

Some aspects of dilepton production in HIC

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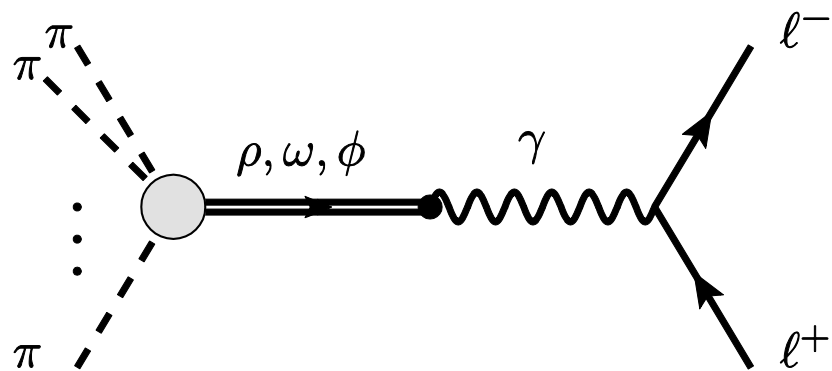
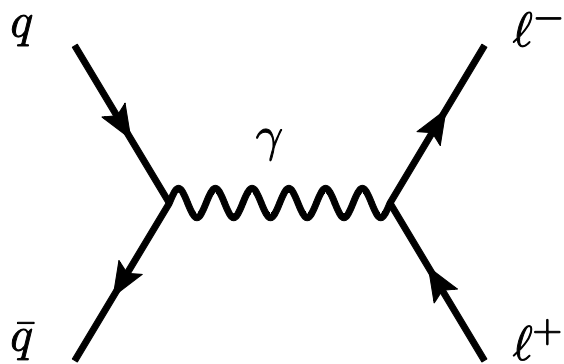
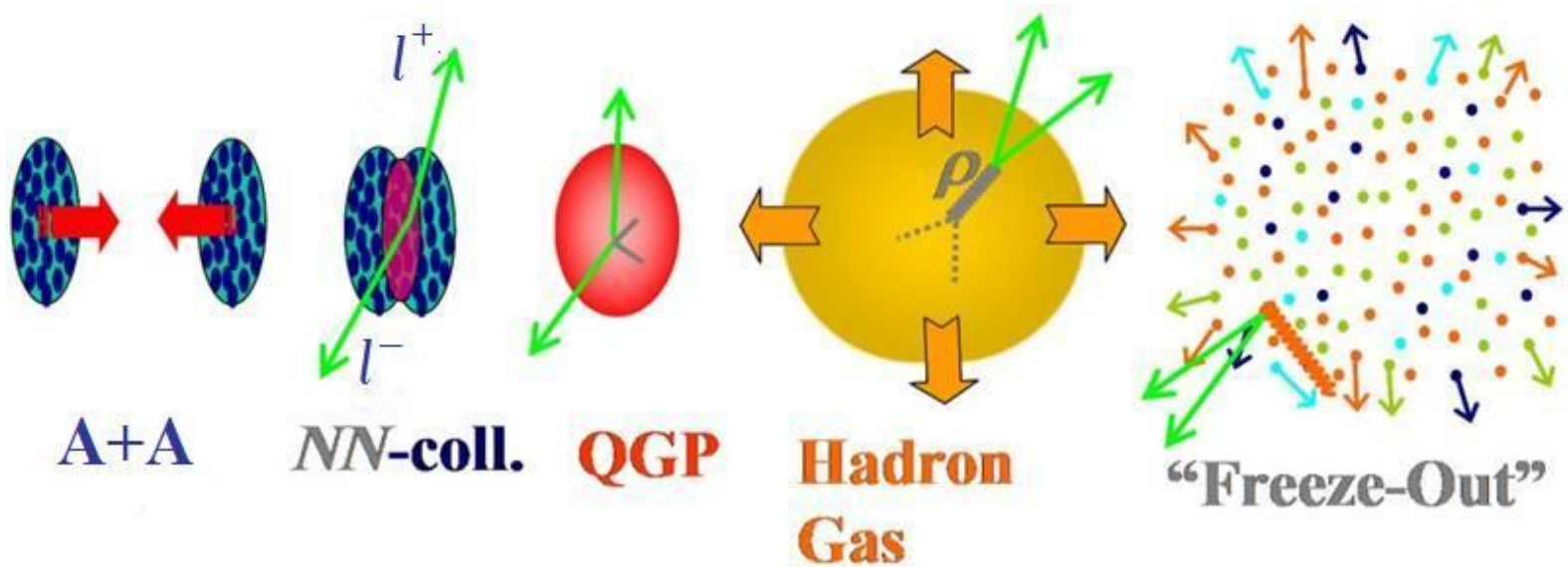
**J.Deng, Q.Wang, N.Xu, P.-f. Zhuang, PLB701,581(2010);
H.-j. Xu, H.-f. Chen, X.Dong, Q.Wang, Y.-f. Zhang,
PRC85, 024906(2012)**

Thermal radiation workshop, RBRC-BNL, Dec 4-7, 2012

Outline

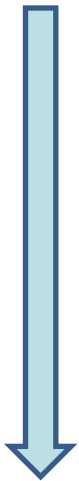
- 1. Introduction: a very brief timeline (history) of dilepton production in HIC (R. Rapp's talk, I. Tserruya's talk)**
- 2. Vector meson contribution via Vector Meson Dominance model (VDM)**
- 3. T_{eff} as probe to EOS of dense matter**
- 4. Comparison with STAR data**
- 5. Summary and conclusion**

Di-lepton sources in HIC



A little history (1)

Why dileptons ?



(1985) McLerran & Toimela, PRD31, 545

(1990) Braaten, Pisarski, Yuan, PRL64, 2242

(1993) Shuryak, Xiong, PRL70, 2241

- Suffer less interaction after their production
- Emitted in the whole space-time volume of the fireball
- Provide crucial information of the medium properties created by HIC

Earlier
experiments

(1997) DLS collaboration, PRL79,1229

(1998) DLS collaboration, PRC57,1865

- Rare production
- Complex background

Ca+Ca, C+C,
C+Ca, d+Ca

$M < 0.7$ GeV

A little history (2)

Low statistics

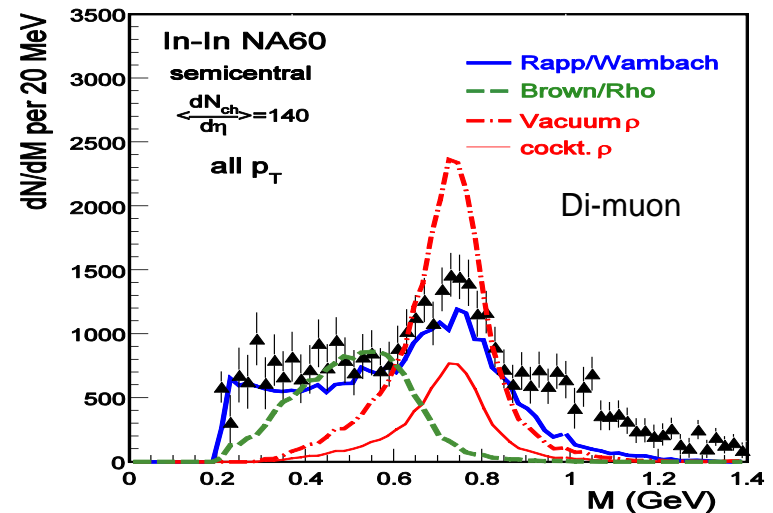
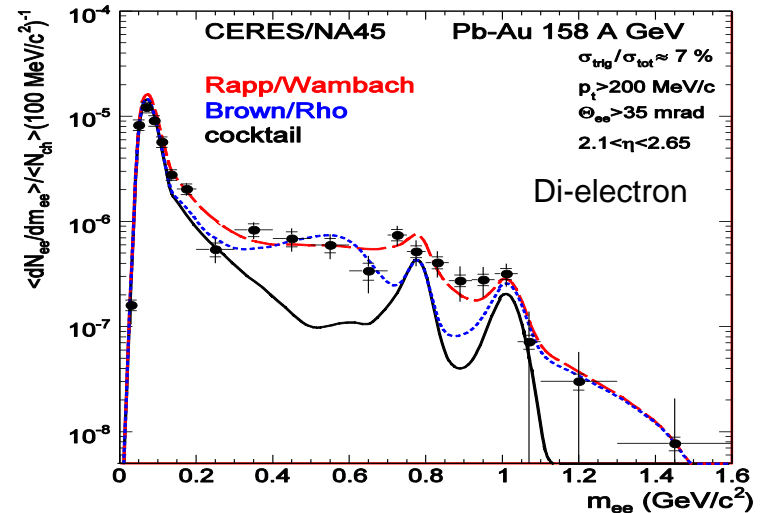
(1998) HELIOS-3, EPJC 5, 63
(1998) CERES/NA45, PLB422,405
(2000) NA50, EPJC 13,69

Low mass
enhancement

High statistics

(2006) NA60, PRL96,162302

In-medium rho
spectra



A little history (3)

