho_{\theta}(q_0, \vec{q}) = \frac{1}{\pi} \text{Im}[G_R^{\mu\mu, \nu\nu}(q_0, \vec{q})] = 9q_0 \frac{\zeta}{\pi} + \frac{9}{\pi} (\epsilon + p) \left(\frac{1}{3} - c_s^2 \right)^2 \frac{q_0 \Gamma_s \vec{q}^4}{(q_0^2 - c_s^2 \vec{q}^2)^2 + (q_0 \Gamma_s \vec{q}^2)^2}$$



bulk viscosity



sound peak

- real photons, sound peak does not contribute:

$$\rho_{\theta}(q_0, \vec{q}) \approx 9q_0 \frac{\zeta}{\pi}$$

$G\tilde{G}$

- Similar calculations can be done for $F\tilde{F}$ $G\tilde{G}$
- Spectral function $G\tilde{G}$ in hydro approximation is defined by sphaleron transition rate and was calculated in pQCD and AdS/CFT.

Bulk viscosity

- first principle Lattice QCD:

H. Meyer SU(3) Yang Mills (YM)

However, there are issues.

- approximations:

$$\zeta = C_\zeta \eta (1/3 - c_s^2)^2 \quad (\text{vs ADS/QCD } \zeta \cong 2 \eta (1/3 - c_s^2))$$

$C_\zeta = 15$ in relaxation time appr. (S. Weinberg '71)

$C_\zeta = 45$ in NLO SU(3) YM (K. Dusling and T. Schafer '11)

$C_\zeta = 2.5-5$ phenomenological constraints

in this talk: conservative $C_\zeta = 2.5-5$

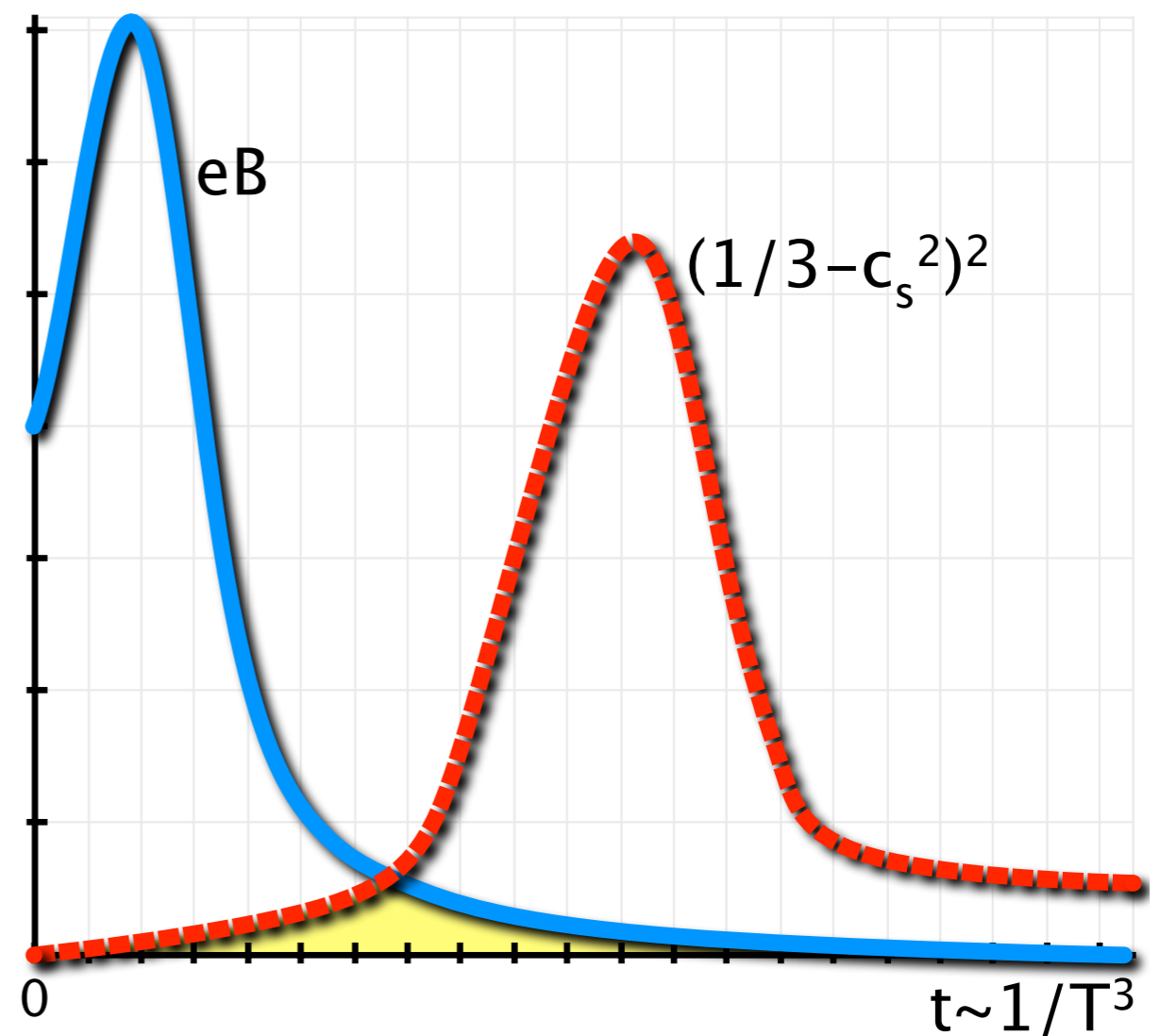
- also conservative $\eta/s = 1/(4\pi)$. Entropy, s , from matrix model fitted to YM SU(3) (R. Pisarski)

So one would expect...

$$q_0 \frac{d\Gamma_B}{d^3q} = 2 \left(\frac{g_{\sigma\gamma\gamma}}{\pi f_\sigma m_\sigma^2} \right)^2 \times \frac{(B_y^2 - B_x^2)q_x^2 + q_\perp^2 B_x^2}{\exp(\beta q_0) - 1} \rho_\theta(q_0 = |\vec{q}|).$$

$$\rho_\theta(q_0, \vec{q}) \approx 9q_0 \frac{\zeta}{\pi} \quad \zeta = C_\zeta \eta (1/3 - c_s^2)^2$$

- negligible contribution from this mechanism?!
eB is non-zero at early stage where $(1/3 - c_s^2)^2$ is small



