

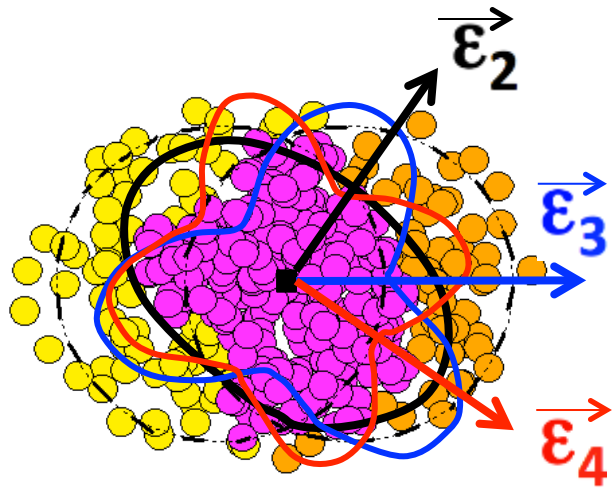


# Event-shape fluctuations and flow correlations

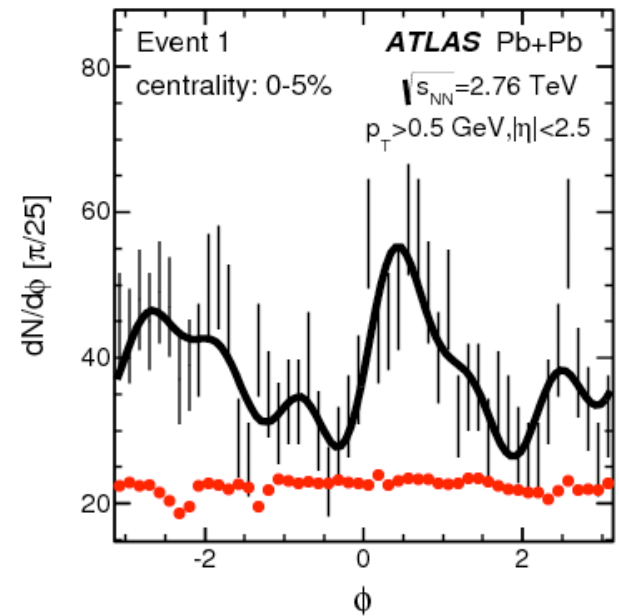
Jiangyong Jia

ATHIC 2014 Conference

# Geometry and harmonic flow



Collective expansion



$$\vec{\epsilon}_n \equiv \epsilon_n e^{in\Phi_n^*} \equiv -\frac{\langle r^n e^{in\phi} \rangle}{\langle r^n \rangle}$$

$$\frac{dN}{d\phi} \propto 1 + 2 \sum_n v_n \cos n(\phi - \Phi_n)$$

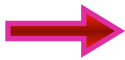
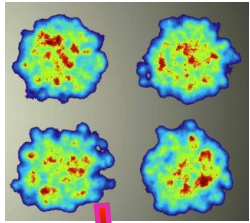
$$\vec{v}_n \equiv v_n e^{in\Phi_n}$$

- Probes: initial geometry and transport properties of QGP
  - How  $(\epsilon_n, \Phi_n^*)$  are transferred to  $(v_n, \Phi_n)$ ?
  - What is the nature of final state (non-linear) dynamics?
  - What is the nature of longitudinal flow dynamics?

# Event-by-event observables

Many little bangs

1104.4740, 1209.2323, 1203.5095, 1312.3572



$$p(v_n, v_m, \dots, \Phi_n, \Phi_m, \dots) = \frac{1}{N_{\text{evts}}} \frac{dN_{\text{evts}}}{dv_n dv_m \dots d\Phi_n d\Phi_m \dots}$$

Moments:  $\langle \cos(n_1 \phi_1 + n_2 \phi_2 \dots + n_m \phi_m) \rangle = \sum n_i = 0$

Examples:  $\langle v_{n_1} v_{n_2} \dots v_{n_m} \cos(n_1 \Phi_{n_1} + n_2 \Phi_{n_2} \dots + n_m \Phi_{n_m}) \rangle$

$$\langle \cos(n\phi_1 - n\phi_2 + n\phi_3 - n\phi_4) \rangle = \langle v_n^4 \cos(n\Phi_n - n\Phi_n + n\Phi_n - n\Phi_n) \rangle = \langle v_n^4 \rangle$$

$$\langle \cos(n\phi_1 - n\phi_2 + m\phi_3 - m\phi_4) \rangle = \langle v_n^2 v_m^2 \cos(n\Phi_n - n\Phi_n + m\Phi_m - m\Phi_m) \rangle = \langle v_n^2 v_m^2 \rangle$$

$$\langle \cos(3\phi_1 + 3\phi_2 - 6\phi_3) \rangle = \langle v_3^2 v_6 \cos 6(\Phi_3 - \Phi_6) \rangle$$

Cumulants obtained by combining with lower order correlators:

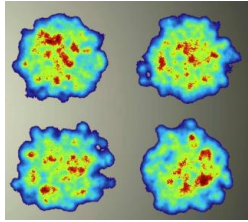
$$\langle \cos(n\phi_1 - n\phi_2 + n\phi_3 - n\phi_4) \rangle_c = \langle v_n^4 \rangle - 2 \langle v_n^2 \rangle^2 \quad \rho(v_n)$$

$$\langle \cos(n\phi_1 - n\phi_2 + m\phi_3 - m\phi_4) \rangle_c = \langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle \quad \rho(v_n, v_m)$$

# Event-by-event observables

Many little bangs

1104.4740, 1209.2323, 1203.5095, 1312.3572

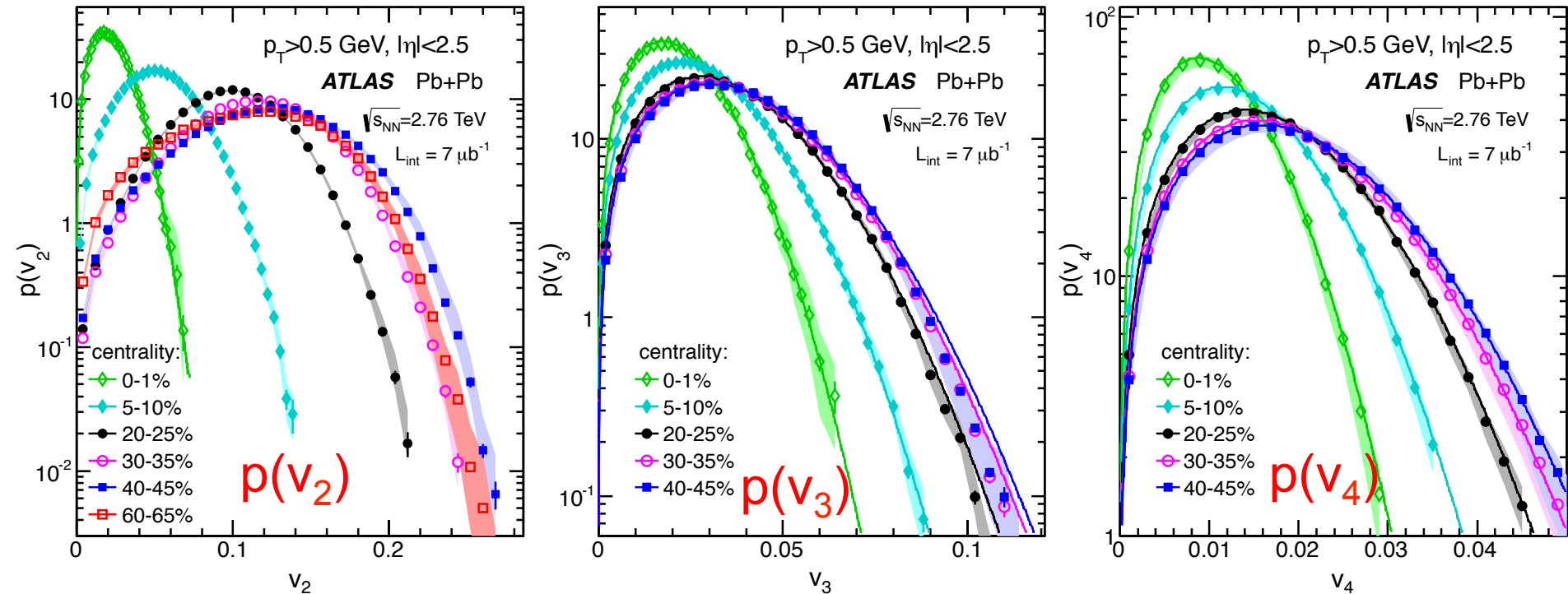


$$p(v_n, v_m, \dots, \Phi_n, \Phi_m, \dots) = \frac{1}{N_{\text{evts}}} \frac{dN_{\text{evts}}}{dv_n dv_m \dots d\Phi_n d\Phi_m \dots}$$

	pdf's	cumulants	event-shape method
	$p(v_n)$	$v_n \{2k\}, k = 1, 2, \dots$	NA
	$p(v_n, v_m)$	$\langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle$	yes
Flow-amplitudes	$p(v_n, v_m, v_l)$	$\langle v_n^2 v_m^2 v_l^2 \rangle + 2\langle v_n^2 \rangle \langle v_m^2 \rangle \langle v_l^2 \rangle - \langle v_n^2 v_m^2 \rangle \langle v_l^2 \rangle - \langle v_m^2 v_l^2 \rangle \langle v_n^2 \rangle - \langle v_l^2 v_n^2 \rangle \langle v_m^2 \rangle$	yes
	...	Obtained recursively as above	yes
EP-correlation	$p(\Phi_n, \Phi_m, \dots)$	$\langle v_n^{c_n} v_m^{c_m} \dots \cos(c_n n \Phi_n + c_m m \Phi_m + \dots) \rangle$ $\sum_k k c_k = 0$	yes
Mixed-correlation	$p(v_l, \Phi_n, \Phi_m, \dots)$	$\langle v_l^2 v_n^{c_n} v_m^{c_m} \dots \cos(c_n n \Phi_n + c_m m \Phi_m + \dots) \rangle - \langle v_l^2 \rangle \langle v_n^{c_n} v_m^{c_m} \dots \cos(c_n n \Phi_n + c_m m \Phi_m + \dots) \rangle$ $\sum_k k c_k = 0$	yes

Flow fluctuation:  $p(v_n)$

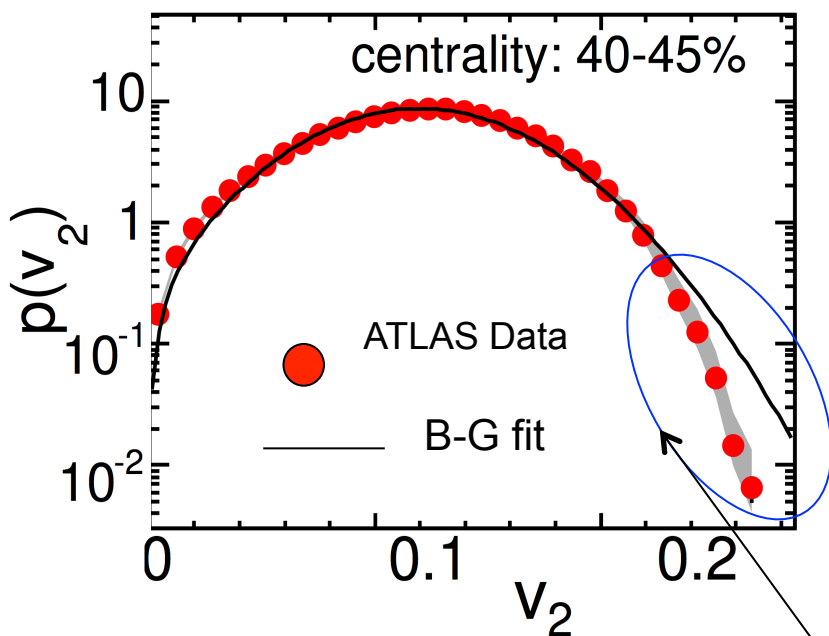
# $p(v_2)$ , $p(v_3)$ and $p(v_4)$ distributions



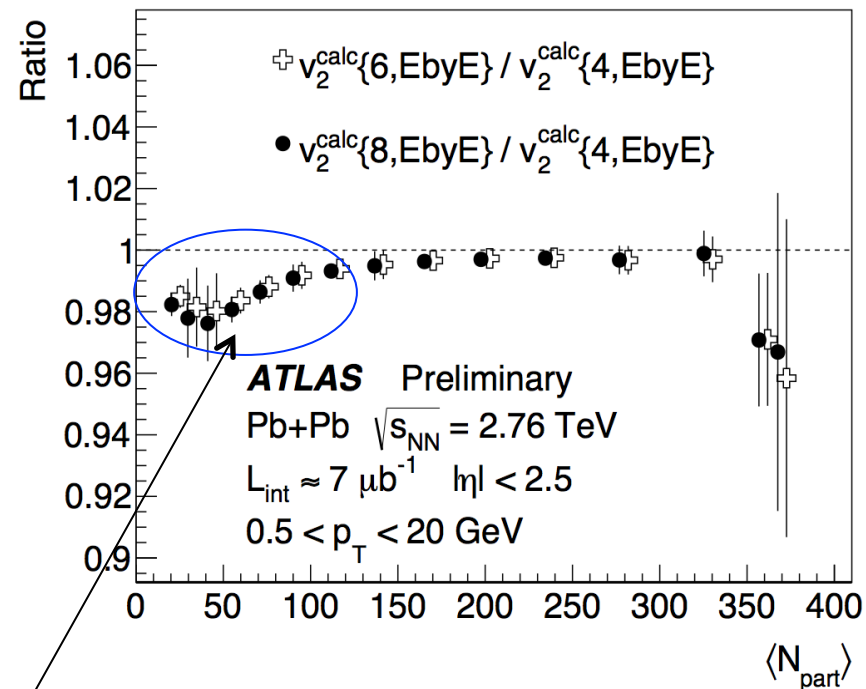
$$v_n \{4\}^4 = 2 \langle v_n^2 \rangle^2 - \langle v_n^4 \rangle \neq 0 \text{ for } n = 2, 3$$

- The non-zero  $v_n \{4,6..\}$  either due to
  - average geometry such as  $v_2^{RP} \neq 0$  or
  - non-Gaussianness in the flow fluctuation

# Cumulants from $p(v_2)$



Non-Gaussian behavior



- Non-gaussianity is reflected by a 1-2% change beyond 4<sup>th</sup> order cumulants

