

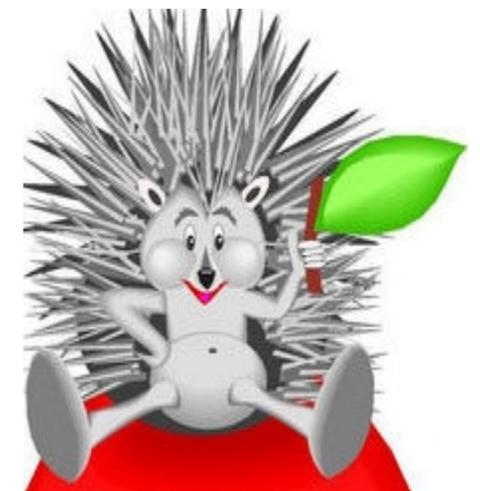
BNL workshop on
E/M probes. Dec.2012



Stony Brook
University

New ideas on the Penetrating Probes

Edward Shuryak
Stony Brook University



outline

time=0



glasma
time



outline

- **“hot glue scenario”** (ES,1992) + **comment on anisotropy vs dilepton polarization** (ES .2012)

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 $\Rightarrow e+e^-$ (Staig,ES,2009)
coherent photon+q $\Rightarrow q + \text{real photon}$
(Liao,ES,in progress)

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- **B+photon \Leftrightarrow 0^+ glue (scale anomaly) inchannel and G^2BA interaction** (Basar,Kharzeev,Skokov 2012)
- **(DIS on GLASMA) or (hydro stress tensor coupled to 2 photons) B A \Leftrightarrow $2^+, 0^+$ glue $O(T_{\mu\nu}AA)$** (Basar,Kharzeev,ES in

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QGP

near T_c

- photon radiation rates \leq quarks and gluons are not the only quasiparticles, **monopoles appear in large number and help to explain large scattering rates** (small viscosity). We need to include those collisions in photon rates
- news on chiral symmetry restoration **hadronic time**

Two-Stage Equilibration in High Energy Heavy Ion Collisions

E. Shuryak

Department of Physics, State University of New York at Stony Brook, Stony Brook, New York 11794

(Received 9 March 1992)

Using the (lowest-order) perturbative QCD, we argue that high energy heavy ion collisions proceed via two stages: *equilibration of gluons* takes time $\tau_g \sim \frac{1}{2}$ fm/c, while *production and equilibration of quarks* needs time at least $\tau_q \sim 2$ fm/c. If so, the initial gluon plasma is much hotter than usually estimated, $T_g \sim 400$ MeV, which leads to enhanced charm production and significant modifications of other proposed signals.

finned by $d\sigma/dt = (\pi\alpha_s^2/s^2)M^2$] are

$$M_{gg \rightarrow gg}^2 = \frac{9}{2} \left(3 - \frac{ut}{s^2} - \frac{us}{t^2} - \frac{st}{u^2} \right),$$

$$M_{gg \rightarrow \bar{q}q}^2 = \frac{1}{6} \frac{(u^2+t^2)}{ut} - \frac{3}{8} \frac{u^2+t^2}{s^2},$$

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$$M_{q_1q_2 \rightarrow q_1q_2}^2 = \frac{4}{9} \frac{s^2+u^2}{t^2}.$$

Spectra of the produced photons and dileptons should also be significantly modified in this scenario: During the "transitory period" ($\tau_g < \tau < \tau_q$) one has *smaller* number of quarks, but those are *hotter*. The reason is again that $gg \rightarrow \bar{q}q$ is dominated by *small angles*, so the produced quarks have the same momentum distribution as gluons. As most photons and dileptons to be observed actually correspond to the *tails of the distribution functions*, it is important that their relaxation happens *from above*.

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charm contribution needs to be subtracted to test early dileptons $M=1..3$ GeV

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In this note we discuss how angular distribution of the dileptons produced in heavy ion collisions at RHIC/LHC energies can provide an information about a degree of local equilibration of the quark-gluon plasma produced at different invariant mass regions.

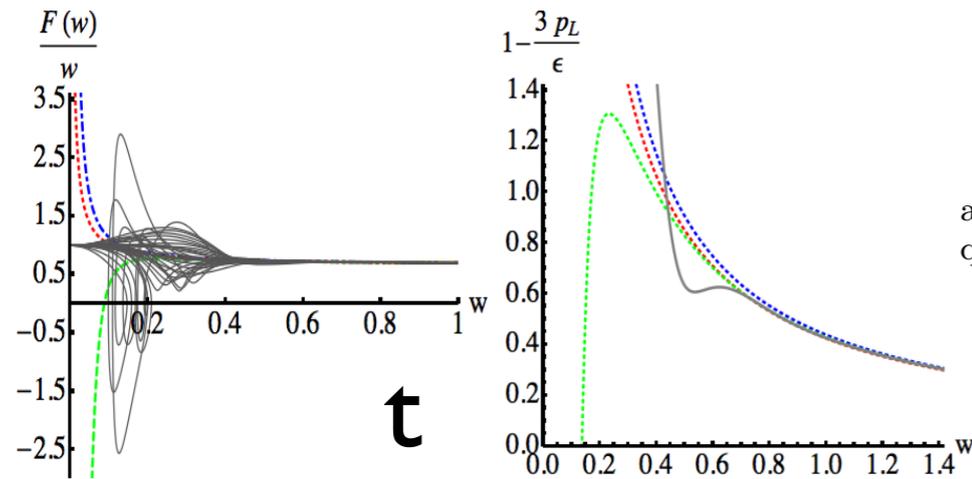


FIG. 1. a) $F(w)/w$ versus w for all 29 initial data. b) Pressure anisotropy $1 - \frac{3p_L}{\epsilon}$ for a selected profile. Red, blue and green curves represent 1st, 2nd and 3rd order hydrodynamics fit.

$w = \tau T_0 = \tau / (.5 \text{ fm})$ at LHC

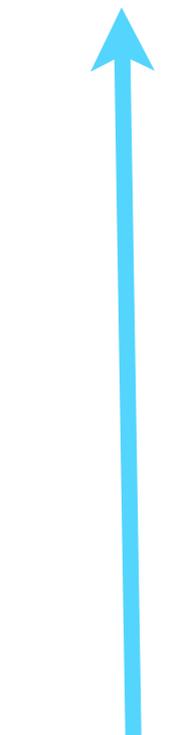
AdS/CFT collisions with various initial conditions converge to the same hydro, and they do so **when anisotropy is still large!**

τ



$$\frac{d\sigma}{d\Omega} \sim (1 + a \cos^2 \theta)$$

$M_{e^+e^-}$



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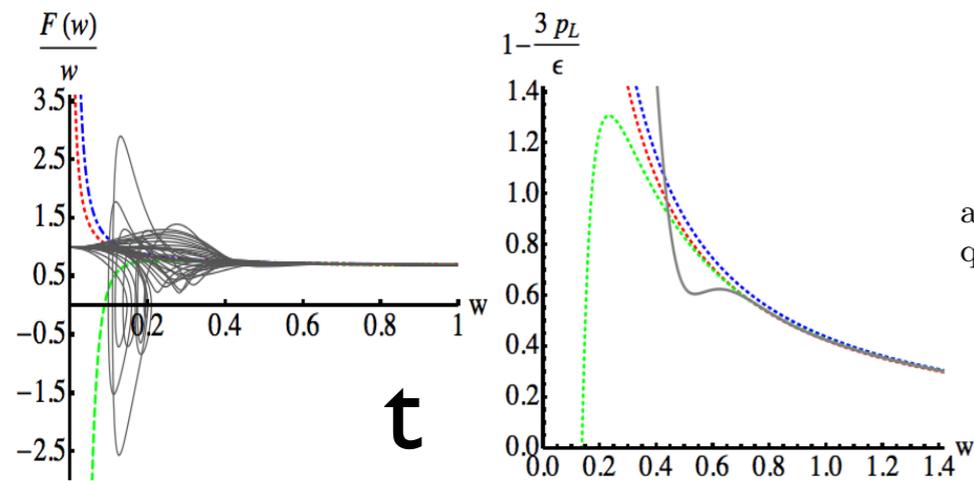


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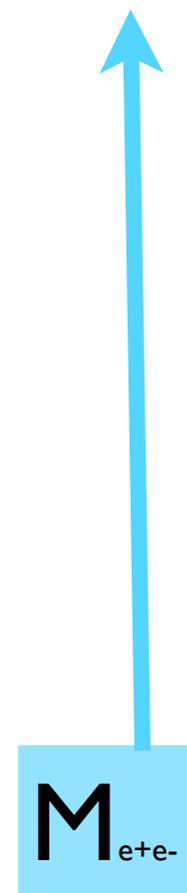
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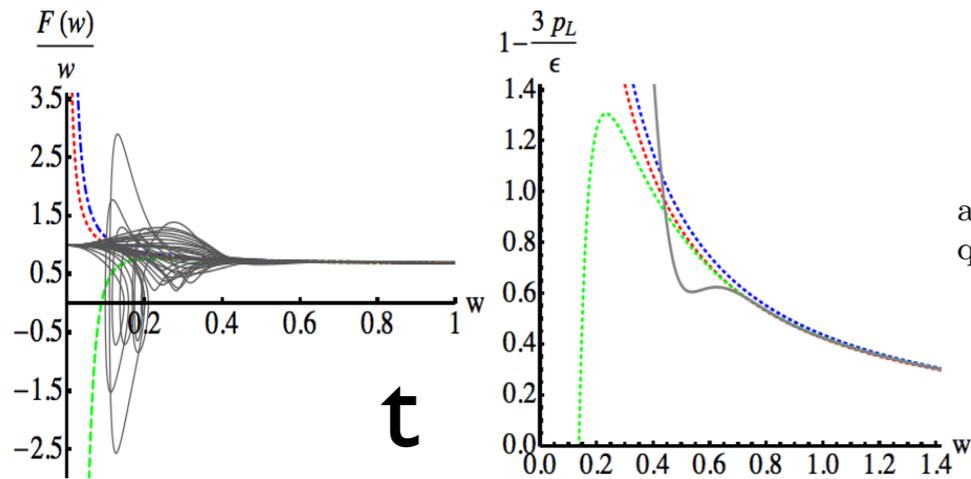