However

- rate is proportional to $(1-c_s^2)^2 T^3 (eB)^2$

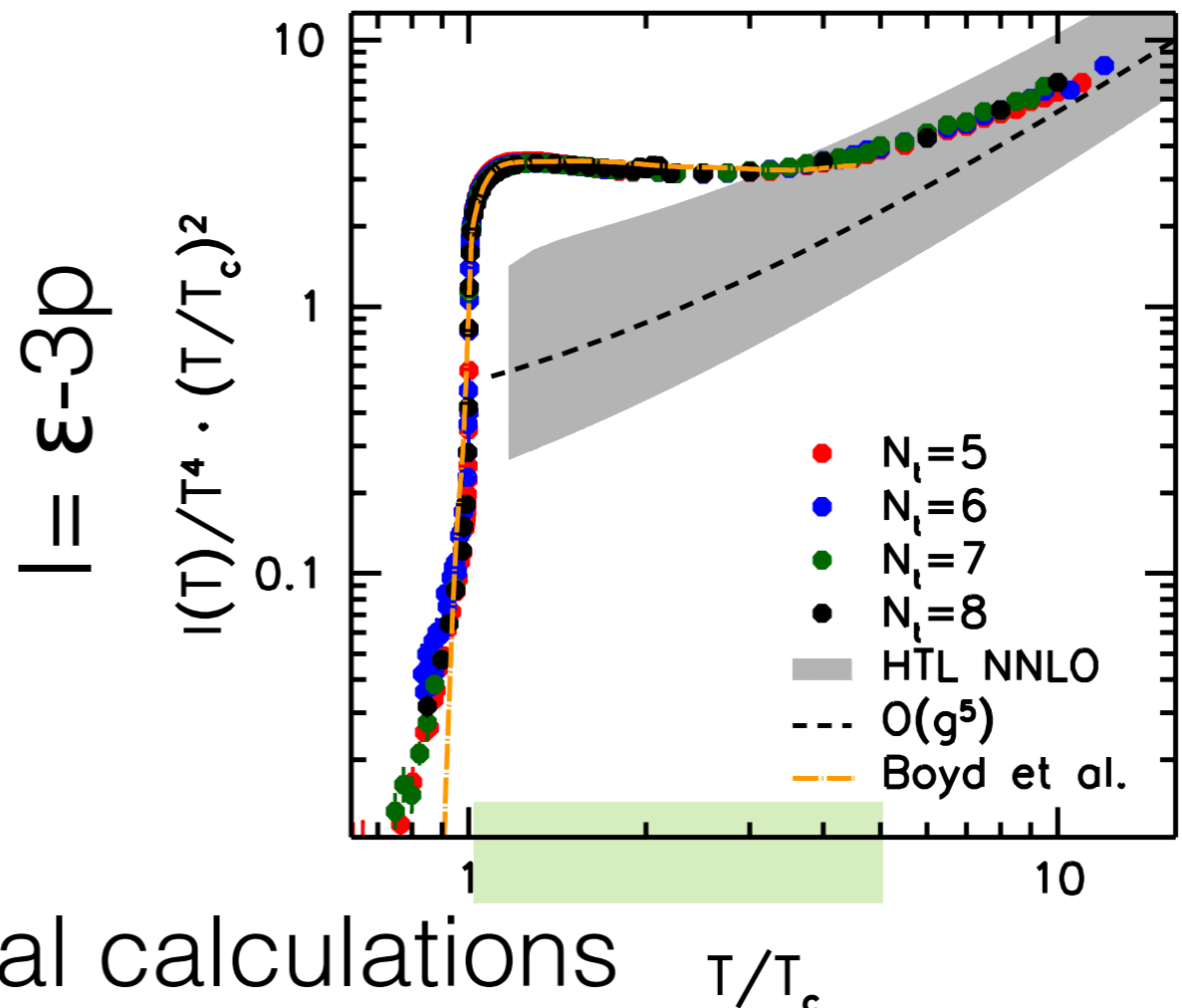
rough estimate at early stage

$(1-c_s^2) \sim (\epsilon - 3p)/T^4 \sim (\text{from LQCD}) \sim \mathbf{1/T^2}$ (talks by R. Pisarski)

$(1-c_s^2)^2 T^3 \sim \mathbf{1/T} \sim (\text{Bjorken expansion}) \sim t^{1/3}$

while $eB \sim 1/(t^2 + \text{const})$

stringy?
fuzzy bag?
quasiparticles?
monopoles?



- rigorous answer: numerical calculations

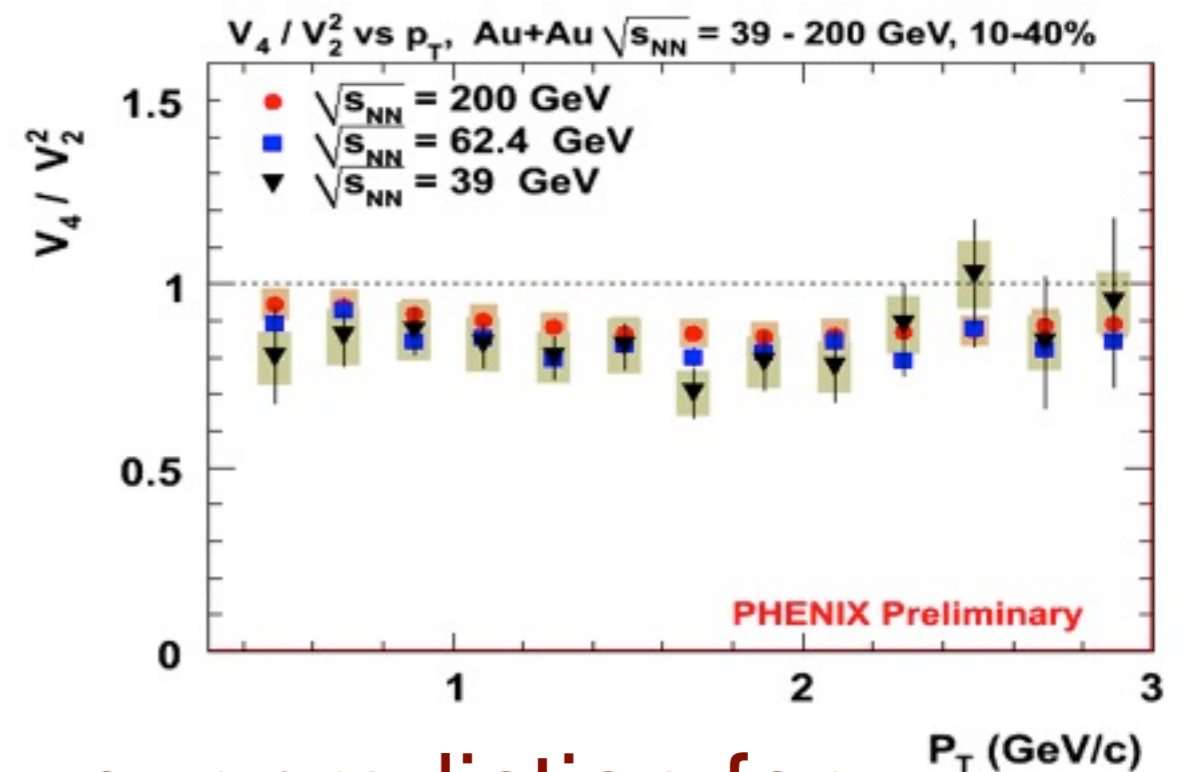
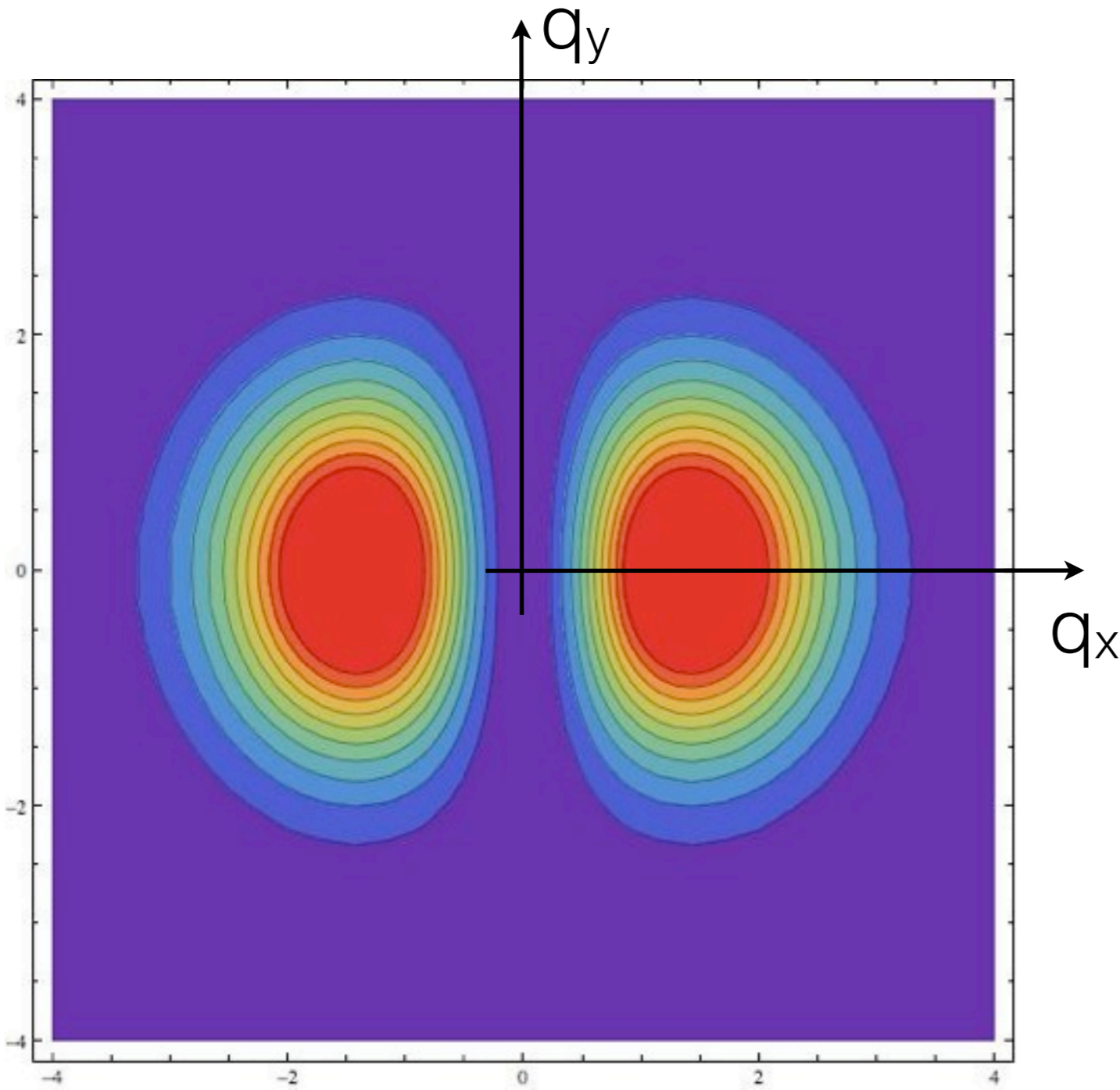
Anisotropy of production rate