# Event-plane correlations $\rho(\Phi_n, \Phi_m \dots)$

1403.0489

$$\langle \cos 4(\Phi_2 - \Phi_4) \rangle$$

$$\langle \cos 8(\Phi_2 - \Phi_4) \rangle$$

$$\langle \cos 12(\Phi_2 - \Phi_4) \rangle$$

$$\langle \cos 6(\Phi_2 - \Phi_3) \rangle$$

$$\langle \cos 6(\Phi_2 - \Phi_6) \rangle$$

$$\langle \cos 6(\Phi_3 - \Phi_6) \rangle$$

$$\langle \cos 12(\Phi_3 - \Phi_4) \rangle$$

$$\langle \cos 10(\Phi_2 - \Phi_5) \rangle$$

$$\langle \cos(2\Phi_2 + 3\Phi_3 - 5\Phi_5) \rangle$$

$$\langle \cos(-8\Phi_2 + 3\Phi_3 + 5\Phi_5) \rangle$$

$$\langle \cos(2\Phi_2 + 4\Phi_4 - 6\Phi_6) \rangle$$

$$\langle \cos(-10\Phi_2 + 4\Phi_4 + 6\Phi_6) \rangle$$

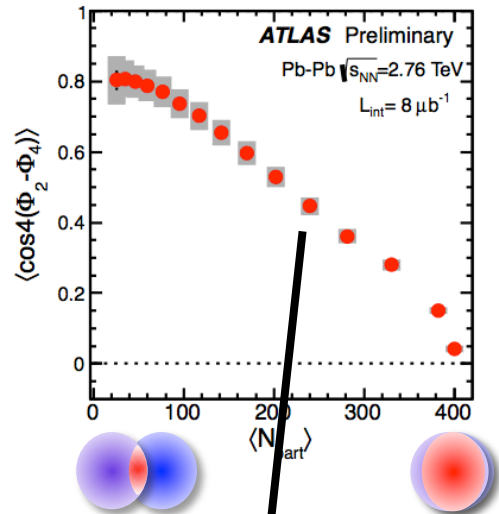
$$\langle \cos(2\Phi_2 - 6\Phi_3 + 4\Phi_4) \rangle$$

$$\langle \cos(-10\Phi_2 + 6\Phi_3 + 4\Phi_4) \rangle$$

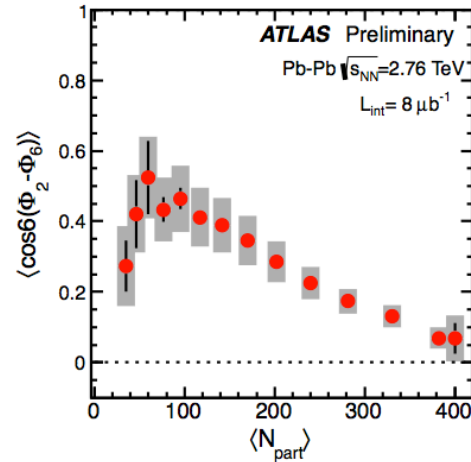


# Event-plane correlation results

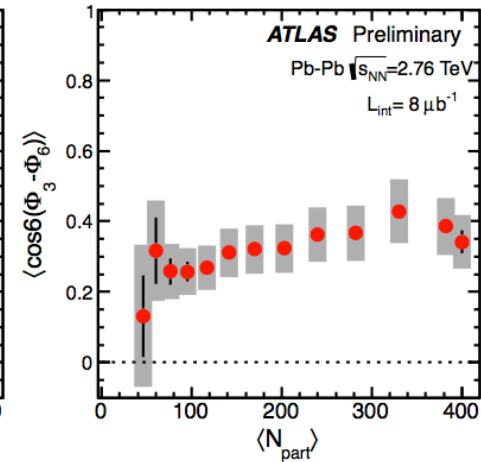
$$\langle \cos 4(\Phi_2 - \Phi_4) \rangle$$



$$\langle \cos 6(\Phi_2 - \Phi_6) \rangle$$

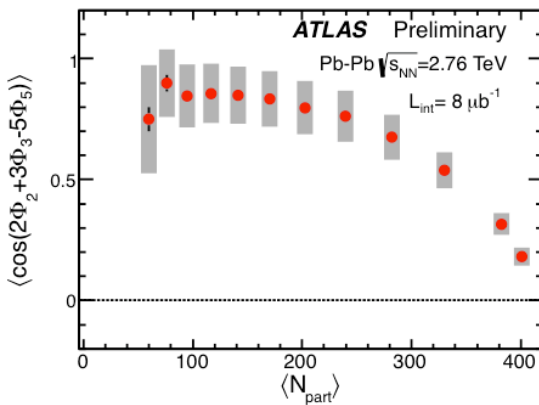


$$\langle \cos 6(\Phi_3 - \Phi_6) \rangle$$



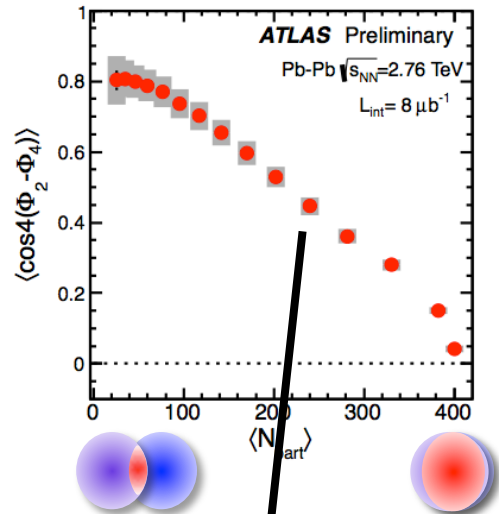
$$v_4 e^{i4\Phi_4} \propto \varepsilon_4 e^{i4\Phi_4^*} + c v_2^2 e^{i4\Phi_2} + \dots$$

$$\langle \cos(2\Phi_2 + 3\Phi_3 - 5\Phi_5) \rangle$$



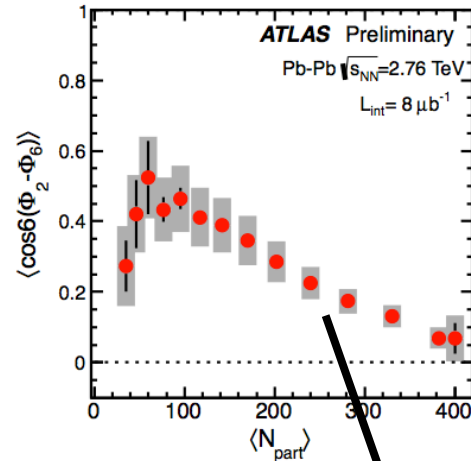
# Event plane correlation results

$$\langle \cos 4(\Phi_2 - \Phi_4) \rangle$$

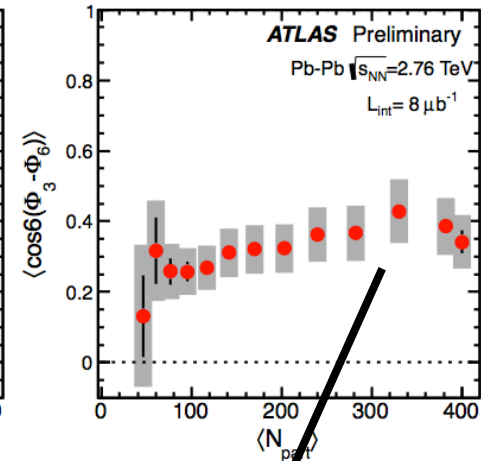


$$v_4 e^{i4\Phi_4} \propto \varepsilon_4 e^{i4\Phi_4^*} + c v_2^2 e^{i4\Phi_2} + \dots$$

$$\langle \cos 6(\Phi_2 - \Phi_6) \rangle$$

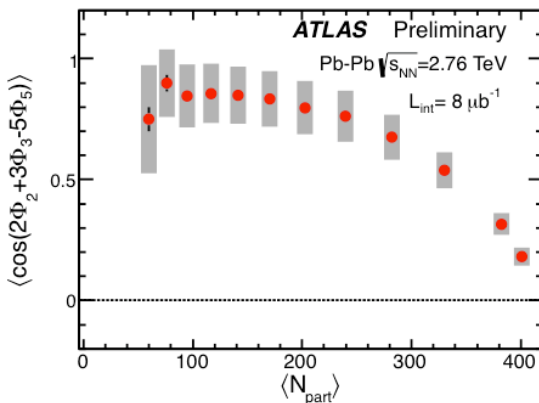


$$\langle \cos 6(\Phi_3 - \Phi_6) \rangle$$



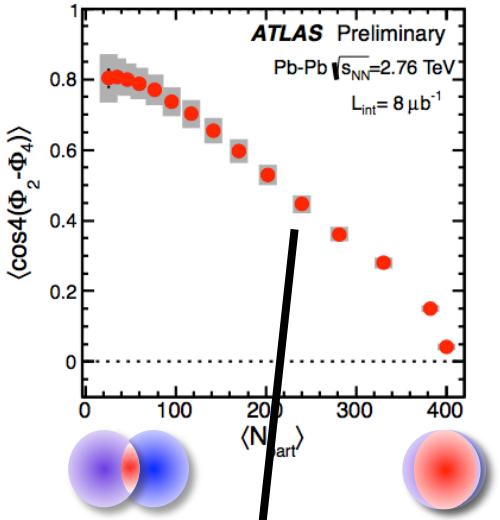
$$v_6 e^{i6\Phi_6} \propto \varepsilon_6 e^{i6\Phi_6^*} + c_1 v_2^3 e^{i6\Phi_2} + c_2 v_3^2 e^{i6\Phi_3} + \dots$$

$$\langle \cos(2\Phi_2 + 3\Phi_3 - 5\Phi_5) \rangle$$



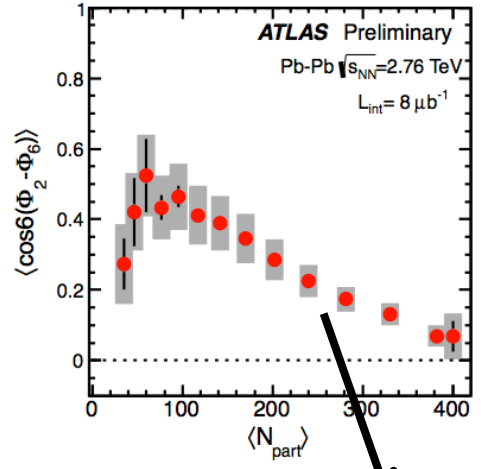
# Event plane correlation results

$$\langle \cos 4(\Phi_2 - \Phi_4) \rangle$$



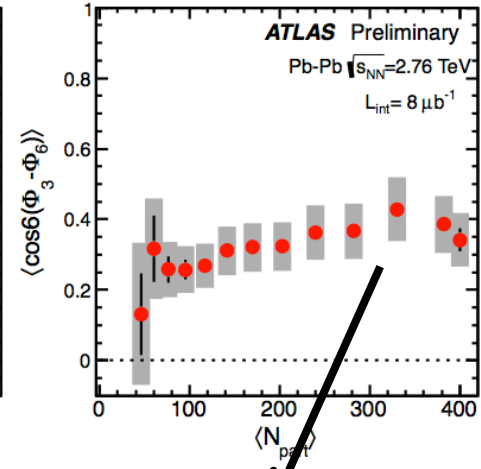
$$v_4 e^{i4\Phi_4} \propto \varepsilon_4 e^{i4\Phi_4^*} + c v_2^2 e^{i4\Phi_2} + \dots$$

$$\langle \cos 6(\Phi_2 - \Phi_6) \rangle$$

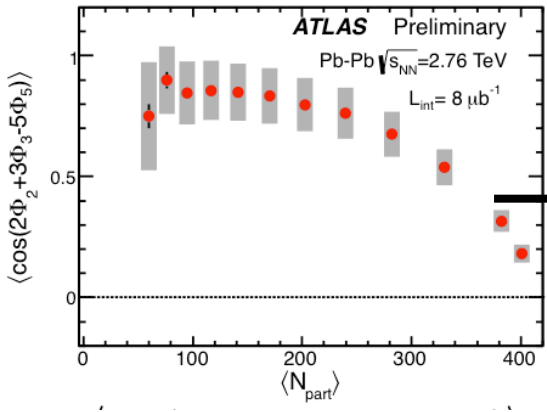


$$v_6 e^{i6\Phi_6} \propto \varepsilon_6 e^{i6\Phi_6^*} + c_1 v_2^3 e^{i6\Phi_2} + c_2 v_3^2 e^{i6\Phi_3} + \dots$$

$$\langle \cos 6(\Phi_3 - \Phi_6) \rangle$$



$$\langle \cos(2\Phi_2 + 3\Phi_3 - 5\Phi_5) \rangle$$



$$v_5 e^{i5\Phi_5} \propto \varepsilon_5 e^{i5\Phi_5^*} + c v_2 v_3 e^{i(2\Phi_2 + 3\Phi_3)} + \dots$$

# How $(\varepsilon_n, \Phi_n^*)$ are transferred to $(v_n, \Phi_n)$ ?

- Flow response is linear for  $v_2$  and  $v_3$ :  $v_n \propto \varepsilon_n$  and  $\Phi_n \approx \Phi_n^*$  i.e.

$$v_2 e^{-i2\Phi_2} \propto \varepsilon_2 e^{-i2\Phi_2^*}, \quad v_3 e^{-i3\Phi_3} \propto \varepsilon_3 e^{-i3\Phi_3^*}$$

- Higher-order flow arises from EP correlations., e.g. :

$$v_4 e^{i4\Phi_4} \propto \varepsilon_4 e^{i4\Phi_4^*} + c v_2^2 e^{i4\Phi_2} + \dots$$

Ollitrault, Luzum, Teaney, Li, Heinz, Chun....

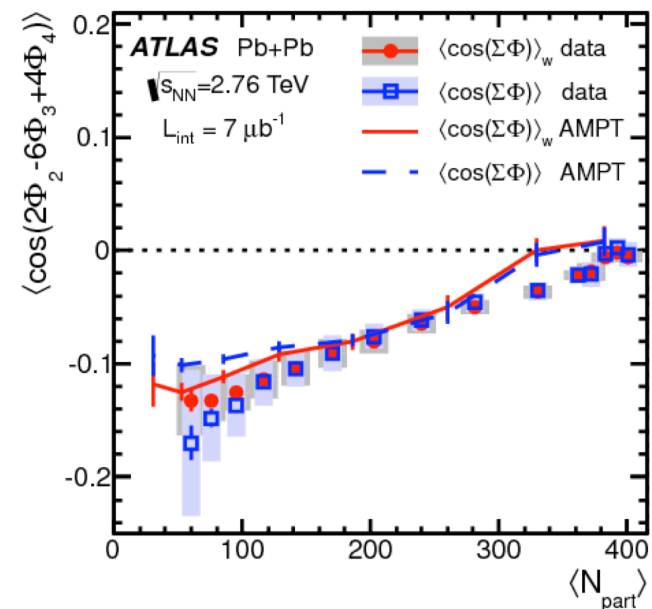
$$v_5 e^{i5\Phi_5} \propto \varepsilon_5 e^{i5\Phi_5^*} + c v_2 v_3 e^{i(2\Phi_2+3\Phi_3)} + \dots$$

$$v_6 e^{i6\Phi_6} \propto \varepsilon_6 e^{i6\Phi_6^*} + c_1 v_2^3 e^{i6\Phi_2} + c_2 v_3^2 e^{i6\Phi_3} + c_3 v_2 \varepsilon_4 e^{i(2\Phi_2+4\Phi_4^*)} \dots$$