Low statistics

(2004) Brown & Rho, Phys.Rept 398, 301

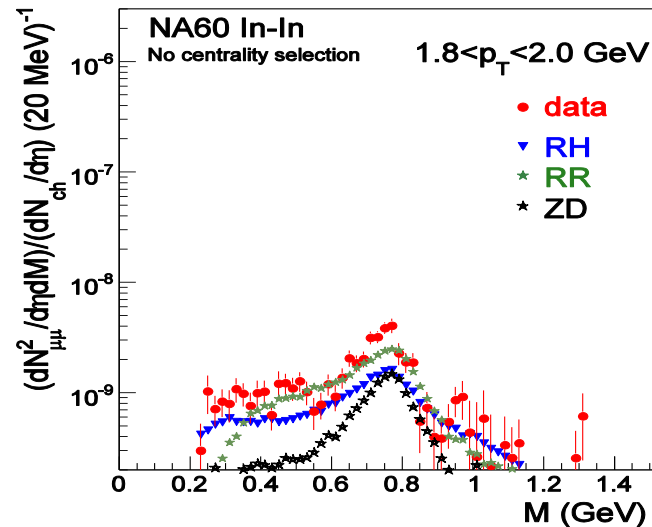
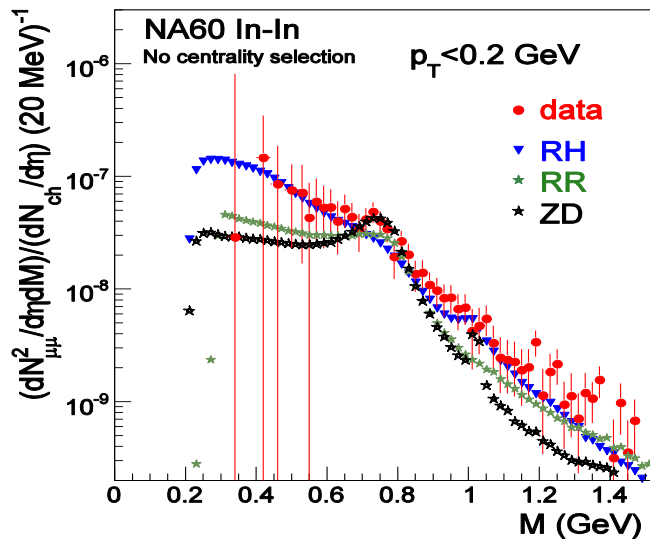
Broadening spectra vs Dropping mass

High statistics

(2006) van Hees & Rapp. PRL97, 102301

(2008) Ruppert, Gale, et al. PRL100, 162301

(2007) Dusling, Teaney & Zahed. PRC75, 024908



A little history (4)

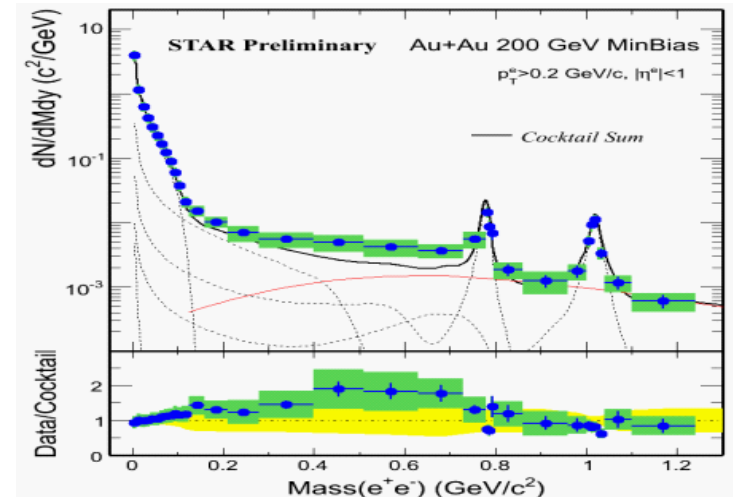
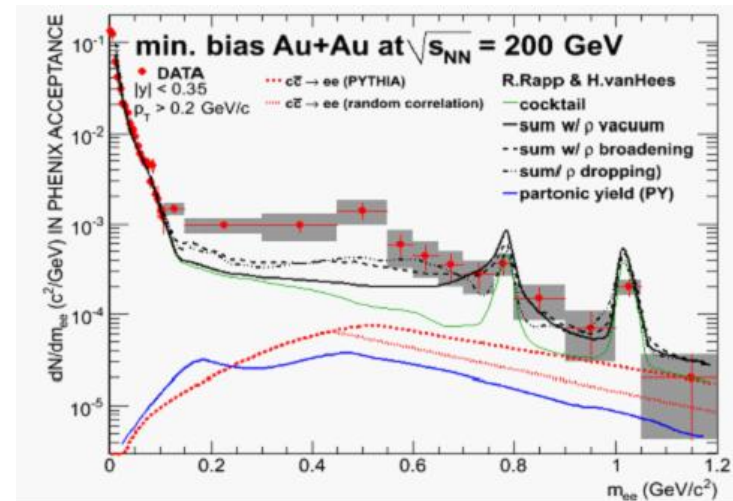
RHIC energy

(2010) PHENIX, PRC81, 034911
(2011) J. Zhao, JPG38, 124134

PHENIX puzzle

Future

- STAR-BES
- LHC
- CMB



A little history (5)

RHIC energy

- (2009) Rapp, Wambach, 0901.3289
- (2012) Linnyk et al, PRC85, 024910
- (2012) Xu, et al, PRC85, 024906

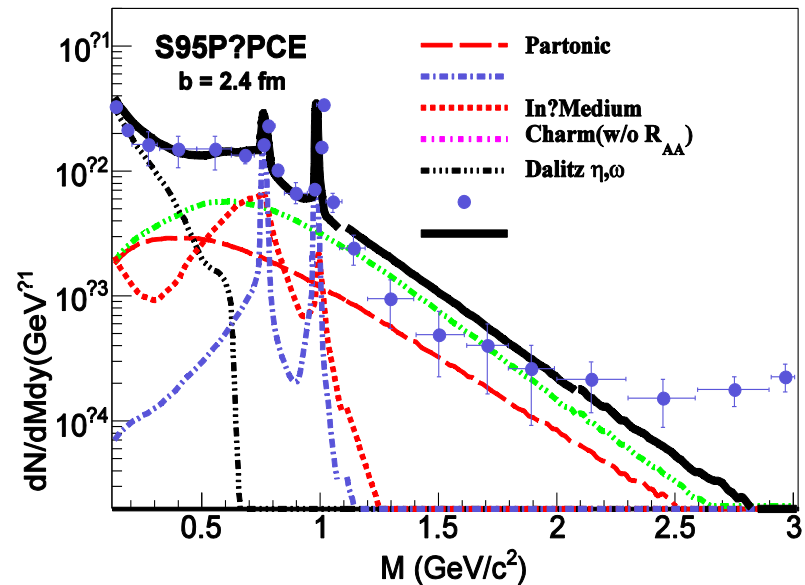


phenomenology study

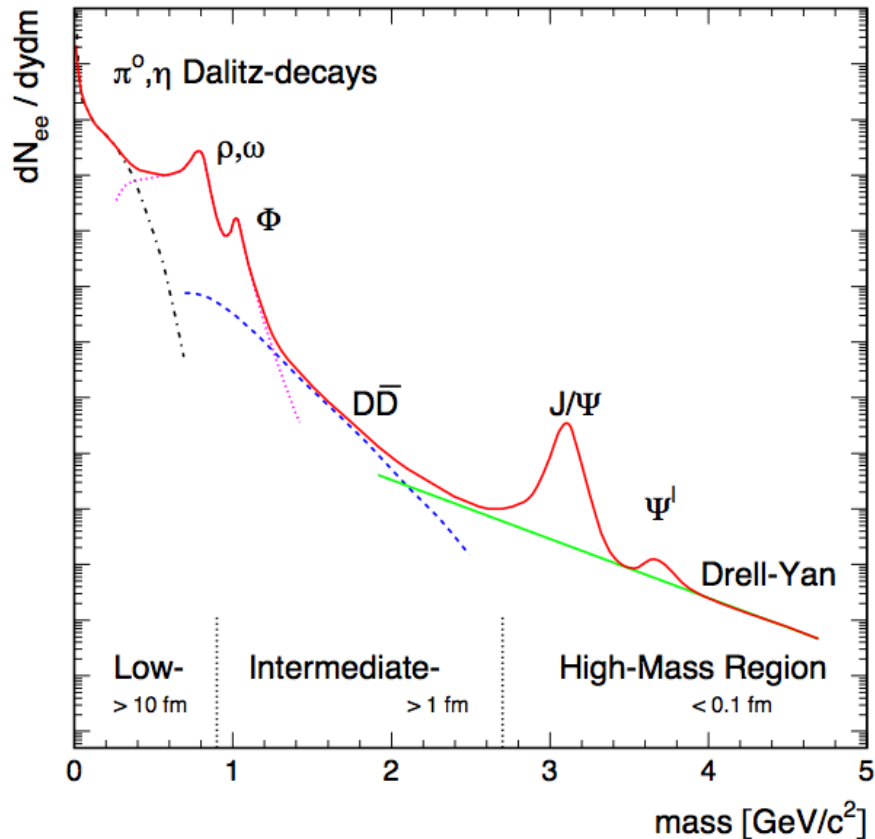
- PHENIX puzzle
- The signals from BES
- v_2 , HBT
- etc.