- in this mechanism:

$$dN/d\boldsymbol{\varphi} \sim q_x^2 = q_T^2 \cos^2(\boldsymbol{\varphi}) = q_T^2 [1 + \cos(2\boldsymbol{\varphi})]/2$$

consequently:

- non-zero v_2
- small v_n , $n=4, \dots$
in contrast to hadronic v_4
PHENIX: $v_4/v_2^2 \sim 1$



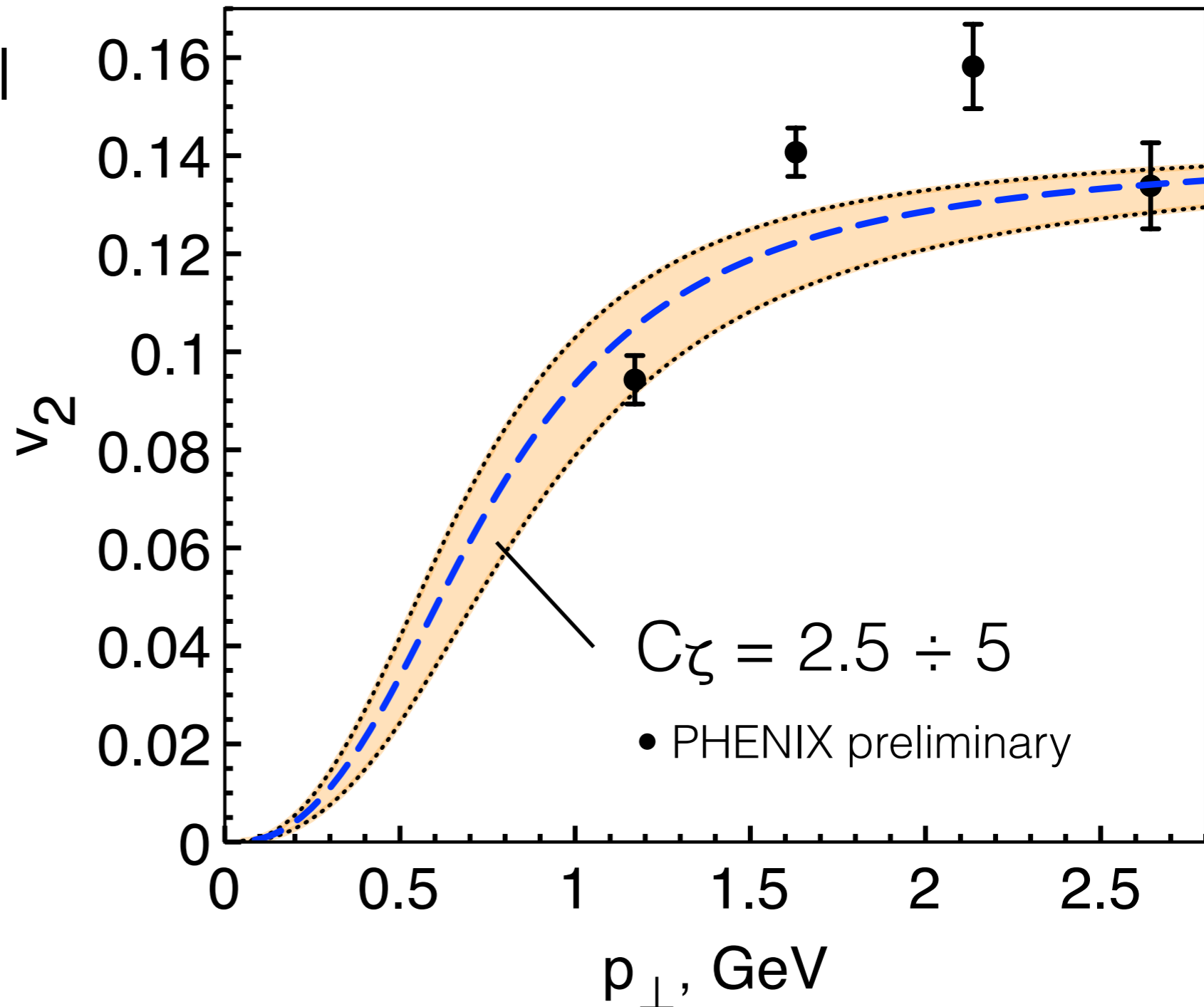
our prediction for
photons: $v_4/v_2^2 \ll 1$