- Some correlators lack intuitive explanation

e.g.  $\langle \cos(2\Phi_2 - 6\Phi_3 + 4\Phi_4) \rangle$  correlation

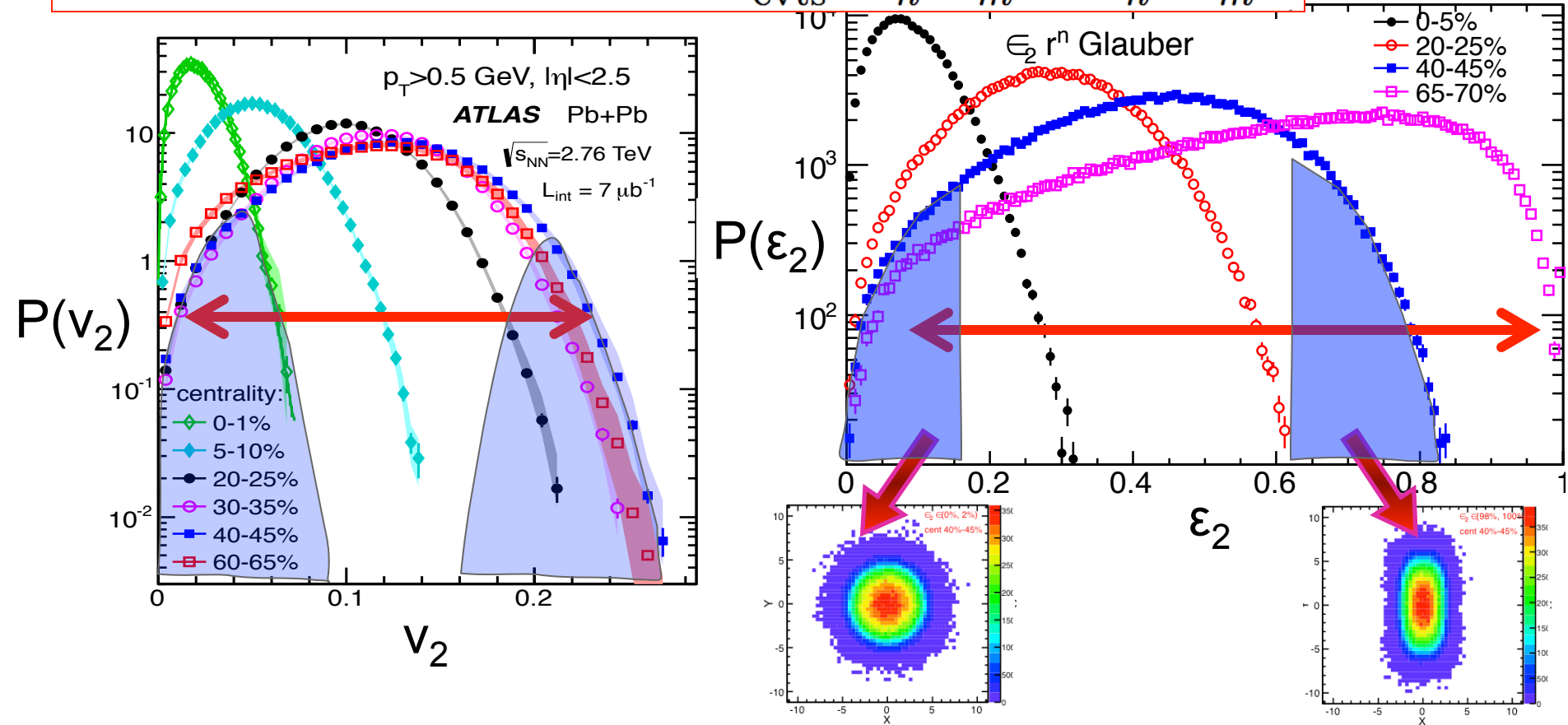
- Although described by EbyE hydro and AMPT



# Event-shape selection technique

# Can we do better?

$$p(v_n, v_m, \dots, \Phi_n, \Phi_m, \dots) = \frac{1}{N_{\text{evts}}} \frac{dN_{\text{evts}}}{dv_n dv_m \dots d\Phi_n d\Phi_m \dots}$$

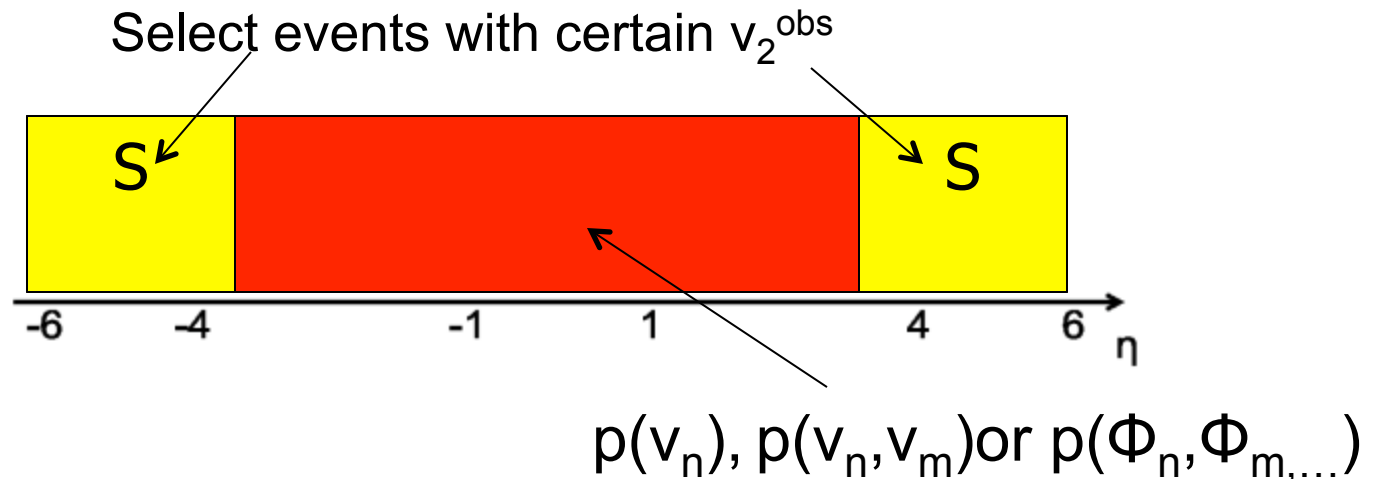


- Study the variation of  $v_n$  at fixed centrality but varying event-geometry: “event-shape-selected  $v_n$  measurements”

# Event-shape selection technique

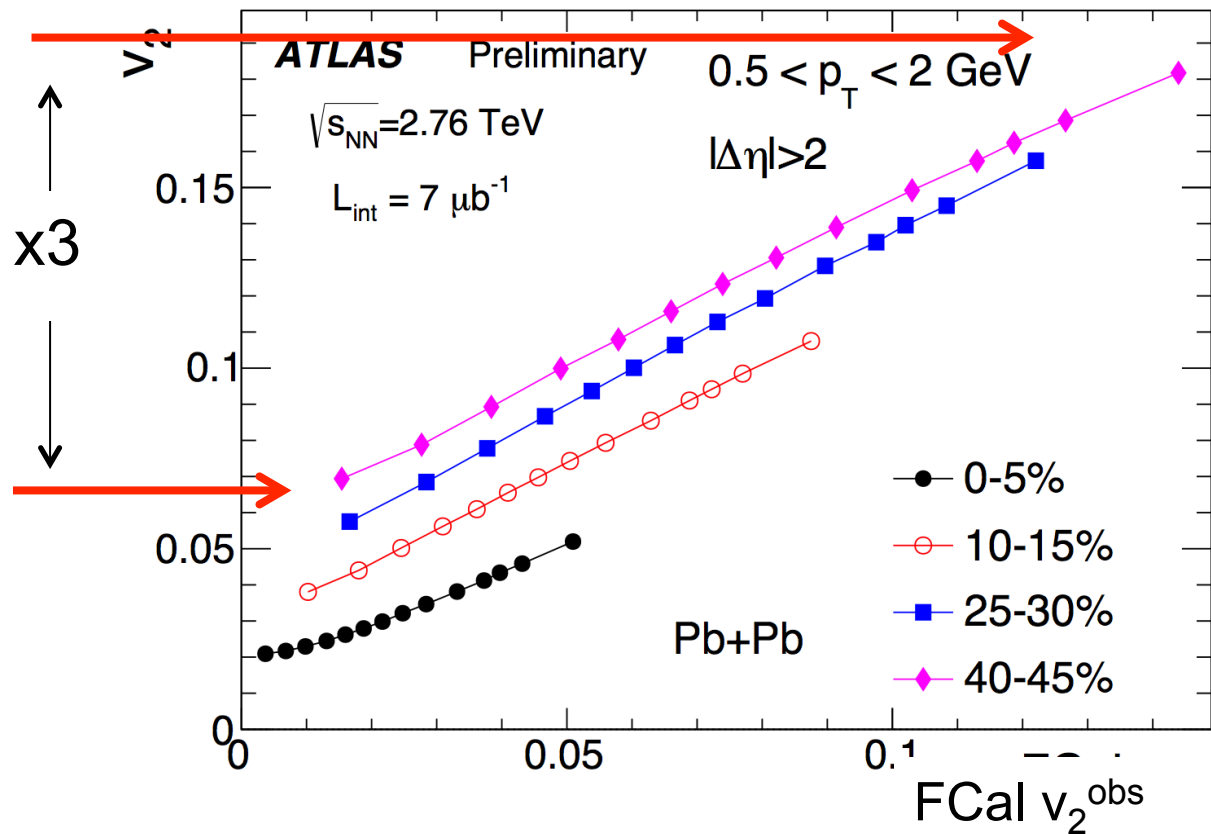
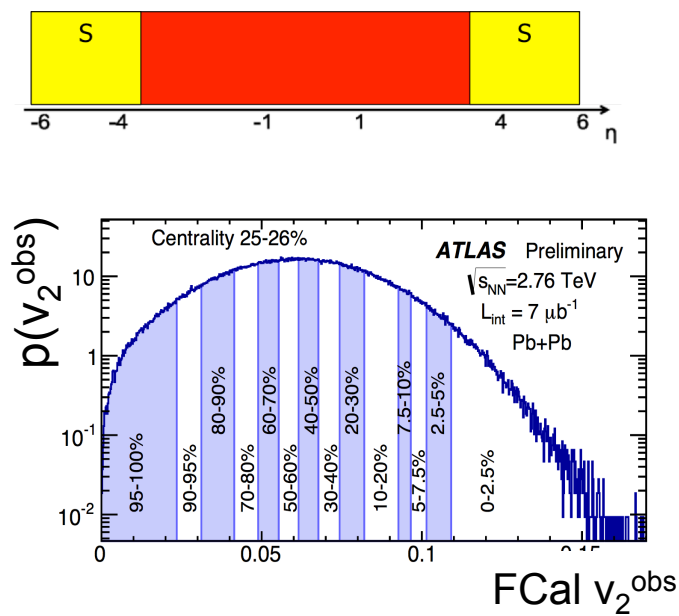
Schukraft, Timmins, and Voloshin, arXiv:1208.4563

Huo, Mohapatra, JJ arxiv:1311.7091



$$\vec{v}_n^{\text{obs}} = \frac{1}{\sum w} (\sum w \cos n\phi_n, \sum w \sin n\phi_n), \quad w = p_T, \quad v_n^{\text{obs}} = \left| \vec{v}_n^{\text{obs}} \right|$$

# More info by selecting on event-shape



Fix centrality then select events with certain  $v_2^{\text{obs}}$  in large rapidity:

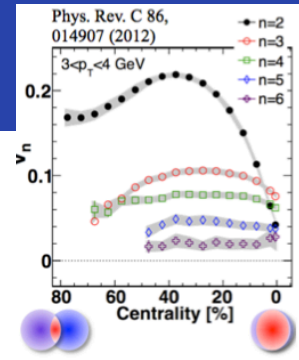
→ measure  $v_n$  via two-particle correlations in  $|\eta| < 2.5$

Vary ellipticity by a factor of 3!

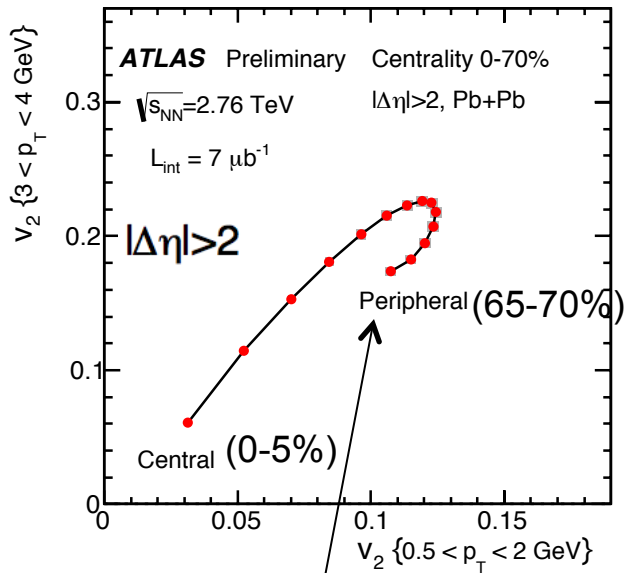


# $v_n - v_2$ correlations: centrality dependence

- First correlation without event  $v_2$ -selection, 5% steps

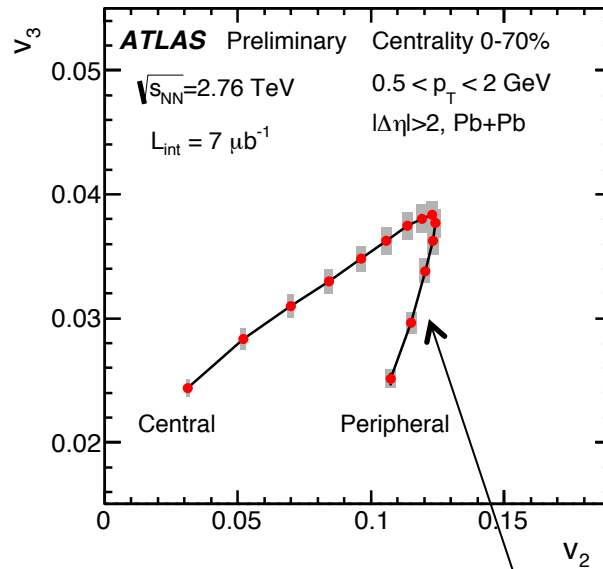


## $v_2(p_T^a, b)$ vs $v_2(p_T^b, b)$



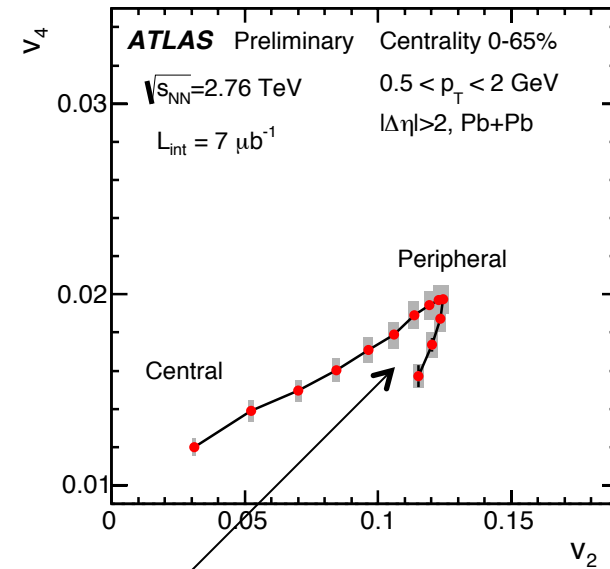
“Boomerang” reflects stronger viscous damping at higher  $p_T$  and peripheral

## $v_3(b)$ vs $v_2(b)$



“Boomerang” reflects reflects different centrality dependence, which is also sensitive to the viscosity effect.