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 $a=1$ for DY (partons come from the beam direction)
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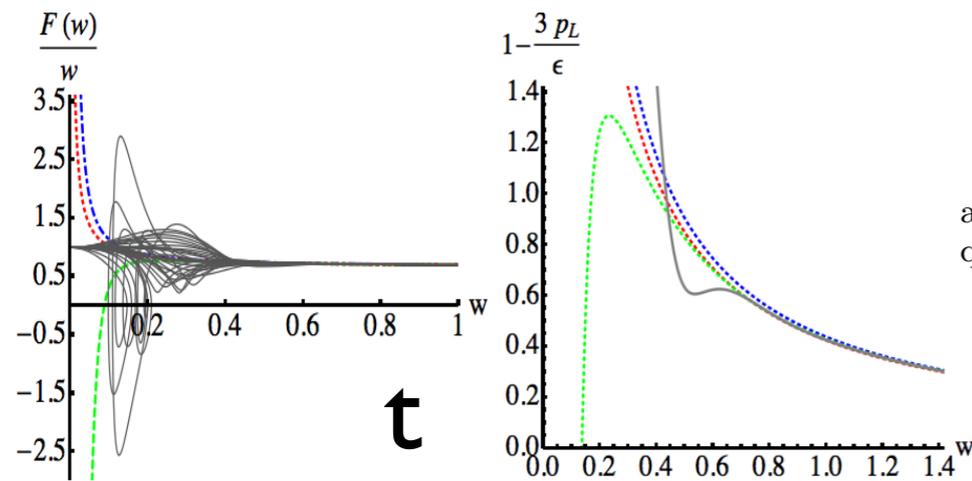


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Perhaps STAR/new PHENIX can catch out-of-equilibrium stage via $a < 0$

Production of soft e^+e^- Pairs in Heavy Ion Collisions at RHIC by Semi-coherent Two Photon Processes

Pilar Staig and Edward Shuryak

Department of Physics and Astronomy

Stony Brook University, Stony Brook NY 11794 USA

(Dated: May 20, 2010)

“exotic”

We calculate the contribution of the two photon production process into e^+e^- spectra, and compare the results with experimental data from the PHENIX detector at RHIC. We study the contribution given by “semi-coherent” kinematics, in which one photon is relatively hard and is incoherently emitted by participating protons, while another can be soft enough to be in a coherent domain.

coherent Z^2

incoherent $P's Z$

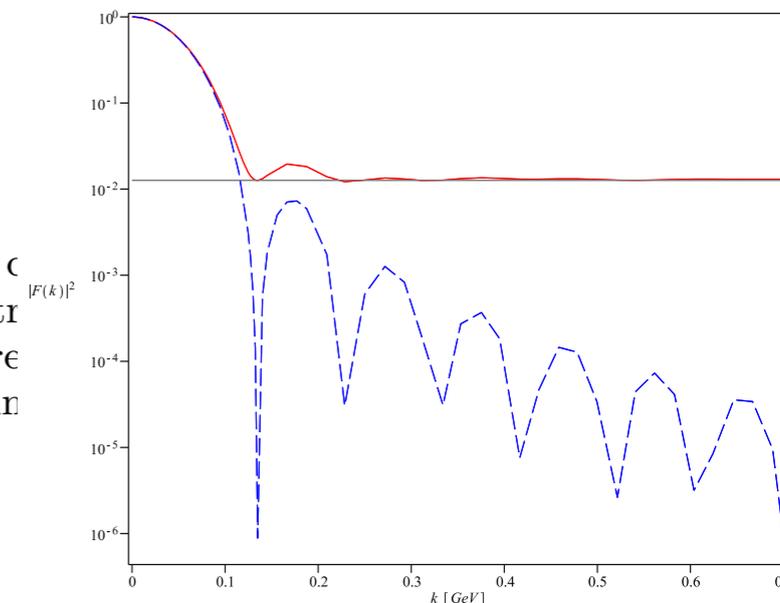
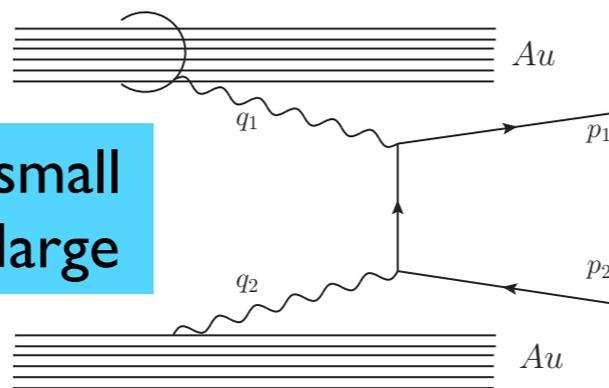


FIG. 3: (Color online) The square of the form factors plotted on a logarithmic scale. The (blue) dashed line corresponds to smooth Woods-Saxon charge distribution, the (red) continuous line corresponds to resolved discrete protons (but not quarks), as explained in the text.)

q_1 small
 q_2 large



- ultra-peripheral is coherent-coherent
- parton processes are incoherent-incoherent
 $gq \Rightarrow \gamma q$
- semi-coherent has coherent field with small q_1 and incoherent one with large q_2 (resolving protons)

arXiv:1005.3531v1 [nucl-th] 19 May 2010

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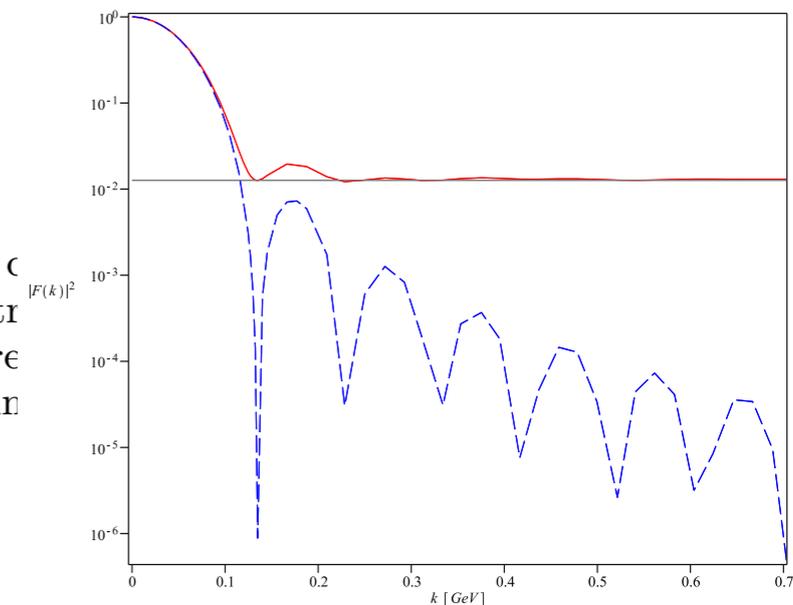


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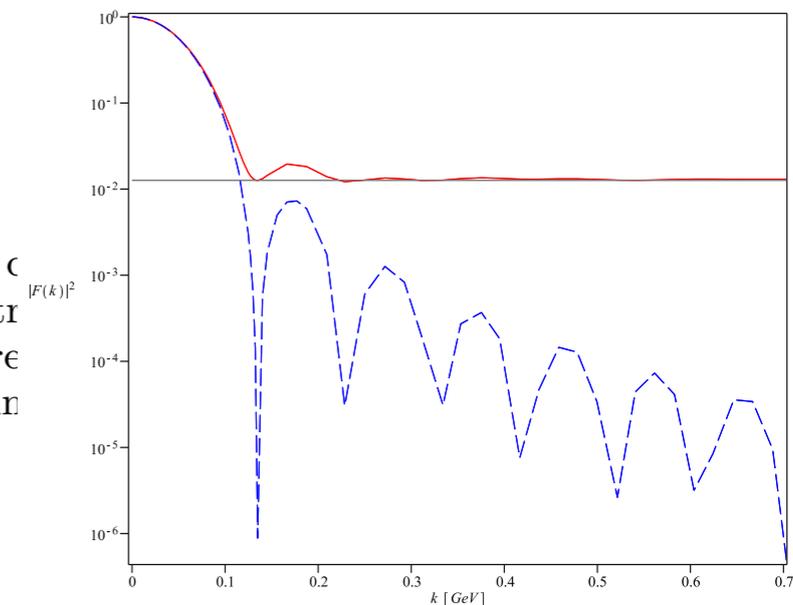


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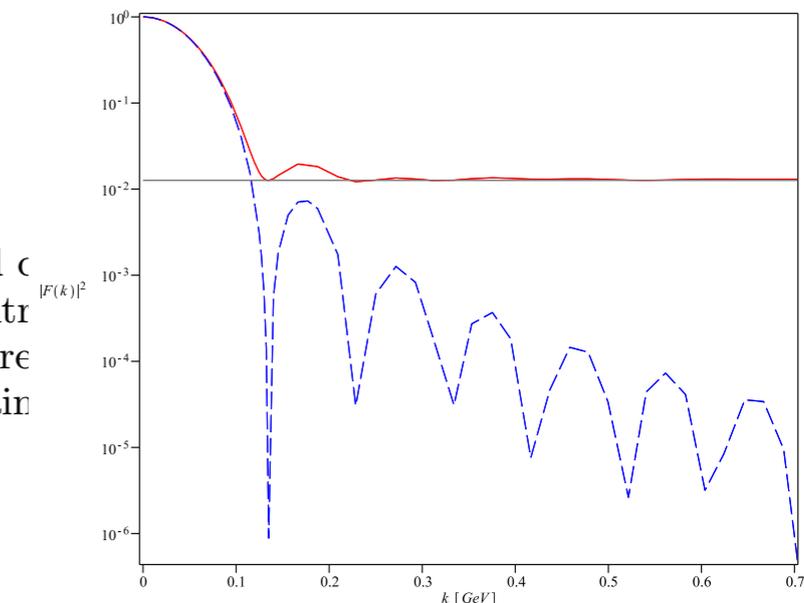
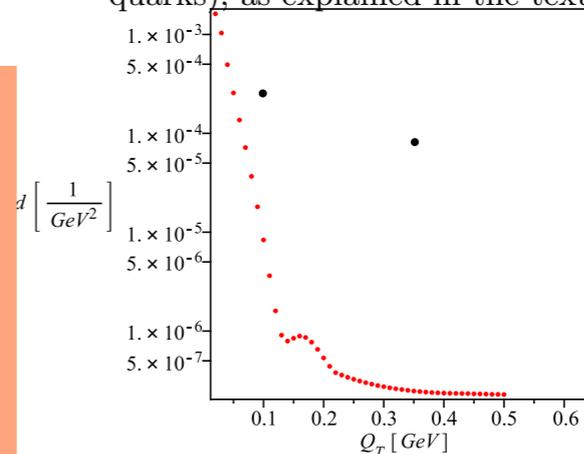


