

# Measuring Histograms of the Number of Pairs of Particles, Constructing Ratios of Histograms, and the relationship to Correlations – PART 1

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## OUTLINE

- Define two-particle correlations and density ratios of the number of pairs of particles.
- Describe sibling and mixed pairs of particles and number of pair densities and histograms.
- Application to  $\eta, \phi$  space
- **Technical issues including:**
  - Kinematic projections
  - Normalization
  - Charge dependence
  - Notation conventions

# Correlations & Density Ratios for the Number of Pairs of Particles:

Two-particle correlation in momentum space

$$C_{(2)}(\vec{p}_1, \vec{p}_2) \equiv \underbrace{\rho_{(2)}(\vec{p}_1, \vec{p}_2)}_{\text{Two-particle density}} - \underbrace{\rho_{(1)}(\vec{p}_1)}_{\text{Single-particle densities}} \underbrace{\rho_{(1)}(\vec{p}_2)}_{\text{Single-particle densities}}$$

Two-particle density

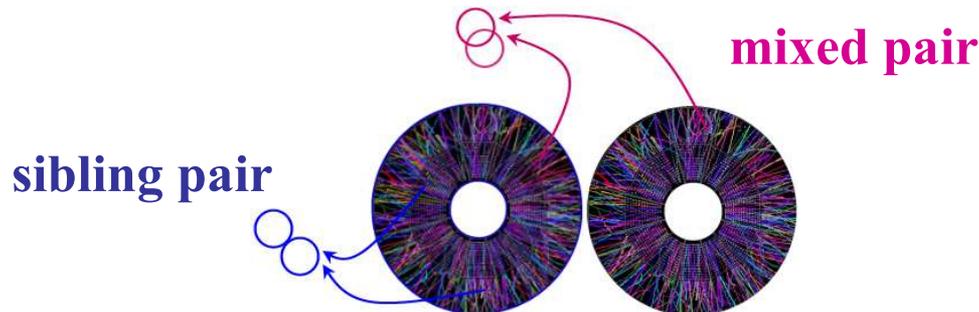
Single-particle densities:  
product equals uncorrelated  
two-particle density

$$\rho_{(2)}(\vec{p}_1, \vec{p}_2) \equiv \rho_{sibling}(\vec{p}_1, \vec{p}_2)$$

$$\rho_{(1)}(\vec{p}_1)\rho_{(1)}(\vec{p}_2) \equiv \rho_{mix}(\vec{p}_1, \vec{p}_2)$$

density of “**sibling**” pairs - each pair of particles is taken from the same event, hence the term *sibling*; pairs from many events are summed.

density of “**mixed**” pairs - one particle for a mixed-event pair is taken from one event and the second particle is taken from a different, but similar event; corresponding pairs from many events are summed.



## Correlations & Density Ratios for the Number of Pairs of Particles: (continued)

*Combining:*  $C_{(2)}(\vec{p}_1, \vec{p}_2) = \rho_{mix}(\vec{p}_1, \vec{p}_2)[r(\vec{p}_1, \vec{p}_2) - 1],$

*Define ratio of densities of the number of pairs of particles:*

$$r(\vec{p}_1, \vec{p}_2) \equiv \frac{\rho_{sibling}(\vec{p}_1, \vec{p}_2)}{\rho_{mixed}(\vec{p}_1, \vec{p}_2)}$$

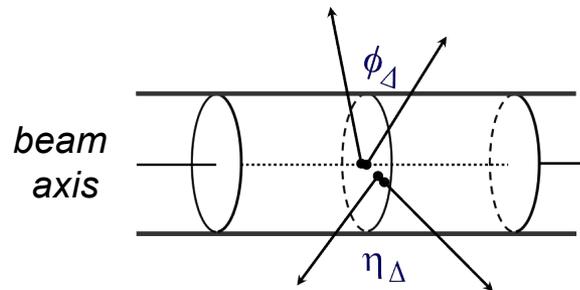
# Measuring the Ratio of Histograms of Number of Pairs of Particles: Projections onto $\eta, \phi$ - or - Axial momentum space

Sort events within a centrality bin according to:

- (1) Location of primary vertex along beam line (z-axis)
- (2) Number of primary particles ( $N_{ch}$ ) within  $p_t, \eta, \phi$  acceptance.

For all accepted particles in this subset of events compute

$(\phi_1 - \phi_2)$  and  $(\eta_1 - \eta_2)$  for every sibling and mixed pair.



