



Beam Energy Dependence of Azimuthal Correlations in Au-Au Collisions at Mid and Forward Rapidity

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Outline

Introduction

- i. QCD Phase Diagram
- ii. STAR Detector
- iii. Correlation function technique

II. Results

- *i.* $v_n p_T$ dependence
- *ii.* $v_n \eta$ dependence
- *iii.* v_n beam energy dependence
- iv. Viscous coefficient

III. Conclusion

QCD Phase Diagram

The BES at RHIC allows the study of a broad domain of (μ_B, T) – plane.





F. Becattini, PoSCPOD07:012,2007

 $\geq \mu_B \& T$ variations via beam energy or rapidity selections.

QCD Phase Diagram

Strong interest in measurements which span a broad (μ_B, T) domain.

 Investigate signatures for the first-order phase transition



PRL 112,162301(2014)

- Investigate transport coefficients as a function of (μ_B, T)
- Possible non-monotonic patterns



PRL 116, 112302 (2016)

 Search for critical fluctuations



PRL 112, 032302 (2014)

STAR Detector at RHIC



TPC detector covers | η | < 1 FTPC detector covers 2.5 < | η | < 4

Correlation function technique

All current techniques used to study v_n are related to the correlation function. Two particle correlation function $C(\Delta \varphi = \varphi_1 - \varphi_2)$ used in this analysis,

$$C(\Delta \varphi) = \frac{dN/d\Delta \varphi(same)}{dN/d\Delta \varphi(mix)} \text{ and } v_n^2 = \frac{\sum_{\Delta \varphi} C(\Delta \varphi) \cos(n \Delta \varphi)}{\sum_{\Delta \varphi} C(\Delta \varphi)}$$

$$PLB \ 708, 249 \ (2012)$$

$$v_n(p_T) = \frac{v_n^2(p_{Tref}, p_T)}{\sqrt{v_n^2(p_{Tref})}} \text{ and } v_n(\eta) = \frac{v_n^2(\eta, \eta_{ref})}{\sqrt{v_n^2(\eta_{ref})}}$$

$$2$$

$$v_n(p_T) = \frac{n(r_Tref^{n-1})}{\sqrt{v_n^2(p_{Tref})}}$$
 and $v_n(\eta) = \frac{v_n(\eta)}{\sqrt{v_n^2(\eta)}}$

Factorization ansatz for v_n verified.

Non-flow signals, as well as some residual detector effects (track merging/splitting) minimized with $|\Delta \eta = \eta_1 - \eta_2| > 0.7$ cut.

Results

$$\succ v_n(pT) = \frac{v_n^2(p)}{\sqrt{v_n^2}}$$

$$\frac{v_{\bar{n}}(p_{T_{ref}},p_{T})}{\sqrt{v_{n}^{2}(p_{T_{ref}})}}$$

$$\checkmark |\eta| < 1$$

$$\succ v_n(\eta) = \frac{v_n^2(\eta, \eta_{ref})}{\sqrt{v_n^2(\eta_{ref})}}$$

$$\sqrt{|\eta_{ref}|} < 1 \text{ and } |\eta| < 4$$

$$\sqrt{0.2} < p_T < 4 \text{GeV/c}$$

$$\sqrt{v_n} (\sqrt{s_{NN}})$$

Data presented with only statistical uncertainties

Viscous coefficient



 $v_n(p_T)$ indicate a similar trend for different beam energies.
 $v_n(p_T)$ decreases with harmonic order n.



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ν₂(η) show similar trends for the respective beam energies.
 ν₂(η) increases with beam energy over the measured η range.



Mid and forward rapidity $v_n(\eta)$ decreases with harmonic order n.



Reasonable agreement between the STAR and PHOBOS measurements.

 $v_n(\sqrt{S_{NN}})$ $|\eta| < 1 \text{ and } |\Delta \eta| > 0.7$ $0.2 < p_T < 4 \text{GeV/c}$



- ≻ Mid rapidity v_n(√s_{NN}) shows a monotonic increase with beam energy.
 > w (√2) decreases with hormonic order p
- $\succ v_n(\sqrt{s_{NN}})$ decreases with harmonic order n.

 $v_2(\sqrt{S_{NN}})$ TPC and FTPC $0.2 < p_T < 4$ GeV/c



- → Mid and forward rapidity $v_2(\sqrt{s_{NN}})$ shows a monotonic increase with beam energy.
- > Forward rapidity $v_2(\sqrt{s_{NN}})$ shows a stronger dependence.

Viscous coefficient

→ Use $v_n(p_T, \text{cent})$ to extract the viscous coefficient as a function of $\sqrt{s_{NN}}$, based on the acoustic ansatz

> ✓ the viscous coefficient encodes the transport coefficient $\frac{\eta}{s}$

Viscous coefficient

The v_n measurements are sensitive to ε_n , transport coefficient η/s and the expanding parameter RT.

Acoustic ansatz

✓ Sound attenuation in the viscous matter reduces the magnitude of v_n .

Anisotropic flow attenuation,

$$\frac{v_n}{\varepsilon_n} \propto e^{-\beta n^2}$$

arXiv:1305.3341

For two different harmonics n and n'(n'=2), $\frac{(v_n)^{\frac{1}{n}}}{(v_{n'})^{\frac{1}{n'}}} \propto c \ e^{-\beta (n-n')} \quad \text{and} \quad \beta \propto \frac{\eta}{s} \frac{1}{RT}, \ c = \frac{(\epsilon_n)^{\frac{1}{n}}}{(\epsilon_{n'})^{\frac{1}{n'}}}$

From macroscopic entropy considerations $(RT)^3 \propto \frac{dN}{dn}$ arXiv:1601.06001

• The viscous coefficient ξ encodes the transport coefficient $\frac{\eta}{s}$,

$$\xi = \left(\frac{dN}{d\eta}\right)^{1/3} \ln\left(\frac{(v_n)^{\frac{1}{n}}}{(v_2)^{\frac{1}{2}}}\right) \propto -(n-2)\frac{\eta}{s} + \left(\frac{dN}{d\eta}\right)^{1/3} \ln(c)$$

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Viscous coefficient $|\eta| < 1$ and $|\Delta \eta| > 0.7$ $0.2 < p_T < 4$ GeV/c



The viscous coefficient ξ shows a non-monotonic behavior with beam energy in both cases, n = 3 and n = 4.

III. Conclusion

Comprehensive set of STAR measurements for $v_n(p_T, \eta, \text{cent}, \sqrt{s_{NN}})$ presented.

> Mid and forward rapidity v_2 shows a monotonic increase with beam energy,

✓ Stronger $\sqrt{s_{NN}}$ dependence for forward v_2 .

- For a given √s_{NN} v_n decrease with the harmonic order.
 ✓ Similar patterns but different magnitude for different √s_{NN}
- The viscous coefficient ξ, which encodes the transport coefficient ^η/_s, indicates a non-monotonic pattern for the beam energy range studied.