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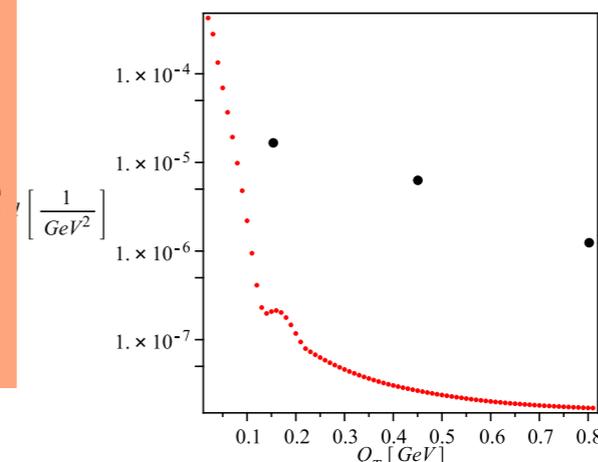
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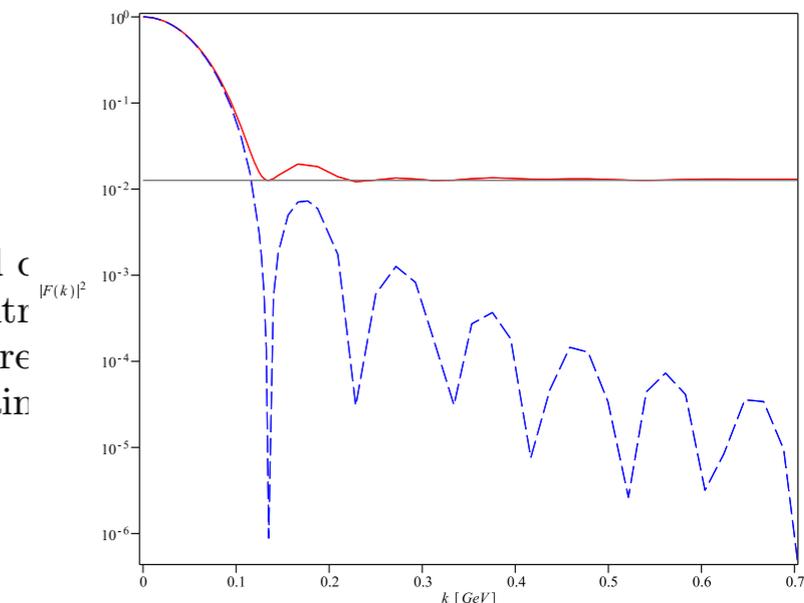
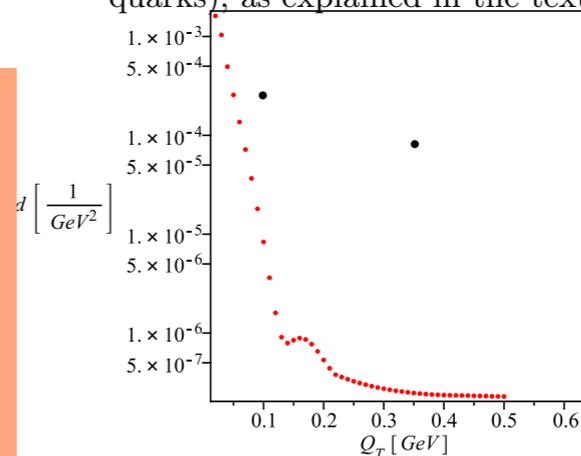


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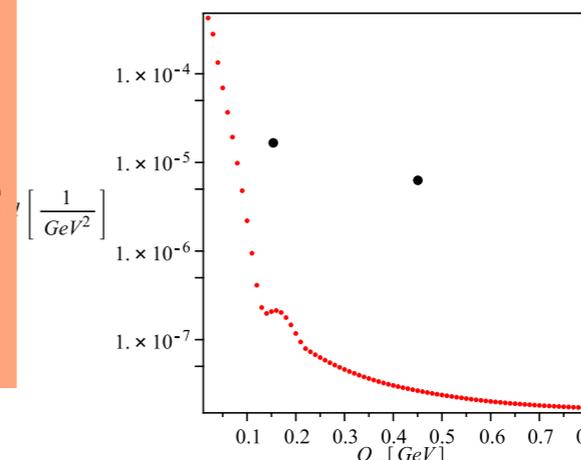
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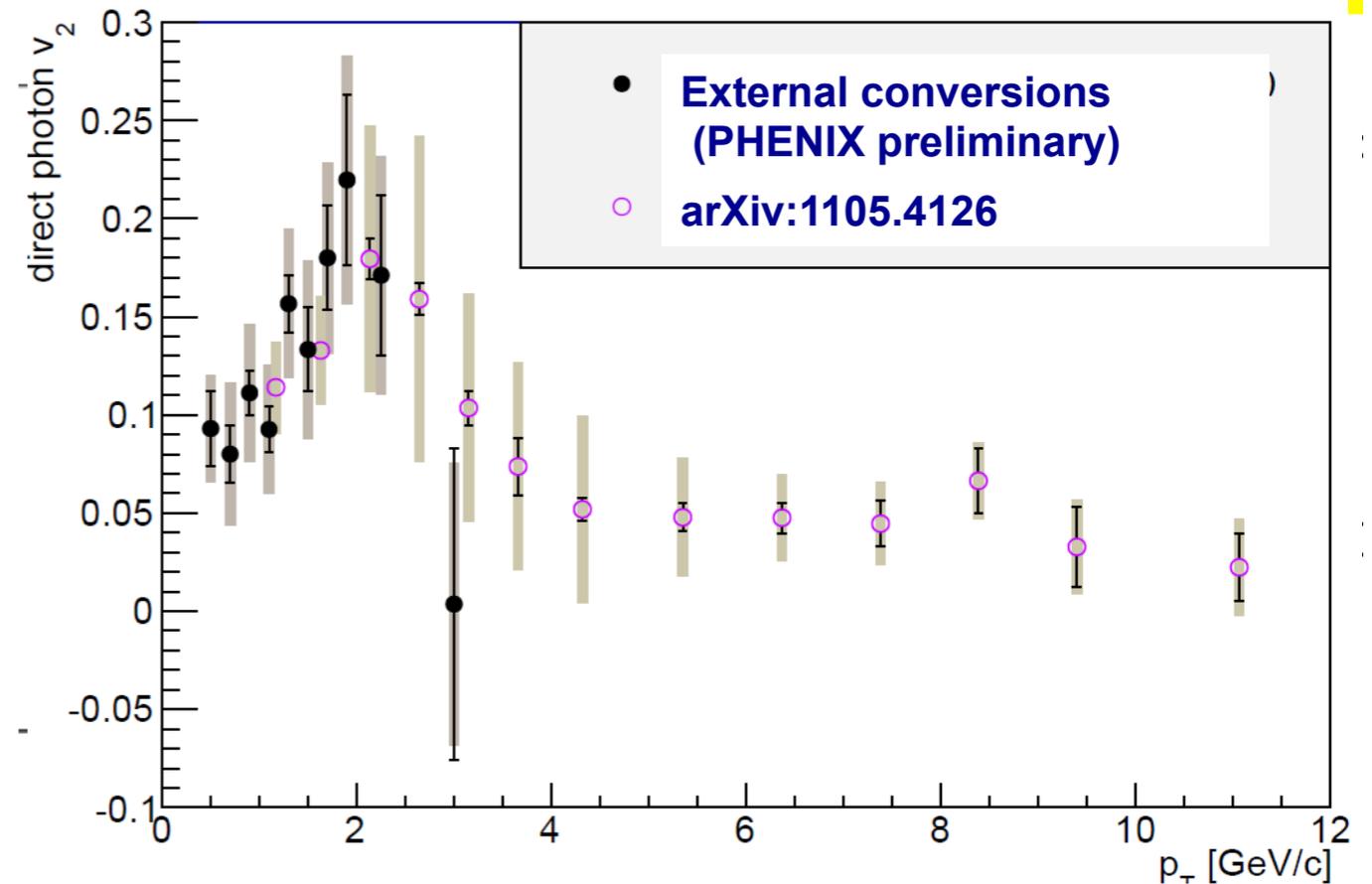


M=0.81..0.99 GeV

arXiv:1005.3531v1 [nucl-th] 19 May 2010

New puzzle: large v_2 of photons (Phenix)

=> v_2 is as large as that for hadrons
=> seems to persist even at large p_t (?)
=> also seen by ALICE



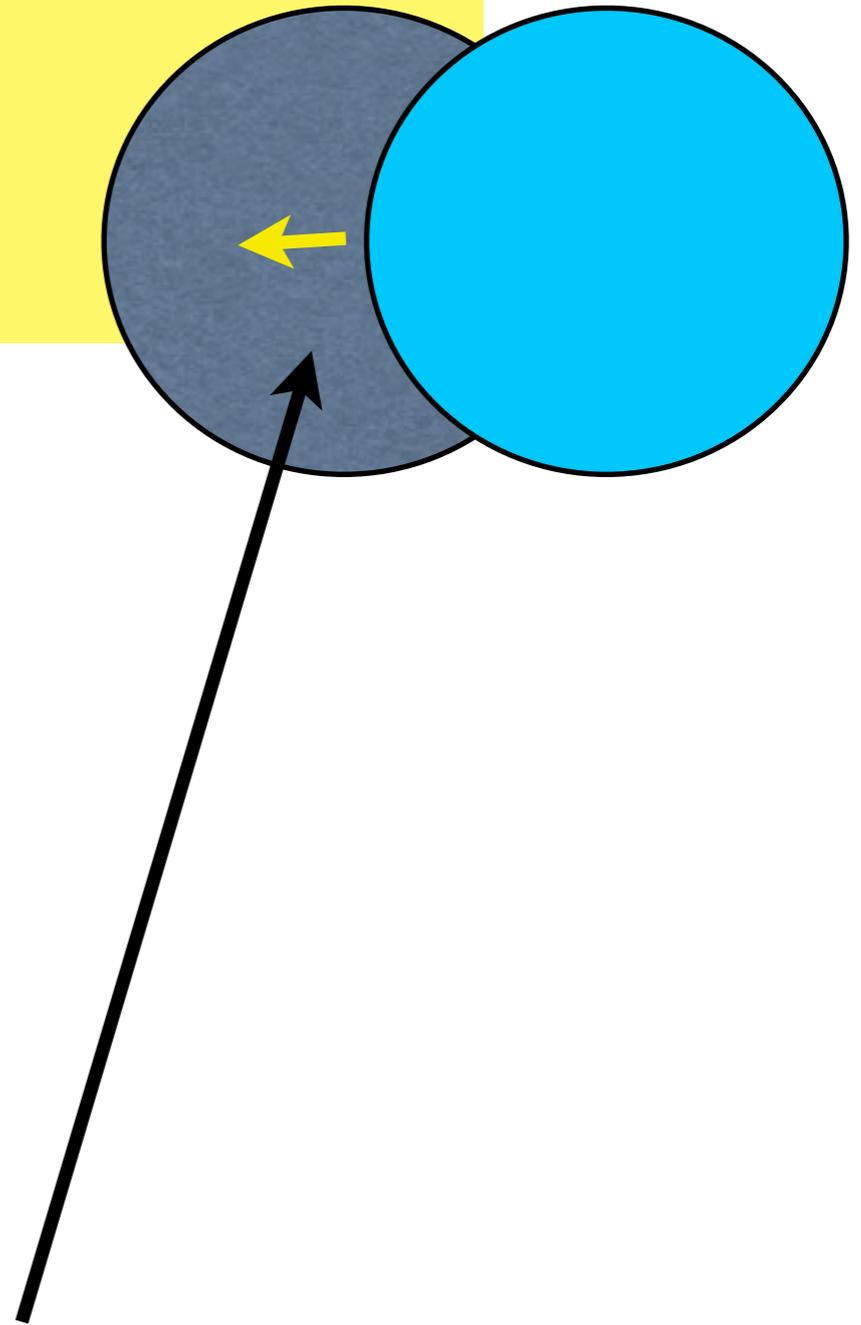
● q^+ (**semi-coherent eff. photon**) \Rightarrow q^+ + gamma