Numerical calculations: v_2

- ingredients: thermal photons and photons from conformal anomaly + eB
- significant contribution to v_2
- higher p_\perp : prompt photons

Numerical calculations: v_2

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- significant contribution to v_2
- higher p_{\perp} : prompt photons



Other parameters

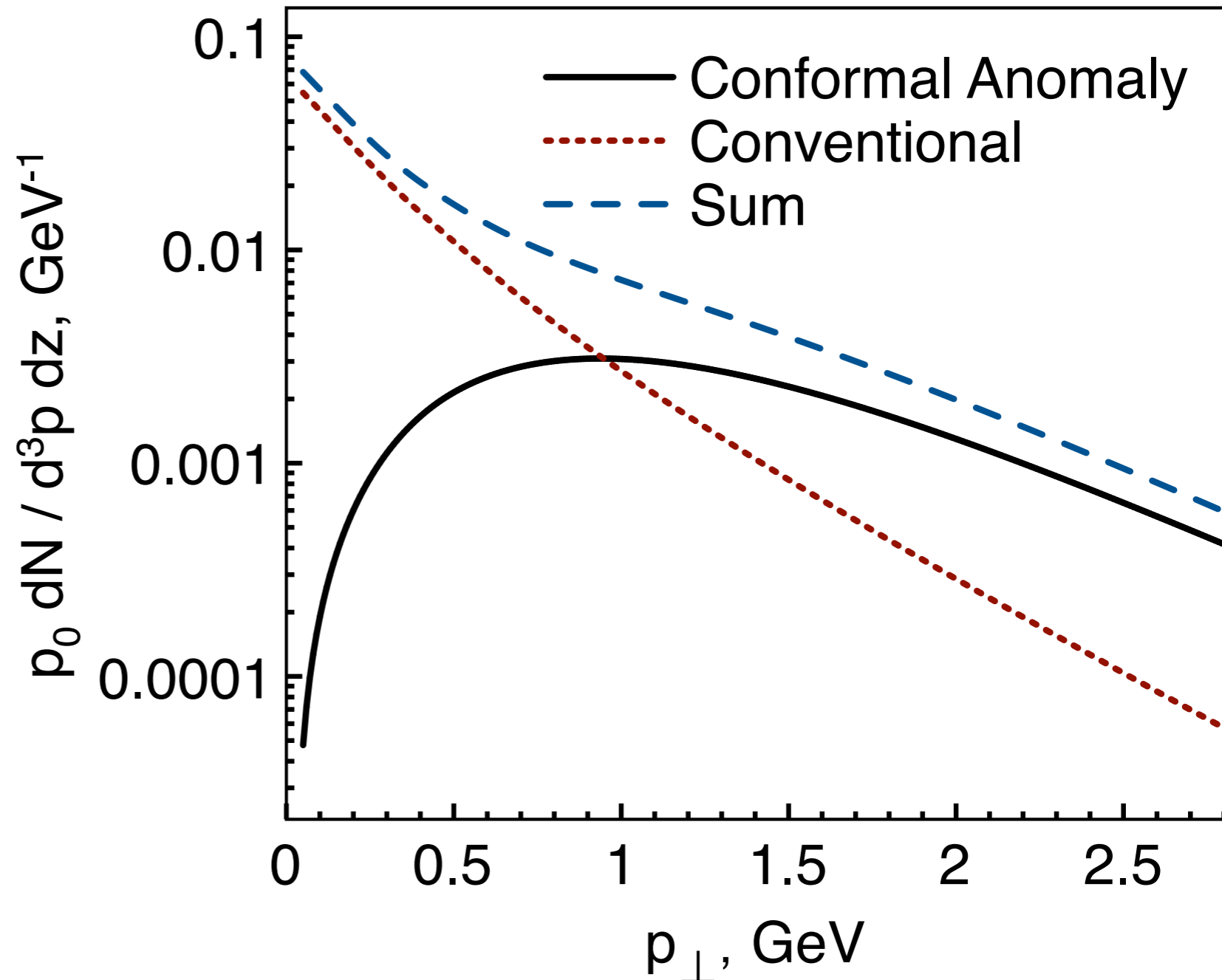
- initial temperature $T = 400 \text{ MeV}$
- initial time $\tau = 0.1 \text{ fm}/c$ (no need for complete equilibrium, is to be discussed later)
- Bjorken expansion for T
- electromagnetic field from spectators only (with fluctuations taken into account). Possible induced magnetic field will only enhance production via this mechanism

Transverse momentum spectra

- conformal anomaly:
 $dN/dp_{\perp} \sim \mathbf{p}_{\perp}^2 / [\exp(p_0/T) - 1]$
similar to effect of direct flow
- higher than thermal photons for $p_{\perp} > 1$ GeV
- higher p_{\perp} : prompt photons

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 $dN/dp_{\perp} \sim \mathbf{p}_{\perp}^2 / [\exp(p_0/T) - 1]$
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Experimental tests I

- **1) magnetic field B** is generated mostly by spectators thus, B is defined by centrality (measured by ZDC), reaction plane
- **2) hadronic flow:** initial eccentricity ϵ
 ϵ depends on details of hadron interaction (Glauber fluctuations, fluctuations of energy deposition); participant plane
- so switch off either **1)** or **2)**

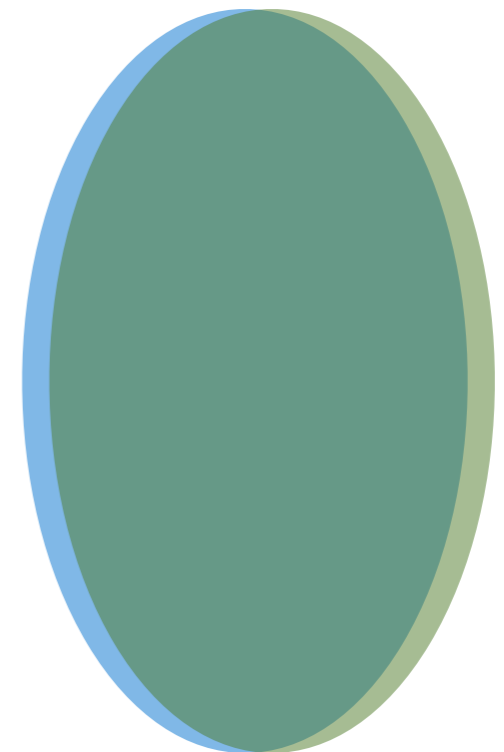
Switching of B

- central U+U collisions

U is deformed ion:

events with (almost) no particles in ZDC: **B=0**, **$\epsilon \neq 0$** ;