- (2006) Chatterjee, et al. PRC75, 054909
- (2008) Dusling et al. NPA809, 246
- (2011) Deng, et al. PLB701,581
- (2011) Mohanty, et al. PRC84, 024903

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Invariant mass spectra



LMR

Chiral symmetry restoration,
Vector meson production,
in-medium effect

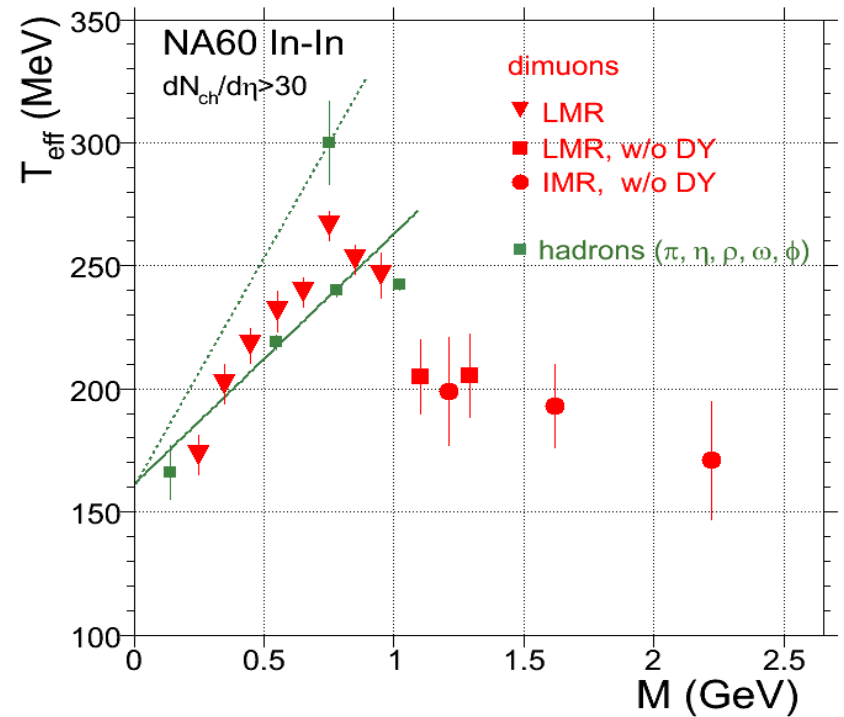
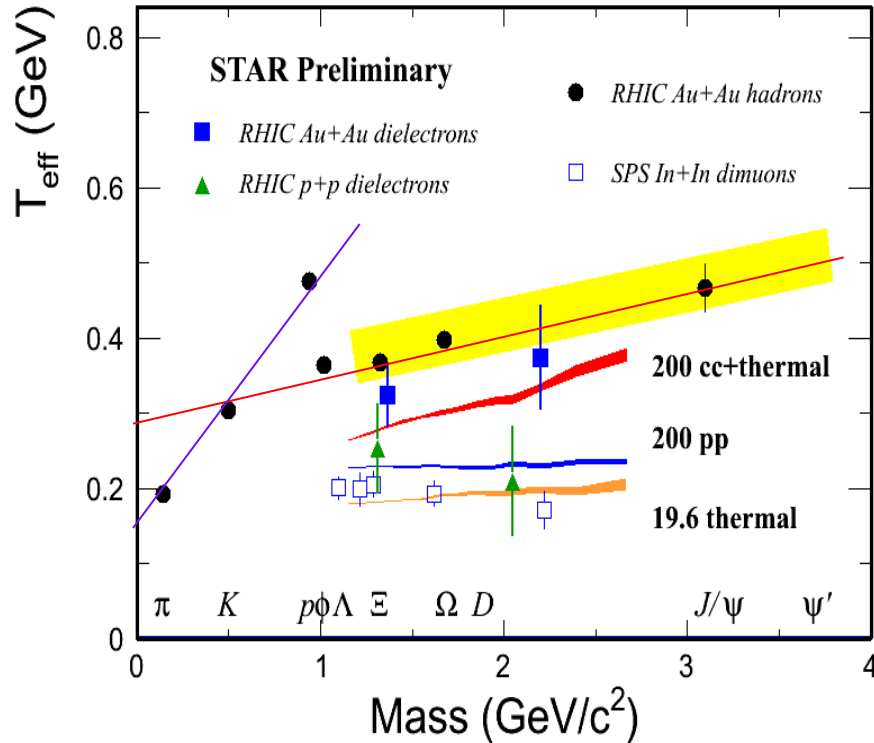
IMR

Heavy quark, QGP thermal
radiation

HMR

Heavy quarkonia, Drell-Yan

Effective temperature for hadrons and dileptons



$$T_{\text{eff}} = T_0 + Mv_T^2$$

The transition tregon may signal a transition from a hadronic source to a partonic source
 NA60, PRL100, 022302(2008); EPJC59, 607(2009)

Strategy

- Dilepton production in Au+Au collisions at 200 GeV in IMR, QGP phase: $q\text{-}q\text{-bar}$ annihilation; Hadron phase: D_ρ with vertices $\rho\pi X$ (X : all mesons below 1300 MeV), and vertices of ρNN^* and $\rho N\Delta^*$ (N^* and Δ^* : baryon resonances); D_ω with vertices $\omega\rho\pi$, $\omega\pi\pi\pi$; and D_ϕ with vertices $\phi KK\text{-bar}$.
- Space-time evolution of medium is described by a 2+1 ideal hydro model, different EOS are used for hydro-simulation.
- Slope parameters and elliptic flows show distinct features from two phases
- Comparison to recent STAR data.

Hydrodynamics for HIC

- Assumption: thermalization, Ideal or Viscous
- Inputs: EOS, initial conditions, freeze-out conditions
- Outputs: space-time evolution
- Comparison with data: pt-spectra, flows, ...
- Further application: fluctuation & correlation, non-equilibrium statistics, ...

Baym, Friman, Blaizot et al 86', Rischke 98', Kolb, Huovinen, Heinz, 00-01', Romatschke, 08', Song, Heinz, 08', Schenke, Jeon, Gale, 10', Pang, Q.W., Wang 12',

Vector Meson Dominance Model (VDM)

- VDM (Kroll, Lee, Zumino, 67'). The Lagrangian for $\rho\pi\gamma$ -system:

$$L = (D_\mu \pi)^* (D^\mu \pi) - m_\pi^2 \pi^* \pi - \frac{1}{4} G_{\mu\nu} G^{\mu\nu} + \frac{1}{2} m_\rho^2 \rho_\mu \rho^\mu + \bar{\psi} (i\gamma_\mu D_A^\mu - m) \psi - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \boxed{\frac{e}{2g} F_{\mu\nu} G^{\mu\nu}} \leftarrow \text{Photon-rho coupling}$$

where

$$D^\mu = \partial^\mu + ig\rho^\mu + ieA^\mu \quad F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$$

$$D_A^\mu = \partial^\mu - ieA^\mu \quad G^{\mu\nu} = \partial^\mu \rho^\nu - \partial^\nu \rho^\mu$$

- EOM for EM field

$$\partial_\sigma F^{\alpha\sigma} \approx e\bar{\psi}\gamma^\alpha\psi - \frac{e}{g} m_\rho^2 \rho^\alpha \quad \longrightarrow \quad \langle J^\alpha(x) J^\beta(y) \rangle \rightarrow \frac{e^2 m_\rho^4}{g^2} \langle \rho^\alpha(x) \rho^\beta(y) \rangle$$

(keep terms linear in e/g)

$$\Pi_\gamma(x, y) \rightarrow \frac{e^2 m_\rho^4}{g^2} D_\rho^{\alpha\beta}(x, y)$$

Dilepton emission rate (1)

In-medium rate

$$\frac{d^8 N}{d^4 x d^4 p} = -\frac{\alpha}{4\pi^4} \frac{1}{M^2} n_B(p \cdot u) L(M) \text{Im}\Pi_\gamma^R(p)$$

$$L(M) = \left(1 + \frac{2m_l^2}{M^2}\right) \sqrt{1 - \frac{4m_l^2}{M^2}}$$

u: fluid velocity
T: temperature
 Π_γ and n_B
 depend on space-time
 via **u** and **T**

Photon selfenergy (VDM + quark) $\text{Im}\Pi_\gamma^R = \begin{cases} \sum_{V=\rho,\omega,\phi} e^2 \left(\frac{m_V^2}{g_V}\right)^2 \text{Im}D_V & \text{hadronic source} \\ -\frac{N_c M^2}{12} \left(1 + \frac{\alpha_S}{\pi}\right) \sum_{i=u,d,s} e_i^2 & \text{quark source} \end{cases}$

Imaginary part of Retarded propagator of rho meson

$$\text{Im}D_V(M, q, T) = \frac{1}{3} \text{Im}D_V^L(M, q, T) + \frac{2}{3} \text{Im}D_V^T(M, q, T)$$

$$\text{Im}D_V^{L,T}(M, q, T) = \frac{\text{Im}\Sigma_V^{L,T}(M, q, T)}{|M^2 - (m_V^0)^2 + \Sigma_V^{L,T}(M, q, T)|^2}$$

Dilepton emission rate (2)

- Freezeout (FO) dilepton rate is related to FO vector meson rate. Most of ρ mesons decay inside medium. But most of ω and φ meson decays take place after FO due to their long life time.

$$\frac{dN_{l\bar{l}}^{fo}}{d^4p} = \frac{\alpha}{3} \left(\frac{e}{g}\right)^2 \frac{m_V}{\Gamma_V} \frac{dN_V^{fo}}{d^4p}$$

$$\frac{dN_V^{fo}}{d^4p} = \frac{g_s^\rho}{4\pi^4} \int_{T_f} d\Sigma_\mu p^\mu \text{Im}D_V n_B(p \cdot u)$$

Freezeout
Emission rate of
vector meson

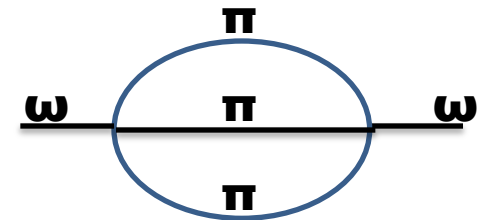
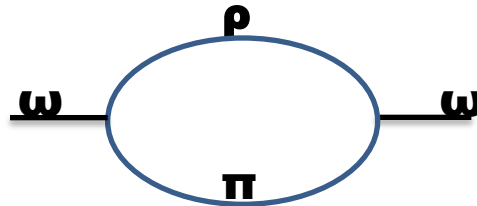
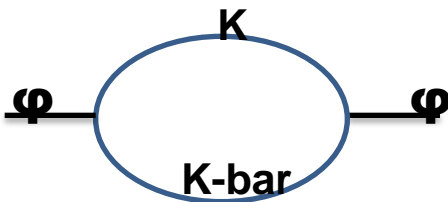
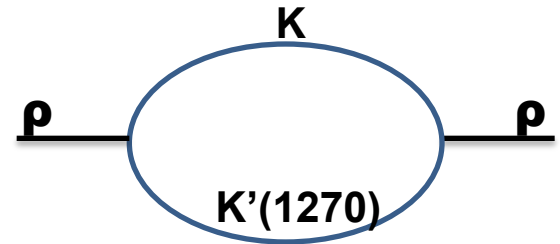
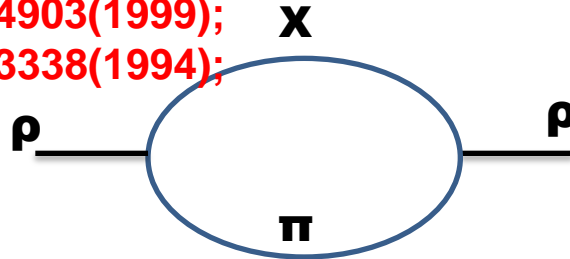
Freezeout
hypersurface