## $v_4(b)$ vs $v_2(b)$

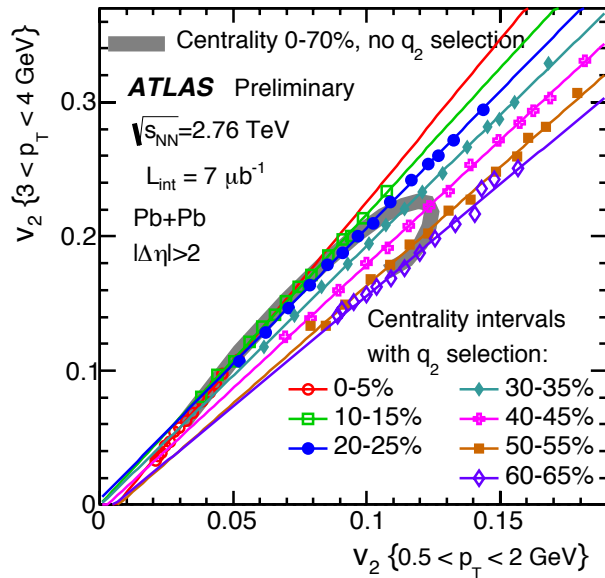


# $v_n$ - $v_2$ correlations: within fixed centrality

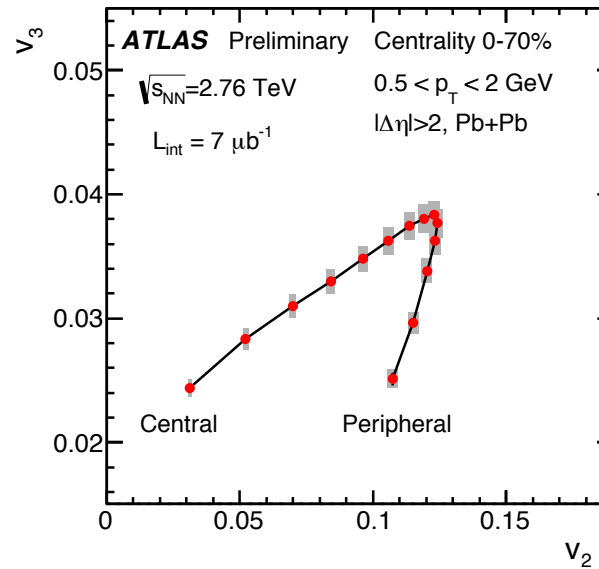
- Fix system size and vary the ellipticity!

Probe  $p(v_n, v_2)$

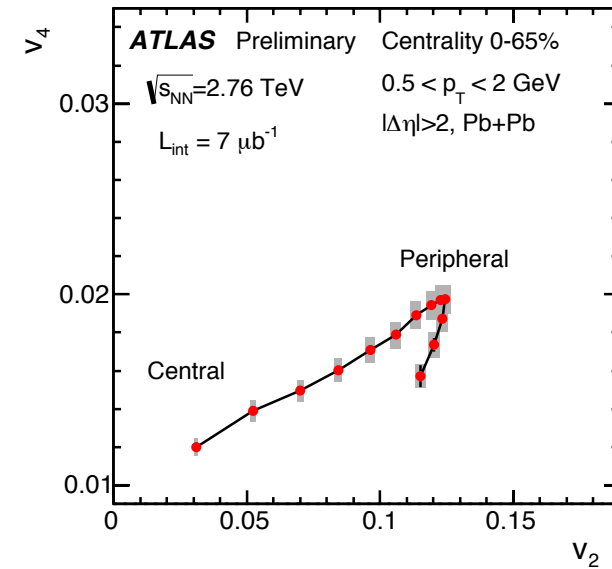
$v_2$  (higher  $p_T$ )



$v_3$



$v_4$



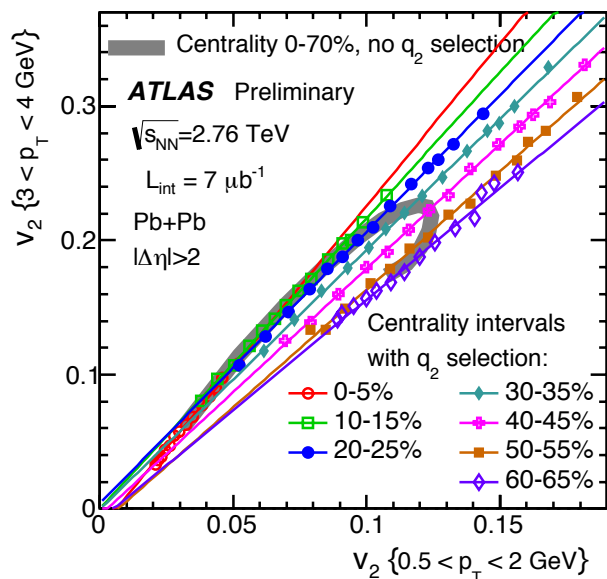
Linear correlation for forward  $v_2$ -selected bin  $\rightarrow$  viscous damping controlled by system size, not shape

# $v_n$ - $v_2$ correlations: within fixed centrality

- Fix system size and vary the ellipticity!

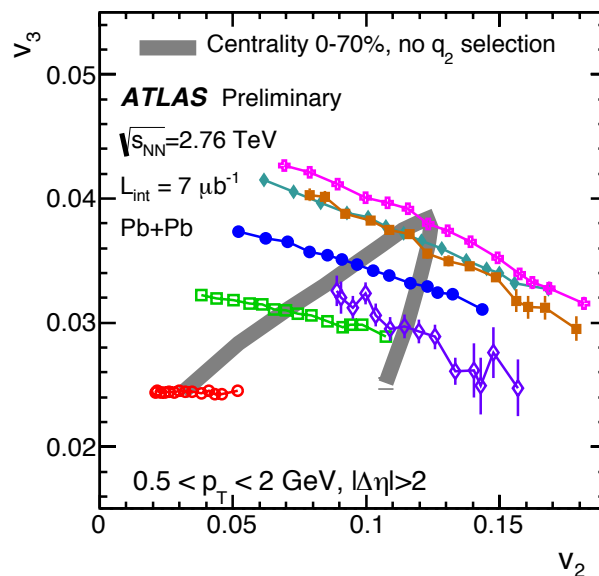
Probe  $p(v_n, v_2)$

$v_2$  (higher  $p_T$ )



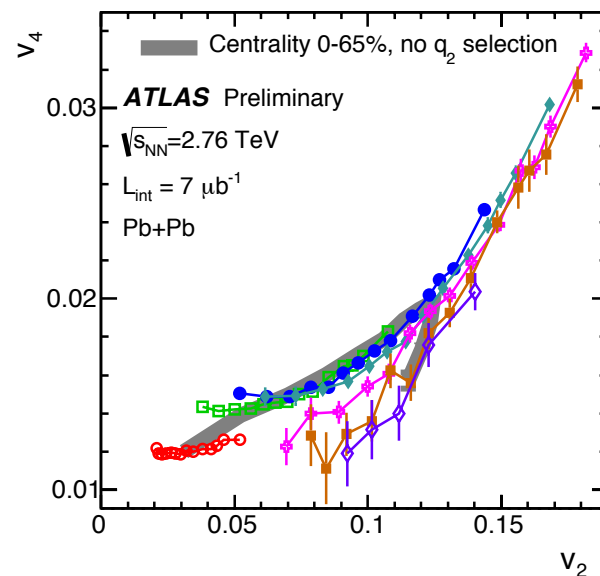
Linear correlation for forward  $v_2$ -selected bin  $\rightarrow$  viscous damping controlled by system size, not shape

$v_3$



Clear anti-correlation,

$v_4$



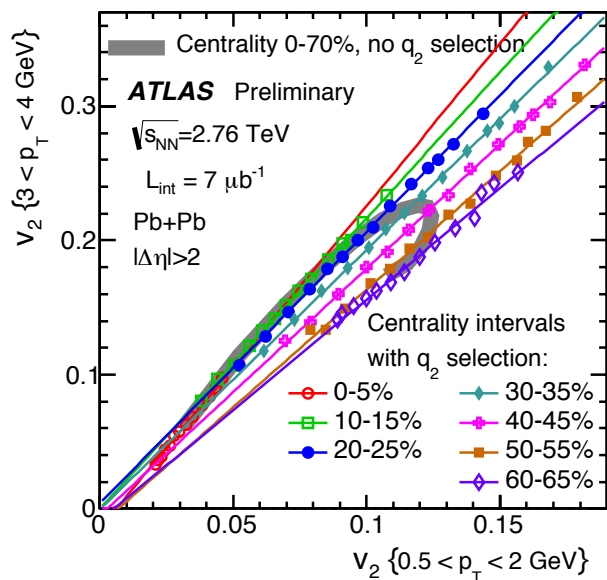
quadratic rise from non-linear coupling to  $v_2^2$

# $v_n$ - $v_2$ correlations: within fixed centrality

- Fix system size and vary the ellipticity!
- Overlay  $\varepsilon_3$ - $\varepsilon_2$  and  $\varepsilon_4$ - $\varepsilon_2$  correlations, rescaled

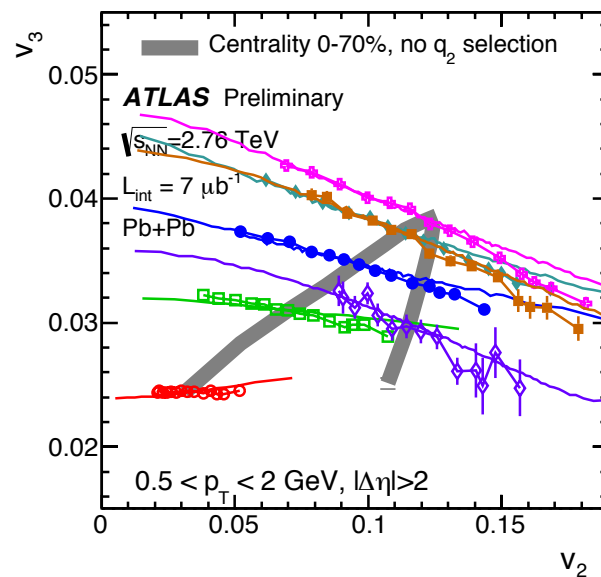
Probe  $p(v_n, v_2)$

## $v_2$ (higher $p_T$ )



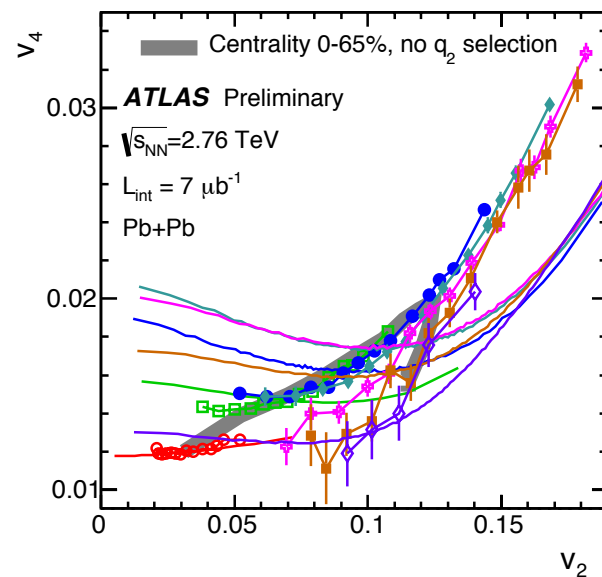
Linear correlation for forward  $v_2$ -selected bin  $\rightarrow$  viscous damping controlled by system size, not shape

## $v_3$



Clear anti-correlation, mostly initial geometry effect!!

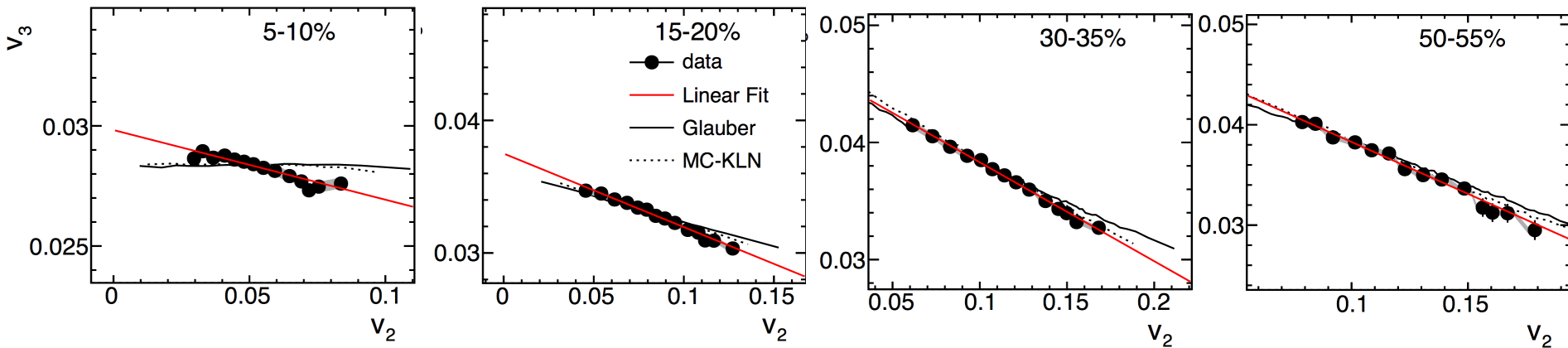
## $v_4$



quadratic rise from non-linear coupling to  $v_2^2$  initial geometry only does not work!!

Initial geometry describe  $v_3$ - $v_2$  but fails  $v_4$ - $v_2$  correlation

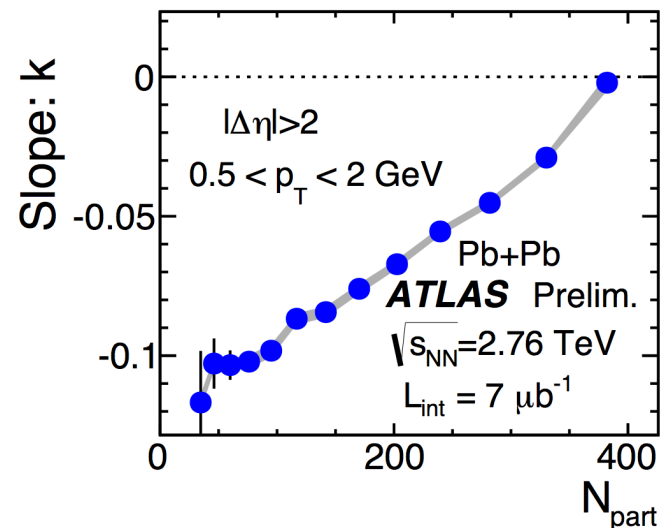
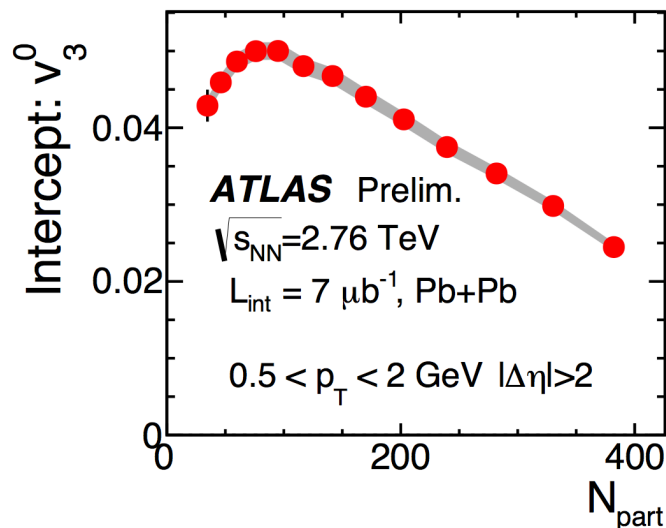
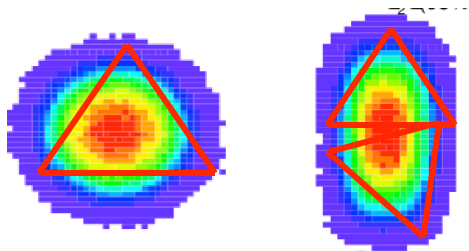
# Anti-correlation between $v_3$ and $v_2$



Can be used to fine tune initial geometry models!