(with J-F. Liao)

(WW approximation)

● q^+ (B-correlated gluon) \Rightarrow q^+ + gamma

$$Z\alpha \sim \alpha_s$$

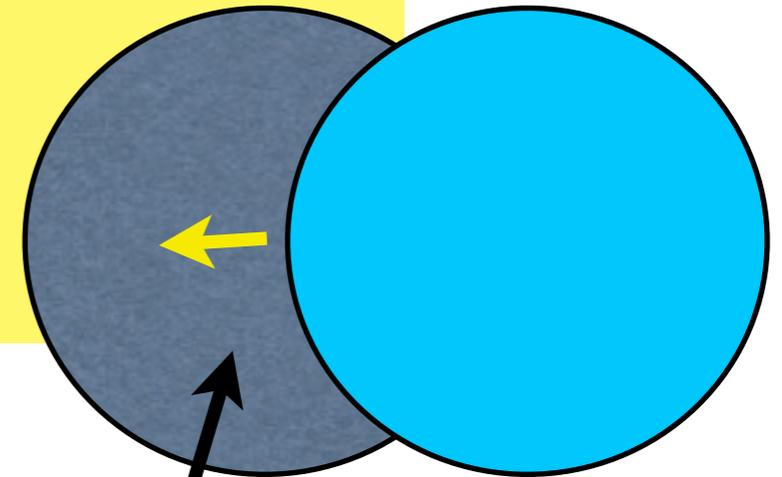


● q^+ (semi-coherent eff. photon) \Rightarrow q^+ + gamma

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while $Z\alpha \sim \alpha_s$

effective gluon density is $O(10)$ larger than equivalent photons,

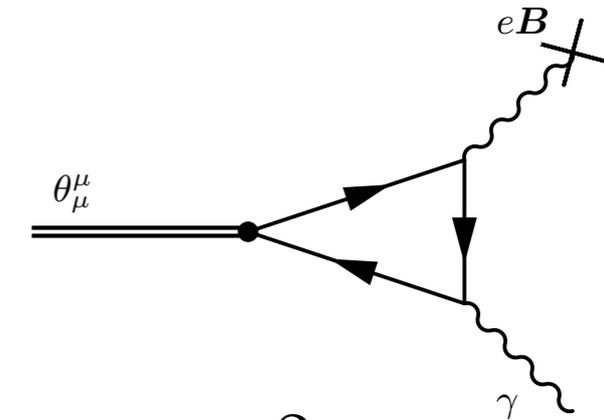
but photon's momentum is **strongly correlated with the impact parameter b**

coherent E/B field of the photon extends beyond the edge of the nuclei

even small effects need to be calculated!

Basar, Kharzeev, Skokov arXiv:1206.1334

- (1) virtual quark loop
- (2) coherent magnetic field
- (3) B+photon mix with scalar G^2
- (4) Its correlators are related to bulk viscosity and correlators

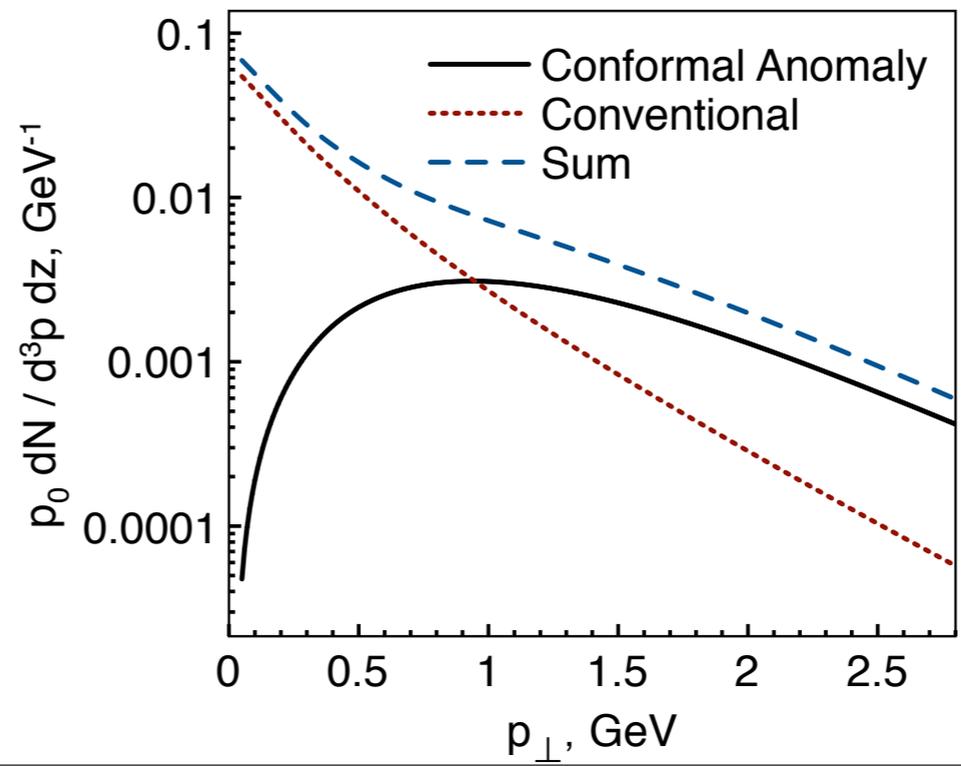
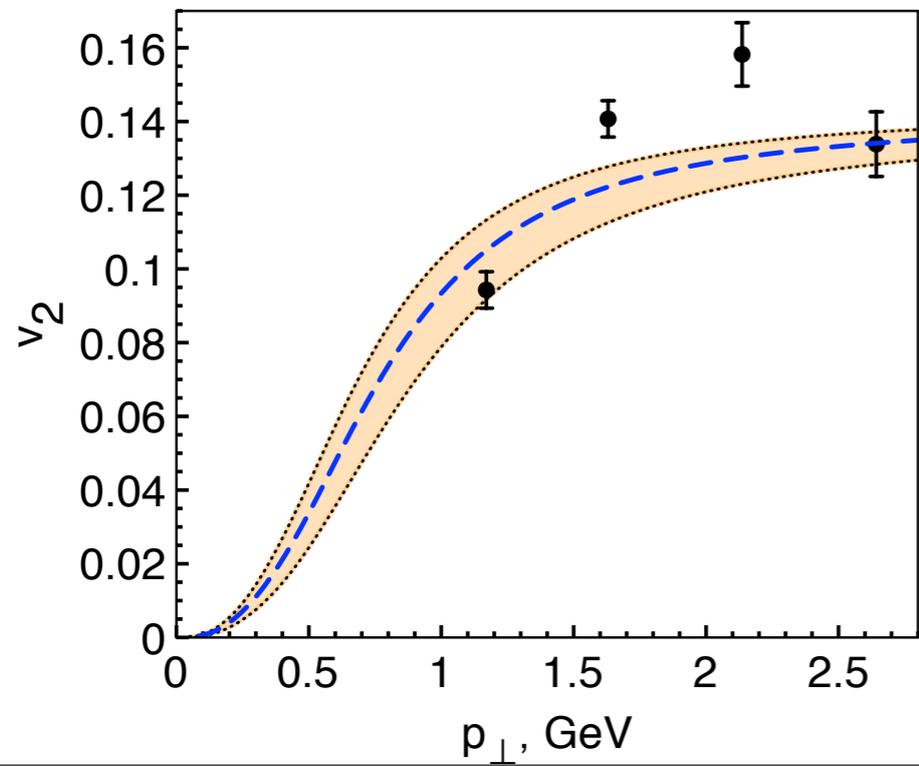


$$\sim (B_y q_x^\gamma)^2$$

Spectral densities known from AdS/QCD

models

no effect along the B field!



DIS on GLASMA

Basar, Kharzeev,
ES, in progress

The tensor channel

$$\vec{E}^2 - \vec{B}^2 \quad \vec{E}^2 + \vec{B}^2$$

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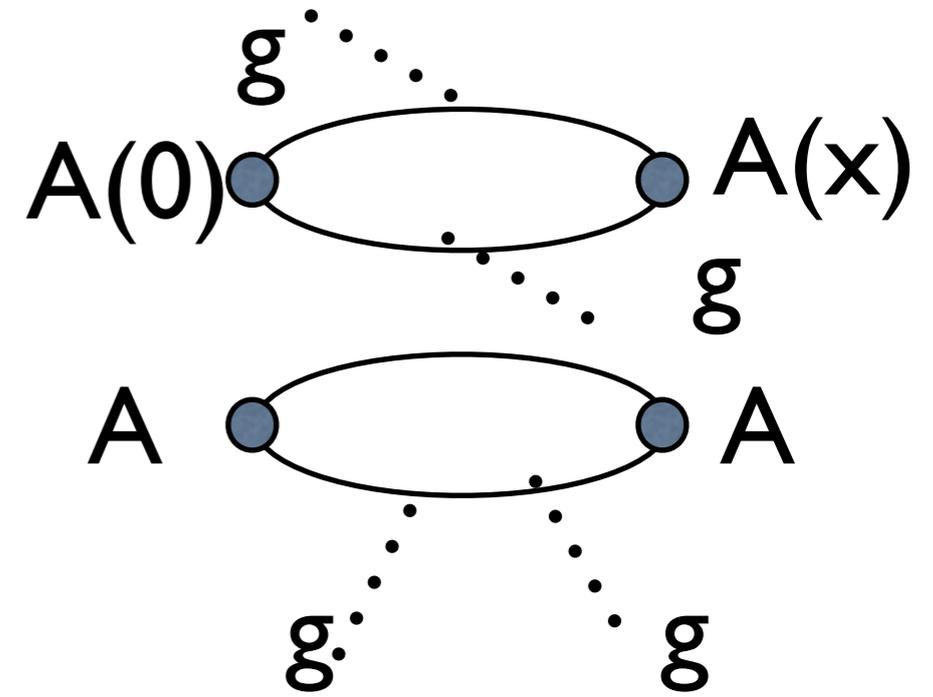
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- 2 OPE cases: (i) $\gamma\gamma$ hard, gg soft; (ii) $\gamma\gamma$ soft, gg hard. Calculation follows [ES, Vainshtein Nucl.Phys. B201 \(1982\) 141](#) and the result is not naive products of two stress tensors