In-vacuum
vector meson
propagator

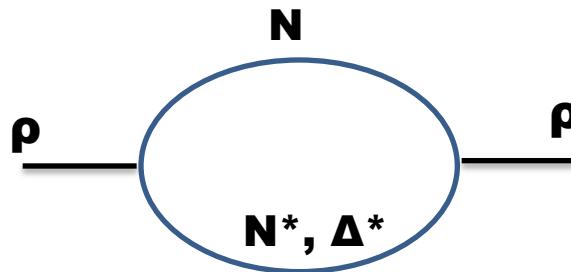
Rho self energy (1)

Chanfray, Schuck, NPA555,329(1993);
 Herrmann, Friman, Noerenberg, NPA560(1993);
 Rapp, Gale, PRC 60,024903(1999);
 Gale, Lichard, PRD 49,3338(1994);

Meson Contribution
 $X = \pi, \omega, h_1(1170)$
 $a_1(1260), \pi'(1300)$



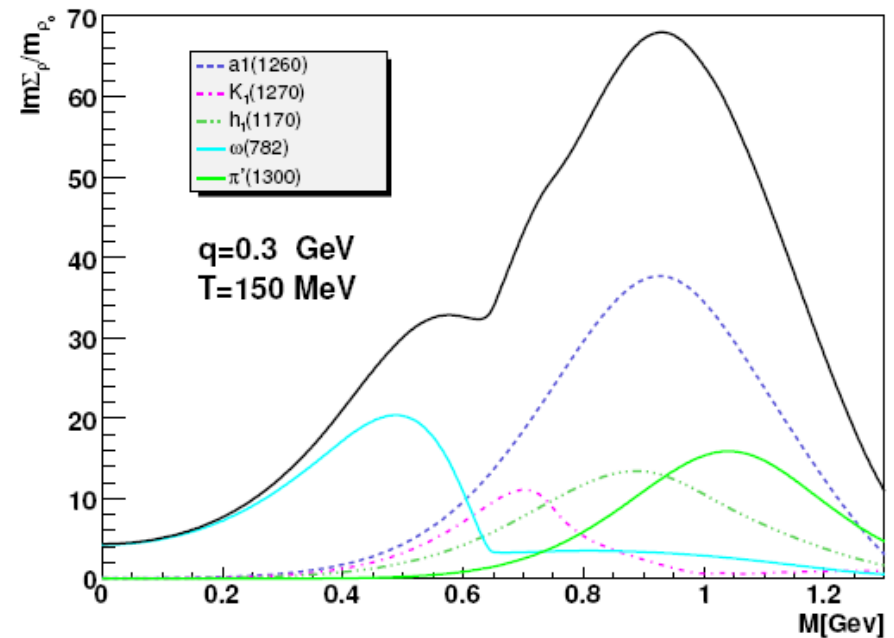
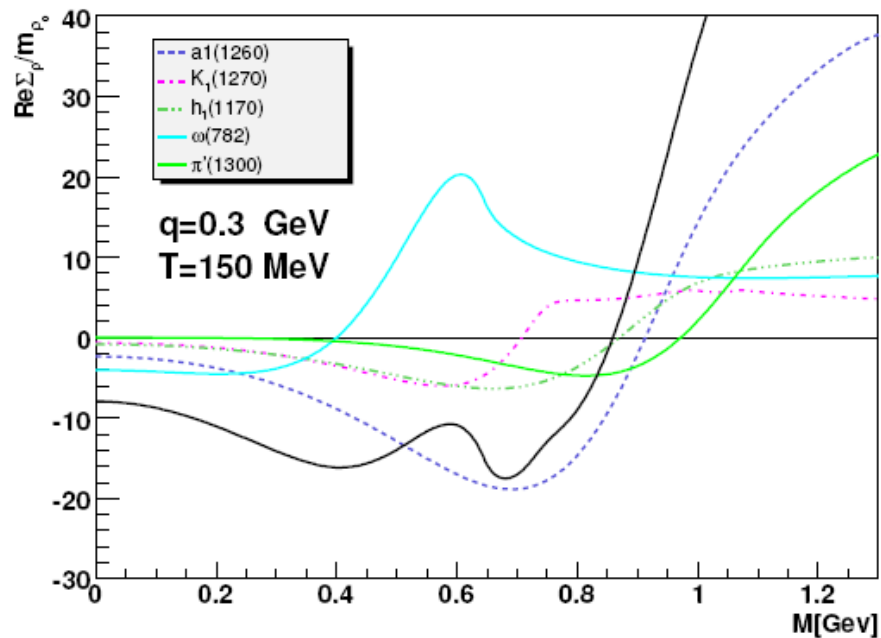
Baryon Contribution:
 $N(1700), N(1720), N(1900)$
 $N(2000), N(2080), N(2090)$
 $N(2100), N(2190),$
 $\Delta(1700), \Delta(1900), \Delta(1905),$
 $\Delta(1940), \Delta(2000)$



Eletsky, Belkacem,
 Ellis, Kapusta,
 PRC 64, 035202(2001);
 Eletsky, Kapusta,
 PRC 59, 2757(1999)

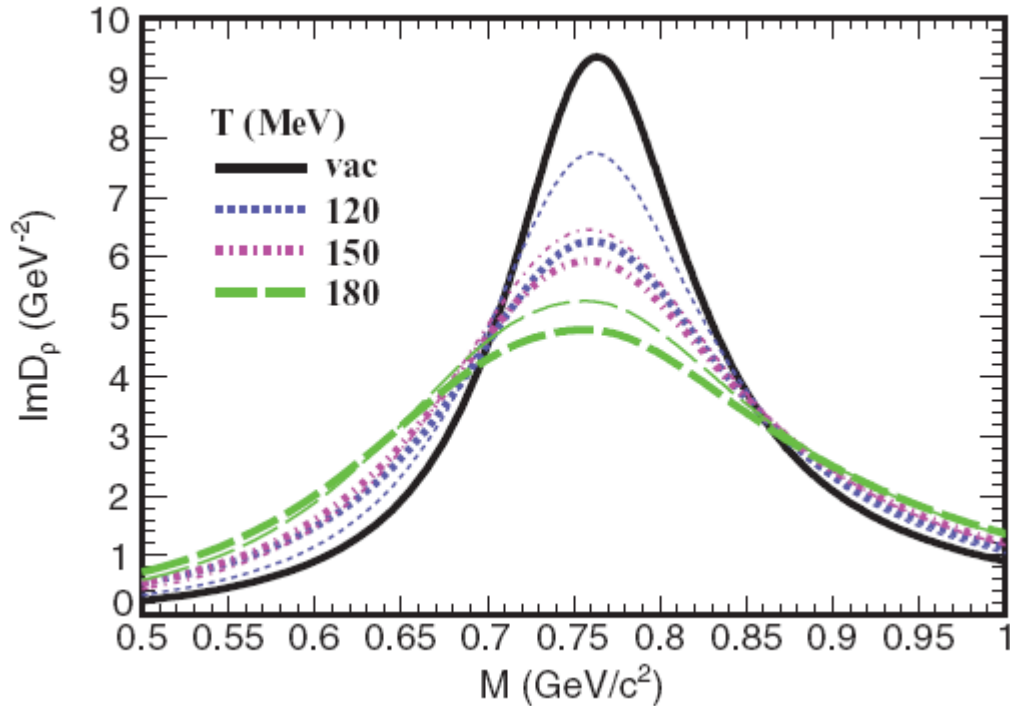
Rho self energy (2)

Re and Im from mesons



Rapp & Gale, PRC 60,024903(1999)

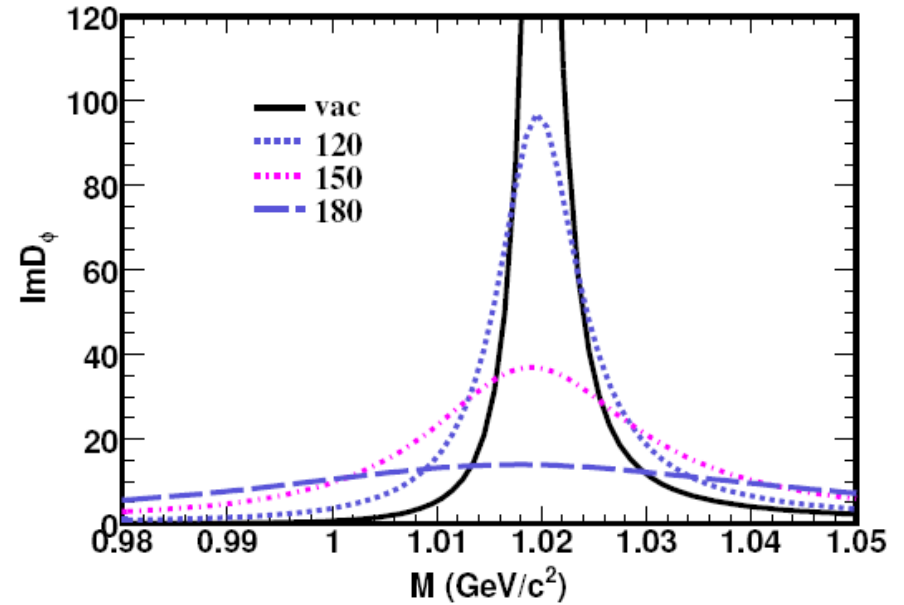
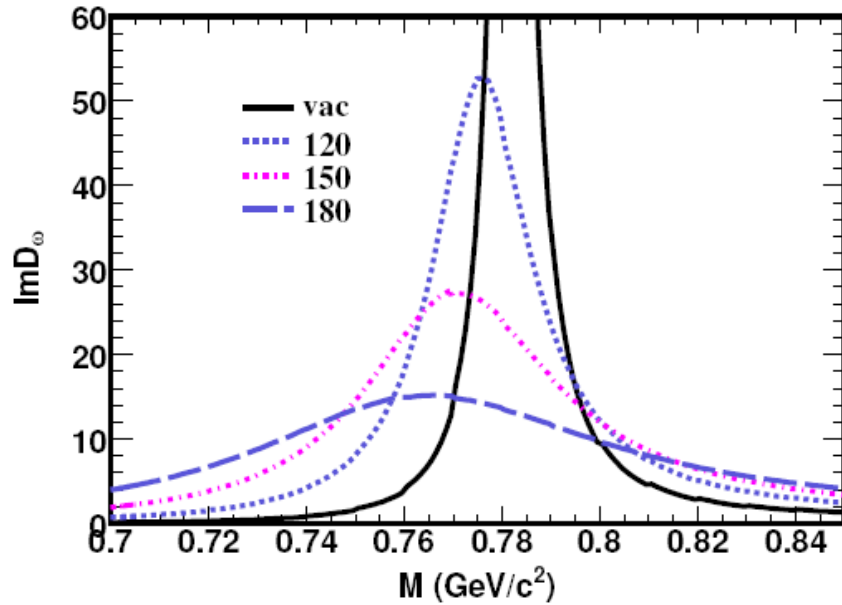
Rho self energy (3): Im D_ρ w/o $NN^*+N\Delta^*$ contribution