if photon v_2 is the same as the one of hadrons,
our mechanism is ruled out



Switching of ϵ

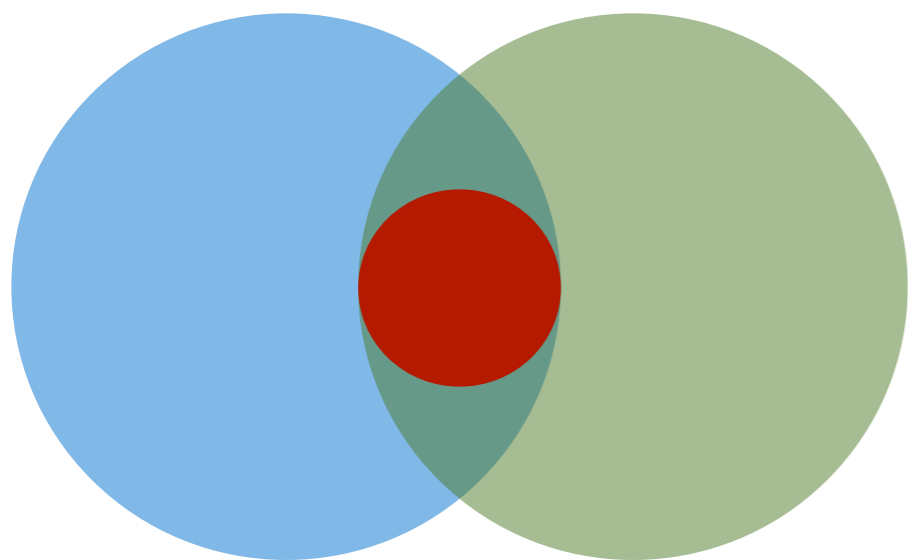
- non-central collisions: fluctuations of eccentricity

in given centrality class (e.g. 45-50% defined by ZDC), **B = const**; while hadronic v_2 fluctuates because of initial eccentricity fluctuations.

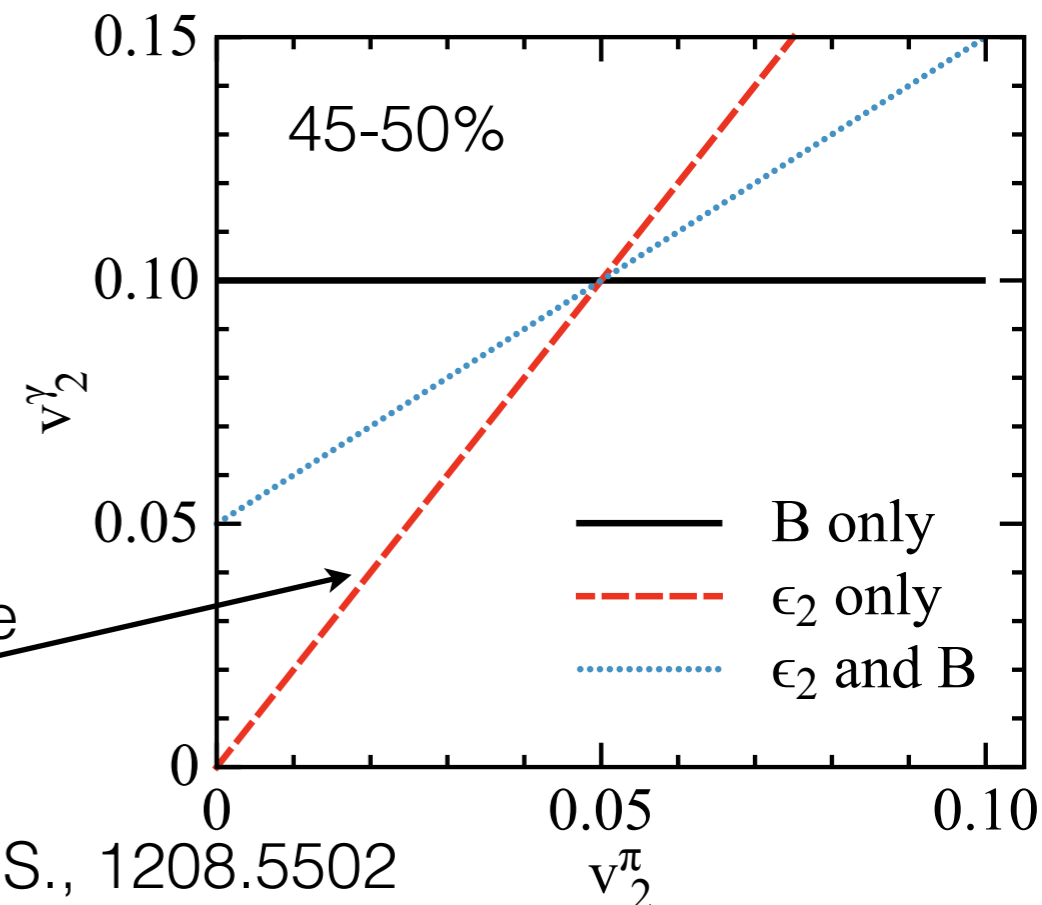
Limiting case:

non-central collisions ($\rightarrow eB \neq 0$) with zero **v_2** .

thus in such events anisotropy of photon production is due to eB .



Our mechanism will be ruled out, if



Experimental tests II

- small v_4 ; violation of scaling $v_4 \sim v_2^2$
- number of photons $N_{\text{in-plane}} > N_{\text{out-plane}}$
($N_{\text{in-plane}} - N_{\text{out-plane}}$) $\sim eB^2$ and thus is quadratic function of impact parameter
- in-plane polarization of photons

Outlook:LHC energies

- Higher initial temperatures \rightarrow lower bulk viscosity

$$T_{ave} = 304 \text{ MeV}$$

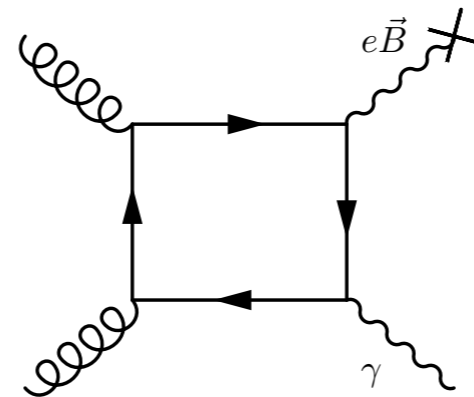
- Large $\gamma \rightarrow$ short time scales for non-zero magnetic field (modulo plasma response)

$$t_{LHC} = t_{RHIC} \Upsilon_{RHIC} / \Upsilon_{LHC} \rightarrow t_{LHC} \propto 0.01 \text{ fm/c vs } t_{RHIC} \propto 0.1 \text{ fm/c}$$

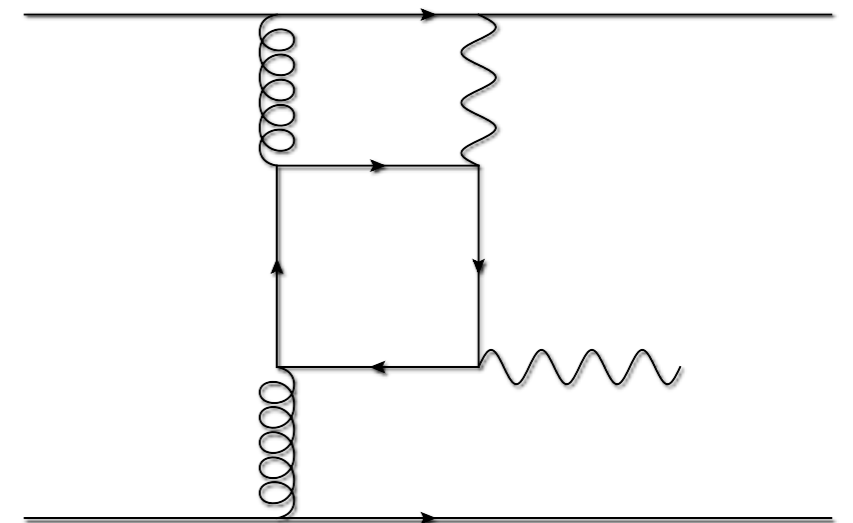
- LHC data can be described by:

$$T_{init}^{LHC} / T_{init}^{RHIC} = Q^{sat}_{LHC} / Q^{sat}_{RHIC} \text{ and } \tau_0^{LHC} / \tau_0^{RHIC} = Q^{sat}_{RHIC} / Q^{sat}_{LHC}$$