$$L_{eff} = \Pi_{\mu\nu} A_\mu A_\nu$$

$$\Pi_{\mu\nu} = Tr(\gamma_\mu S(q, G) \gamma_\nu S^+(q, G))$$



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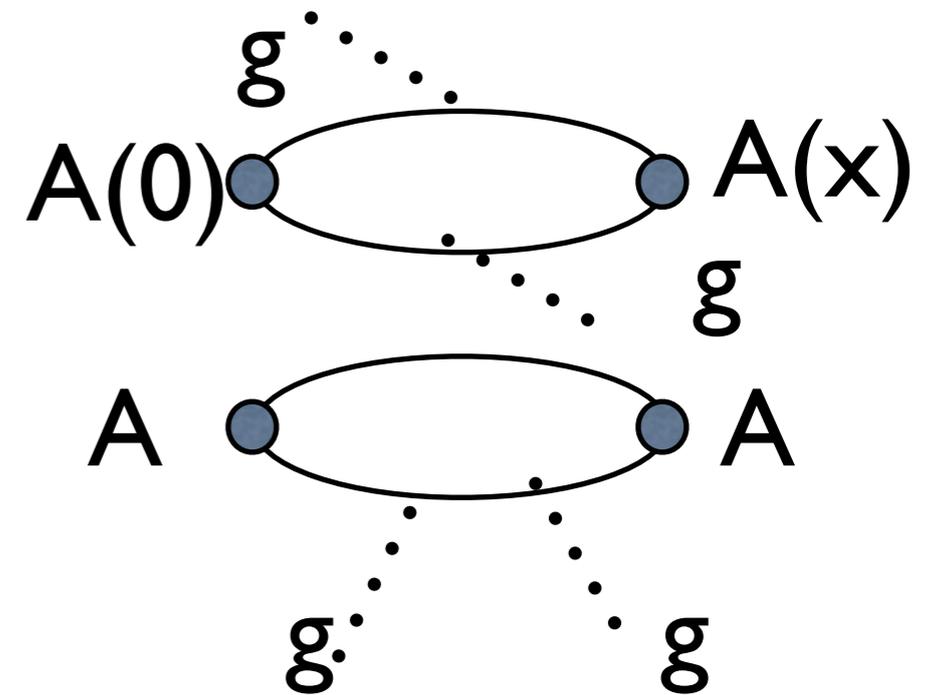
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the usage of the so called fixed-point gauge

$$x_\mu A_\mu(x) = 0 \quad (12)$$

invented by Fock, Schwinger and perhaps others. In this gauge $A_\mu(0) = 0$ and next order terms in x expansion can be written as covariant derivatives of the field strength

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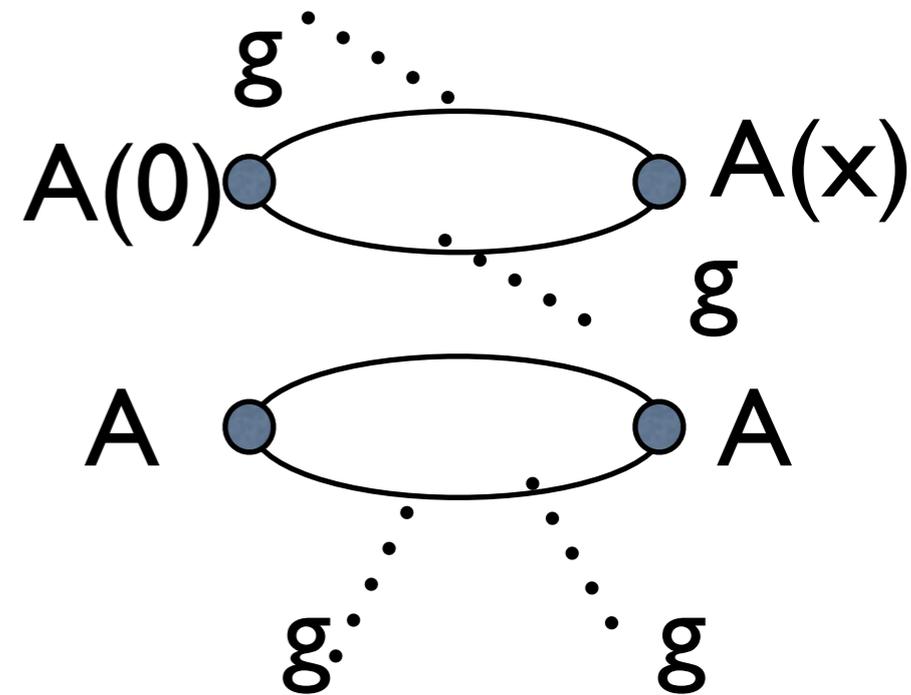
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$$S(q) = \frac{1}{\hat{q}} - \frac{g}{2q^4} q_\alpha \tilde{G}_{\alpha\beta\gamma\beta\gamma 5}$$

$$- \frac{g^2}{2q^8} \hat{q} q_\alpha G_{\alpha\beta} G_{\beta\gamma} q_\gamma + \frac{g^2}{4q^8} q^2 q_\alpha \{G_{\alpha\beta}, G_{\beta\gamma}\} + \gamma_\gamma$$

$$- \frac{g^2}{4q^8} q^2 q_\alpha [G_{\alpha\beta}, G_{\beta\gamma}] - \gamma_\gamma + O\left(\frac{1}{q^6}\right) \quad (15)$$



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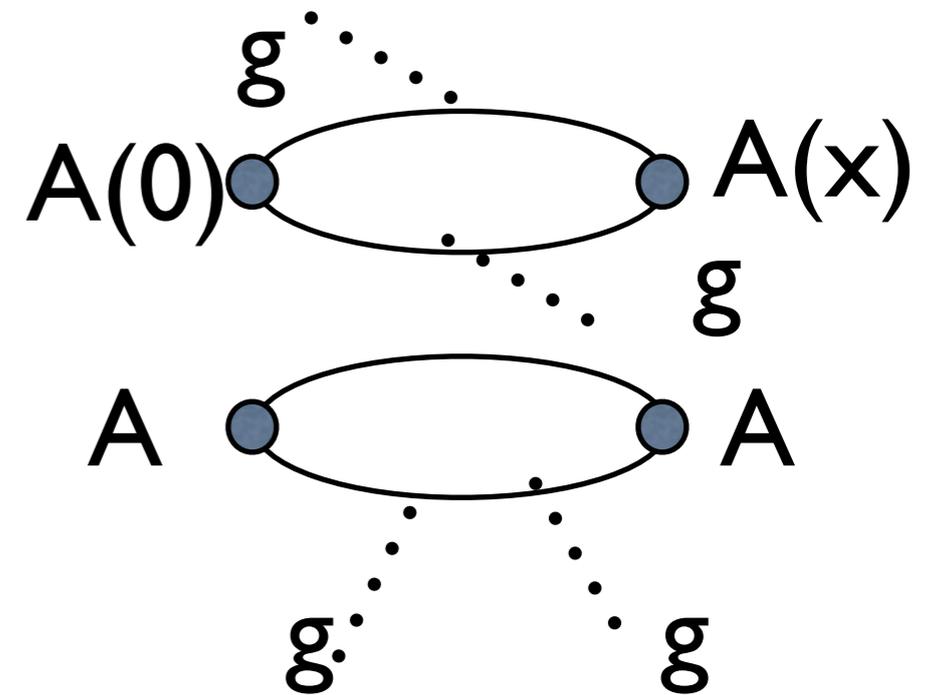
$$A_\mu(x) = \sum_{k=0} \frac{1}{k!(k+2)} x_\nu x_{\alpha_1} \dots x_{\alpha_k} (D_{\alpha_1} \dots D_{\alpha_k} G_{\mu\nu}(0))$$

$$S(q) = \frac{1}{\hat{q}} - \frac{g}{2q^4} q_\alpha \tilde{G}_{\alpha\beta\gamma\beta\gamma 5}$$

$$- \frac{g^2}{2q^8} \hat{q} q_\alpha G_{\alpha\beta} G_{\beta\gamma} q_\gamma + \frac{g^2}{4q^8} q^2 q_\alpha \{G_{\alpha\beta}, G_{\beta\gamma}\} + \gamma_\gamma$$

$$- \frac{g^2}{4q^8} q^2 q_\alpha [G_{\alpha\beta}, G_{\beta\gamma}] - \gamma_\gamma + O\left(\frac{1}{q^6}\right) \quad (15)$$