- Quantified by a linear fit and extract the intercept and slope

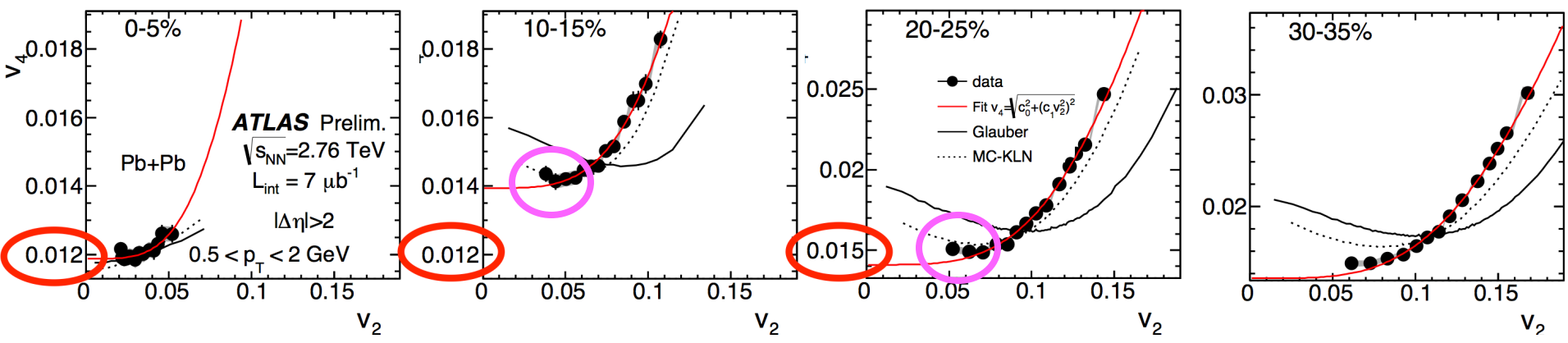
$$v_3 = kv_2 + v_3^0$$



Events with zero  $\varepsilon_2$  has larger average  $\varepsilon_3 \rightarrow$  larger  $v_3$ .

# linear ( $\epsilon_4$ ) and non-linear ( $v_2^2$ ) component of $v_4$

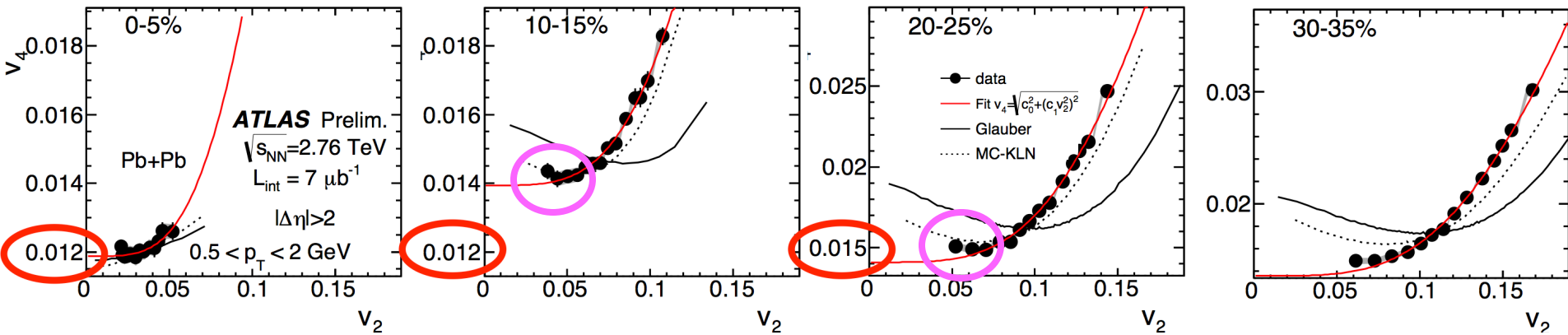
- $v_4$ - $v_2$  correlation for fixed centrality bin  $v_4 e^{i4\Phi_4} = c_0 e^{i\Phi_4} + c_1 (v_2 e^{i2\Phi_2})^2 \Rightarrow$  Fit by  $v_4 = \sqrt{c_0^2 + c_1^2 v_2^4}$



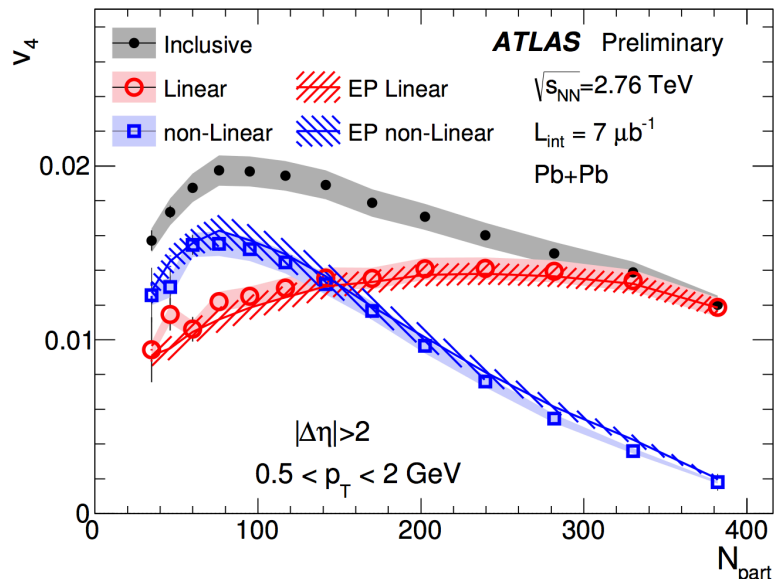
- Fit  $v_4 = \sqrt{c_0^2 + c_1^2 v_2^4}$  to separate linear ( $\epsilon_4$ ) and non-linear ( $v_2^2$ ) component

# linear ( $\varepsilon_4$ ) and non-linear ( $v_2^2$ ) component of $v_4$

- $v_4$ - $v_2$  correlation for fixed centrality bin  $v_4 e^{i4\Phi_4} = c_0 e^{i\Phi_4} + c_1 (v_2 e^{i2\Phi_2})^2 \Rightarrow$  Fit by  $v_4 = \sqrt{c_0^2 + c_1^2 v_2^4}$



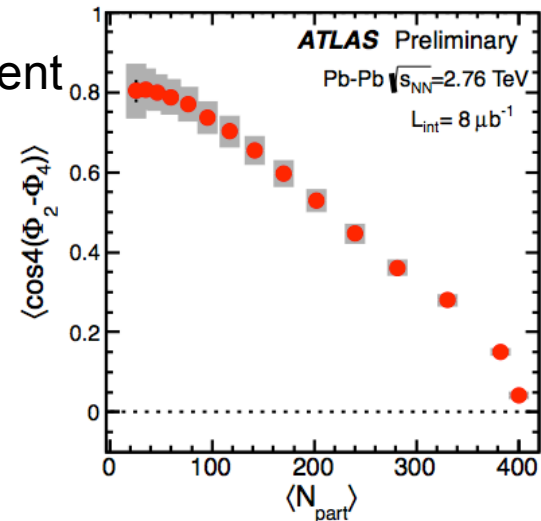
- Fit  $v_4 = \sqrt{c_0^2 + c_1^2 v_2^4}$  to separate linear ( $\varepsilon_4$ ) and non-linear ( $v_2^2$ ) component



predict L and NL component from EP correlations:

$$v_4^{NL} = v_4 \langle \cos 4(\Phi_2 - \Phi_4) \rangle$$

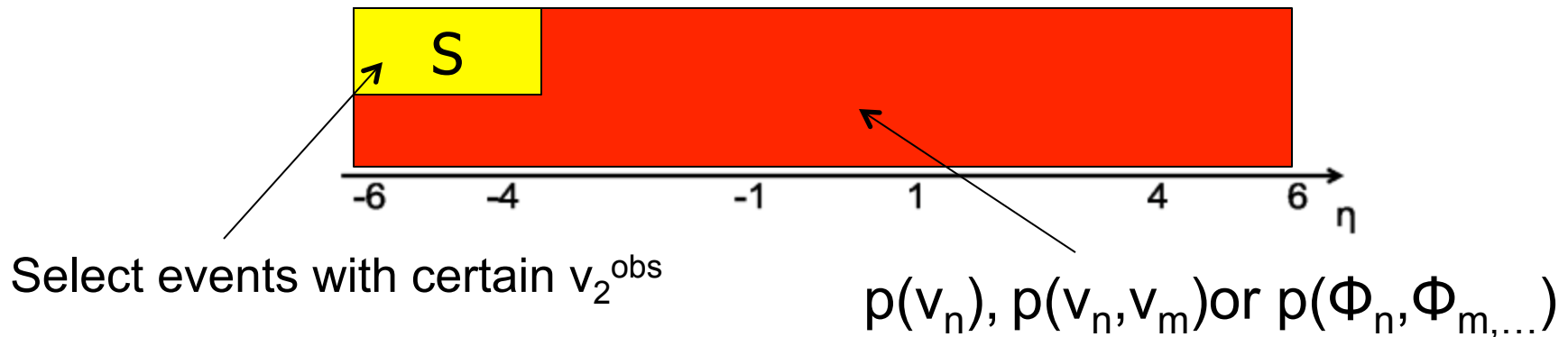
$$v_4^L = \sqrt{v_4^2 - (v_4^{NL})^2}$$



# What about select on one side?

Schukraft, Timmins, and Voloshin, arXiv:1208.4563

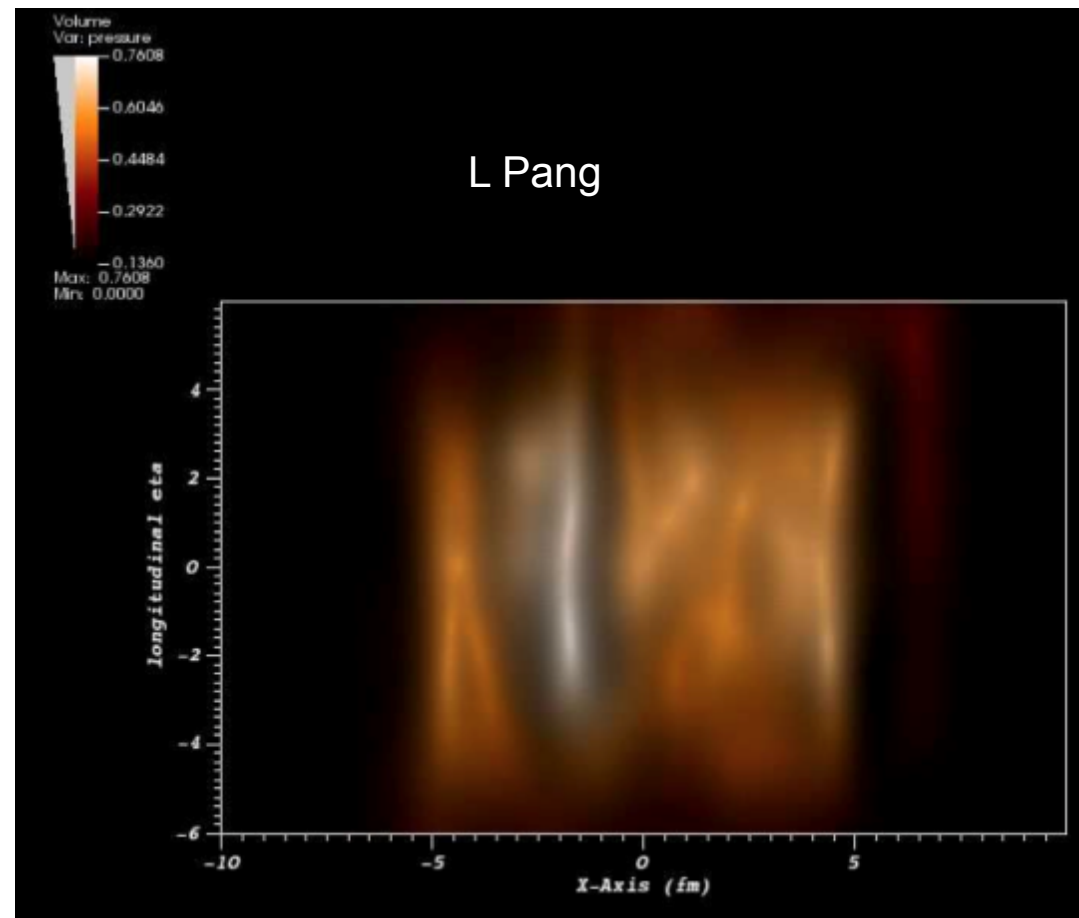
Huo, Mohapatra, JJ arxiv:1311.7091





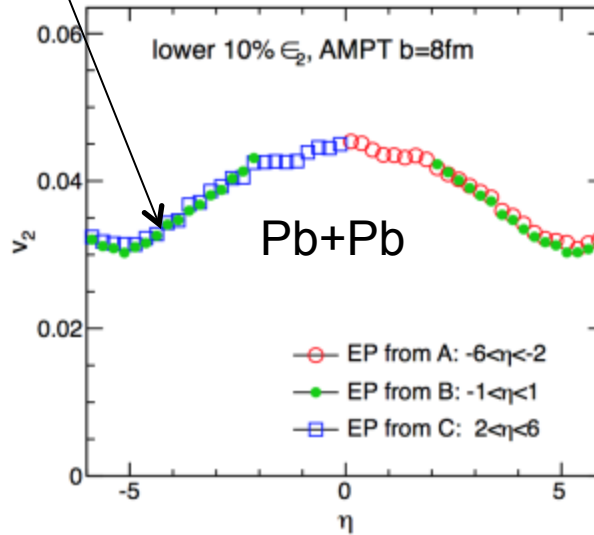
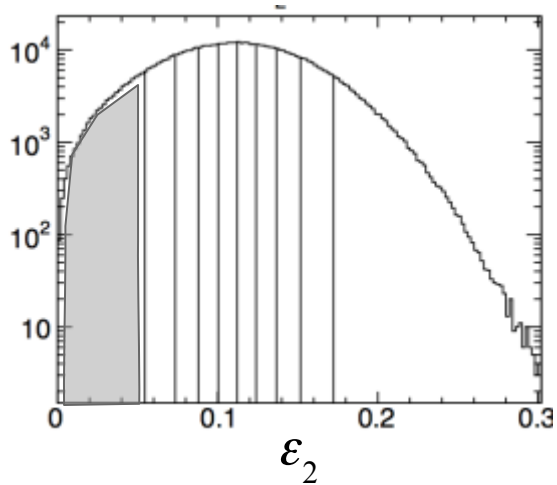
# AMPT model

- AMPT model: Glauber+HIJING+transport
  - Has **fluctuating geometry** and **collective flow**
  - **Longitudinal fluctuations** and **initial flow**



# $v_2(\eta)$ : select on $\epsilon_2$

Flow suppressed



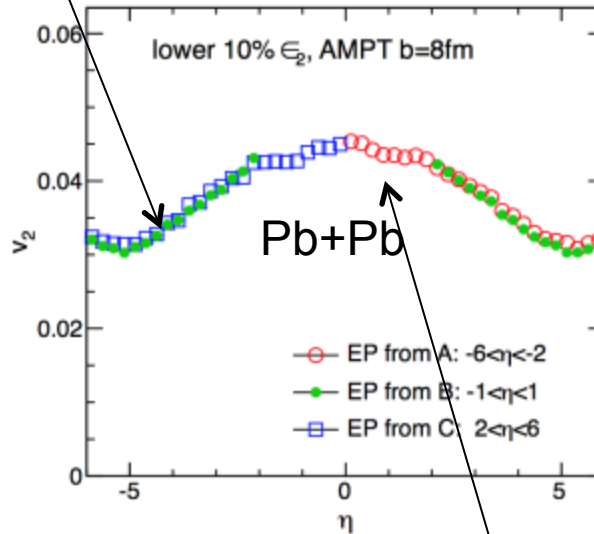
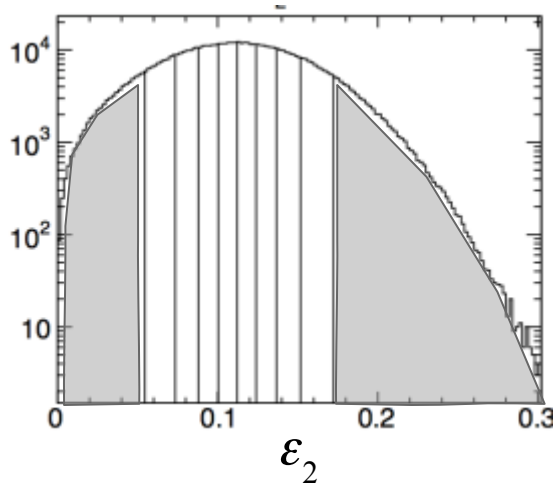
$v_2(\eta)|_{\eta>0}$  when EP in  $-6 < \eta < -2$