- No need for equilibrium
 - production from Glasma:



- production from CGC:



Outlook: RHIC low energy scan

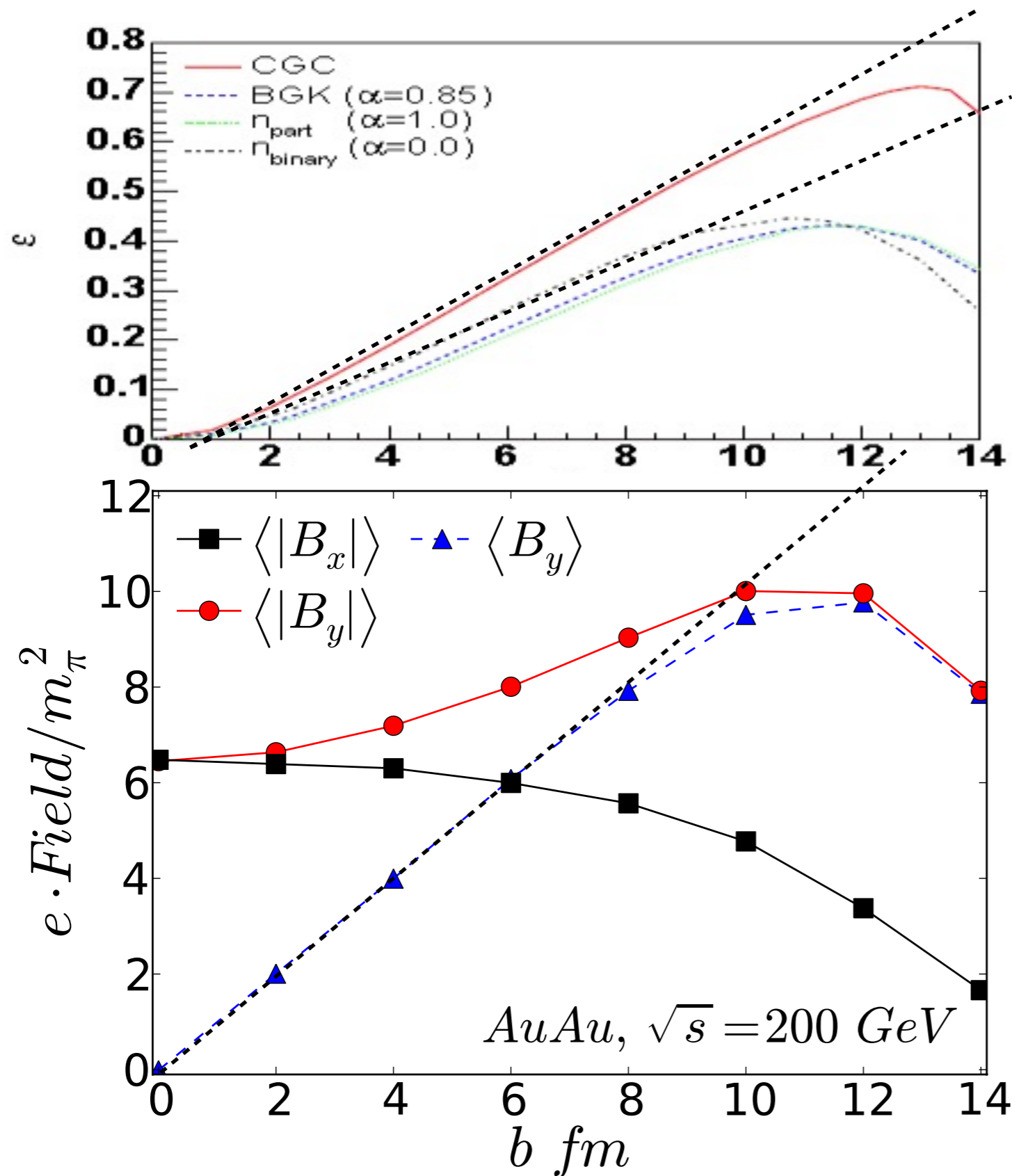
- lower energies \rightarrow lower eB , but longer time scales
- in equilibrium: bulk viscosity, ζ , is divergent at critical point (CP)
- in reality (HIC): at CP $\zeta \sim \xi^{2.8}$, ξ is correlation length
- on $O(4)$ line: $\zeta \sim \xi^2$ (while shear viscosity is finite)
see E. Nakano, V.S. and B. Friman, arXiv:1109.6822
- rather speculative: small eB can compensate large ζ

MAGNETO-HYDRO

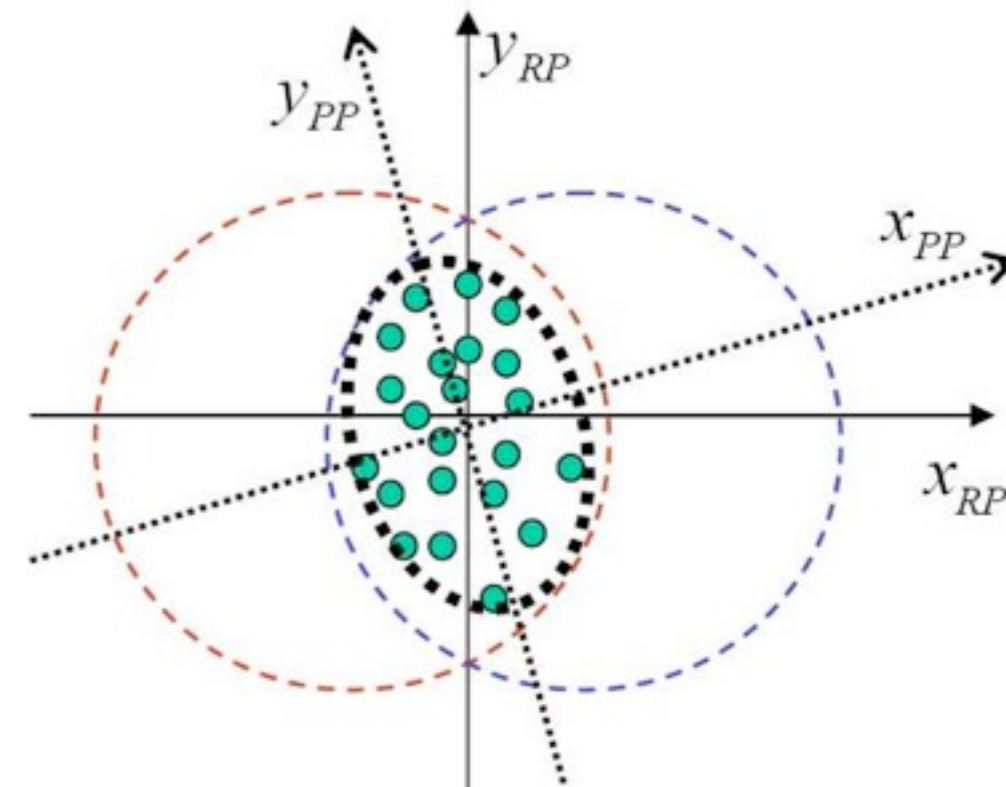
Summary

- photon v_2 puzzle:
 - hadronic physics?!
 - or effect of non-zero magnetic field?!
- there are ways to discriminate between hadronic and magnetic field related mechanisms of v_2 !
- axial anomaly: similar effect, but unknown μ_5 (similar to CME);
synchrotron radiation: similar effect, but unknown quark distribution in initial state