The imaginary parts of the in-medium ρ meson propagators (or in-medium spectral functions) with (thick lines) and without (thin lines) baryonic contributions. The chemical potentials in the PCE EOS are used.

The imaginary part of the propagator is sensitive to temperature, but insensitive to its momentum.

Im of ω and ϕ propagator



**T_{eff} as probe to
EOS of dense matter**

Transverse flow: slope parameter

$$\frac{d^4 N}{p_T dp_T M dM d\phi_p} \approx -\frac{\alpha}{2\pi^4} \frac{1}{M^2} \left(1 + \frac{2m^2}{M^2}\right) \sqrt{1 - \frac{4m^2}{M^2}} \times \int d^4 x \exp \left[\frac{1}{T} \gamma_T M_T v_T p_T \cos(\phi_v - \phi_p) \right] K_0 \left(\frac{\gamma_T M_T}{T} \right) \text{Im}\Pi^R$$

$M_T = \sqrt{M^2 + p_T^2}$

differential rate

space-time integral

transverse fluid velocity

azimuthal angle of fluid

spectra in transverse momentum and invariant mass

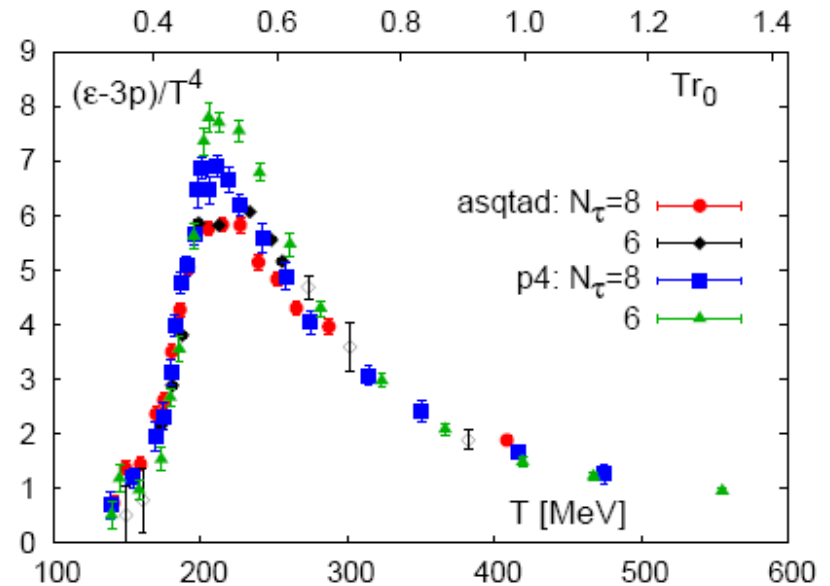
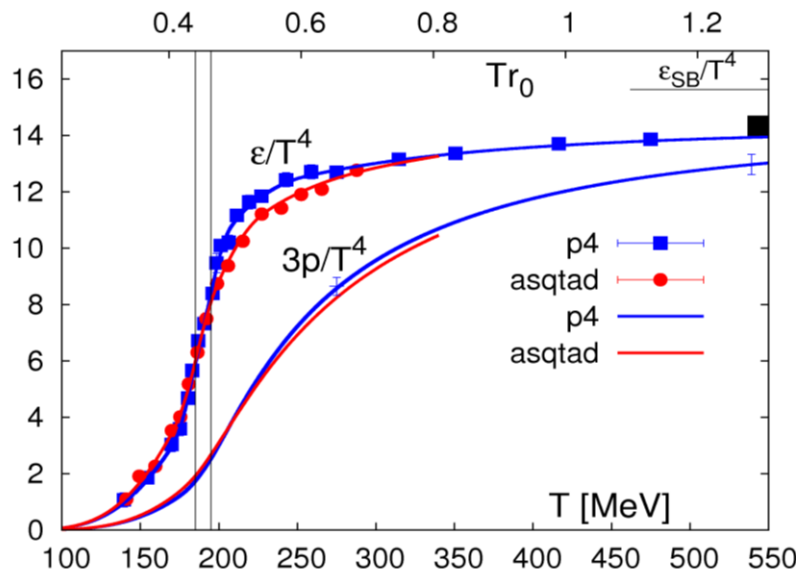
$$\frac{d^2 N}{m_T dm_T M dM} \sim \sqrt{\frac{\bar{T}}{\bar{\gamma}_T}} \frac{\sqrt{m_T + M}}{m_T} \exp\left(-\frac{m_T + M}{T_{eff}}\right)$$

slope parameter

$$T_{eff} \sim \begin{cases} \bar{T} + M^* \bar{v}_T^2, & \text{for } p_T \ll M \\ \bar{T} \sqrt{\frac{1+\bar{v}_T}{1-\bar{v}_T}}, & \text{for } p_T \gg M \end{cases}$$

$m_T \equiv M_T - M$

Dense or hot QCD matter EOS



*Bernard et al, (MILC) PRD 75 (07) 094505,
Cheng et al, (RBC-Bielefeld) PRD 77, 014511(2008);*

Bazavov et al, (HotQCD), Phys.Rev.D80, 014504(2009).

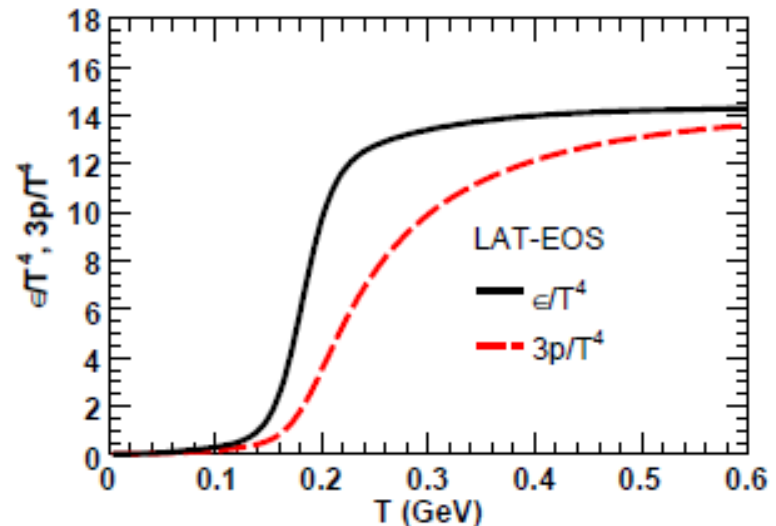
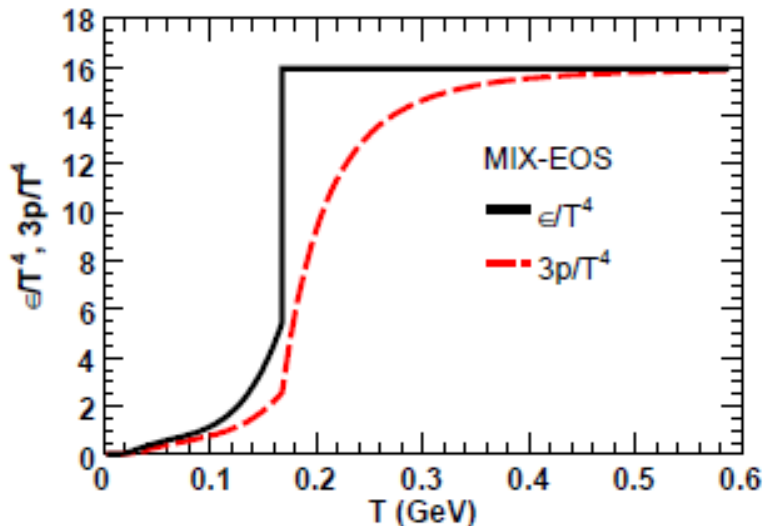
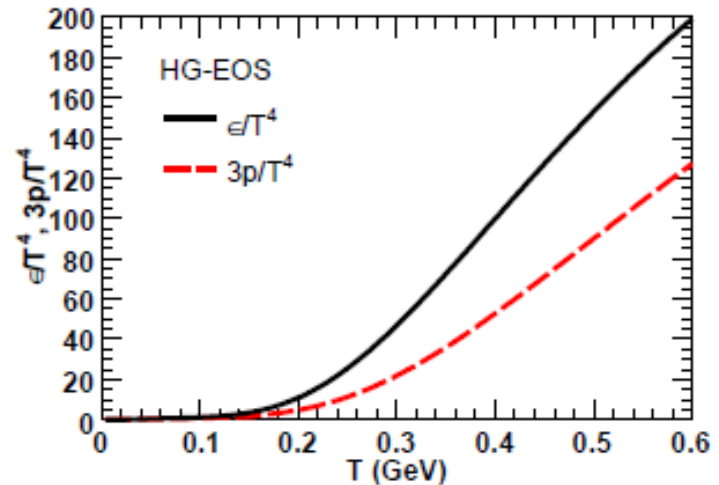
Four equations of state (EOS)