$v_2(\eta)|_{\eta<0}$  when EP in  $2 < \eta < 6$

$v_2(\eta)|_{|\eta|>2}$  when EP in  $|\eta| < 1$

# $v_2(\eta)$ : select on $\epsilon_2$

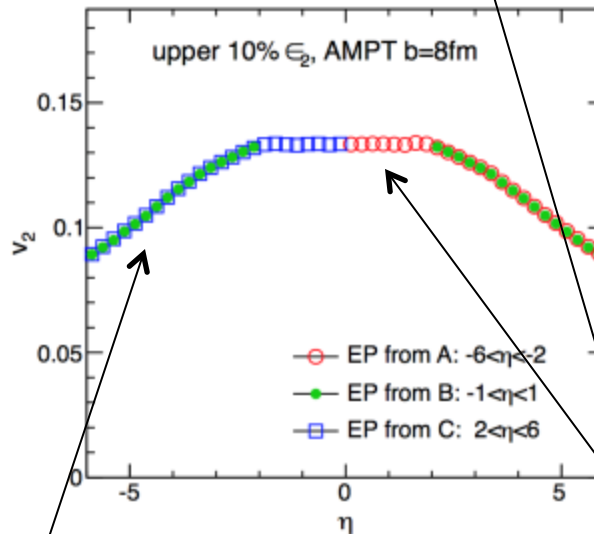
Flow suppressed



$v_2(\eta)|_{\eta > 0}$  when EP in  $-6 < \eta < -2$

$v_2(\eta)|_{\eta < 0}$  when EP in  $2 < \eta < 6$

$v_2(\eta)|_{|\eta| > 2}$  when EP in  $|\eta| < 1$

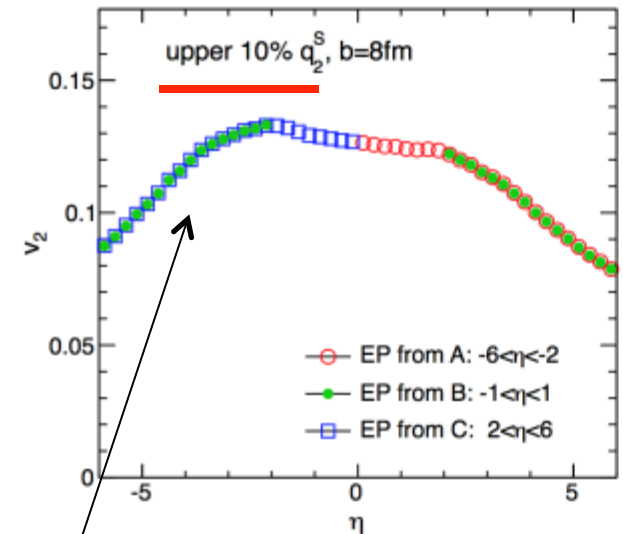
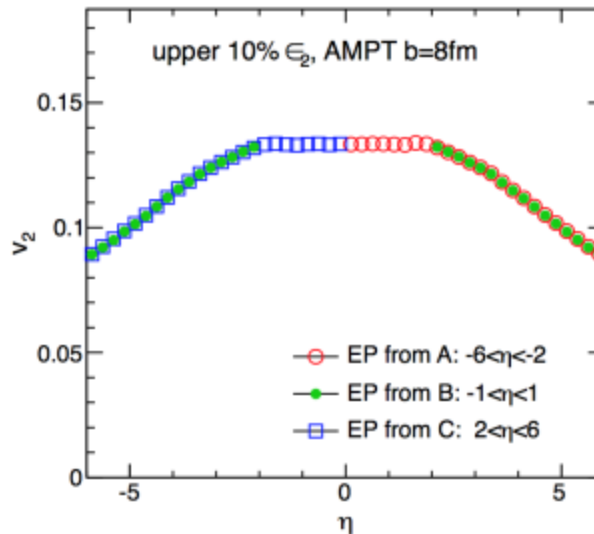
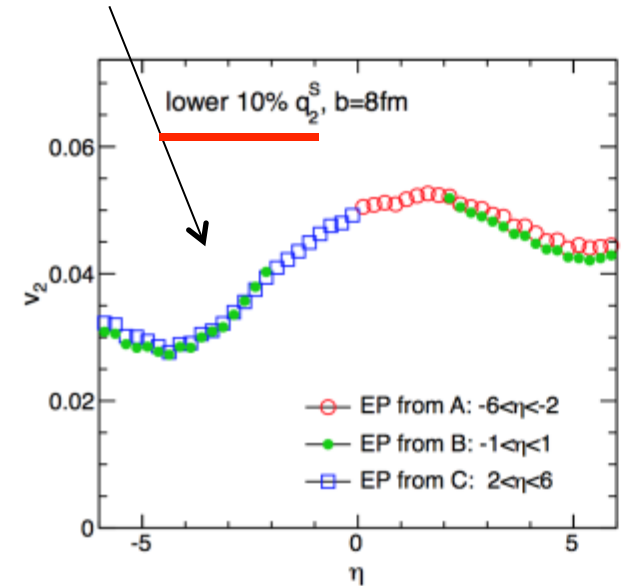
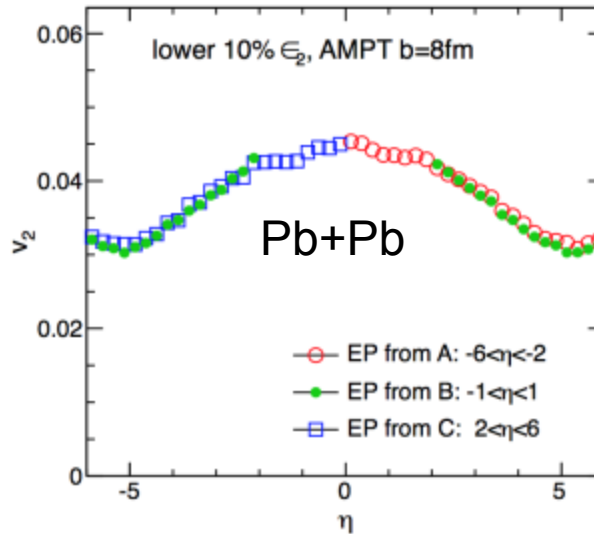
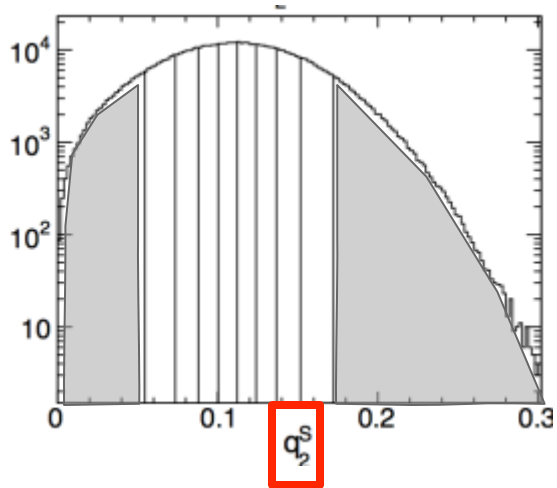


Flow enhanced

Symmetric distribution expected

# $v_2(\eta)$ : compare with selection on $q_2$

Suppression of flow in the selection window



$v_2(\eta)|_{\eta>0}$  when EP in  $-6 < \eta < -2$

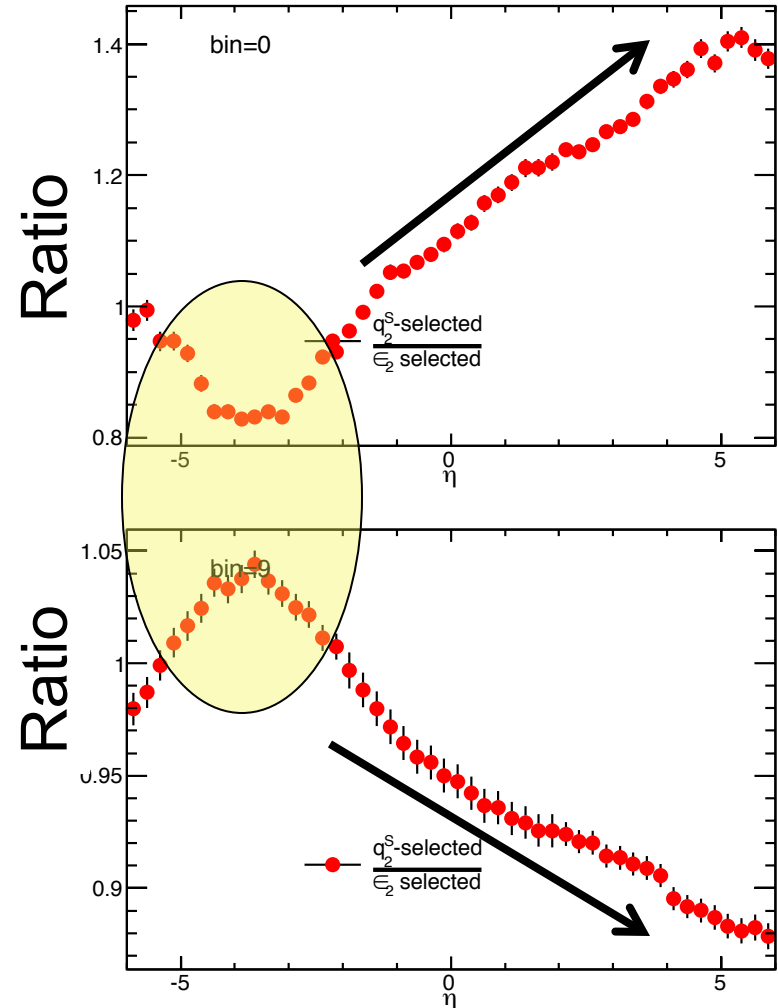
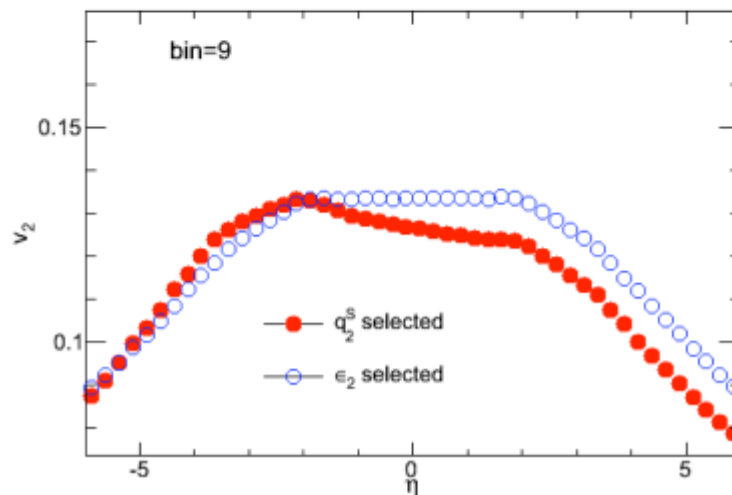
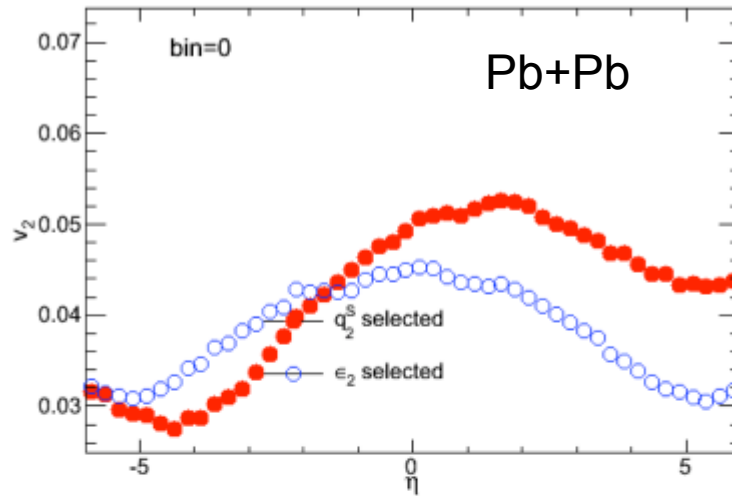
$v_2(\eta)|_{\eta<0}$  when EP in  $2 < \eta < 6$

$v_2(\eta)|_{|\eta|>2}$  when EP in  $|\eta| < 1$

enhancement of flow in the selection window

# What is the origin of $v_2(\eta)$ asymmetry?

- Suppression/enhancement of flow in the selected window
- Decreasing response to flow selection outside the selection window



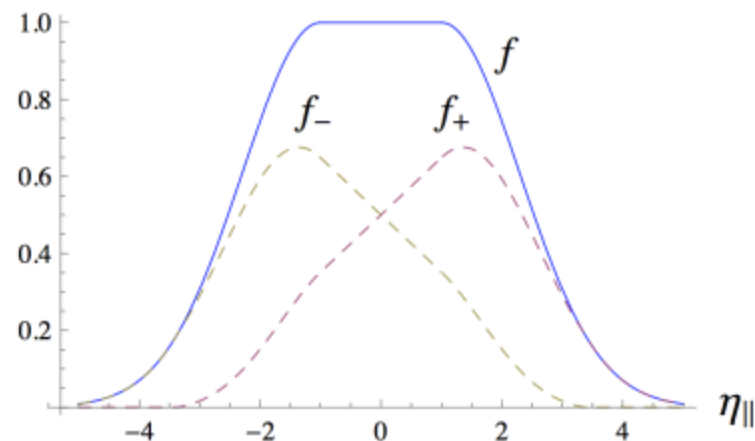
# Longitudinal particle production

wounded nucleon model

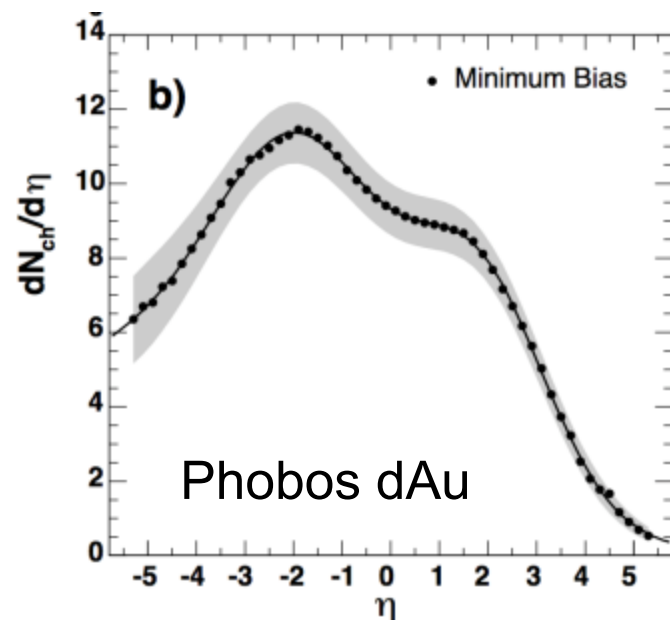
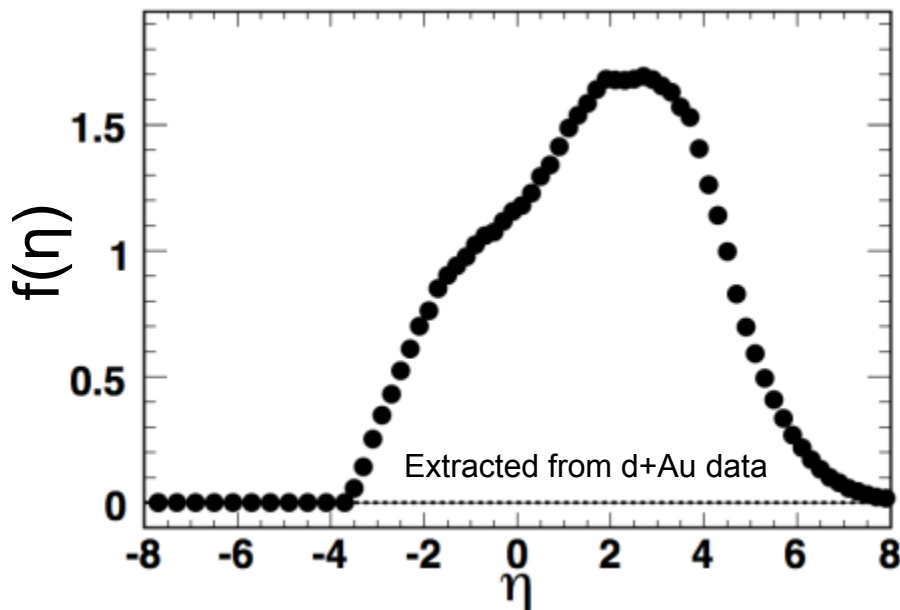
Bialas, Bzdak, Zalewski, Wozniak.... STAR/PHOBOS