**=> to x space
=> square (no
integrals!)**



$$L_{eff} = \Pi_{\mu\nu} A_\mu A_\nu$$

$$\Pi_{\mu\nu} = Tr(\gamma_\mu S(q, G) \gamma_\nu S^+(q, G))$$

the usage of the so called fixed-point gauge

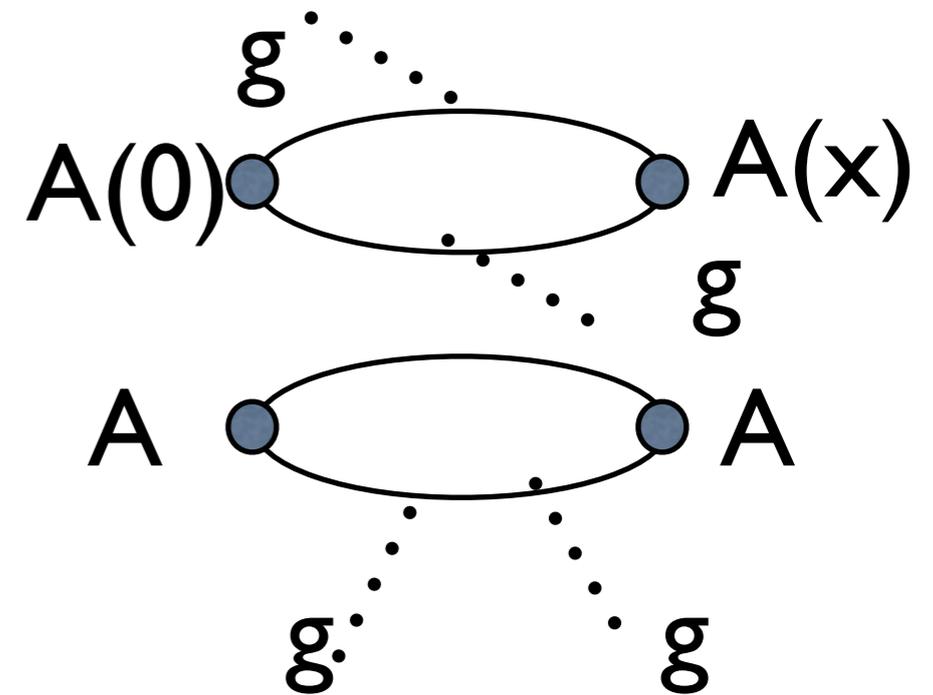
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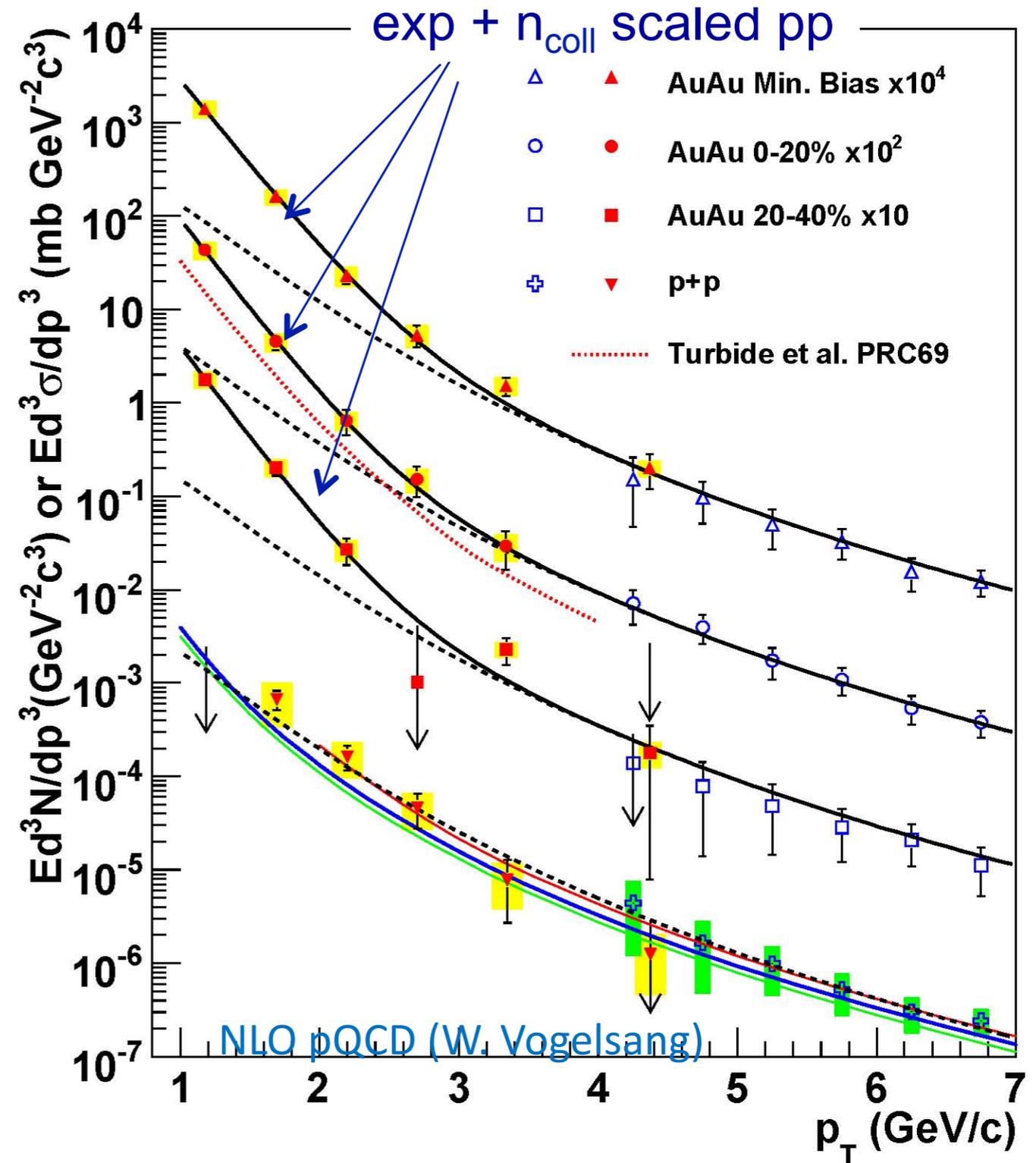
we have local effective interaction GG AA/q4 which of course preserve both gauge invariances



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PHENIX photons

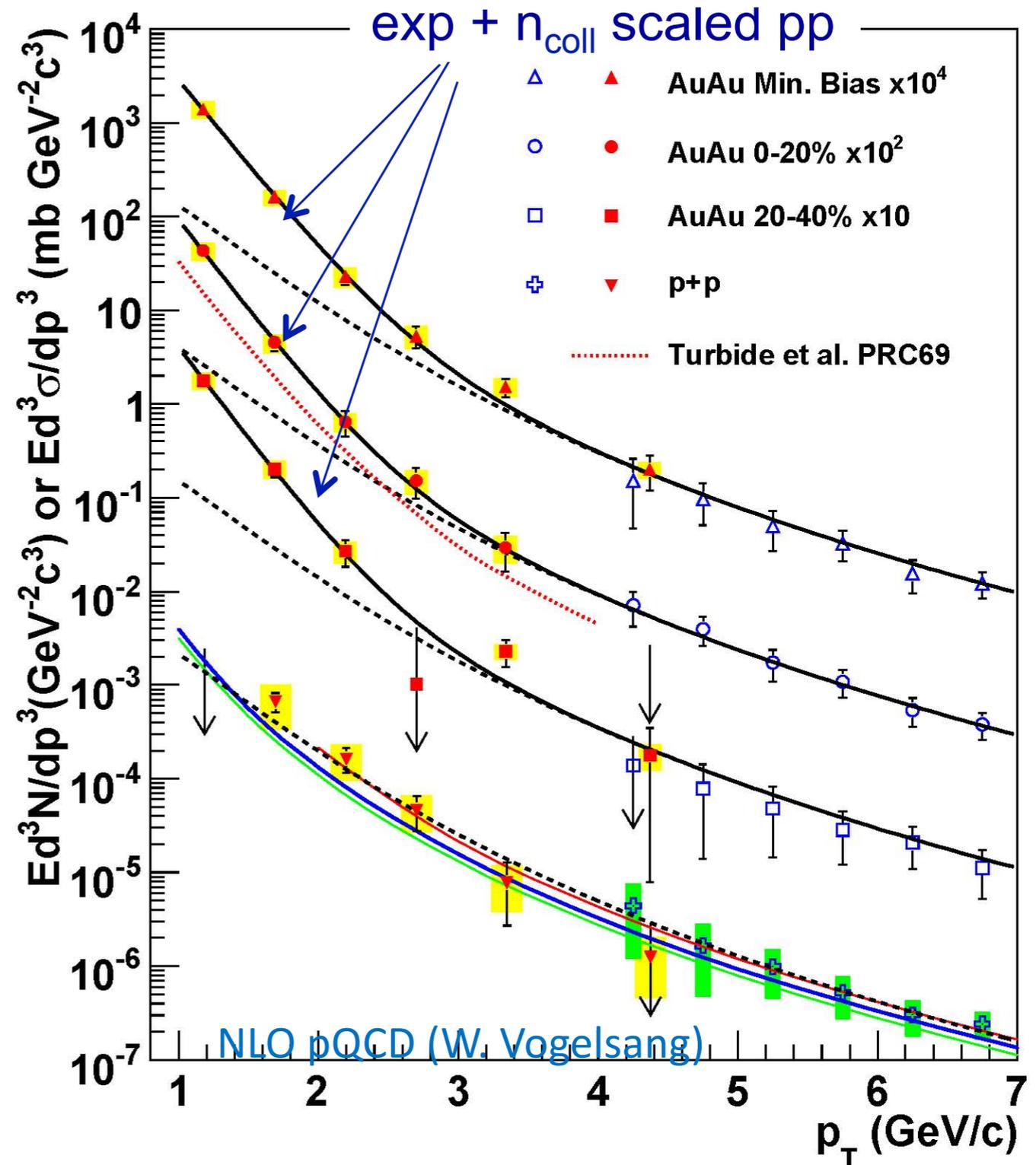
pQCD describes the pp case



PHENIX photons

pQCD describes the pp case

Note that Turbide et al miss a factor 2..3 in the rate.
 Hadronic rate is calculated from a gas of pions, K, rho, K*, A1 plus baryons (rho-N resonances)
 QGP rate is the HTL-corrected QCD Compton



Are the rates too small?

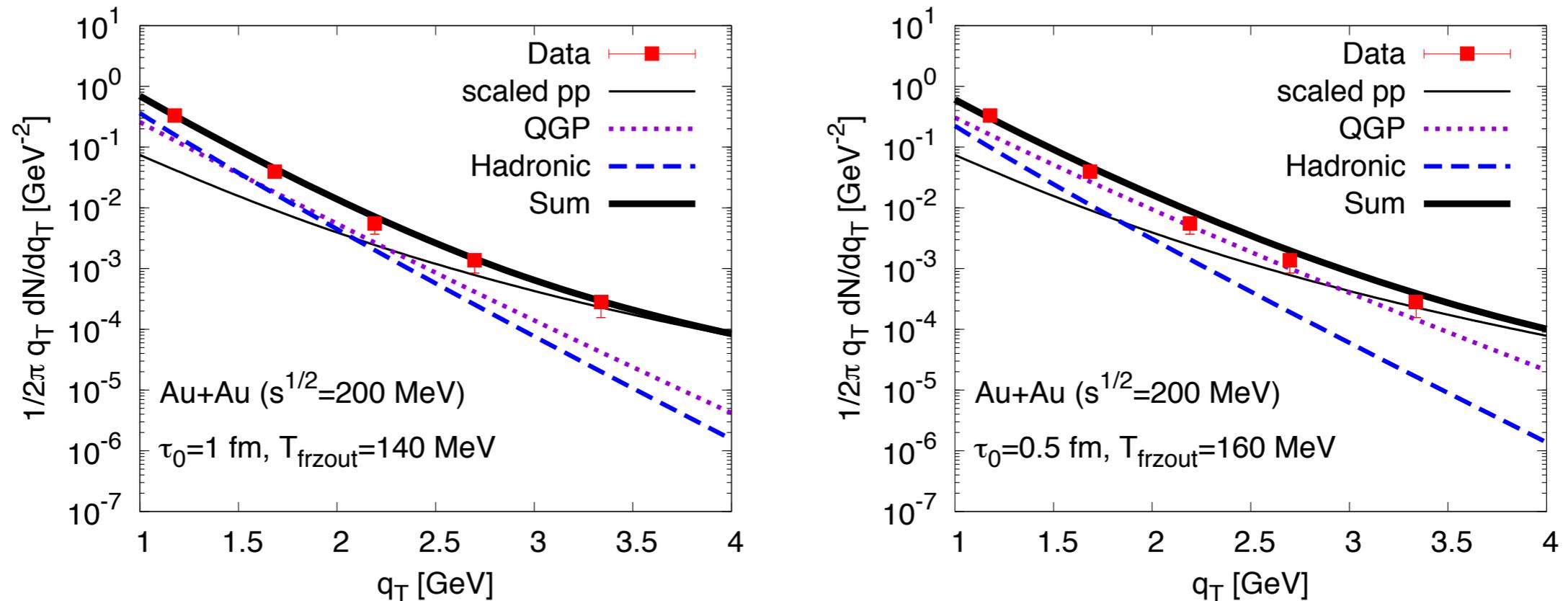


FIG. 9: Photon spectra at RHIC compared to the recent PHENIX data. The left plot is the evolution set RHIC 1 and the right is for RHIC 2.