Massless ideal QGP

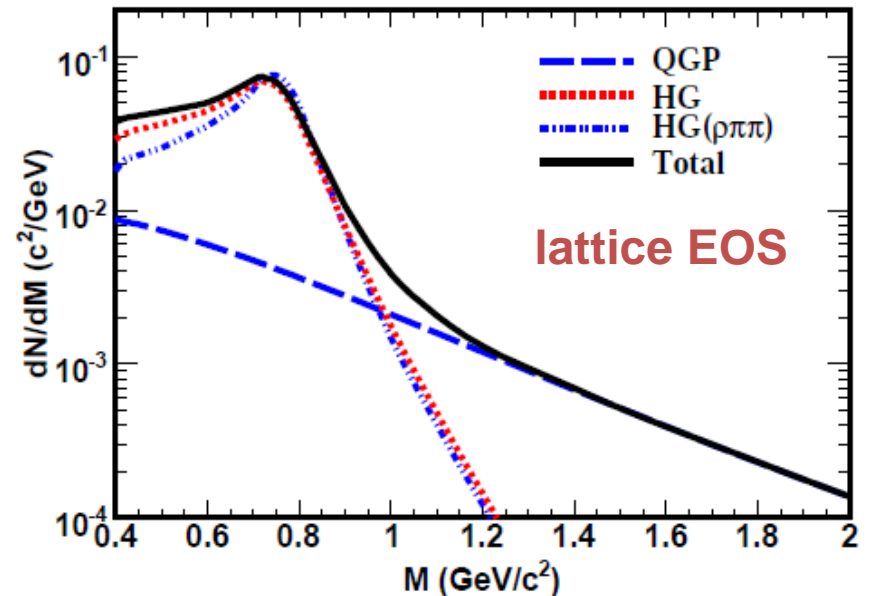
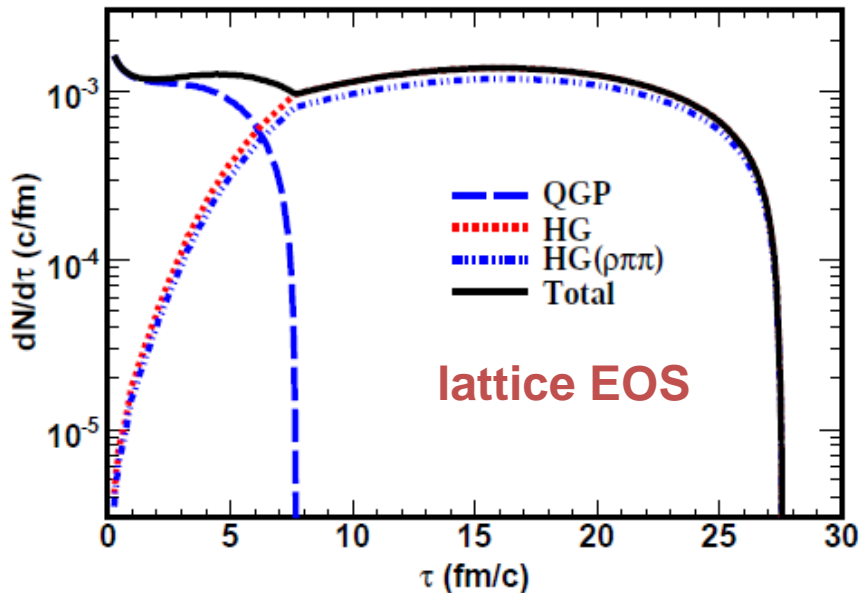
$$\epsilon = 3p = 16T^4$$

Resonance Hadron gas

[Braun-Munzinger, Redlich,
Stachel, nucl-th/0304013]

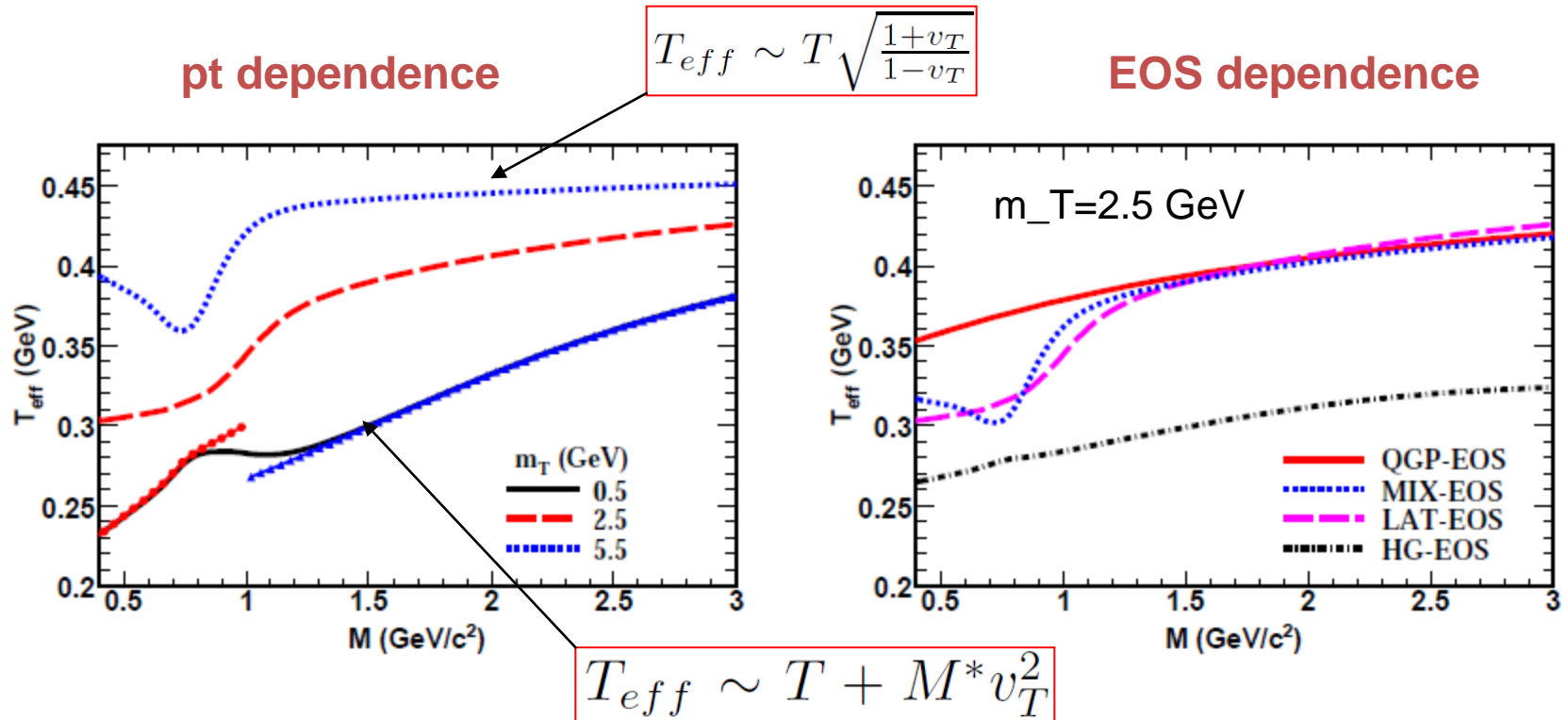


Time evolution of the rate



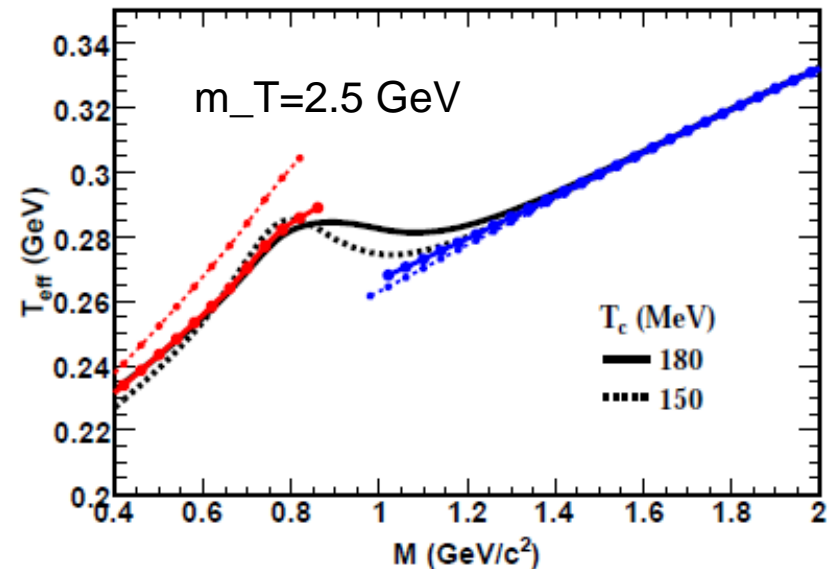
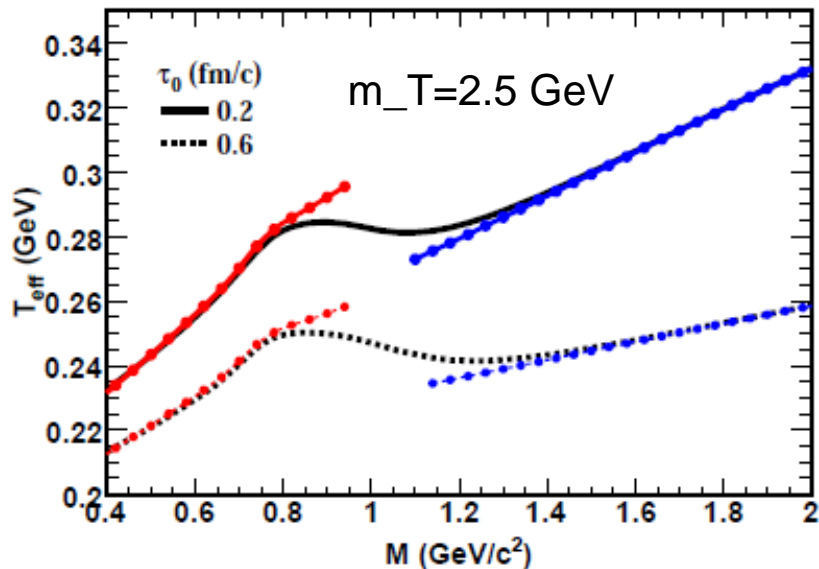
Differential multiplicity as functions of the dilepton invariant mass and proper time. The lattice EOS is used. The unit is arbitrary. The contributions from QGP and HG are shown in dashed and dotted lines.