- Assumes that after the collision of two nuclei, the secondary particles are produced by independent fragmentation of wounded nucleons

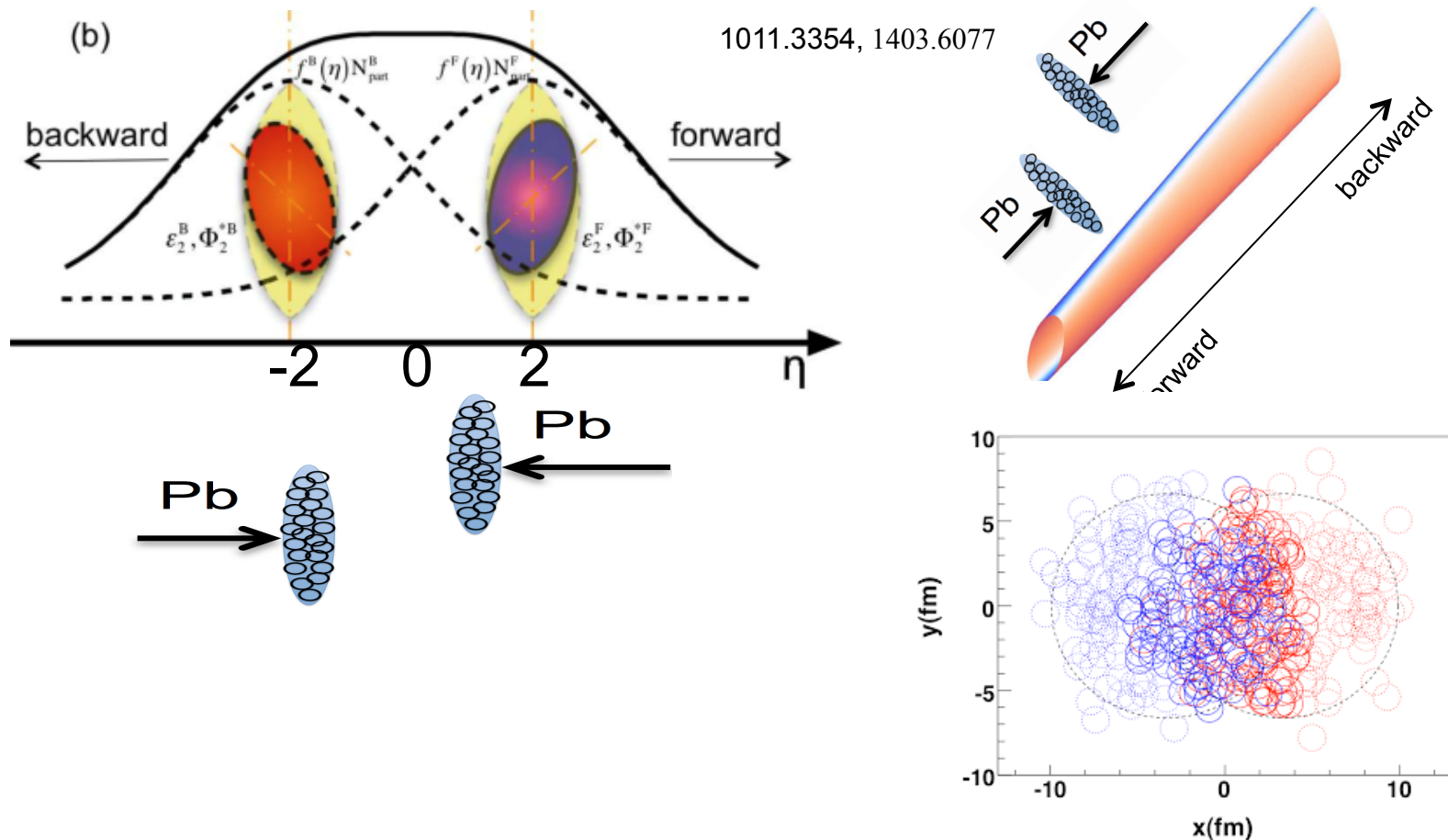
$$dN/d\eta \propto f^F(\eta)N_{\text{part}}^F + f^B(\eta)N_{\text{part}}^B$$



Emission function of one wounded nucleon



# Flow longitudinal dynamics

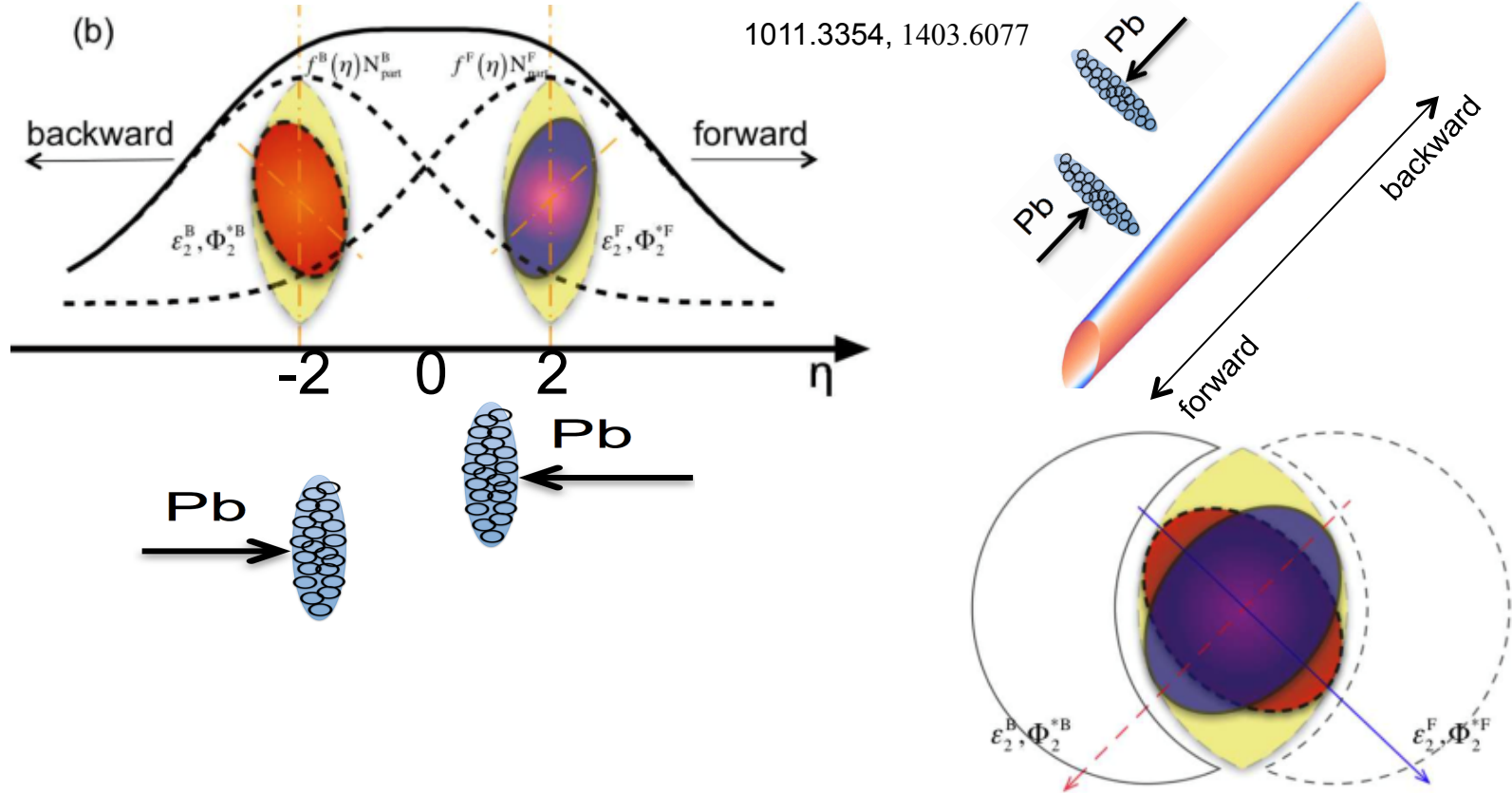


- Shape of participants in two nuclei not the same due to fluctuation

$$\varepsilon_m^F, \Phi_m^{*F} \quad \varepsilon_m^B, \Phi_m^{*B} \quad \varepsilon_m, \Phi_m^* \quad N_{part}^F, N_{part}^B, N_{part} \quad \varepsilon_n^F, \Phi_n^{*F} \neq \varepsilon_n^B, \Phi_n^{*B}$$

- Particles are produced by independent fragmentation of wounded nucleons, emission function  $f(\eta)$  not symmetric in  $\eta \rightarrow$  Wounded nucleon model

# Flow longitudinal dynamics



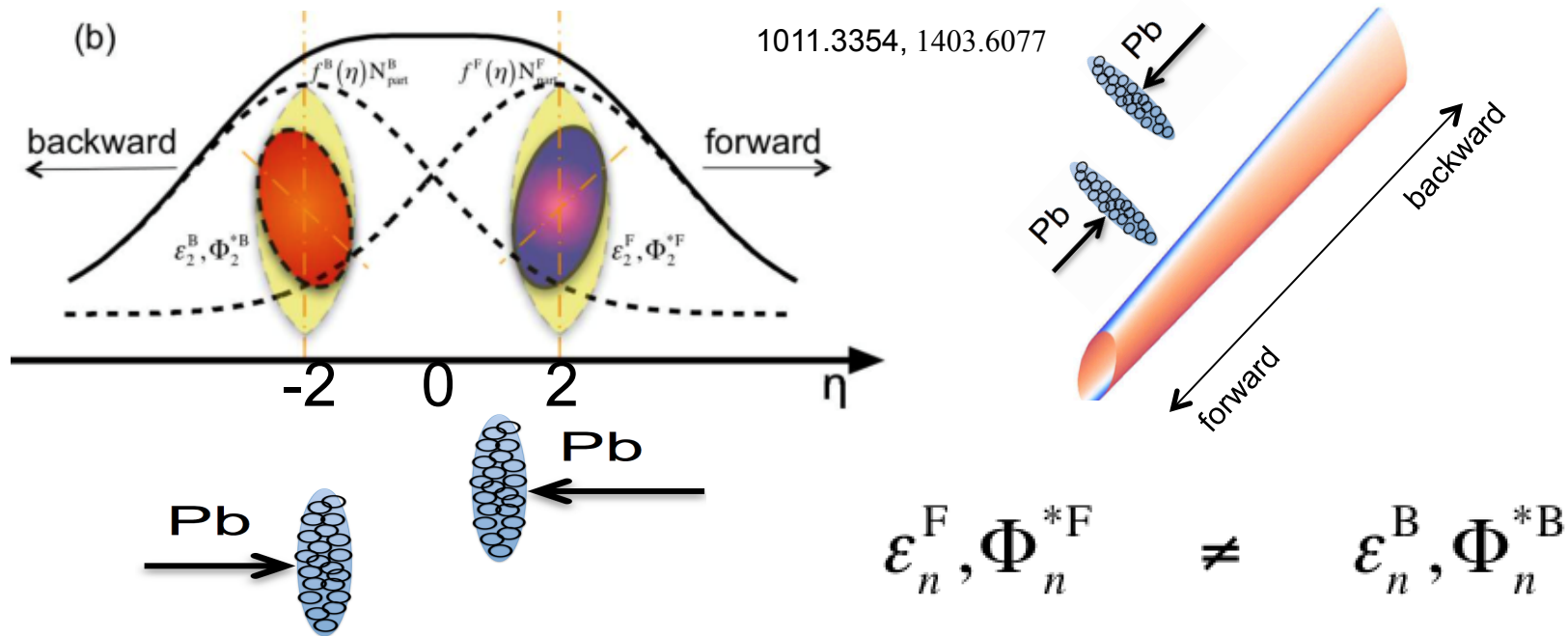
- Shape of participants in two nuclei not the same due to fluctuation

$$\varepsilon_m^F, \Phi_m^{*F} \quad \varepsilon_m^B, \Phi_m^{*B} \quad \varepsilon_m, \Phi_m^* \quad N_{part}^F, N_{part}^B, N_{part} \quad \varepsilon_n^F, \Phi_n^{*F} \neq \varepsilon_n^B, \Phi_n^{*B}$$

- Particles are produced by independent fragmentation of wounded nucleons, emission function  $f(\eta)$  not symmetric in  $\eta \rightarrow$  Wounded nucleon model



# Flow longitudinal dynamics



- Eccentricity vector interpolates between  $\vec{\epsilon}_n^F$  and  $\vec{\epsilon}_n^B$

$$\vec{\epsilon}_n^{\text{tot}}(\eta) \approx \alpha(\eta)\vec{\epsilon}_n^F + (1 - \alpha(\eta))\vec{\epsilon}_n^B \equiv \epsilon_n^{\text{tot}}(\eta)e^{in\Phi_n^{*\text{tot}}(\eta)}$$