**seems OK in this work, but
this is for $\alpha_s=0.75$!**

photons from QGP

(my original QGP paper, 1978) + HTL completion to
it (Kapusta et al, PRD44, 1991)

$gq \rightarrow q\gamma, qq \rightarrow g\gamma$

$$q_0 \frac{dR_\gamma}{d^3q} = \frac{6}{9} \frac{\alpha\alpha_s}{2\pi^2} T^2 e^{-q_0/T} \ln \left(1 + \frac{2.912}{4\pi\alpha_s} \frac{q_0}{T} \right),$$

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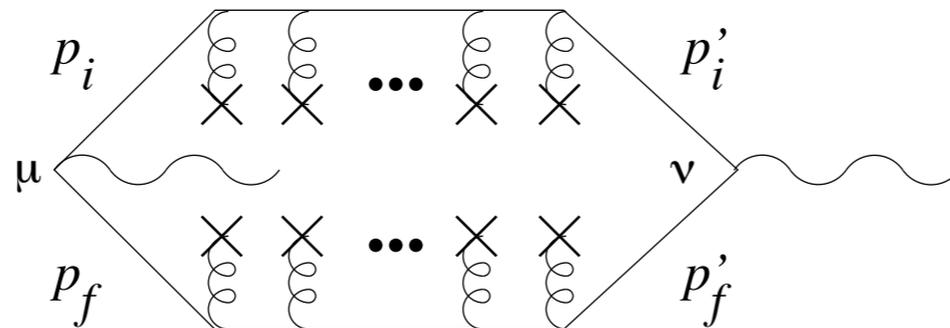
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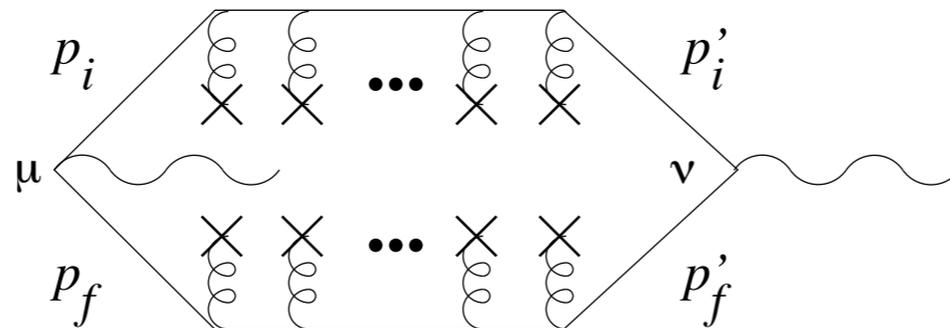


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may work at large p_t ,
otherwise pQCD series does not converge
perhaps, strong coupling methods needed

should there be extra contribution to the
perturbative photon radiation rates?

(one expects it for e.g. **viscosities**)

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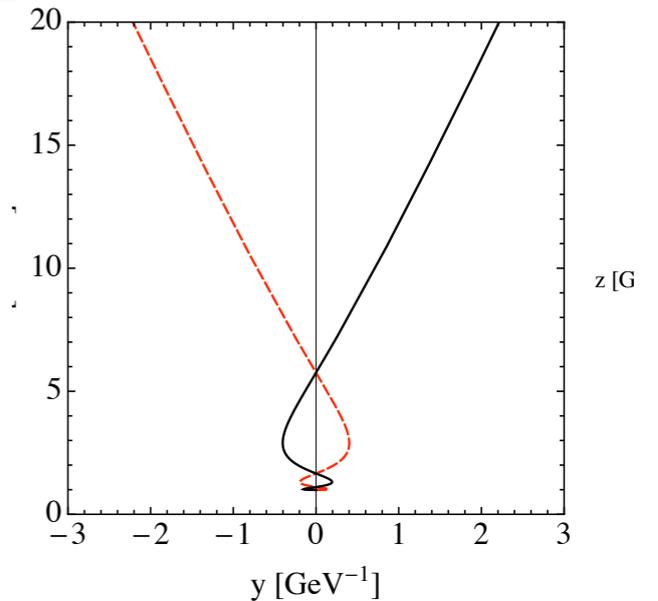
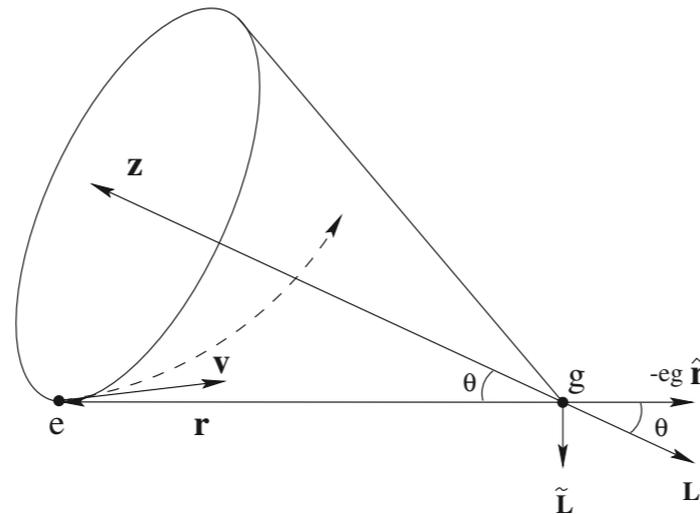
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- Same with electromagnetic rates

monopoles contribute to viscosities, via large angle scattering

$$\vec{B} = g \frac{\vec{r}}{r^3}.$$

$$m \frac{d^2 \vec{r}}{dt^2} = e \vec{v} \times \vec{B} = \frac{eg}{r^3} \frac{d\vec{r}}{dt} \times \vec{r};$$

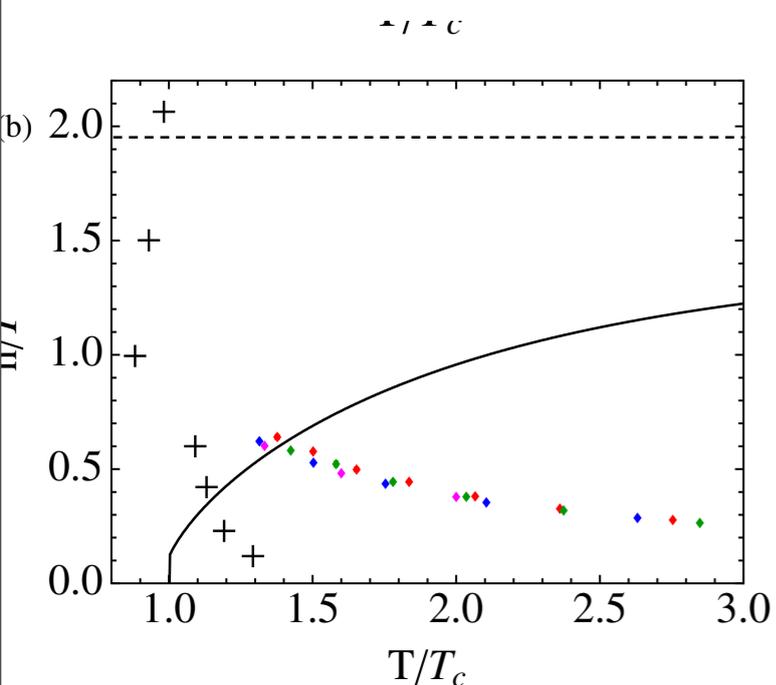


PHYSICAL REVIEW D **80**, 034004 (2009)

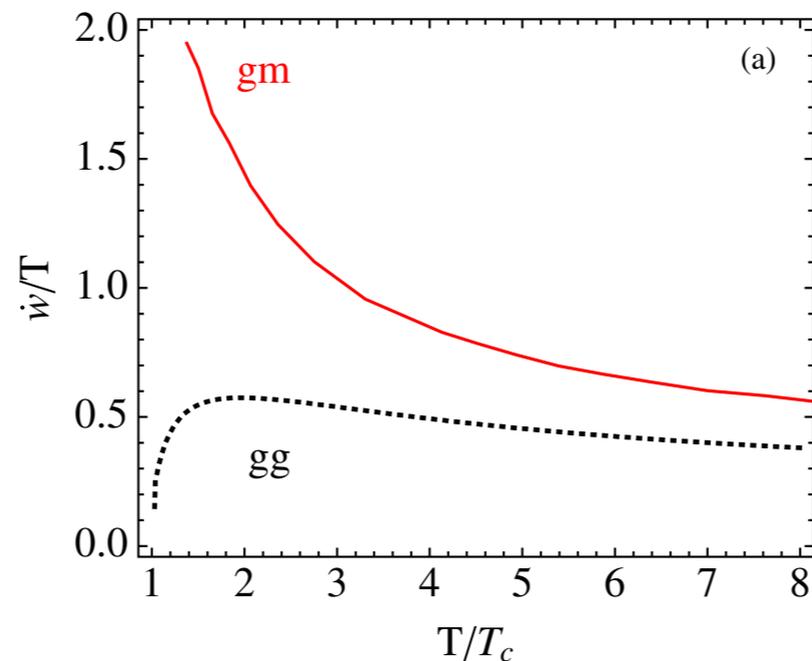
Role of monopoles in a gluon plasma The motion of an electric charge in the field of a magnetic monopole.