Slope parameter: pt and EOS dependence



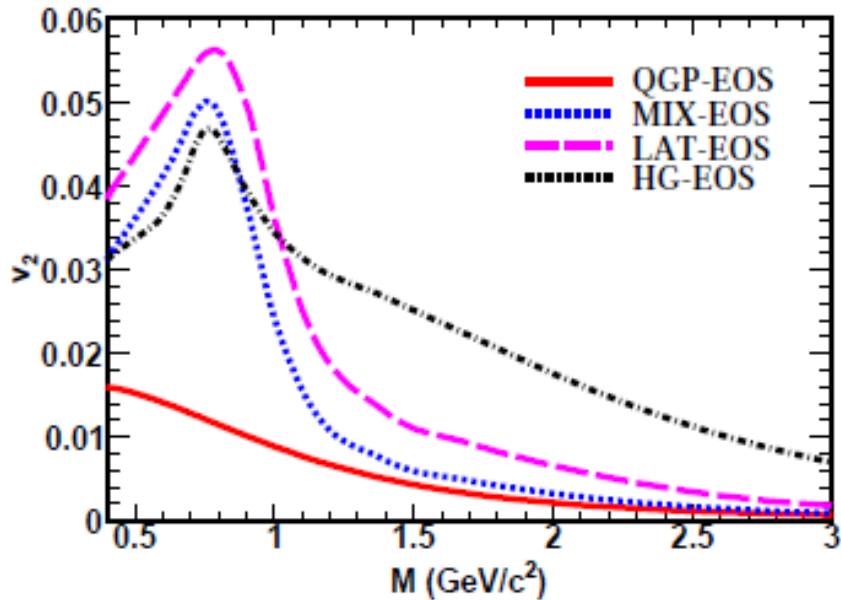
Slope parameter as functions of M for the mixed phase (left panel) and the lattice (right panel) EOS. The results for different values of m_T are shown in the solid, dashed and dotted lines. The lines with hollowed circles/triangles are extracted from the HG/QGP components.

Slope parameter: parameter dependence

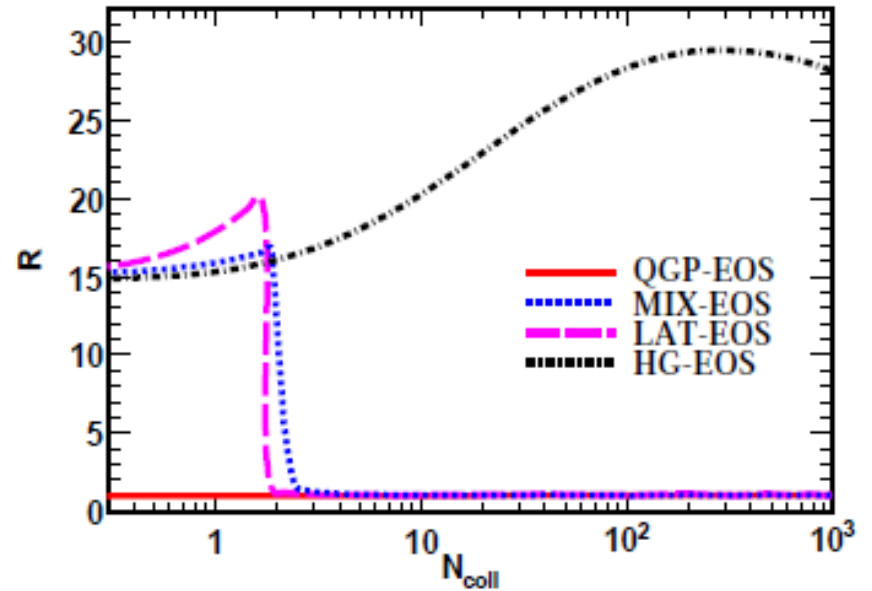


Parameter dependences of the slope parameter with the lattice EOS. Left panel: the initial time for the hydrodynamic evolution $\tau_0 = 0.2; 0.6$ fm/c. Right panel: the phase transition temperature $T_c = 180, 150$ MeV. $m_T = 2.5$ GeV.

Elliptic flow v_2 and R



$$v_2(p_T, M) = \langle \cos(2\phi_p) \rangle$$

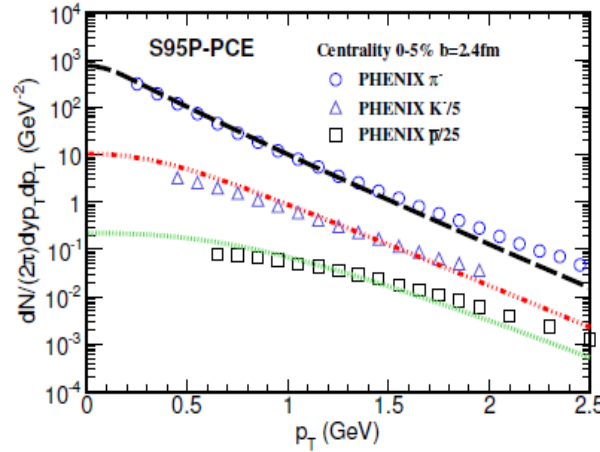
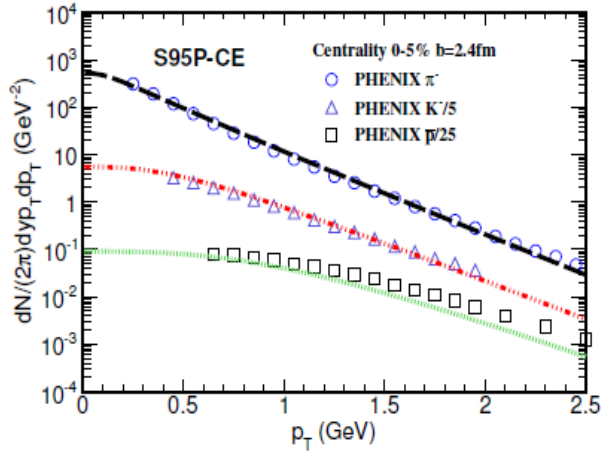


$$R = \frac{dN/dM_T^2 dp_T^2 dy|_{M_T=2.6\text{GeV}, p_T=2\text{GeV}}}{dN/dM_T^2 dp_T^2 dy|_{M_T=2.6\text{GeV}, p_T=0\text{GeV}}}$$

Asakawa, Ko, Levai, PRL70, 398(1993).

Comparison to STAR dilepton data

Transverse momentum spectra



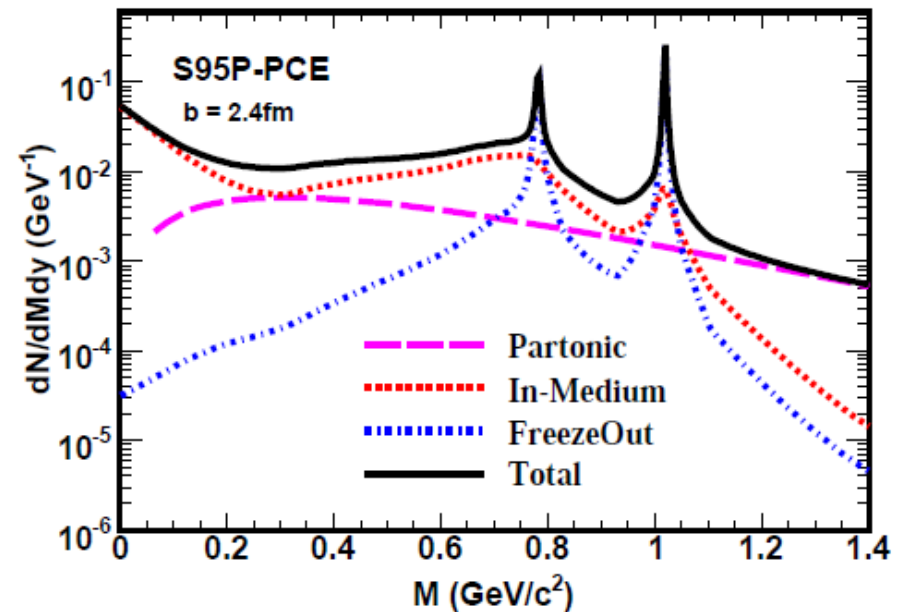
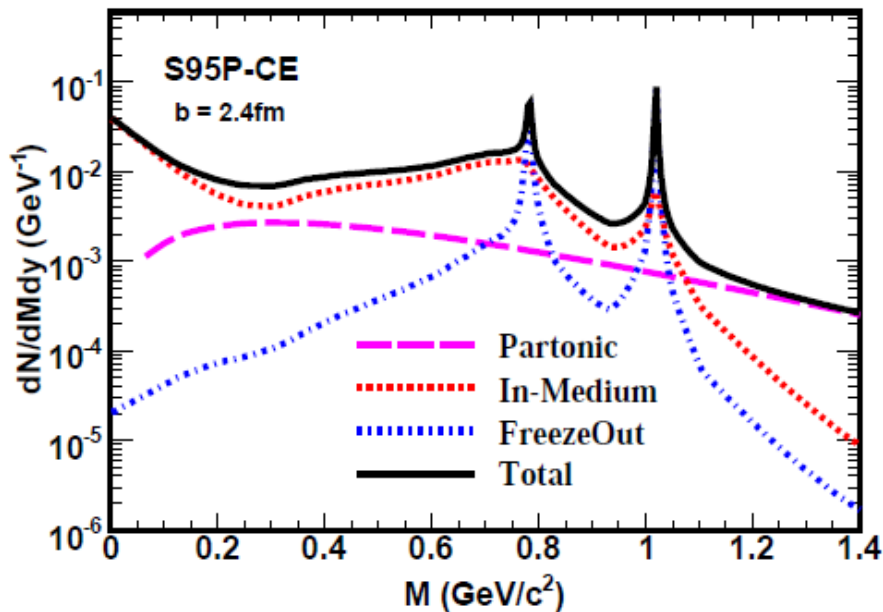
Parameters