Where we define:

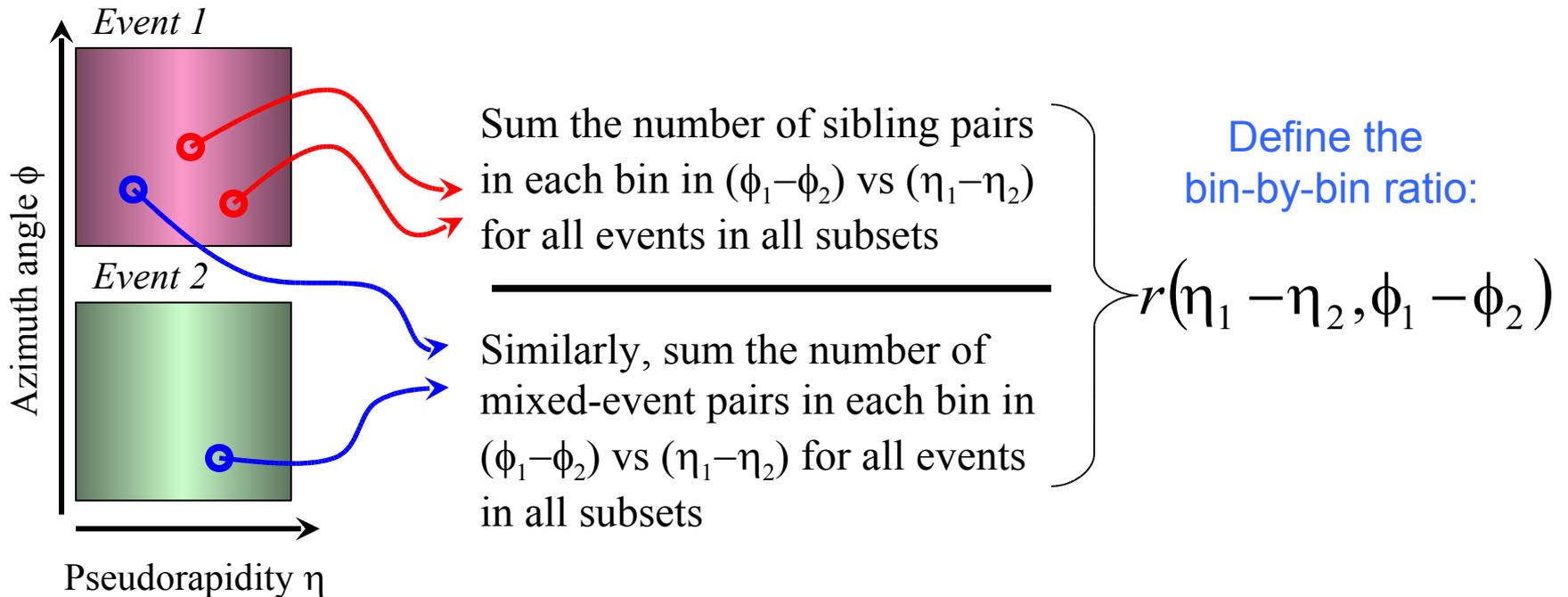
$$\phi_\Delta = \phi_1 - \phi_2$$

$$\eta_\Delta = \eta_1 - \eta_2$$

(see last slide in this tutorial)

Compute histograms for the number of sibling and mixed pairs of particles using only those subsets of events within 7.5 cm along the z-axis and  $N_{ch}$  within 50. Repeat for all event subsets in a centrality bin.

# Measuring the Ratio of Histograms of Number of Pairs of Particles: Projections onto $\eta, \phi$ – or - Axial momentum space (continued)



This is an example of projecting the six-dimensional correlation distribution onto two-dimensional, relative pair-wise momentum coordinates.

We refer to such projections as the **joint** (*i.e.* two things) – **autocorrelation** (*i.e.* on relative coordinates)

# Normalized Ratio of Histograms - $\hat{r}$

$$\hat{r}(\eta_1 - \eta_2, \phi_1 - \phi_2) \equiv \frac{N_{mixed-pairs}}{N_{sibling-pairs}} r(\eta_1 - \eta_2, \phi_1 - \phi_2)$$

where:

$N_{sibling-pairs} =$  Total number of sibling pairs for all events in the centrality bin.

$N_{mixed-pairs} =$  Total number of mixed-event pairs used among the “similar” events within each  $Z_{vertex}$  and  $N_{ch}$  bin, summed over all such subsets of events in the centrality bin.

# Charge Independent and Charge Dependent Ratios:

Define normalized ratios of histograms for the four combinations of charge signs ( $++$ ,  $--$ ,  $+-$ ,  $-+$ ), for like-sign (LS) pairs and unlike-sign (US) pairs, and for all charge-sign pairs as in the preceding discussion.

We have verified that  $\hat{r}_{++} = \hat{r}_{--}$  and  $\hat{r}_{+-} = \hat{r}_{-+}$  within statistics.