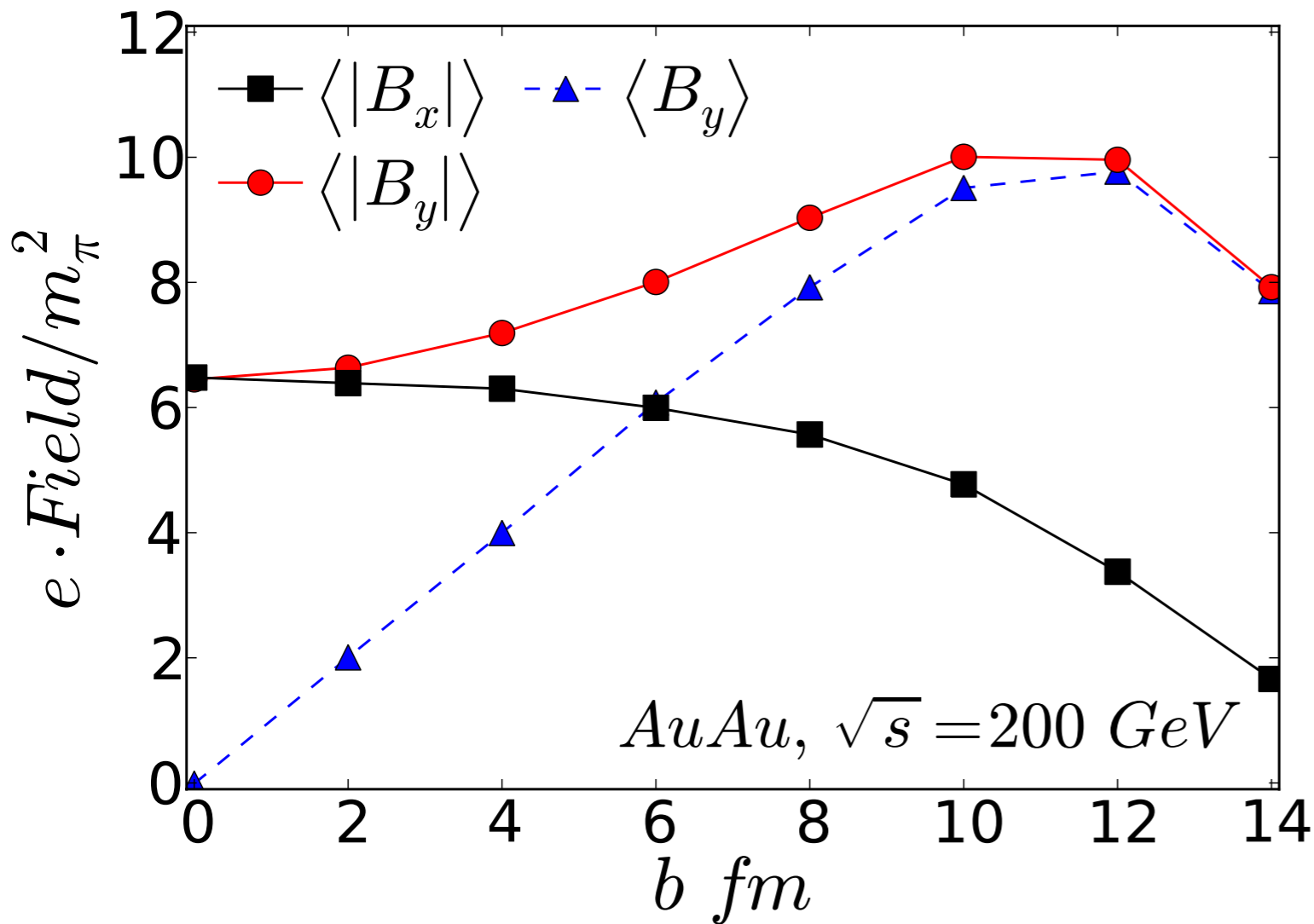
Backup



*Drescher Dumitru Hayashigaki Nara,
nucl-th/0605012*



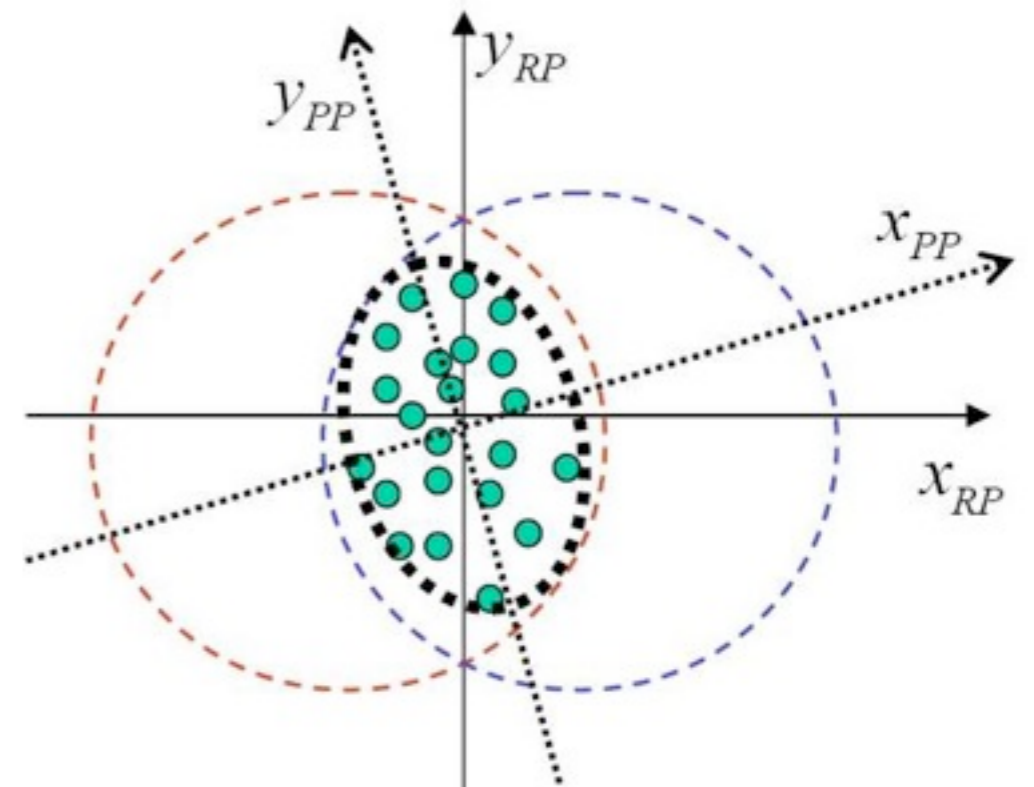
Fluctuations of eB



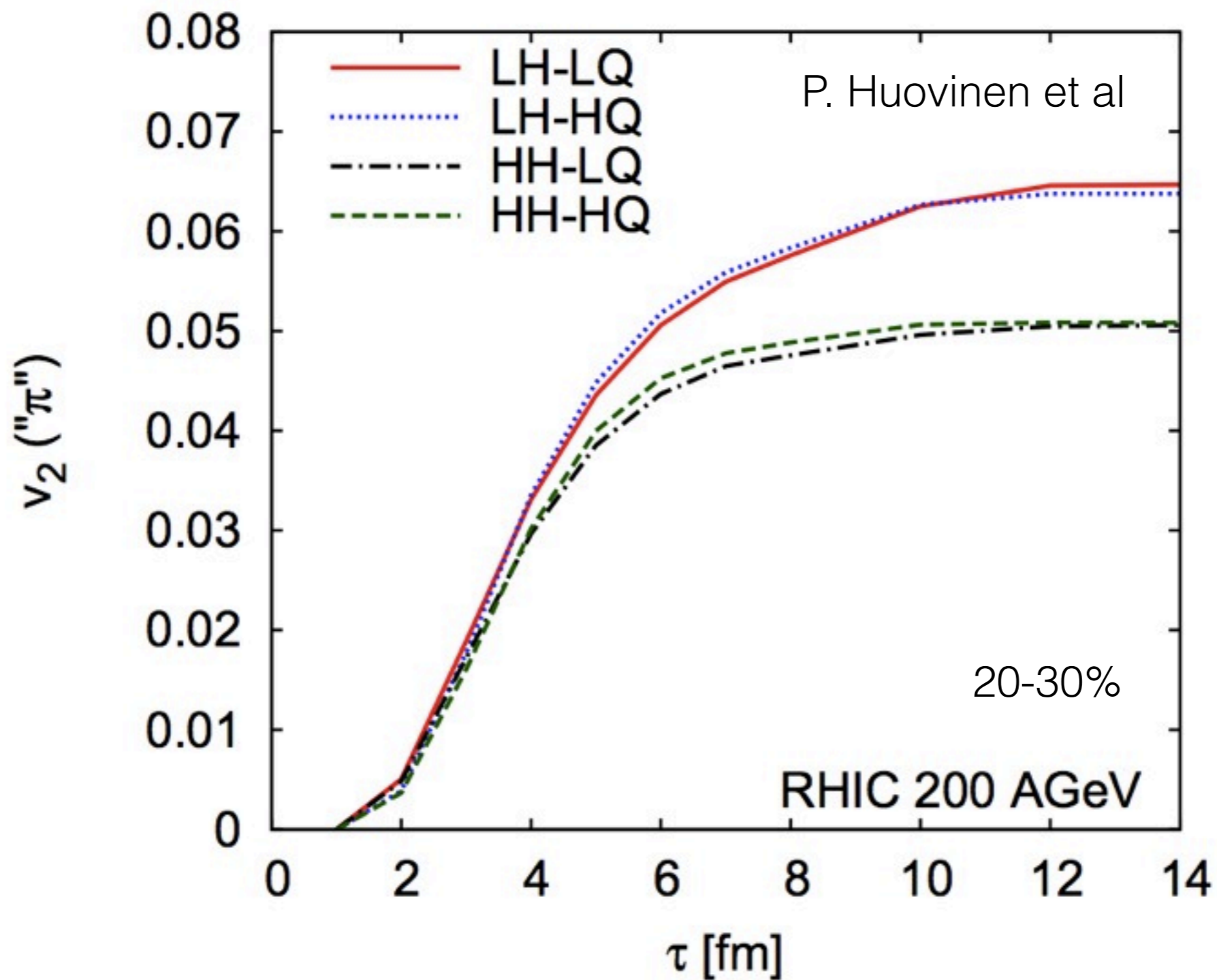
- for observables $\langle eB \rangle$ is not as significant as as $\langle |eB_y - eB_x| \rangle$