Claudia Ratti and Edward Shuryak*

Physics Department, State University of New York at Stony Brook, Stony Brook, NY 11794-3840
(Received 4 February 2009; published 5 August 2009)



ROLE OF MONOPOLES IN A GLUON PLASMA



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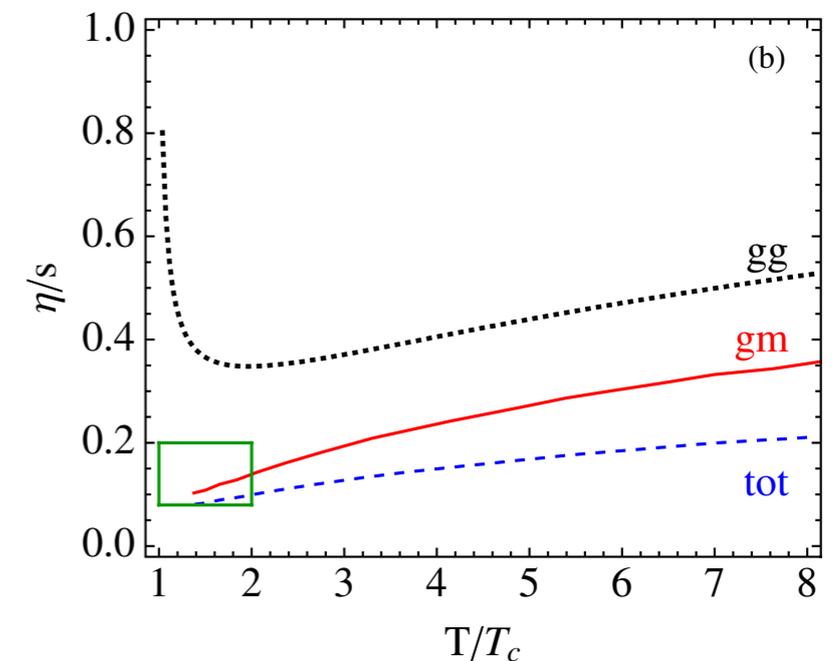
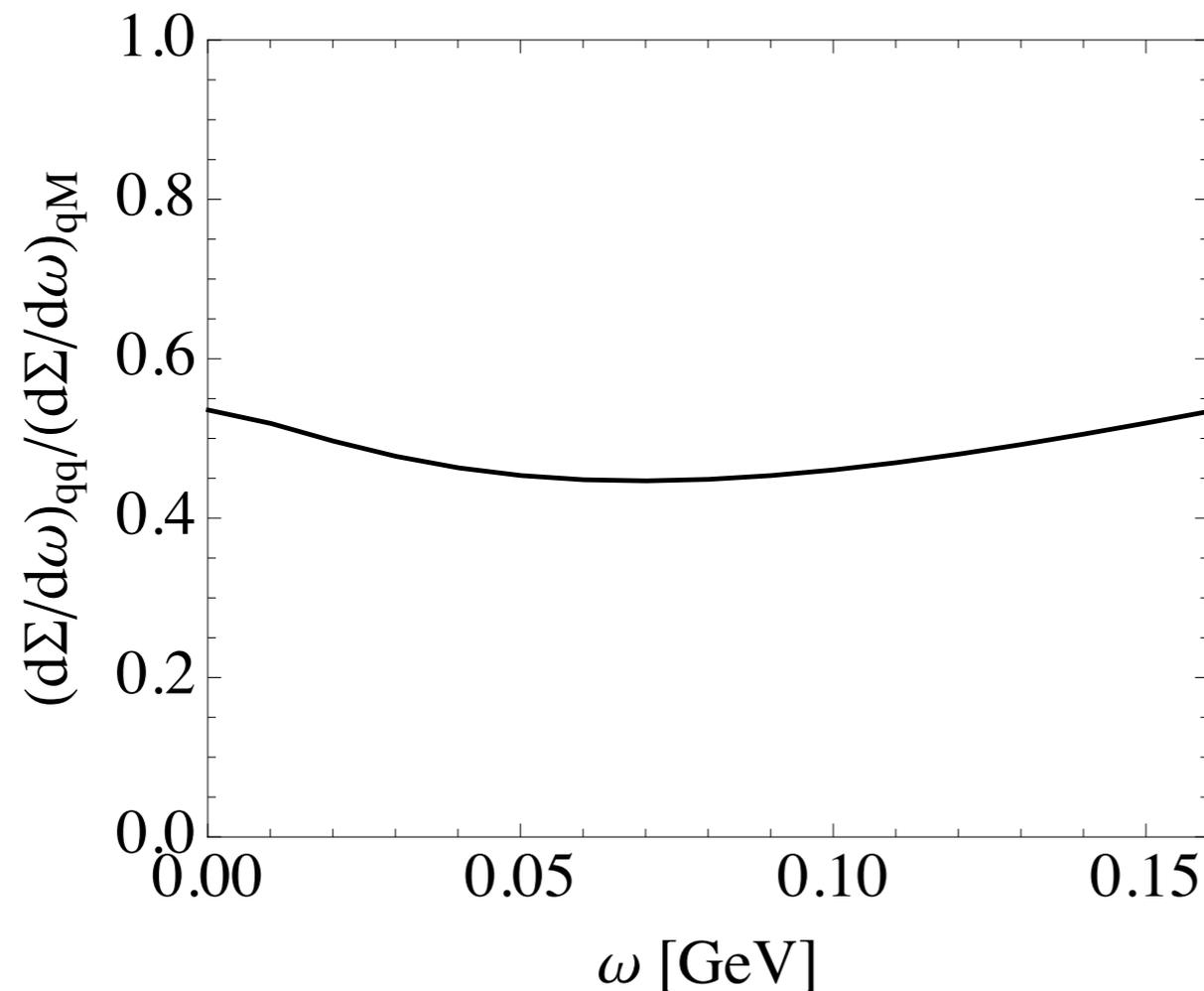


FIG. 15 (color online). (a) gluon-monopole and gluon-gluon scattering rates. (b) gluon-monopole and gluon-gluon viscosity over entropy ratio, η/s . The blue, dashed curve is the total η/s , which is evaluated from the gg and gm contributions. The green box represents the present estimate of η/s in the RHIC temperature regime.

monopoles contribute to the photon radiation:

Radiation of an electric charge in the field of a magnetic monopole

Michael Lublinsky^{a,b}, Claudia Ratti^{a,c}, Edward Shuryak^a



disclaimer:
only soft photons
calculated classically
from the trajectories

qM rate in RHIC condition
is about twice that for qq
(=> qq gamma)

comments about chiral symmetry restoration

- Kapusta+ES, 1993 Weiberg-type sum rules for $\langle VV-AA \rangle$ correlator (Rapp revived recently)
- yet it was unclear how pion, rho, A1, rho' move as T grows (hard on the lattice as Matsubara box shrinks), e.g.
rho \Rightarrow pi, rho' \Rightarrow A1 (Brown-Rho);
or **no pion, rho \Rightarrow A1**; or nobody moves and **all melt (???)**
- lattice thermodynamics, especially mu-derivatives suggest BARYONS get heavier. The **mass LR term disappears** but **energy (LL+RR) appear** and compensates

chiral breaking is due to small subset of states, ZMZ
and its width is small (ES, 1982)

given by the magnitude of the hopping
from one instanton to the next

$$T_{I\bar{I}} \sim \frac{\rho^2}{R^3} \sim \frac{(0.3 \text{ fm})^2}{(1 \text{ fm})^3} \sim 20 \text{ MeV}$$

- people found it on the lattice and showed pions are completely described by ZMZ
- that is why quark mass dependence is nontrivial, and chiral perturbation

recently the opposite exercise was done by the Graz group

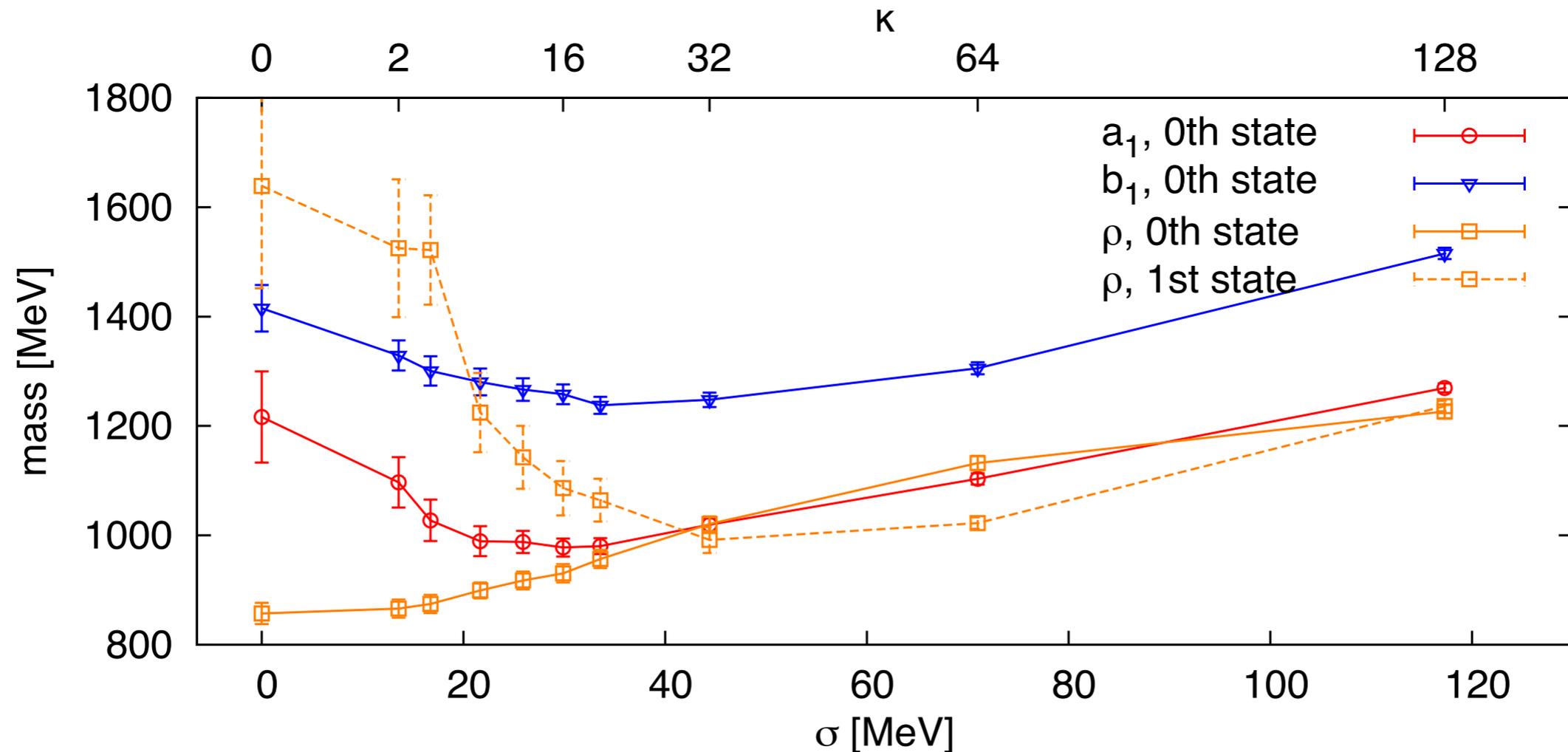
Symmetries of hadrons after unbreaking the chiral symmetry

L. Ya. Glozman,* C. B. Lang,[†] and M. Schröck[‡]

Institut für Physik, FB Theoretische Physik, Universität Graz, A-8010 Graz, Austria

By eliminating narrow band of modes 1/10000 one finds that **A1 moves down**
while rho only slightly go up

near-perfect chiral pairs are left, nearly the same in average => thus resonance gas works



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- photons: small (nonpert) viscosity=> large (nonpert) gamma radiation rates.
- $v_2 \leq$ need to calculate all fermion loop effects
- example: interplay of strong **coherent QED fields** with partonic reactions, including **virtual quark loops**