$\alpha(\eta)$  determined by  $f(\eta)$

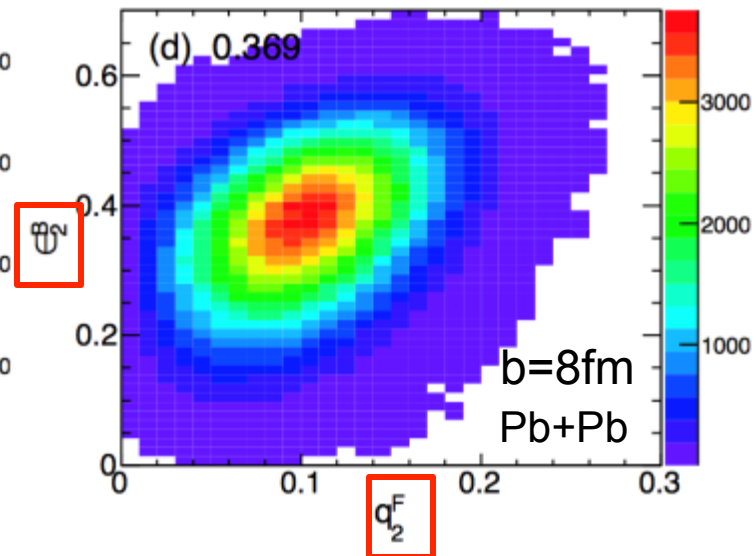
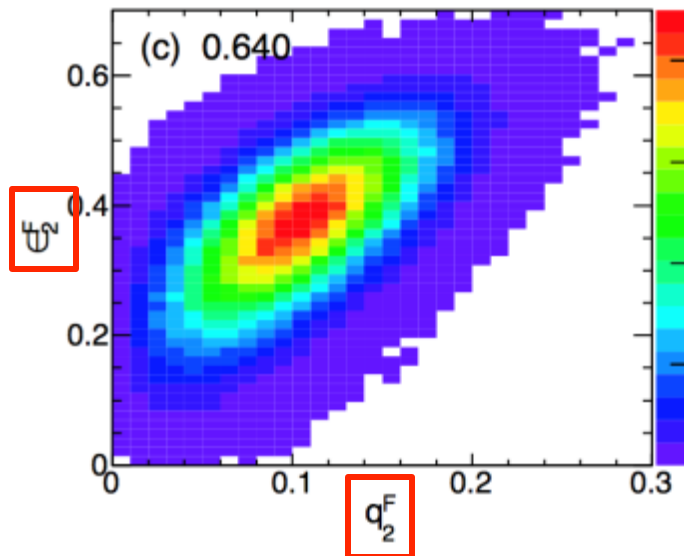
Asymmetry:	$\epsilon_n^F \neq \epsilon_n^B$
Twist:	$\Phi_n^{*F} \neq \Phi_n^{*B}$

- Hence  $\vec{v}_n(\eta) \approx c_n(\eta) [\alpha(\eta)\vec{\epsilon}_n^F + (1 - \alpha(\eta))\vec{\epsilon}_n^B]$  for  $n=2,3$

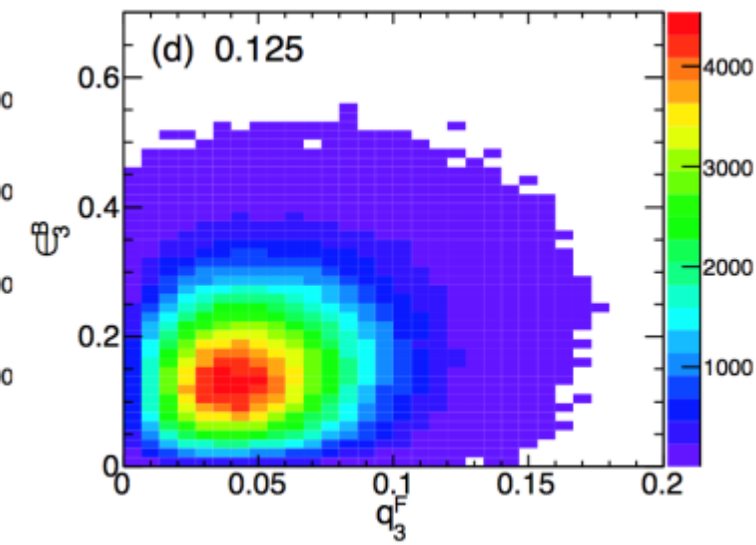
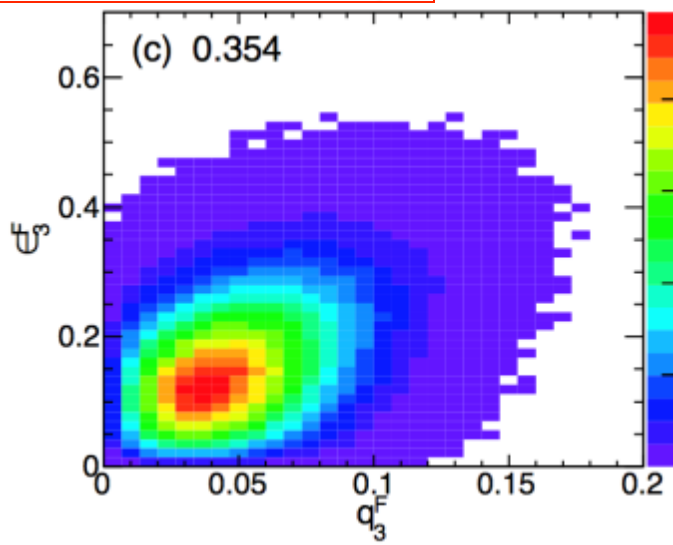
- Picture verified in AMPT simulations, magnitude estimated 1403.6077

# What AMPT tell us?

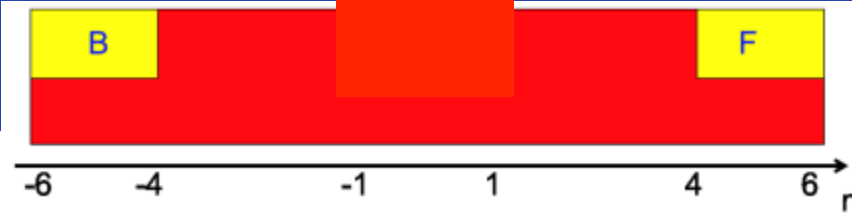
$\varepsilon_2^F$  more correlated with  $q_2^F$  than  $q_2^B$



$\varepsilon_3^F$  more correlated with  $q_3^F$  than  $q_3^B$



FB asymmetry survives



# What AMPT tell us?

- Twist in initial geometry appears as twist in the final state flow

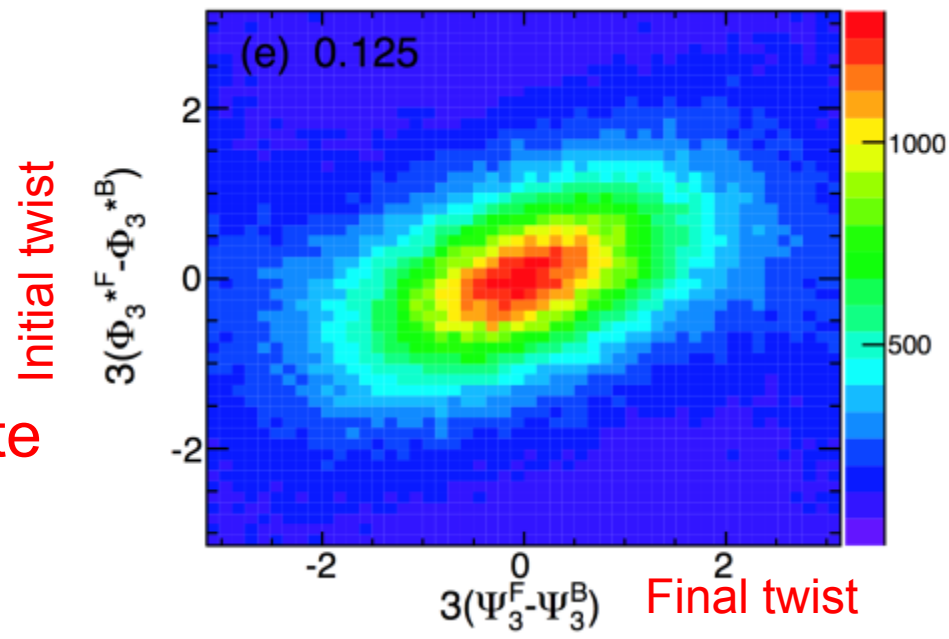
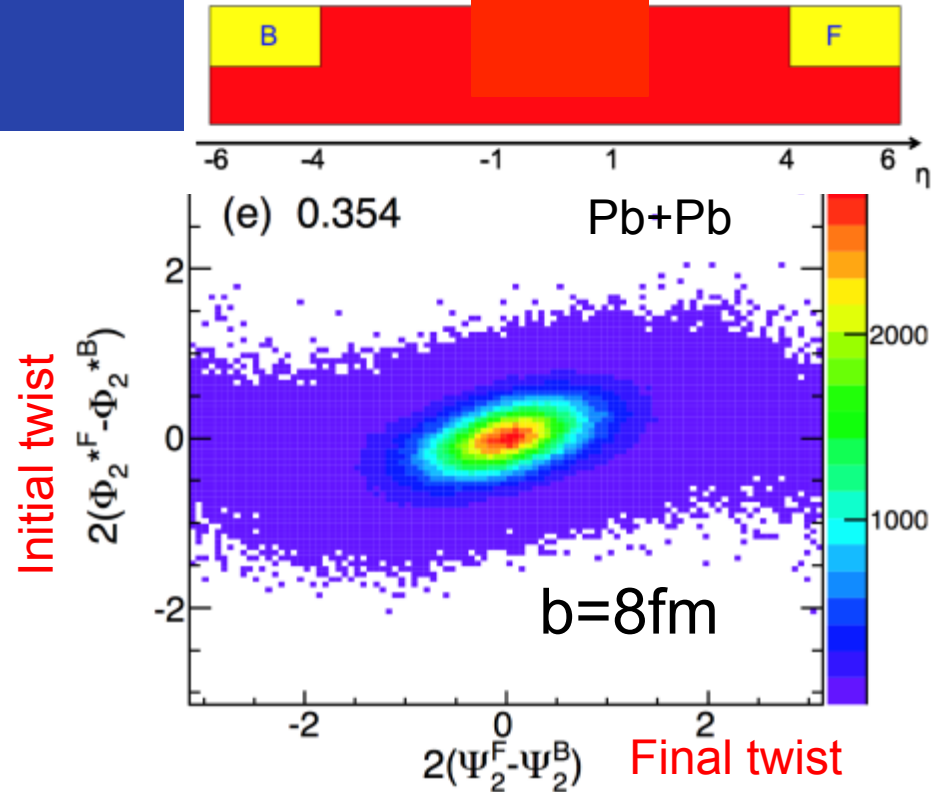
- Participant plane angles:

$$\Phi_n^{*F} \quad \Phi_n^{*B}$$

- Final state event-plane angles

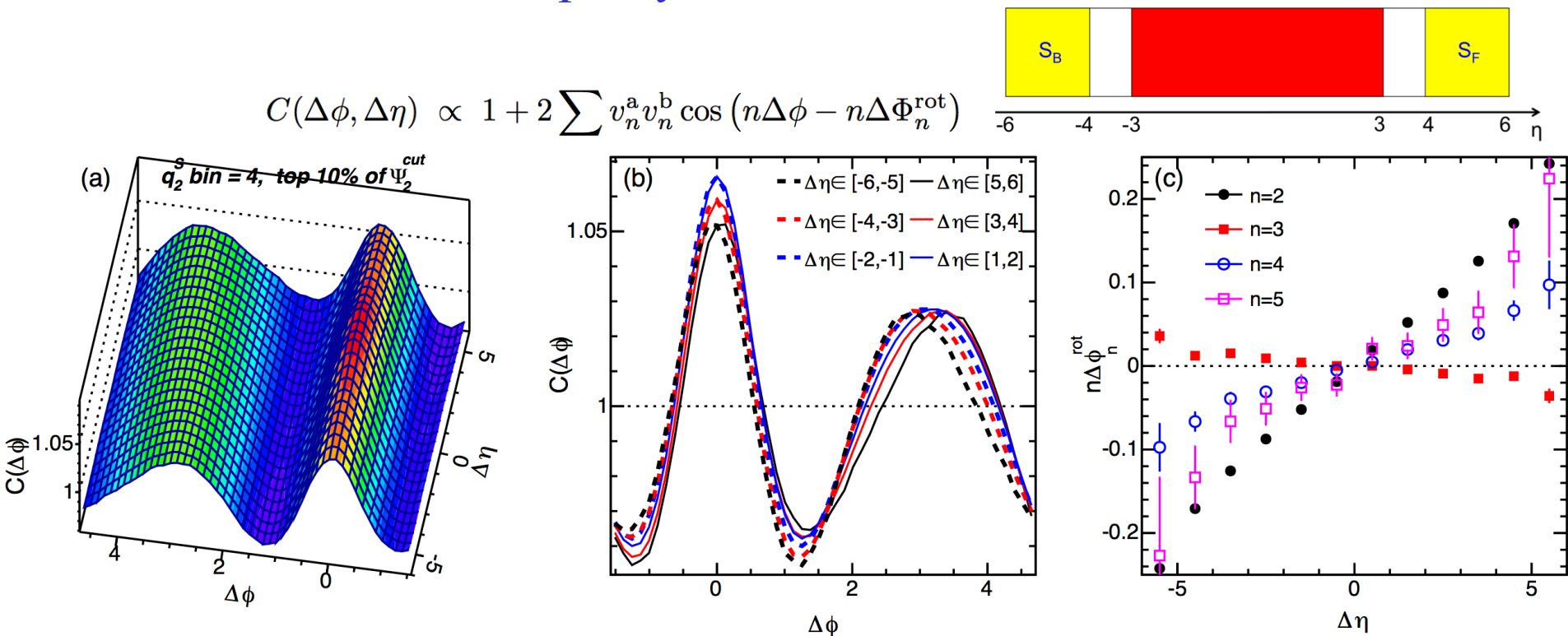
$$\Psi_n^F \quad \Psi_n^B$$

Initial twist survives to final state



# Twist seen in simple 2PC analysis

- NO event-plane determination! Just select twist in large  $\eta$  and check correlation at center-rapidity.



- Though twist is enforced on  $q_2$ , twist also seen for higher order  $v_n$
- Non-linear mixing to the higher order harmonics!! .

# Summary-I

- Event-shape fluctuations contains a lot of information

$$p(v_n, v_m, \dots, \Phi_n, \Phi_m, \dots) = \frac{1}{N_{\text{evts}}} \frac{dN_{\text{evts}}}{dv_n dv_m \dots d\Phi_n d\Phi_m \dots}$$

- Three complementary methods: **Strong fluctuation within fixed centrality!**

	pdf's	cumulants	event-shape method
Flow-amplitudes	$p(v_n)$	$v_n\{2k\}, k = 1, 2, \dots$	NA
	$p(v_n, v_m)$	$\langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle$	yes
	$p(v_n, v_m, v_l)$	$\langle v_n^2 v_m^2 v_l^2 \rangle + 2\langle v_n^2 \rangle \langle v_m^2 \rangle \langle v_l^2 \rangle -$ $\langle v_n^2 v_m^2 \rangle \langle v_l^2 \rangle - \langle v_m^2 v_l^2 \rangle \langle v_n^2 \rangle - \langle v_l^2 v_n^2 \rangle \langle v_m^2 \rangle$	yes
	...	Obtained recursively as above	yes
EP-correlation	$p(\Phi_n, \Phi_m, \dots)$	$\langle v_n^{c_n} v_m^{c_m} \dots \cos(c_n n \Phi_n + c_m m \Phi_m + \dots) \rangle$ $\sum_k k c_k = 0$	yes
Mixed-correlation	$p(v_l, \Phi_n, \Phi_m, \dots)$	$\langle v_l^2 v_n^{c_n} v_m^{c_m} \dots \cos(c_n n \Phi_n + c_m m \Phi_m + \dots) \rangle -$ $\langle v_l^2 \rangle \langle v_n^{c_n} v_m^{c_m} \dots \cos(c_n n \Phi_n + c_m m \Phi_m + \dots) \rangle$ $\sum_k k c_k = 0$	yes

# Summary-II

- Rich patterns forward/backward EbyE flow fluctuations:

$$\vec{v}_n(\eta) \approx c_n(\eta) [\alpha(\eta)\vec{\epsilon}_n^F + (1 - \alpha(\eta))\vec{\epsilon}_n^B]$$

Event-shape  
selection and event  
twist techniques

- New avenue to study initial state fluctuations, particle production and collective expansion dynamics.