A. Bzdak and V.S., 1111.1949

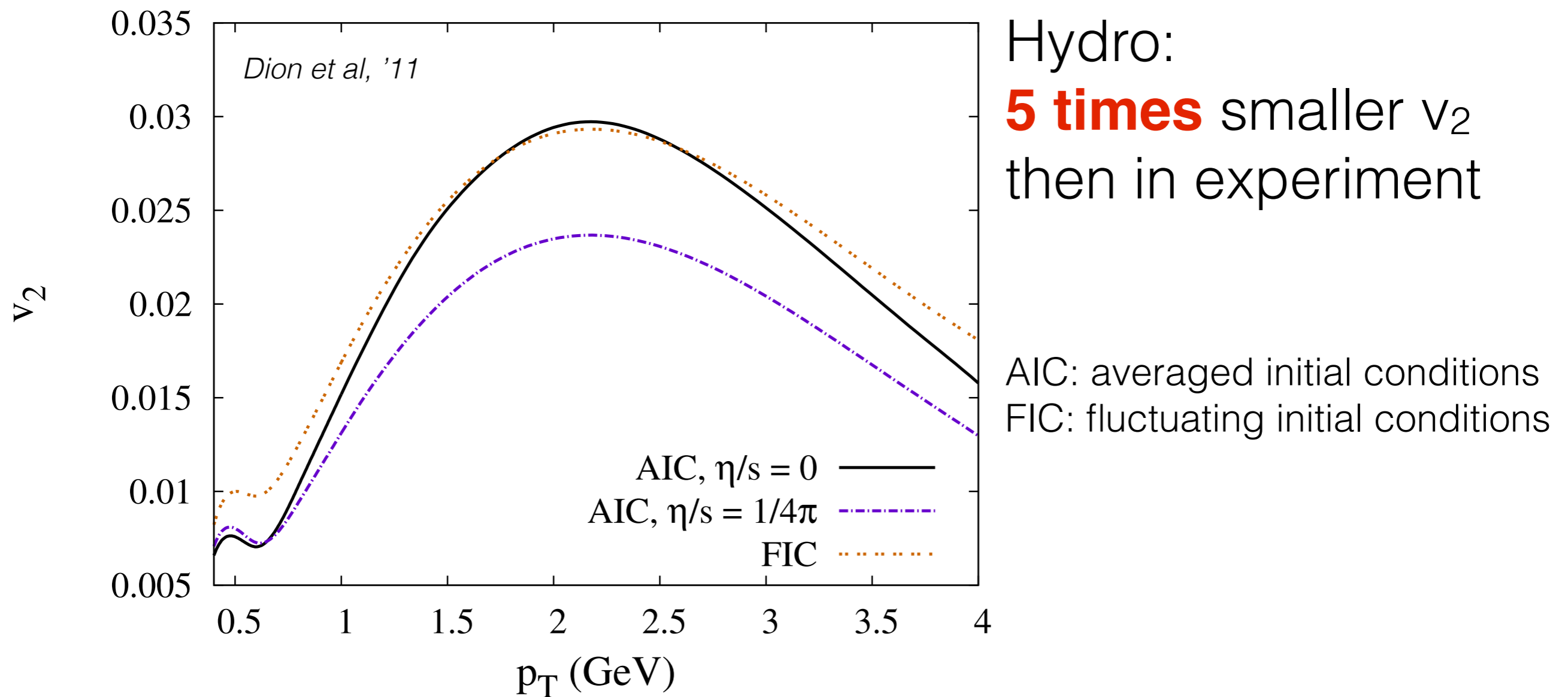
- azimuthal fluctuations of eB relative to orientation of participant plane: J. Błochyński et al 1209.6594



Development of hadronic flow



Theory: Hydrodynamics II



Large photon v_2 is difficult to explain with dominant QGP source