Charge independent (CI) normalized ratios of histograms are computed exactly as described in the preceding slides using all charged particles.

$$\hat{r}^{CI}(\eta_{\Delta}, \phi_{\Delta}) \equiv \frac{N_{mixed-pairs}}{N_{sibling-pairs}} r(\eta_{\Delta}, \phi_{\Delta})$$

Charge dependent (CD) normalized ratios of histograms are computed by applying the steps described in the preceding slides for all like-sign pairs of charged particles ( $++$  and  $--$ ) and for all unlike-sign pairs of charged particles ( $+-$  and  $-+$ ) and forming the difference:

$$\hat{r}^{CD}(\eta_{\Delta}, \phi_{\Delta}) \equiv \hat{r}_{LS} - \hat{r}_{US}$$

# Why does the Event Structure PWG use $\mathbf{x}_\Delta$ , rather than the more common $\Delta\mathbf{x}$ for difference quantity $\mathbf{x}_1 - \mathbf{x}_2$ ?

The reason is related to the fact that we often refer to rotated coordinate systems which take advantage of the observed symmetries in the two-point correlation data.

Rotated axes lie along the sum  $(\mathbf{x}_1 + \mathbf{x}_2)$  and difference  $(\mathbf{x}_1 - \mathbf{x}_2)$  directions.

Notation  $\mathbf{x}_\Sigma, \mathbf{x}_\Delta$  follow naturally.

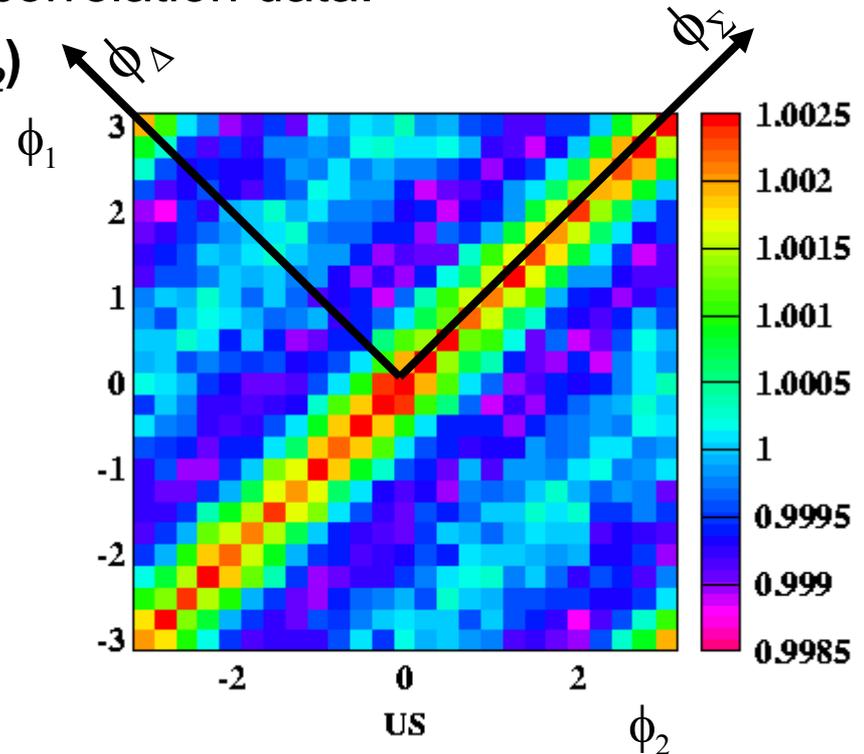
On the other hand definition

$\Delta\mathbf{x} = \mathbf{x}_1 - \mathbf{x}_2$  calls for the correspond-

ing notation “ $\Sigma\mathbf{x}$ ” which would be confused with summation. For

example the  $\phi_1$  versus  $\phi_2$  two-point ratio data are shown here along with

the  $\phi_\Sigma, \phi_\Delta$  coordinate system.



## Application of these methods in STAR papers:

- **“Hadronization geometry and charge-dependent number autocorrelations on axial momentum space in Au-Au collisions at  $\sqrt{s_{NN}} = 130$  GeV,”**  
*Submitted June 30, 2004*  
e-Print Archives (nucl-ex/0406035).
- **“Minijet deformation and charge-independent particle correlations on axial momentum space in Au-Au collisions at  $\sqrt{s_{NN}} = 130$  GeV,”**  
Presently on starpapers-I (*August, 2004*).
- **“Transverse momentum correlations and minijet dissipation in Au-Au collisions at  $\sqrt{s_{NN}} = 130$  GeV,”**  
*Submitted August 11, 2004*  
e-Print Archives (nucl-ex/0408012)