- S95-CE:
 $e_0=40 \text{ GeV}/\text{fm}^3$,
 $\tau_0=0.4 \text{ fm}$,
 $T_0=375 \text{ MeV}$,
 $T_f=128 \text{ MeV}$

- S95-PCE:
 $e_0=60 \text{ GeV}/\text{fm}^3$,
 $\tau_0=0.4 \text{ fm}$,
 $T_0=413 \text{ MeV}$,
 $T_f=140 \text{ MeV}$

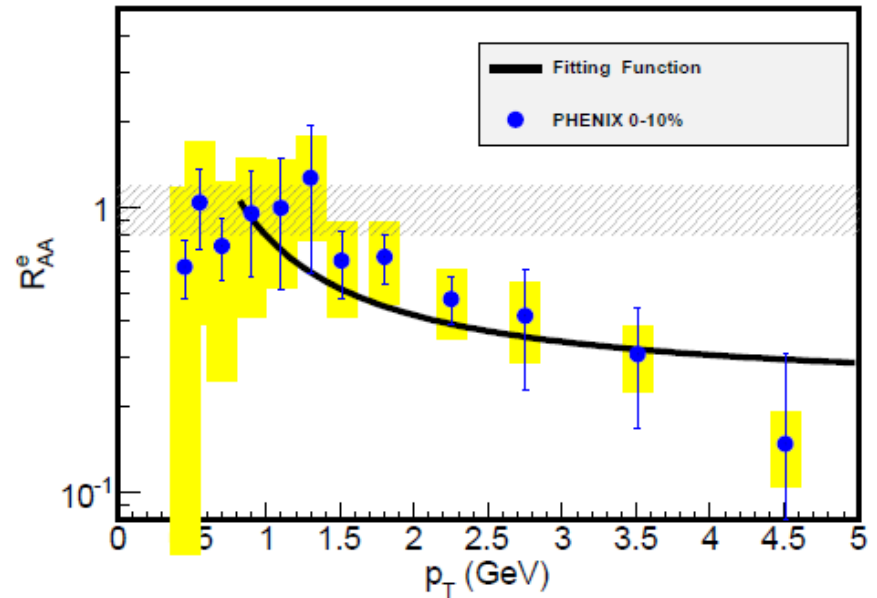
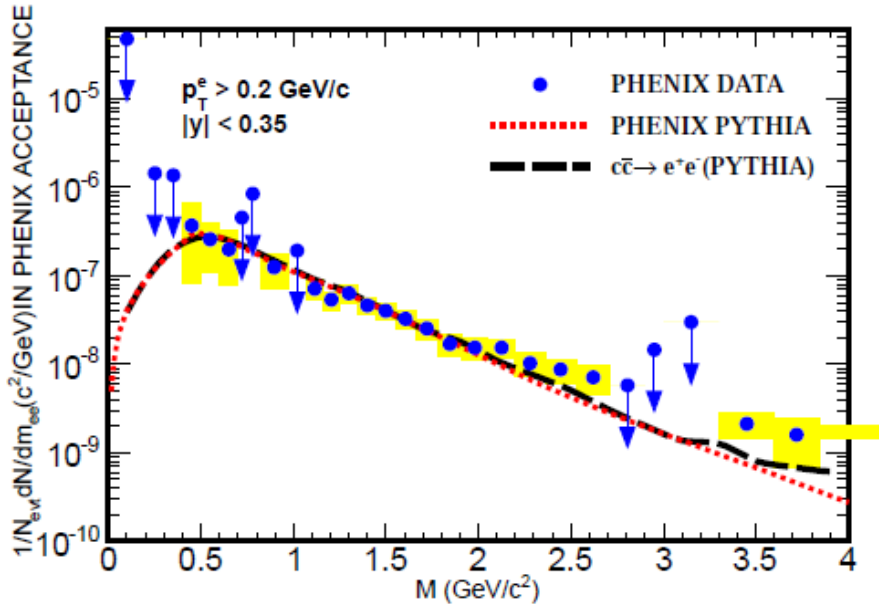
	π	ρ/π	ω/π	ϕ/π
CE	312.65	5.95×10^{-2}	5.7×10^{-2}	1.27×10^{-2}
PCE	349.47	7.8×10^{-2}	9×10^{-2}	2.8×10^{-2}
PHENIX	281.8	1.03×10^{-1}	8.98×10^{-2}	2.14×10^{-2}
STAR	-	1.69×10^{-1}	-	2.65×10^{-2}

Cocktail of in-medium and freezeout contribution in full space



Total = Partonic + In-medium hadronic + Freezeout

Semi-leptonic decays of charm hadrons

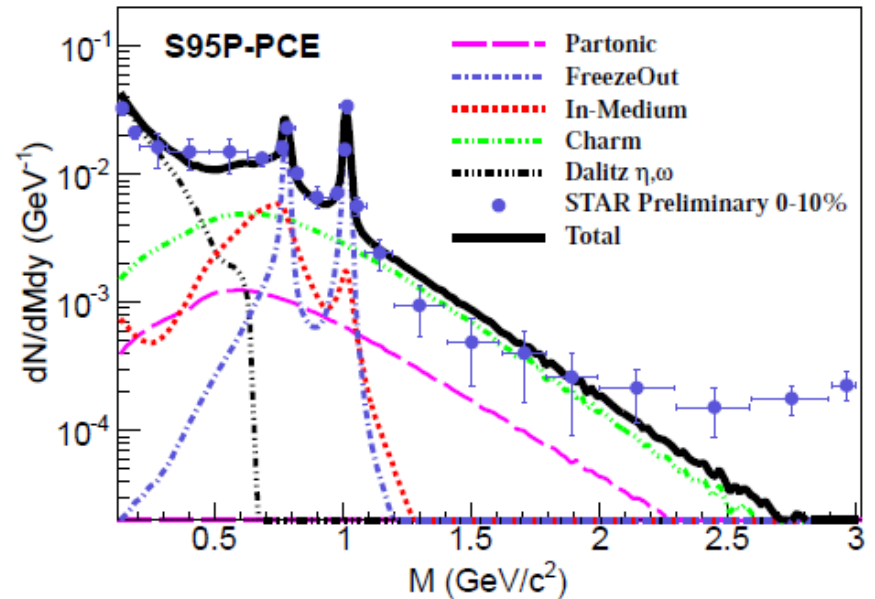
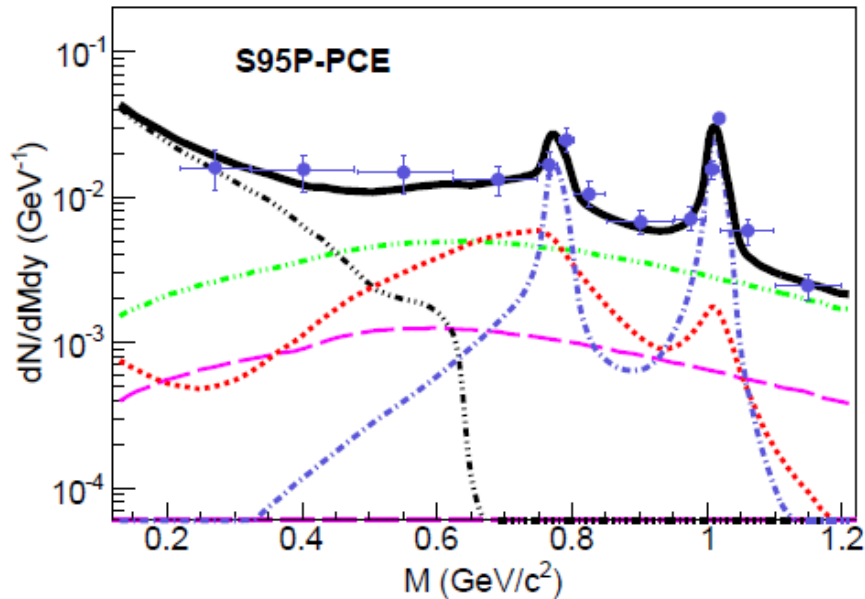


Re-scaled di-electron cross section from charm mesons of semi-leptonic decays in p+p collisions by PYTHIA.

Branch ratios from PDG: $D^\pm, D^0, D^s, \Lambda^c$

the nuclear modification factor for nonphotonic electrons in central Au+Au collisions from the PHENIX.

Comparison with STAR preliminary data



Comparison with STAR preliminary data in most central (0-10%) Au+Au collisions with the STAR acceptance.

$$\text{STAR acceptance: } |\eta^e| < 1, p_T^e > 0.2 \text{ GeV}, |y^{ee}| < 1$$

Summary and Conclusion

- T_{eff} of di-lepton can serve as a probe to EOS of the dense matter in high energy HIC
- Rho meson self-energy from meson resonances below 1300 MeV and baryon resonances (from $\rho NN^* + \rho N\Delta^*$ couplings) are taken into account
- In-medium and freezeout contributions are identified, open charm contribution is modeled
- Comparison with STAR data is made with good agreement

Thanks!